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**exoALMA. XVII. Characterizing the Gas Dynamics around Dust Asymmetries**

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Abstract

The key planet-formation processes in protoplanetary disks remain an active matter of research. One promising mechanism to radially and azimuthally trap millimeter-emitting dust grains, enabling them to concentrate and grow into planetesimals, is anticyclonic vortices. While dust observations have revealed crescent structures in several disks, observations of their kinematic signatures are still lacking. Studying the gas dynamics is, however, essential to confirm the presence of a vortex and understand its dust trapping properties. In this work, we make use of the high-resolution and sensitivity observations conducted by the exoALMA large program to search for such signatures in the ¹²CO and ¹³CO molecular line emission of four disks with azimuthal dust asymmetries: HD 135344B, HD 143006, HD 34282, and MWC 758. To assess the vortex features, we constructed an analytical vortex model and performed hydrodynamical simulations. For the latter, we assumed two scenarios: a vortex triggered at the edge of a dead zone and of a gap created by a massive embedded planet. These models reveal a complex kinematical morphology of the vortex. When compared to the data, we find that none of the sources show a distinctive vortex signature around the dust crescents in the kinematics. HD 135344B exhibits a prominent feature similar to the predictions from the simulations, thus making this the most promising target for sensitive follow-up studies at higher resolution and in particular with less abundant molecules at higher resolution and sensitivity to trace closer to the disk midplane.

Unified Astronomy Thesaurus concepts: [Protoplanetary disks \(1300\)](#); [Planet formation \(1241\)](#); [Planetary-disk interactions \(2204\)](#); [Submillimeter astronomy \(1647\)](#)

1. Introduction

In recent years, high-resolution and sensitivity observations of the dust and gas material in protoplanetary disks with facilities such as the Atacama Large Millimeter/submillimeter Array (ALMA; ALMA Partnership et al. 2015), the Spectro-Polarimetric High-contrast Exoplanet Research (SPHERE; J. L. Beuzit et al. 2019), or the Gemini Planet Imager (B. Macintosh et al. 2014)

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have revealed that various substructures are ubiquitous in disks and may indicate planet–disk interactions. In the dust, this includes spiral arms, gaps or cavities, rings, and azimuthal asymmetries (e.g., review in J. Bae et al. 2023). The latter two are associated with local pressure maxima, where dust grains can be radially and/or azimuthally trapped and prevented from undergoing processes such as inward drift or fragmentation (e.g., T. Birnstiel et al. 2010; P. Pinilla et al. 2012a, 2012b). Dust traps can further enable the streaming instability (e.g., S. M. Stammer et al. 2019), thus allowing dust to grow to larger sizes and form planetesimals (e.g., X.-N. Bai & J. M. Stone 2010a, 2010b; N. Raettig et al. 2015).

One possible origin of pressure bumps in disks is embedded planets, which—if massive enough—are capable of opening a gap in the disk and form symmetric dust rings (e.g., B. A. Ayliffe et al. 2012; G. Dipierro et al. 2015; G. P. Rosotti et al. 2016; R. Dong et al. 2017a; G. Dipierro et al. 2018; R. Dong et al. 2018a). In low-viscosity regions, the edges of such gaps or cavities can then become unstable against the Rossby wave instability (RWI; R. V. E. Lovelace et al. 1999; H. Li et al. 2000), resulting in the formation of an anticyclonic vortex where dust can be trapped radially and azimuthally, possibly giving rise to the observed crescent-shaped asymmetries (e.g., T. Birnstiel et al. 2013; W. Fu et al. 2014; J. Bae et al. 2016). Evidence of azimuthal dust trapping has been found in a number of sources with ALMA, seen as a large horseshoe structure (e.g., HD 142527, S. Casassus et al. 2013; IRS 48, N. van der Marel et al. 2013; AB Aur, Y.-W. Tang et al. 2012, 2017) or as a concentration of dust within a dust ring (e.g., HD 135344B, SR 21 L. M. Pérez et al. 2014; V1247 Ori S. Kraus et al. 2017; HD 34282 G. van der Plas et al. 2017; MWC 758 Y. Boehler et al. 2018; HD 143006, L. M. Pérez et al. 2018), with high-resolution observations revealing more complex substructures in some of these sources (e.g., N. van der Marel et al. 2016a; P. Cazzoletti et al. 2018; G. A. Muro-Arena et al. 2020; Y. Yang et al. 2023).

Alternatively to dust trapping, azimuthal dust asymmetries may be caused in a circumbinary disk by a sufficiently massive companion carving an eccentric cavity, which results in a gas overdensity at the cavity edge and a horseshoe structure (e.g., E. Ragusa et al. 2017; D. J. Price et al. 2018). N. van der Marel et al. (2021) compare several disk observables in a number of asymmetric disks to distinguish between vortices, horseshoes created by binary companions, and spiral density waves, concluding that the current data do not allow for a clear classification. Apparent dust asymmetries may further arise in inclined sources due to optically thick emission at the hot inner rim of a disk cavity (O. M. Guerra-Alvarado et al. 2024; Á. Ribas et al. 2024).

The edge of a disk cavity or gap favors vortex formation regardless of the underlying process creating the structure, not limited to planets. The RWI can further be triggered by infalling material (J. Bae et al. 2015; A. Kuznetsova et al. 2022) or at the edge of a dead zone (e.g., M. Flock et al. 2015, 2017). Moreover, there are other instabilities that can lead to the formation of vortices: the vertical shear instability (S. Richard et al. 2016; M. Flock et al. 2020; N. Manger et al. 2020; T. Pfeil & H. Klahr 2021), the convective overstability (W. Lyra 2014; N. Raettig et al. 2021), or the zombie vortex instability (e.g., P. S. Marcus et al. 2013, 2015, 2016). To differentiate the mechanisms and understand the process of

dust trapping in vortices, it is important to study their gas dynamics.

The gas in protoplanetary disks is studied through a range of molecular lines, tracing different disk regions and probing various physical and chemical processes. While it is easier to observe substructures in the dust component of the disk due to the higher achieved sensitivity, an increasing number of substructures have also been observed in the gas (e.g., S. Bruderer et al. 2014; K. Zhang et al. 2014; N. van der Marel et al. 2016b; R. Dong et al. 2017b; C. J. Law et al. 2021; L. Wölfer et al. 2023; R. Teague et al. 2025). Similar to the dust, this includes spirals, rings, and gaps or cavities. Narrowband line emission brings the advantage of providing information about the gas dynamics. Studying the velocity field of the rotating gas and identifying deviations from purely Keplerian rotation yields information about the underlying processes that are shaping the rotation pattern (for a review, see C. Pinte et al. 2023). The different disk processes leave individual fingerprints in the kinematics, but it remains difficult to disentangle them. Yet, observations of localized (e.g. C. Pinte et al. 2018, 2019, 2020; R. Teague et al. 2018, 2019a; S. Casassus & S. Pérez 2019; G. P. Rosotti et al. 2021; H. Yu et al. 2021; F. Alarcón et al. 2022; S. Casassus et al. 2022; A. F. Izquierdo et al. 2022) and more extended (e.g., non-Keplerian spirals; R. Teague et al. 2019b, 2021, 2022; G. P. Rosotti et al. 2020; S. Casassus et al. 2021; H. Garg et al. 2021; L. Wölfer et al. 2021; J. Stadler et al. 2023) kinematical deviations from recent years are generally consistent with the presence of massive embedded planets.

Identifying vortex signatures in the gas dynamics is crucial to not only confirm their presence but also to study their dust-trapping properties and understand the formation of the dust concentrations in disks. As shown by C. M. T. Robert et al. (2020), vortices produce signatures potentially discernible in the line-of-sight velocity through observations with current facilities. This includes the anticyclonic motions around the density maximum of a vortex as well as vortex-driven spirals, with the latter being comparable to those excited by a planet with a mass of a few to a few tens of Earth masses (P. Huang et al. 2019). H.-W. Yen & P.-G. Gu (2020) indeed report kinematical deviations in the HD 142527 disk, which are consistent with a local pressure bump and colocated with the large horseshoe structure seen in the dust disk, yet their observations do not allow them to distinguish if the pressure bump is formed by a vortex or massive planet. Y. Boehler et al. (2021) compare CO emission lines of the same disk with a vortex model and find kinematic signatures consistent with the presence of a large vortex; however, due to limitations in the spatial resolution of the data, artificial velocity deviations generated by beam smearing cannot be ruled out.

In this work, we make use of the unprecedented sensitivity at high resolution of the exoALMA data (R. Teague et al. 2025) to study the kinematical features seen around the azimuthal asymmetries of four disks in the sample and search for a kinematical vortex counterpart. In that regard, we compare the observations with both an analytical model of vortices and hydrodynamical simulations of a vortex triggered at the edge of a dead zone and gap created by a massive planet, respectively. The Letter is structured as follows: In Section 2 we describe the source sample and present the observations. Our models are outlined in Section 3 and 4 and compared to the data in Section 5. Our results are summarized in Section 6.

