



Universiteit
Leiden
The Netherlands

The large magellanic cloud revealed in gravitational waves with LISA

Keim, M.A.; Korol, V.; Rossi, E.M.

Citation

Keim, M. A., Korol, V., & Rossi, E. M. (2023). The large magellanic cloud revealed in gravitational waves with LISA. *Monthly Notices Of The Royal Astronomical Society*, 521(1), 1088-1098. doi:10.1093/mnras/stad554

Version: Publisher's Version

License: [Creative Commons CC BY 4.0 license](https://creativecommons.org/licenses/by/4.0/)

Downloaded from: <https://hdl.handle.net/1887/3719208>

Note: To cite this publication please use the final published version (if applicable).

The large magellanic cloud revealed in gravitational waves with LISA

Michael A. Keim¹, Valeriya Korol^{2,3,4} and Elena M. Rossi²

¹Department of Astronomy, Yale University, PO Box 208101, New Haven, CT 06520-8101, USA

²Leiden Observatory, Leiden University, PO Box 9513, RA Leiden 2300, The Netherlands

³Max-Planck-Institut für Astrophysik, Karl-Schwarzschild-Straße 1, Garching 85741, Germany

⁴Institute for Gravitational Wave Astronomy & School of Physics and Astronomy, University of Birmingham, Birmingham B15 2TT, UK

Accepted 2023 February 4. Received 2023 February 4; in original form 2022 July 28

ABSTRACT

The Laser Interferometer Space Antenna (LISA) will unveil the non-transient gravitational wave sky full of inspiralling stellar-mass compact binaries within the local universe. The Large Magellanic Cloud (LMC) is expected to be prominent on the LISA sky due to its proximity and its large population of double white dwarfs (DWD). Here, we present the first dedicated study of the LMC with gravitational wave sources. We assemble three LMC models based on: (1) the density distribution and star formation history from optical wavelength observations, (2) a detailed hydrodynamic simulation, and (3) combining the two. Our models yield a hundred to several hundred detectable DWDs: indeed, the LMC will be a resolved galaxy in the LISA sky. Importantly, amongst these we forecast a few tens to a hundred double degenerate supernovae type Ia progenitors, a class of binaries which have never been unambiguously observed. The range in the number of detections is primarily due to differences in the LMC total stellar mass and recent star formation in our models. Our results suggest that the total number, periods, and chirp masses of LISA sources may provide independent constraints on both LMC stellar mass and recent star formation by comparing LISA observations with the models, although such constraints will be highly model-dependent. Our publicly available model populations may be used in future studies of the LMC, including its structure and contribution to LISA confusion noise.

Key words: gravitational waves – binaries (including multiple): close – white dwarfs – Magellanic Clouds – galaxies: stellar content.

1 INTRODUCTION

The milli-Hertz frequency band offers an opportunity to study persistent non-transient gravitational wave (GW) sources such as inspiralling stellar-mass compact object binaries in the local universe. Space-based GW observatories like the European Space Agency-led Laser Interferometer Space Antenna (LISA, Amaro-Seoane et al. 2017, see also TianQin Luo et al. 2016 and Taiji Ruan et al. 2018 missions) in the 2030s will detect such sources in abundance (for a review see Amaro-Seoane et al. 2022a). In particular, double white dwarfs (DWD) – being the most common stellar binary remnants – will make it possible to map our Galaxy with GWs (Adams, Cornish & Littenberg 2012; Korol, Rossi & Barausse 2019; Wilhelm et al. 2021; Georgousi et al. 2022). Next to the Milky Way, the Large Magellanic Cloud (LMC) will be the biggest and likely the only extended structure visible on the LISA sky (Korol, Koop & Rossi 2018; Korol et al. 2020; Roebber et al. 2020).

As the most massive satellite of our Milky Way, the Large Magellanic Cloud (LMC) is a prime science target for LISA. Its close proximity (at ~ 50 kpc, Pietrzyński et al. 2013; de Grijs, Wicker & Bono 2014) allows the detection of DWD GW signals with LISA for periods shorter than ~ 20 min (Korol et al. 2018). In addition, the LMC is an active site of star formation (Harris & Zaritsky 2009), which has been shown to have a critical effect on

the detectable DWD population (Korol et al. 2020). As a result, the LMC has been predicted to house $\sim \mathcal{O}(10^2)$ LISA-detectable DWDs, the most of any Milky Way satellite (Korol et al. 2020; Roebber et al. 2020; possibly including a few – several double neutron star binaries Seto 2019; Andrews et al. 2020, Lau et al. 2020). GW observations will therefore extend studies of the LMC to stellar populations which are generally inaccessible with electromagnetic observatories (e.g. Rebassa-Mansergas et al. 2019). Still, earlier studies have neither explored the full range of observed LMC stellar masses, nor the spatial distribution of the LMC that may play a relevant role given its proximity (e.g. Storck & Church 2022, for double neutron stars in the Milky Way). Including such factors and the star formation history inferred from electromagnetic observations will give a realistic picture of the LMC's DWD population and help us to understand LMC properties that LISA can explore.