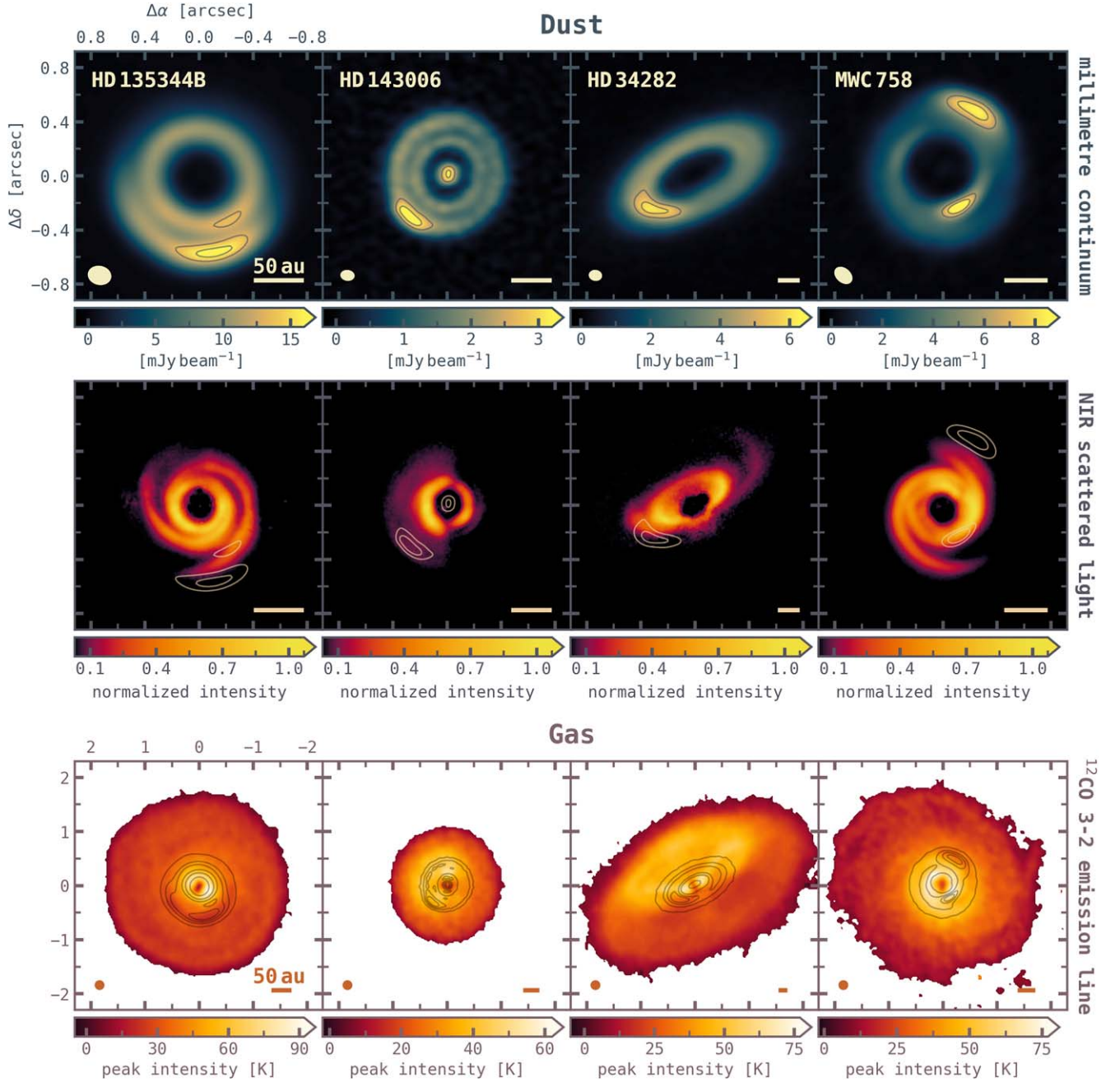


Figure 1. Overview of the four sources studied in this work, shown on the same spatial scale for the dust (top two rows) and gas (bottom row). Top: millimeter dust continuum as observed within the exoALMA program (P. Curone et al. 2025). Middle: micron-sized dust observed through NIR scattered light with SPHERE (T. Stolker et al. 2017; J. de Boer et al. 2021; B. B. Ren et al. 2023). The SPHERE images are normalized to the peak value and shown with a logarithmic normalization of the color map to emphasize the structures in the outer disk. To highlight the location of the crescents, two continuum contours around the peak of the dust asymmetries are overlaid on the dust images. Bottom: peak intensity of the $^{12}\text{CO } J=3-2$ emission line with overlaid continuum emission (five contour levels equally placed between 3σ and the peak flux). Note the larger spatial scale of the gas vs. dust images.

2. Observations

Our source sample consists of four sources—namely HD 135344B, HD 143006, HD 34282, and MWC 758—which were observed as part of the exoALMA large program (R. Teague et al. 2025). They are among the six most asymmetric disks in continuum (P. Curone et al. 2025) in the exoALMA sample, and all show evidence of azimuthal dust traps in their millimeter continuum (compare top row of Figure 1), possibly caused by the presence of a vortex. Three of the sources—HD 135344B, HD 34282, and MWC 758—are

further marked by spiral structures in near-infrared (NIR) scattered light (compare middle row of Figure 1; T. Stolker et al. 2017; J. de Boer et al. 2021; B. B. Ren et al. 2023). A summary of each source is provided in Appendix A. The inclination and position angle of the targets, obtained with DISCMINER (A. F. Izquierdo et al. 2025), are summarized in Table 1.

To study the gas features seen around the dust azimuthal asymmetries in our sample, we make use of the fiducial exoALMA images (R. Teague et al. 2025) of the ^{12}CO and $^{13}\text{CO } J=3-2$ emission lines ($0''.15$; 100 m s^{-1} , continuum

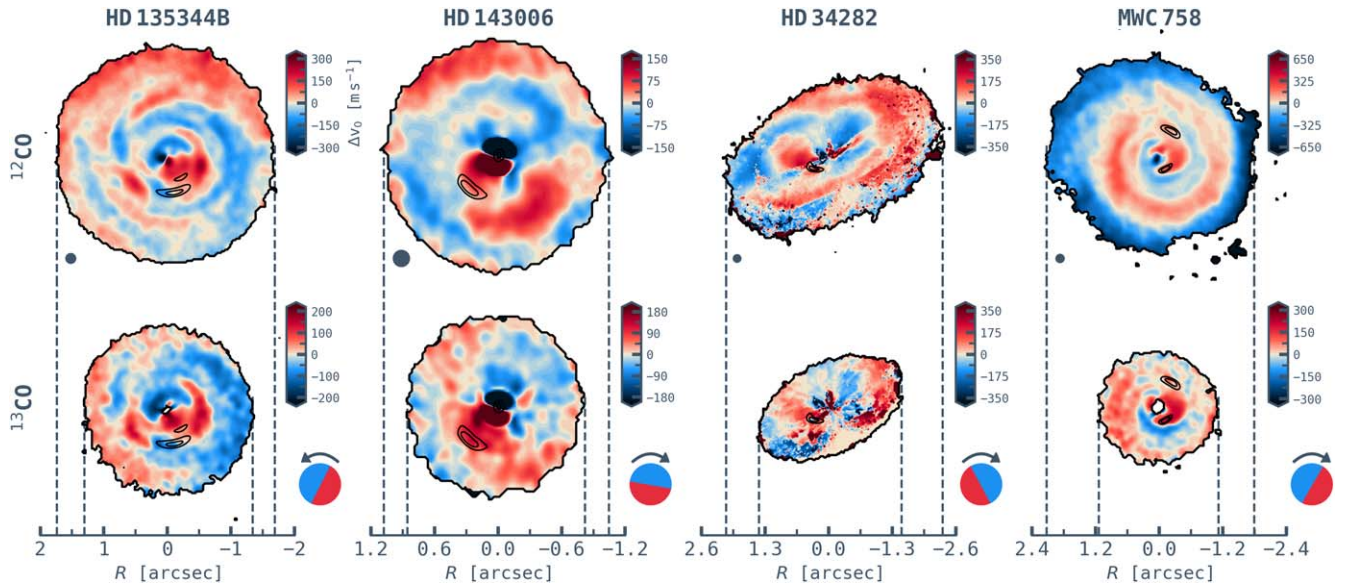


Figure 2. Velocity residuals Δv_0 after subtraction of the Keplerian model from the data, shown for ^{12}CO and ^{13}CO $J = 3-2$ and all four targets. For HD 34282, only emission from the top side of the disk was considered when making the v_0 map, following the procedure described in A. F. Izquierdo et al. (2025). The beam is shown in the bottom left corner of the top row. The red-blue circle and arrow in the bottom right corner of the second row indicate the underlying rotation pattern of each disk. To highlight the location of the crescents, two continuum contours around the dust asymmetries are overlaid on the images.

Table 1

Inclination and Position Angle of the Sources, Obtained by Modeling the ^{12}CO Emission Line with DISCMINER (A. F. Izquierdo et al. 2025)

Source	HD 135344B	HD 143006	HD 34282	MWC 768
i (deg)	16	17	58	19
PA	243	168	117	240

Note. The position angle is measured counterclockwise from the north to the redshifted axis of the disk.

subtracted). The CS $J = 7-6$ images are excluded from this work due to their much lower signal-to-noise ratio (SNR) at the same resolution. For the full details of the imaging techniques, we refer the reader to R. Loomis et al. (2025). For both lines, the molecular emission channels were modeled with the DISCMINER package (A. F. Izquierdo et al. 2021), which returns a smooth Keplerian model of the line intensity profile. To identify perturbations in the gas dynamics, it is then necessary to extract observables—such as the centroid velocity, line peak intensity, and line width—in a postprocessing step and to compute deviations between the data and model (residuals). Here, this is done by fitting both the data and model line profiles with a Gaussian function, except for the HD 34282 disk for which a double-peaked bell profile is adopted. While for low-inclination disks, a single-peak profile is sufficient to represent the morphology of the profile, higher-inclination disks such as HD 34282 are better reproduced with a double-peak profile since emission from the back side of the disk is visible. This approach then allows us to separate the contribution from the front and back sides of the disk. Details of the modeling procedures and choice of extraction methods can be found in A. F. Izquierdo et al. (2025).

In the bottom row of Figure 1, a gallery of the ^{12}CO peak intensity is shown for the four targets studied here. Emission below 5σ is masked. The continuum emission (top row) is

overlaid as contours, placed equally between 3σ and the continuum peak, highlighting the location of the potential vortices at relatively small radii compared to the large extents of the gas disk. The latter exceeds the dust disk by a factor of roughly 2–3, which is expected from processes such as radial drift or the difference between dust and gas opacities (S. Facchini et al. 2017; L. Trapman et al. 2019). Given that ^{12}CO emission is optically thick, the peak intensity traces the temperature of the gas, with the observed temperatures of up to ~ 100 K being as expected in the upper disk layers (e.g., S. Bruderer 2013; S. Bruderer et al. 2014; M. Leemker et al. 2022). A thorough analysis of the two-dimensional temperature structure of the disks can be found in M. Galloway-Sprietsma et al. (2025).

2.1. Velocity Residuals

In Figure 2, we present the velocity centroid residuals of our targets after subtraction of the Keplerian DISCMINER model from the ^{12}CO (top row) and ^{13}CO (bottom row) data (M. Fukagawa et al. 2025, in preparation), respectively. For the HD 34282 disk, the contribution from the back side has been subtracted in the computation of the residuals, as discussed in A. F. Izquierdo et al. (2025). The other three disks are close to face-on and do not require a removal of the lower surface.

The HD 135344B kinematics are marked by prominent asymmetric arc- and spiral-like structures, seen in both emission lines, as well as by features resembling a Doppler flip. Similarly, MWC 758 shows a large spiral arm, covering more than one full azimuth. For both disks, similar features are seen in the peak intensities and line widths (compare Figure 7). Altogether, this suggests strong perturbations in these disks. Interestingly, they are also the two disks in our sample that exhibit extended spirals in the NIR scattered light (see Appendix B). The patterns seen for the HD 143006 disk are less clear or consistent, but a prominent redshifted arc can be distinguished in the kinematics traced by both lines (even though at an offset), which may connect to the extended

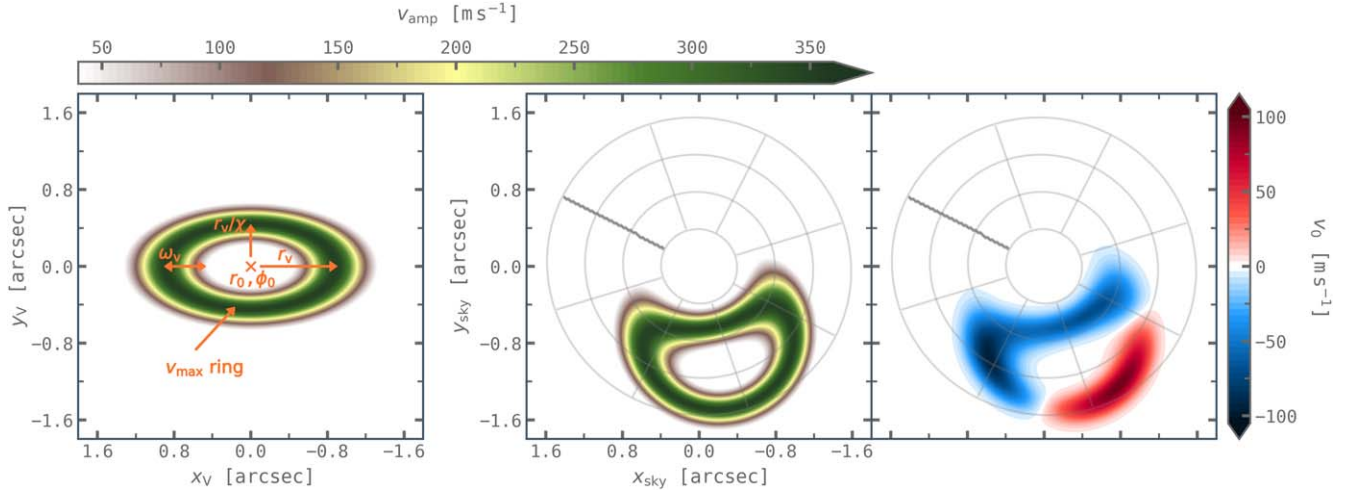


Figure 3. Analytical model of a vortex, described as elliptic streamlines of constant velocity with a Gaussian velocity profile. The first two panels show the absolute velocity of the vortex in both the vortex frame and sky frame. The last panel depicts the velocity projected along the line of sight, showing a complex morphology purely resulting from the projection. In the last two panels, the underlying disk geometry is overlaid as gray contours, corresponding to a position angle of 243° and inclination of 16° (same as for HD 135344B).

redshifted arc in the outer disk. The residuals of the HD 34282 disk are particularly difficult to interpret due to the high-inclination nature of the source. Even though the DISCMINER model and subsequent fit with a double-peaked bell profile account for both the upper and lower disk surface, contributions from the lower side seem to still be visible in the upper surface residuals, and the vertical structure off the disk is seen as a butterfly pattern in the inner disk regions.

3. Methods I: Analytical Vortex Model

In the following sections, we model the kinematical signature of a vortex to compare with the deviations seen in the residuals of our sample. For that purpose, we first construct a simple analytical model similar to Y. Boehler et al. (2021), which serves as a guideline for understanding the general kinematical patterns triggered by a vortex and helps to gain intuition for the more complex residuals. Subsequently, we perform two-dimensional hydrodynamical simulations (see Section 4). By doing so, we are mostly interested in comparing the kinematical patterns of a vortex to the data rather than a full representation of the vortex properties, which remains challenging (P. Huang et al. 2019; C. M. T. Robert et al. 2020).

3.1. Model Description

We follow Y. Boehler et al. (2021) in setting up the vortex flow as streamlines of constant velocity around a central position (r_0, ϕ_0) in polar disk coordinates, with the rotation velocity of the vortex varying as a Gaussian in the radial direction. In contrast to Y. Boehler et al. (2021), the streamlines are elliptical in the vortex frame in Cartesian coordinates, analogous to a shearing box. Our parameterization is depicted in the first panel of Figure 3. Here, the vortex is shown in the Cartesian vortex frame (subscript v) and is thus centered around the origin of this frame. The velocity in the vortex increases to a maximum velocity v_{\max} at the radius r_v . The vortex is further described by an aspect ratio χ , the ratio of the major axis in the azimuthal direction to the minor axis in the radial direction. The width of the Gaussian profile is given by ω_v .

The Cartesian vortex coordinates (x_v, y_v) can be related to the polar disk coordinates (r, ϕ) through

$$x_v = r(\phi - \phi_0) \quad \& \quad (1)$$

$$y_v = r - r_0, \quad (2)$$

where (r_0, ϕ_0) describes the center of the vortex. These can further be transformed into polar vortex coordinates (r', ϕ') via

$$r' = \sqrt{x_v^2 + (\chi \cdot y_v)^2} \quad \text{and} \quad (3)$$

$$\phi' = \tan^{-1} \left(\frac{x_v}{\chi \cdot y_v} \right). \quad (4)$$

The absolute rotation velocity of the vortex is then given as

$$|v(x_v, y_v)| = v_{\max} \cdot \exp \left[\left(\frac{r' - r_v}{\omega_v} \right)^2 \right]. \quad (5)$$

In Figure 3, the absolute orbital velocity is shown in the vortex frame (first panel) alongside the sky frame (second panel), using an inclination of 16° and position angle of 243° for the projection (same as for HD 135344B). The velocity projected along the line of sight is then given as

$$v_0 = |v(x, y)| \cdot \cos(\phi' + \phi) \cdot \sin i, \quad (6)$$

where i represents the disk inclination. The projected velocity of the vortex is presented in the last panel of Figure 3 in the sky frame, showing a clear red-blue pattern with a complex morphology, which purely results from projection effects.