In this work we construct a realistic population of DWDs in the LMC and model their GW emission with the aim of expanding the LISA science case to the Milky Way's neighbourhood. We estimate the number of DWDs detectable within the nominal 4 yr duration of the mission and we show that our results strongly depend on the assumed LMC star formation history and total stellar mass. The ultimate goal is to determine whether LISA can spatially resolve the LMC. In Section 2 we assemble various models of the DWD population in the LMC. These are designed to provide different total stellar masses, star formation histories, and spatial distributions of stars. We populate these models with DWDs from a fiducial binary

* E-mail: michael.keim@yale.edu

population synthesis simulation of Korol et al. (2020) and we model DWDs' GW signals using the pipeline presented in Karnesis et al. (2021) to assess their detectability with LISA. In Section 3 we describe the resulting LMC binary populations detectable by LISA, and we show how they can constrain overall LMC properties. Finally, discuss our results in Section 4.

2 METHODS

In this section, we first summarize LMC features known from electromagnetic radiation that are relevant for our study. Next, we construct three different LMC models: an observation-driven model ('Model 1'), a simulation-driven model ('Model 2'), and a combined model ('Model 3'). We then populate our LMC models with DWDs from a fiducial binary population synthesis simulation. Finally, we calculate the DWDs' GW signals and evaluate their detectability with LISA.

2.1 The large magellanic cloud at electromagnetic wavelengths

The LMC is part of the magellanic system consisting of the Large and Small Magellanic Clouds (SMCs). It is linked to the SMC by a largely gaseous magellanic bridge, it is trailed by a gaseous magellanic stream, and extends towards the Milky Way via a gaseous leading arm (Brüns et al. 2005). The origin of these features has been attributed to either strong interaction or collision between the LMC and SMC (Yoshizawa & Noguchi 2003; Besla et al. 2012), as have the LMC's subdominant arms (Besla et al. 2016; Belokurov & Erkal 2019), warping of the LMC's disk (Choi et al. 2018; Mackey et al. 2018) and central bar (Subramaniam 2003; Besla et al. 2012; Haschke, Grebel & Duffau 2012), the accretion of SMC stars onto the LMC (Olsen et al. 2011; Besla et al. 2012; De Leo et al. 2020), and the star formation activity of both the LMC and SMC (Yoshizawa & Noguchi 2003; Harris & Zaritsky 2009; Indu & Subramaniam 2011).

A wide array of studies has been dedicated to constraining the LMC mass. Estimates from mass-to-light ratio suggest stellar masses of $1.5 \times 10^9 M_\odot$ (McConnachie 2012) to $2.7 \times 10^9 M_\odot$ (van der Marel et al. 2002), whereas rotation curve modelling suggests a neutral gas mass of $5.2 \times 10^8 M_\odot$ and a stellar population mass of $2.0 \times 10^9 M_\odot$ within a radius of 4 kpc (Kim et al. 1998). The LMC's dark matter dominated total (or dynamical) mass has been estimated using kinematics of stars (Schommer et al. 1992; Kim et al. 1998; van der Marel & Kallivayalil 2014; Cullinane et al. 2020; Wan et al. 2020) and its interaction with surrounding galaxies and stellar streams (Kallivayalil et al. 2013; Peñarrubia et al. 2016; Laporte et al. 2018; Shao et al. 2018; Erkal et al. 2019; Erkal & Belokurov 2020) with estimates arriving up to $3.1 \times 10^{11} M_\odot$. The LMC is an active topic of research and investigations have so far ultimately relied on electromagnetic observations only. With this paper we start exploring GWs as a promising avenue for synergistic and independent studies of the LMC's mass content and star formation history, whose results would be based on completely different selection effects.

2.2 Large magellanic cloud models

Model 1: observation driven. We first consider an observation driven model based on Harris & Zaritsky (2009)'s analysis of the 20 million LMC stars observed by the Magellanic Clouds Photometric Survey (Zaritsky et al. 2004). To derive a star formation history, the authors generate artificial colour-magnitude diagrams corresponding to stellar populations with 16 different ages between 6.3 Myr and 15.8 Gyr, sampled in log-space. They then fit these as a weighted

sum to over a thousand different two-dimensional (2D) sky regions of the LMC (see Fig. 4 of Harris & Zaritsky 2009). Since they apply a completeness correction in order to account for faint, unobserved stars, the resulting fit represents not only the star formation history, but also the overall mass distribution as well.