3.2. Parameter Exploration

In Figure 4, we present the line-of-sight velocity for different disk geometries as well as parameters of the analytical vortex model. In this context, the parameters given on top of the figure correspond to the base model shown in the first panel, while the individual parameters changed compared to this model are given in each panel. This exercise highlights that while generally, a similar red-blue pattern is visible in all cases, the exact morphology and strength of the pattern vary substantially for different combinations of parameters. We further note that

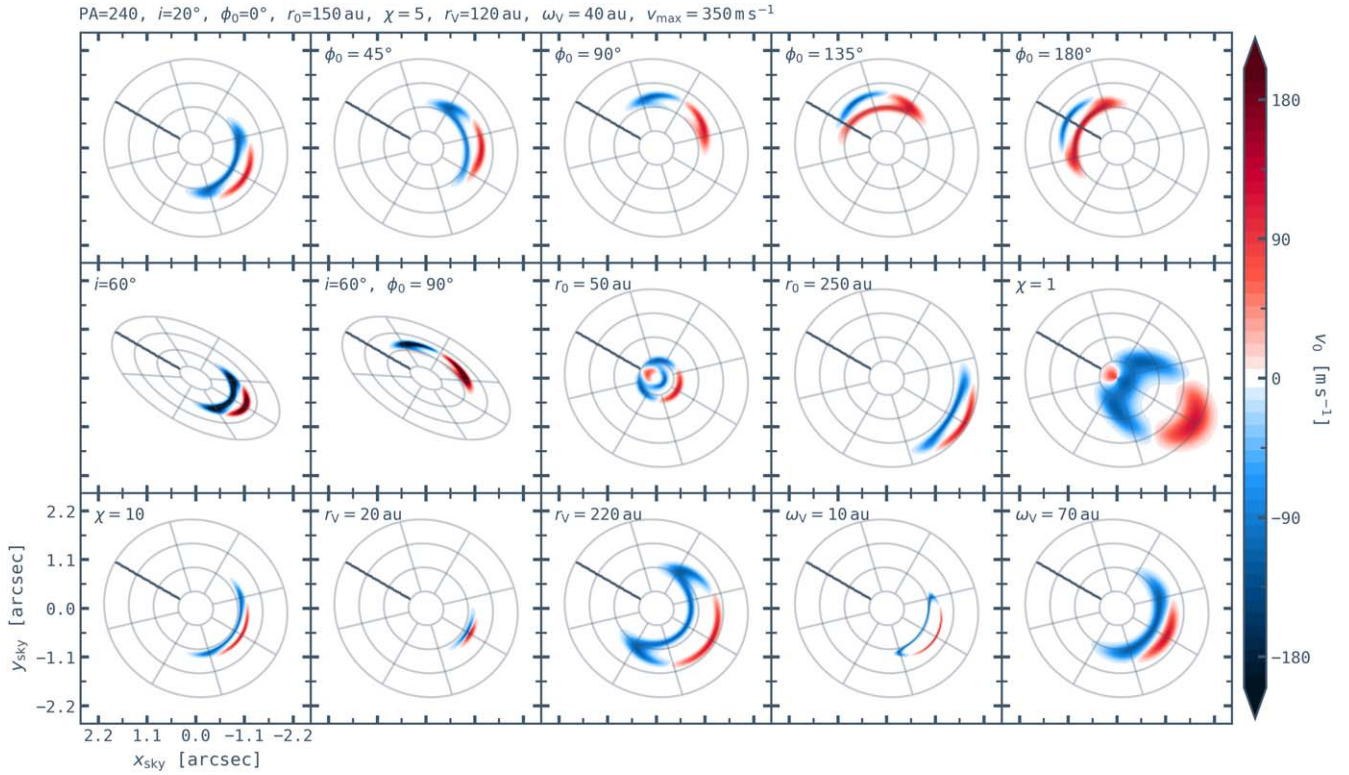


Figure 4. Line-of-sight velocity of the analytical vortex model, shown for different disk geometries and vortex parameters. The parameters given on top of the figure correspond to the morphology shown in the first panel. In the top of each individual panel, the parameters changed compared to this base vortex model are given. The underlying disk geometry is overlaid as gray contours.

the vortex edges have strong radial components, while along the “arcs,” the deviations are dominated by azimuthal motions.

We attempted to constrain the basic properties of the dust crescents seen in our targets by implementing the simple vortex model in the eddy code (R. Teague 2019), which allows us to fit a Keplerian profile to a rotation map. In this context, we first conducted a run to find a good representation of the main rotation pattern and in subsequent runs fixed those values to only fit for the vortex parameters described in Section 3.1, fitting different disk regions. However, none of these runs converged and resulted in very degenerate distributions. Given the complex kinematical patterns seen in our targets (see Figure 2), we conclude that a simple model is not sufficient to constrain the properties of the vortex. We are further limited by the low inclination of three of the four sources, making it hard to detect azimuthal motions projected into the line of sight. This is different for vertical motions. The importance of vertical motions is unclear at this point (e.g., H. Meheut et al. 2010, 2012; M.-K. Lin 2012; S. Richard et al. 2013), but they may dominate the vortex signature in some cases, which cannot be captured by our simple 2D model. Moreover, other disk processes such as planet–disk interactions can create dynamical patterns that overlay the vortex signatures and make it challenging to isolate them (see also M. Barraza-Alfaro et al. 2024). Nevertheless, the analytical model is useful to gain intuition for the general patterns created by a vortex and can help to interpret both the simulations and observations (see Section 4).

4. Methods II: Simulations

We explore the kinematic structure from a vortex (and vortex +planet) by conducting 2D numerical simulations with the grid-

based code PLUTO (A. Mignone et al. 2007), version 4.4. In this context, we do not aim for a comprehensive parameter study but intend to expand beyond the relatively simple analytical model to try to capture some of the complexity seen in the data. We follow the disk evolution, solving the Navier–Stokes equations of fluid dynamics (for details, see Section 2 in M. Barraza-Alfaro et al. 2024) for two scenarios: a vortex triggered by the RWI at the inner edge of a magnetorotational instability (MRI) dead zone, and RWI triggered at the edge of a gap opened by a Jupiter-mass planet. Even though the inner edge of the dead zone—in contrast to the outer edge—is unlikely to be relevant on exoALMA scales, we set up the simulation in this way to trigger a vortex at a particular location. The mechanism generating the vortex is the same for both the inner and outer edges but in the latter case results in a less clean vortex and uncertain location due to radial migration. Here we are mainly interested in the morphology of the vortex, and creating a vortex in the same location for both the dead zone and embedded planet simulation makes it easier to compare the two cases. In both cases, a vortex spontaneously forms within a few hundred orbits, computed at 1 code unit length in the dead-zone scenario and corresponding to planetary orbits in the planetary case.

We use a PLUTO numerical setup with LINEAR reconstruction method, second-order Runge–Kutta time stepping, the Linearized Roe Riemann solver (P. L. Roe 1981), and we include the disk viscosity as a diffusion term integrated with super-time-stepping technique. We set the Courant number to 0.4.

4.1. Disk Setup

We consider a two-dimensional disk in polar coordinates, including the evolution of the gas only. The disk initial

conditions are based on the numerical setup presented in R. Miranda et al. (2017). Our simulations follow a locally isothermal equation of state $P = c_s^2 \Sigma$, where P is the vertically integrated gas pressure, and the disk gas surface density Σ_g follows

$$\Sigma_g = \Sigma_0 \left(\frac{r}{r_0} \right)^{-1}, \quad (7)$$

with r_0 the code unit of length and Σ_0 the gas surface density at $r = r_0$. The disk sound speed c_s has a radial dependency $c_s = c_{s,0} (r/r_0)^{-0.25}$. The value of $c_{s,0}$ is set so that $H(r_0)/r_0 = 0.05$, with H the disk pressure scale height and $H = c_s/\Omega_K$, where $\Omega_K = \sqrt{GM_*/r^3}$ is the Keplerian orbital frequency. Our disk model results in a flared disk, with $H \propto r^{1.25}$.

The simulation domain covers 0.3–3.0 code units of length in the radial direction and 2π rad in azimuth. The simulations were run with a grid of 1024 and 2048 cells in r and ϕ , respectively, with a mesh uniformly spaced in both the radial and azimuthal directions. The grid uniform radial spacing allows a better resolved dead-zone edge transition and outer planetary gap edge relative to a logarithmic mesh. The effects of disk self-gravity are not taken into account in our models.

4.1.1. Dead-zone Edge

To mimic an MRI-dead-zone transition, we follow R. Miranda et al. (2017); however, for the inner edge of a dead zone instead of the outer one. We implement a transition between two levels of viscosity described by a radially dependent α -viscosity model, $\nu = \alpha c_s H$ (N. I. Shakura & R. A. Sunyaev 1973), with α following

$$\alpha(r) = \alpha_0 \left(1 - \frac{1}{2} \left(1 - \frac{\alpha_{\text{DZ}}}{\alpha_0} \right) \left[1 + \tanh \left(\frac{r - r_{\text{DZ}}}{\Delta_{\text{DZ}}} \right) \right] \right), \quad (8)$$

with $\alpha_{\text{DZ}} = 10^{-5}$ and $\alpha_0 = 10^{-3}$ being the α -viscosity parameters in the MRI-dead and MRI-active regions, respectively. The location of the transition between the active and dead zones is put as 1.2 code units of length ($r_{\text{DZ}} = 1.2r_0$), and Δ_{DZ} sets the width of the transition. We set $\Delta_{\text{DZ}} = H(r_{\text{DZ}})$, corresponding to the standard run from R. Miranda et al. (2017; see their Table 1).

4.2. Planet

In a second setup, we include a $1 M_J$ planet and the star as the gravitational potential of points of mass and the indirect acceleration term

$$\Phi = \frac{GM_*}{r} - \frac{Gm_p}{\sqrt{|r - r_p|^2 + s^2}} + \frac{Gm_p}{r_p^2} r \cos \phi, \quad (9)$$

where r_p is the distance from the star to the planet and s is a potential smoothing length set to 60% of the local disk pressure scale height (see, e.g., F. S. Masset 2002; T. W. A. Müller et al. 2012; P. Weber et al. 2019). We fix the planet orbital distance to 1 code unit of length (or $r_p = r_0$). In addition, we include a tapering such that the planet mass smoothly reaches its final

value in 30 planetary orbits. Finally, for the planet simulation, we use a constant α -viscosity value of 10^{-5} .

In Figure 5, we show the density (left panel), radial velocity (middle panel), and azimuthal velocity relative to a sub-Keplerian background (right panel) resulting from the simulations of a vortex triggered by the RWI at the edge of a dead zone (top row) and gap created by a massive planet (bottom row), respectively. We have used a scaling radius of 70 au to convert from code units to physical units (thus $r_{\text{DZ}} = 84$ au and $r_p = 70$ au). To highlight the patterns created by the vortex, we have subtracted the density field and sub-Keplerian rotation velocity of a smooth disk from the perturbed density and velocity fields. Since in this work we are primarily interested in the patterns created by a vortex and not the exact magnitude of the perturbation, we have further normalized the quantities by the local sub-Keplerian velocity. The simulation snapshots depicted correspond to 400 orbits measured at 1 code unit of length (dead zone) and 580 planetary orbits, respectively.

5. Results

Similar to previous works (e.g., P. Huang et al. 2019; C. M. T. Robert et al. 2020), it is apparent that the vortex creates a prominent azimuthally asymmetric feature in the gas dynamics, seen as a clear red-blue dipole pattern, as well as spiral structures (Figures 4 and 5). The vortex motion looks similar in both the simulations and analytical model, but the spiral features cannot be captured by the latter. If a planet is included in the disk, it creates its own Doppler-flip signature and much stronger spiral features, which overlay the pattern associated with the vortex, thus making it much harder to distinguish.

To compare the vortex signatures predicted by the simulations with the features seen in the velocity residuals, we postprocessed the dead zone simulation with RADMC-3D (C. P. Dullemond et al. 2012) and computed mock images of a flat disk model. In this context, we took the geometry of the sources into account and then matched the radial and azimuthal location of the vortex with the dust asymmetries seen in the data. To directly obtain the expected residuals from the simulated vortex dynamics, we subtracted the sub-Keplerian background from the azimuthal velocity used as an input in the radiative transfer model. We then created synthetic ^{12}CO channel maps via ray tracing with the same spectral resolution as the data. As a final step, we convolved the synthetic data cubes with a $0''.15$ Gaussian beam and computed the line-of-sight velocity maps using BETTERMOMENTS (R. Teague & D. Foreman-Mackey 2018).

The result of this test is shown in Figure 6, where the first two rows represent the ^{12}CO and ^{13}CO residuals, and the bottom row represents the dead-zone vortex simulation. The continuum emission is overlaid on the data to highlight the location of the dust crescents. For MWC 758, two representations are included since this disk is marked by several azimuthal asymmetries. The unconvolved model images are shown in Appendix C.

Comparing the data and model, it becomes apparent that most disks do not show a clear vortex kinematic signature around the dust crescents as predicted by the simulation. In HD 135344B, a similar red-blue pattern is traced by both emission lines, yet it is shifted with respect to the dust asymmetry. It is further overlaid by the strong spiral features, and it is nearby the disk minor axis, where a change of sign due

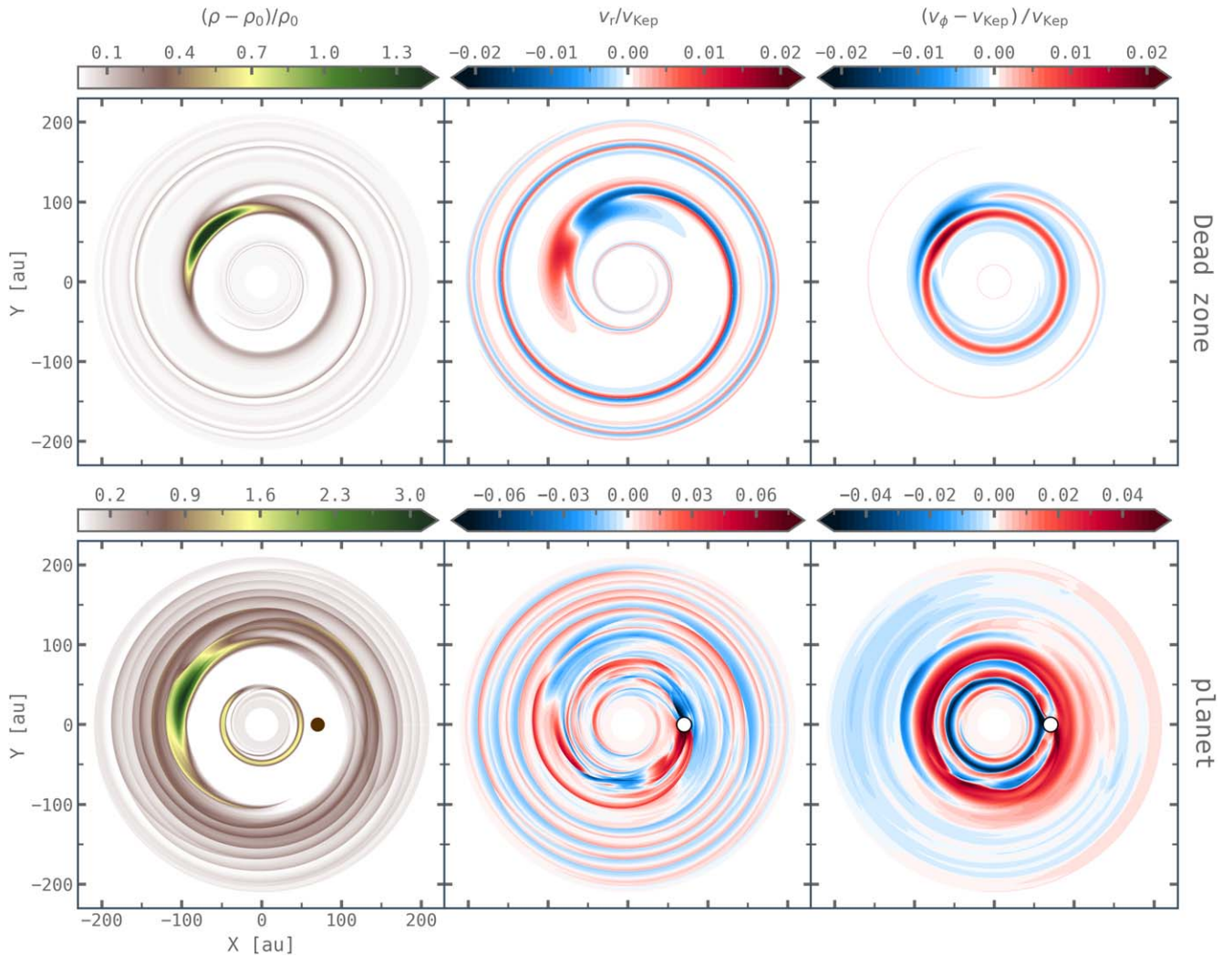


Figure 5. PLUTO simulations of a vortex triggered by the RWI at the edge of a dead zone (top row) and a Jupiter-mass planet (bottom row), using a scaling radius of 70 au to convert from code units to physical units. Shown are the gas density (left panels), radial velocity (middle panels), and azimuthal velocity (right panels). The density field and sub-Keplerian rotation velocity of a smooth disk have been removed (in case of v_r this is zero), and the quantities are normalized. The location of the planet is marked with a circle in the bottom row panels.

to coherent azimuthal variations is expected (Figure 5 in R. Teague et al. 2019b). Even though the vortex also triggers a similar spiral pattern, it is generally not expected to be strong enough to distinguish it in the kinematics (see also P. Huang et al. 2019). Close to the inner crescent in MWC 758, a red-blue flip is present, which resembles the vortex model, but it is part of the red and blue spiral arms and less distinct than the red-blue pair spanning the dust asymmetry in HD 135344B. Among our sample, this makes HD 135344B the most convincing example of a vortex and the best target to pursue with even deeper observations at very high angular resolution. Connected to that, when comparing the size of the dust asymmetries with the size of the beam, it becomes clear that they are very similar; thus, the vortex signatures may not be resolved and therefore do not clearly show up in the kinematics.

Aside from planet-disk interactions, other processes such as (magneto)hydrodynamical instabilities and the gravitational instability (C. Hall et al. 2020; M. Barraza-Alfaro et al. 2025; M. Barraza-Alfaro et al. 2024), stellar interactions including binaries (J. Calcino et al. 2020; B. J. Norfolk et al. 2022) and flybys (N. Cuello et al. 2020), or magnetic winds

(M. Galloway-Sprietsma et al. 2023) can contribute to the kinematics and may obscure the signatures from a vortex, resulting in complex patterns that are hard to disentangle. E. Ragusa et al. (2024) make predictions for the kinematics expected from an eccentric circumbinary disk, finding that they result in a single lobed pattern at the cavity edge. Compared to our data, we cannot distinguish such a pattern though this may result from resolution effects.

6. Summary and Conclusions

In this work, we have studied four disks within the exoALMA sample that show evidence of dust trapping in a vortex to search for a kinematic counterpart of the latter. Our main results are summarized as follows.

1. The kinematics of our targets, which all show azimuthal dust asymmetries, are marked by complex patterns, including spirals and Doppler-flip features (Figure 2). Two of the sources, HD 135344B and MWC 758, show prominent spirals in both the NIR scattered light and all three residuals, with the substructures partially coinciding (Appendix B).

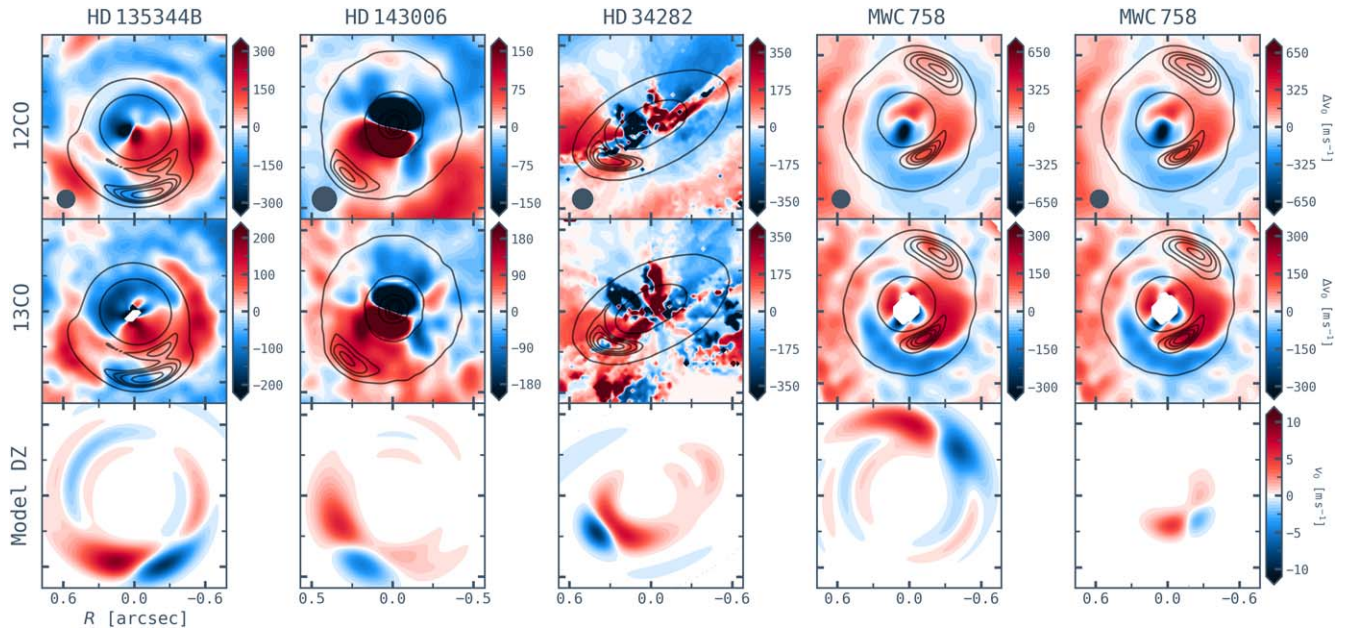


Figure 6. Velocity residuals of the data (top two rows) and dead-zone vortex model (bottom row). The model was postprocessed with RADMC-3D and convolved with the beam of the observation. For each disk, the model has been adapted such that the geometry and vortex location match the data. For MWC 758, two representations are included to account for both vortex features. In the first two rows, the continuum observations are overlaid, highlighting the dust asymmetries. The beam of the observation is shown in the lower left corner of each panel in the first row.

2. An analytical model of the velocity perturbation from a vortex reveals a distinct red-blue pattern but shows different morphologies depending on the underlying disk geometry and exact vortex parameters (Figure 3, 4). This simple model is, however, insufficient to explain the complexity of structure seen in the data.
3. Hydrodynamical simulations of a vortex and vortex +planet also reveal a clear red-blue kinematical pattern in addition to spiral features. In the planetary scenario, the signatures of the vortex are mixed with the strong perturbations from the planet, making it much more difficult to distinguish them (Figure 5). Other disk processes perturbing the gas dynamics are expected to further complicate this picture.
4. Comparing the kinematical patterns of the data and model shows that the features seen around the dust crescents cannot unambiguously be linked to a vortex (Figure 6). Aside from the complications arising due to different physical processes impacting the kinematics, our results may also be limited by the spatial resolution in comparison to the size of the dust asymmetries. While fully resolving the vortex is not necessarily needed to trace the red-blue flip, it is essential if we are to parameterize it and trace the pressure profile across the vortex. As shown by J. Stadler et al. (2025), the pressure morphology of dust rings can only be inferred if they are spatially resolved and we are faced with same challenge in the context of vortices.

While it is necessary to aim for higher spatial resolution to unambiguously detect the kinematic signature of a vortex, it can also be useful to study emission lines that trace closer to the midplane, where vortices reside, to be more sensitive to the vortex dynamics (compare C. Pinte et al. 2019, showing how the strength of a planet’s perturbation diminishes with height). This is the case for CS $J = 7-6$ (compare M. Galloway-Sprietsma et al. 2025), which was observed within the exoALMA program (G. Cataldi et al. 2025, in preparation). Both the SNR at high

spatial resolution and likely the resolution itself are, however, insufficient to draw conclusions about the kinematical pattern around the dust asymmetries but could be studied in a future work. Getting as close as possible to the midplane may increase our chance of detecting the vortex in the gas dynamics significantly.

The vortex predictions presented in this work rely on 2D models. However, to reveal the full dynamics of a vortex and in particular understand the role of vertical motions, global 3D simulations are essential. These simulations are challenging, and the formation and stability of three-dimensional vortices remain an active matter of debate. Earlier works have addressed the question of 2D versus 3D vortices triggered by the RWI but have resulted in different conclusions: while the general vortex patterns are consistent between 2D and 3D simulations, some studies suggest vertical motions to be negligible (M.-K. Lin 2012; S. Richard et al. 2013); others show that they can be more significant in certain cases (H. Meheut et al. 2010, 2012). More detailed theoretical studies are needed to understand how the strength of features varies with height and if they may be reliably detectable in tracers such as ^{12}CO and ^{13}CO .

Within the exoALMA sample, the most promising source for follow-up observations to detect the kinematic signature of an anticyclonic vortex is HD 135344B, as it is marked by a relatively large dust crescent, and the kinematics show a pattern that resembles the signatures of a vortex. In terms of the expected strength of the kinematic signature from a vortex, HD 34282 may be another useful target to follow up with higher spatial resolution and less abundant tracers at higher sensitivity due to its large inclination. The dust asymmetry in this source is, however, relatively small and located closer to the star, making such observations challenging.

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Appendix A Summary of the Sources

A.1. HD 135344B

HD 135344B—or SAO 206462—is a young (8.9 Myr; R. Asensio-Torres et al. 2021) Herbig F-type star located in the Upper Centaurus–Lupus star-forming region at a distance of ~ 135 pc (Gaia DR3; Gaia Collaboration et al. 2023). It is part of a visual binary system (separation $21''$, i.e., ~ 2800 au; B. D. Mason et al. 2001) and surrounded by a low-inclination transition disk that has been extensively studied at different wavelengths.

The disk was first resolved with the Hubble Space Telescope (C. A. Grady et al. 2009), with subsequent scattered-light observations revealing substructures such as an inner dust cavity (~ 28 au) and two prominent spiral arms on a scale of ~ 0.2 – $0.6''$ (T. Muto et al. 2012; A. Garufi et al. 2013; T. Stolker et al. 2016, 2017), one being marked by a brightness asymmetry and pitch-angle change, possibly caused by the spiral moving through a high-density vortex region (J. Bae et al. 2016). Moreover, the scattered-light images exhibit shadow features that may be related to misaligned inner disk regions (T. Stolker et al. 2016, 2017).

A large dust cavity (~ 50 au), originally suggested through spectral energy distribution (SED) modeling by J. M. Brown et al. (2007), was later confirmed in the millimeter continuum with the Submillimeter Array (SMA; J. M. Brown et al. 2009; S. M. Andrews et al. 2011) and ALMA (L. M. Pérez et al. 2014; N. van der Marel et al. 2016a; P. Cazzoletti et al. 2018). Additionally, these data revealed an asymmetric crescent located farther out (~ 80 au), which is marked by an azimuthal peak shift with wavelength and matches the predictions of a dust-trapping vortex (N. van der Marel et al. 2016a; P. Cazzoletti et al. 2018). The two dust rings are connected by a filament, which colocates with the deviation in the spiral arm and may be tracing a planetary wake crossing the dust gap (S. Casassus et al. 2021).

Observations of CO gas in the disk revealed an inner gas cavity peaking inside the dust cavity (~ 30 au) and a spiral feature in the kinematics (N. van der Marel et al. 2016a; S. Casassus et al. 2021). Similar to the NIR data, shadows are seen in archival $C^{18}O$ and exoALMA CS data (G. Cataldi et al. 2025, in preparation). Several works have shown that the

different observations may be explained by one or more massive embedded planets (e.g., T. Muto et al. 2012; A. Garufi et al. 2013; J. Fung & R. Dong 2015; J. Bae et al. 2016; T. Stolker et al. 2016; R. Dong & J. Fung 2017; S. Casassus et al. 2021; A. F. Izquierdo et al. 2025). While such companions were previously not found in a variety of high-contrast imaging programs (e.g., A. L. Maire et al. 2017; G. Cugno et al. 2019; A. Zurlo et al. 2020; R. Asensio-Torres et al. 2021; K. B. Follette et al. 2023), a potential candidate located at ~ 300 au in the disk was identified by G. Cugno et al. (2024) in JWST data.

A.2. HD 143006

The HD 143006 system, located at a distance of ~ 167 pc (Gaia Collaboration et al. 2023) in the Upper Scorpius star-forming region, consists of a G-type T Tauri star of age 7–12 Myr (S. M. Andrews et al. 2018; A. Garufi et al. 2018) and a disk seen close to face-on.

Imaging of the disk’s infrared scattered light revealed various brightness asymmetries, including a broad dark region covering the west side of the disk, two narrow shadow lanes, and a prominent overbrightness located in the southeast of a ring at ~ 18 – 30 au (M. Benisty et al. 2018). A similar, yet less pronounced, overbrightness is also present in the outer disk ring at ~ 50 – 83 au.

The millimeter continuum, on the other hand, is marked by three concentric dust rings (8, 40, and 64 au) and a bright crescent-like structure in the southeast, just outside the outer dust ring at 74 au (L. M. Pérez et al. 2018). The latter is further substructured into three different peaks and generally consistent with a vortex formed at the edge of a cavity, with its morphology, however, pointing toward a second-generation vortex (L. M. Pérez et al. 2018). Moreover, the bright arc coincides with the outer overbrightness in the scattered-light emission but is less extended, which matches predictions of dust trapping in vortices (C. Baruteau & Z. Zhu 2016).

Observations of CO gas emission in HD 143006 exhibit less structures but indicate gas depletion in the very inner disk and the presence of a kink feature in some redshifted channels (L. M. Pérez et al. 2018; C. Pinte et al. 2020). Both the scattered-light and millimeter observations support the scenario of a misaligned inner disk (M. Benisty et al. 2018; L. M. Pérez et al. 2018), with the morphology possibly being explained by the combination of a misaligned binary and an embedded aligned planet in the outer disk (G. Ballabio et al. 2021).

A.3. HD 34282

The Herbig Ae star HD 34282—or V1366 Ori—is about 8.9 Myr old (R. Asensio-Torres et al. 2021) and located at a distance of ~ 309 pc (Gaia Collaboration et al. 2023). A highly inclined circumstellar disk around the source was resolved in the millimeter emission by V. Piétu et al. (2003). A large inner dust cavity in the disk was suggested through modeling of the SED and mid-infrared Q-band images (B. Acke et al. 2009; S. Khalafinejad et al. 2016) and later confirmed (~ 80 au) in ALMA data by G. van der Plas et al. (2017). These images further revealed that the ~ 300 au wide dust ring contains an azimuthal asymmetry in the southeast. The authors suggest a massive brown dwarf companion ($\sim 50 M_J$) to be responsible for both observed substructures.

Scattered-light observations allowed J. de Boer et al. (2021) to identify two inclined rings and a tightly wound single-arm spiral, which is generally consistent with a Jupiter-mass planet and coincides with the millimeter crescent. The latter may be the signature of a vortex in the disk, which, as shown by M. Marr & R. Dong (2022), can resemble a spiral arm in a highly inclined source. So far, no direct evidence of a companion in HD 34282 has been found (R. Asensio-Torres et al. 2021; Y. Boehler et al. 2021; J. Quiroz et al. 2022). The study of A. J. Bohn et al. (2022), who compare position angle and inclination of the inner disk measured with VLTI/GRAVITY and the outer disk measured with ALMA, suggests a misalignment between the inner and outer disks of the system, yet no clear shadow signatures are visible in the scattered-light images.

A.4. MWC 758

Located in the Taurus star-forming region at a distance of 156 pc (Gaia Collaboration et al. 2023), the young (~ 8.9 Myr; R. Asensio-Torres et al. 2021) Herbig Ae star MWC 758—also known as HD 36112—is surrounded by a low-inclination transition disk that has been the subject of a range of high-resolution studies, revealing a complex asymmetric morphology.

NIR scattered-light observations unveiled two prominent spiral arms (C. A. Grady et al. 2013; M. Benisty et al. 2015) and other nonaxisymmetric arc-like features (M. Benisty et al. 2015). The detection of a third spiral arm and potential point-like source at ~ 17 au were reported by M. Reggiani et al. (2018). Both gravitational instability (R. Dong et al. 2015a) and massive planets (R. Dong et al. 2015b; J. Bae & Z. Zhu 2018; C. Baruteau et al. 2019; J. Calcino et al. 2020) have been suggested as the driving mechanism of the spirals, with the latter being favored by an imaging study of the spiral pattern over 5 yr (B. Ren et al. 2020).

No fully depleted cavity was found in the scattered-light data, but the millimeter continuum of the disk is marked by a large and eccentric cavity (~ 40 au; R. Dong et al. 2018b), possibly carved by an embedded planet (A. Isella et al. 2010; S. M. Andrews et al. 2011; S. Marino et al. 2015; Y. Boehler et al. 2018). Additionally, two bright emission clumps inside a double-ring structure are present at radii of ~ 47 and 82 au (S. Marino et al. 2015; Y. Boehler et al. 2018; R. Dong et al. 2018b; S. Casassus et al. 2019) and may be related to dust-trapping pressure maxima in a vortex.

Several spirals have been identified in both the continuum emission (Y. Boehler et al. 2018; R. Dong et al. 2018b; B.-T. Shen et al. 2020) and CO data (Y. Boehler et al. 2018), which are partly colocated with the spirals observed in the infrared data. C. Baruteau et al. (2019) performed hydrodynamical simulations showing that the spirals, the eccentric and asymmetric ring structure, and the crescent-shaped features may all be explained by two Jupiter-mass planets, one located inside and one outside the spirals. High-contrast imaging searches have reported the detection of such planets (M. Reggiani et al. 2018; K. Wagner et al. 2019, 2023), but they remain to be confirmed.

Appendix B Spiral Features

In Figure 7, we present the polar-deprojected CO residuals for the centroid velocity, peak intensity, and line width alongside the polar-deprojected scattered-light emission for

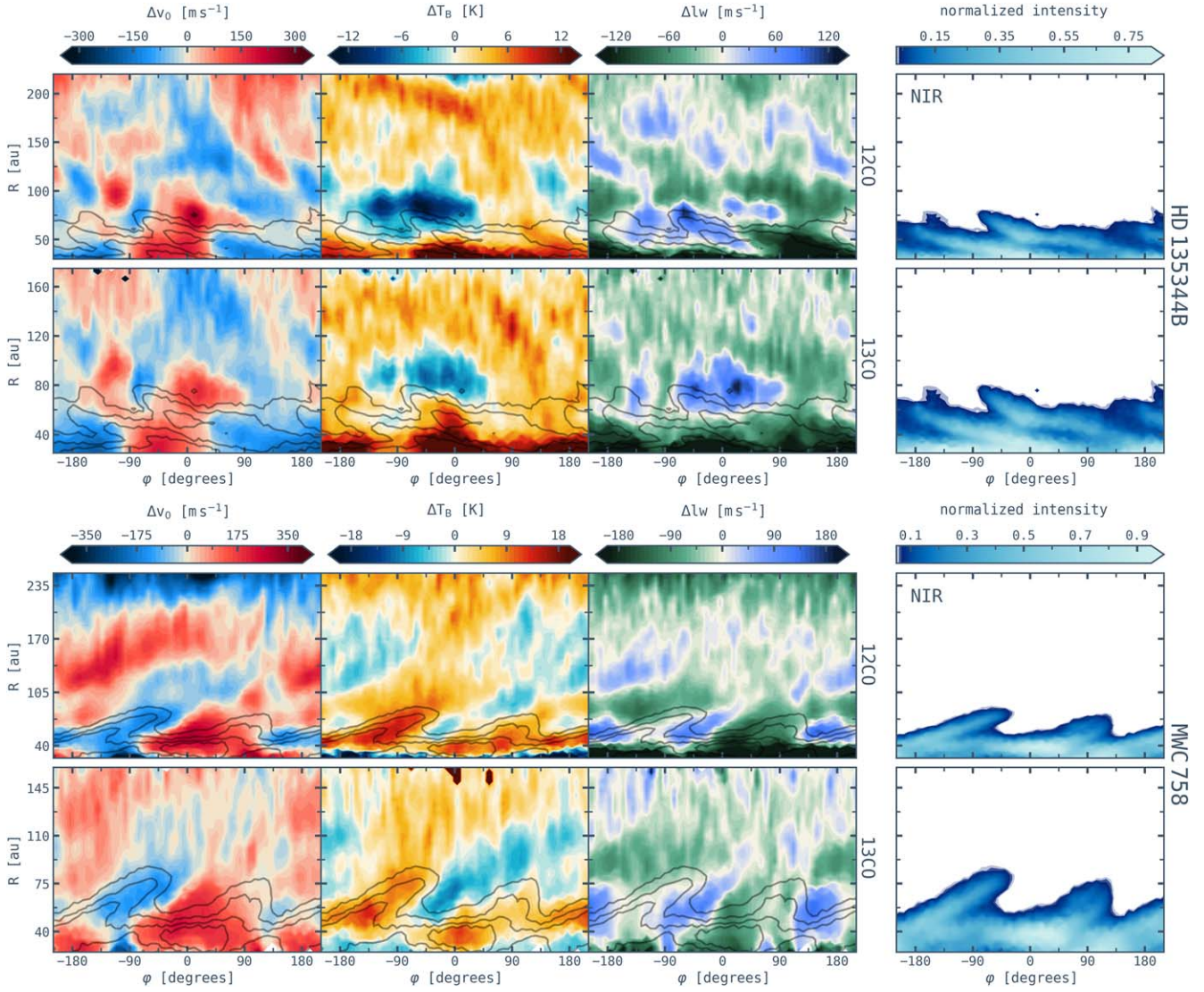


Figure 7. Polar-deprojected residuals of the centroid velocity, peak intensity, and line width, shown for both CO lines and the HD 135344B (top) and MWC 758 (bottom) disks. The deprojection of the NIR scattered-light emission is included in the fourth panel of each row (T. Stolker et al. 2017; B. B. Ren et al. 2023), highlighting the observed spiral features that are overlaid as contours on the other three panels. The SPHERE images are normalized to the peak value and shown with a logarithmic normalization of the color map. Note the different spatial scale of the rows.

the two disks that show prominent spiral features: HD 135344B and MWC 758. For comparison, the spirals found in the NIR are further overlaid as contours on the gas emission residuals. The deprojection is performed by using the inclination and position angle (see Table 1; A. F. Izquierdo et al. 2025) of the individual sources. This is sufficient due to the face-on nature of the disks.

For both disks, the inner spirals seen in the velocity residuals appear to align somewhat with the NIR spirals. In HD 135344B, both NIR spiral arms are located just inside the redshifted velocity arms. In MWC 758, one of the NIR spirals (-90° to 90°) has a redshifted counterpart in the velocity residuals, while the other has a blueshifted counterpart (-180° to -90°). For the peak intensity, on the other hand, the spiral patterns seen for MWC 758 align very well with the NIR spirals, whereas this is not the case for HD 135344B. We note that the strong temperature decrement seen around 80 au in this source is the result of continuum subtraction. In the line width, spiral patterns are mostly visible in the negative residuals, which again seem to somewhat align with the NIR spirals for

MWC 758 but not HD 135344B (for a detailed discussion, see M. Benisty et al. 2025, in preparation).

Appendix C Unconvolved Model Images

In Figure 8, we show the unconvolved model images of the dead-zone simulation alongside the convolved images ($0''.15$ beam). The continuum observations are overlaid as contours to highlight that the model has been adapted to each disk, such that the vortex is present at the same location as the dust asymmetries. In this work, we are mainly interested in the kinematical patterns of a vortex, and we note that the exact strength of these features depends on various factors. Moreover, we assume a flat disk and the magnitude of the vortex may be significantly underestimated. In that regard, the unconvolved images nevertheless emphasize the need for higher spatial resolution. In particular, an inclined source such as HD 34282 may show very strong vortex features if resolved at very high resolution.

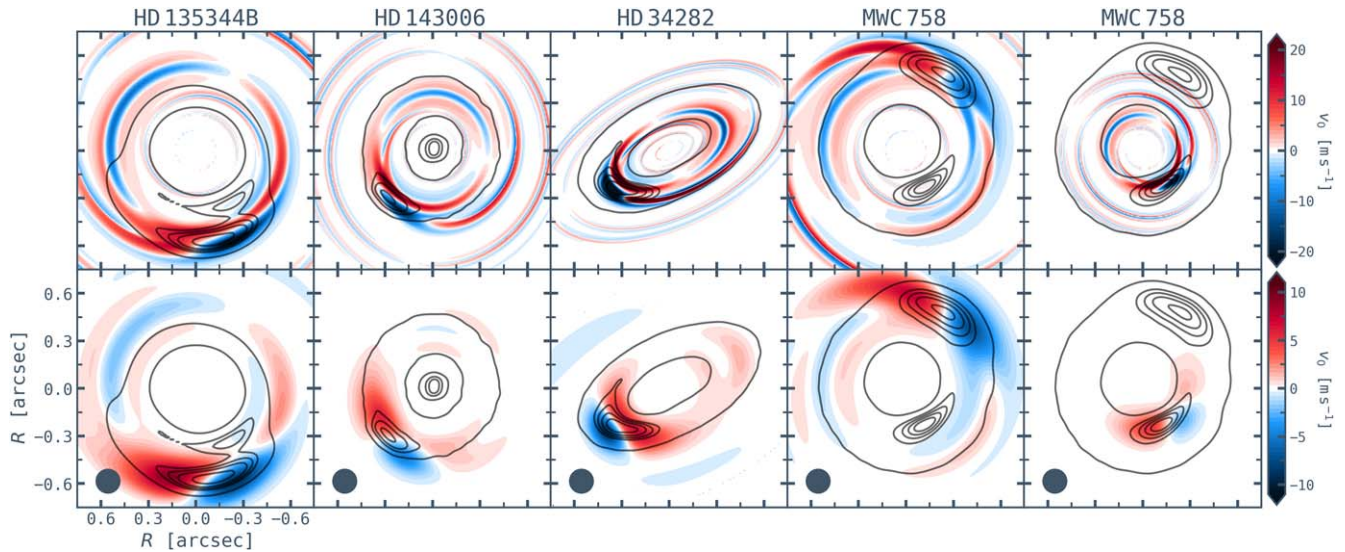


Figure 8. Velocity perturbations of the dead-zone vortex simulation, postprocessed with RADMC-3D and adapted to each dust asymmetry observed in the data. Shown are the unconvolved models (top row) and models convolved with the beam of the observation (bottom row). The continuum observations are overlaid in all panels to highlight the location of the dust asymmetries.

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