First, for each 2D sky region, we fit the star formation history with a spline and apply a global normalization factor so that the resulting star formation rate represents that of main-sequence, main-sequence (MS+MS) binaries which are the progenitors of DWD binaries. Such systems have initial MS binary components with ranging from $0.14 M_\odot$ to $11 M_\odot$, though these masses will evolve subsequent to the zero age main sequence due to stellar processes and binary interactions resulting in mass loss and mass transfer. Analytically, this normalization constant c is obtained by solving

$$c \sum_{i=1}^{1376} \int_{t_0=10 \text{ Gyr}}^{t_f=0 \text{ Gyr}} \text{SFR}_i(t) dt = \sum_{j=1}^N M_{\text{MS+MS},j}, \quad (1)$$

where 1376 is the total number of regions in Harris & Zaritsky (2009), SFR_i is the region i 's spline-fit star formation rate, $M_{\text{MS+MS},j}$ is the initial mass of each MS+MS progenitor binary, and N is the total number of binaries. Given that most LMC stars were born in the last 10 Gyr with earlier times being largely quiescent (e.g. Butcher 1977; Harris & Zaritsky 2009; Ruiz-Lara 2019), on the left hand side of equation (1) we integrate in lookback time from $t_0 = 10$ Gyr until the present day $t_f = 0$ Gyr. The density distribution of binaries across the LMC regions is then determined based on each 2D sky region's total mass in binaries,

$$M_i = \int_{t_0=10 \text{ Gyr}}^{t_f=0 \text{ Gyr}} c \text{SFR}_i(t) dt. \quad (2)$$

We assign formation times by dividing the total masses of binaries by the relevant region's spline-fit star formation rate and advancing in time by $\Delta t = M_{\text{MS+MS},j} / \text{SFR}_i(t)$. To account for periods of quiescence in a region which would make Δt an over estimate, if Δt exceeds 100 Myr we instead evaluate the formation time over finite time steps dt , continually recalculating $dm_k = \text{SFR}_i(t_k) \times dt$ until reaching $\sum dm_i = M_{\text{MS+MS},j}$.

Since the mass distribution found by integrating LMC regions only accounts for 2D sky distribution, to find the distance from the Sun we adopt a Hernquist potential (Hernquist 1990). We consider the transform (e.g. Weinberg & Nikolaev 2001)

$$\begin{aligned} x &= -\cos(\delta) \sin(\alpha - \alpha_{\text{LMC}})d \\ y &= [\sin(\delta) \cos(\delta_{\text{LMC}}) - \sin(\delta_{\text{LMC}}) \cos(\delta) \cos(\alpha - \alpha_{\text{LMC}})]d \\ z &= d_{\text{LMC}} - [\sin(\delta) \sin(\delta_{\text{LMC}}) + \cos(\delta) \cos(\delta_{\text{LMC}}) \cos(\alpha - \alpha_{\text{LMC}})]d \end{aligned} \quad (3)$$

to convert right ascension (RA) α , declination (DEC) δ , distance d coordinates to the (x, y, z) frame for the LMC, where $(\alpha_{\text{LMC}}, \delta_{\text{LMC}}, d_{\text{LMC}})$ is set to $(80^\circ, -69^\circ, 50 \text{ kpc})$. The distance to the equatorial plane ($z = 0$) for each LMC region is then given by

$$d_{z=0} = \frac{d_{\text{LMC}}}{\sin(\delta) \sin(\delta_{\text{LMC}}) + \cos(\delta_{\text{LMC}}) \cos(\delta) \cos(\alpha - \alpha_{\text{LMC}})}. \quad (4)$$

For a region with a given (RA, DEC), equations (3) and (4) give the corresponding (x, y) coordinates of that region's binaries. Binaries are then assigned a z value in the interval $[-4, 4] \text{ kpc}$ using a Hernquist profile, taking the mass density as a probability density so that

$$P(z) \propto \begin{cases} r^{-1}(z) \times (r(z) + r_{\text{hern}})^{-3} & \text{for } r(z) > r_{\text{min}}, \\ r_{\text{min}}^{-1} \times (r_{\text{min}} + r_{\text{hern}})^{-3} & \text{for } r(z) \leq r_{\text{min}}, \end{cases} \quad (5)$$

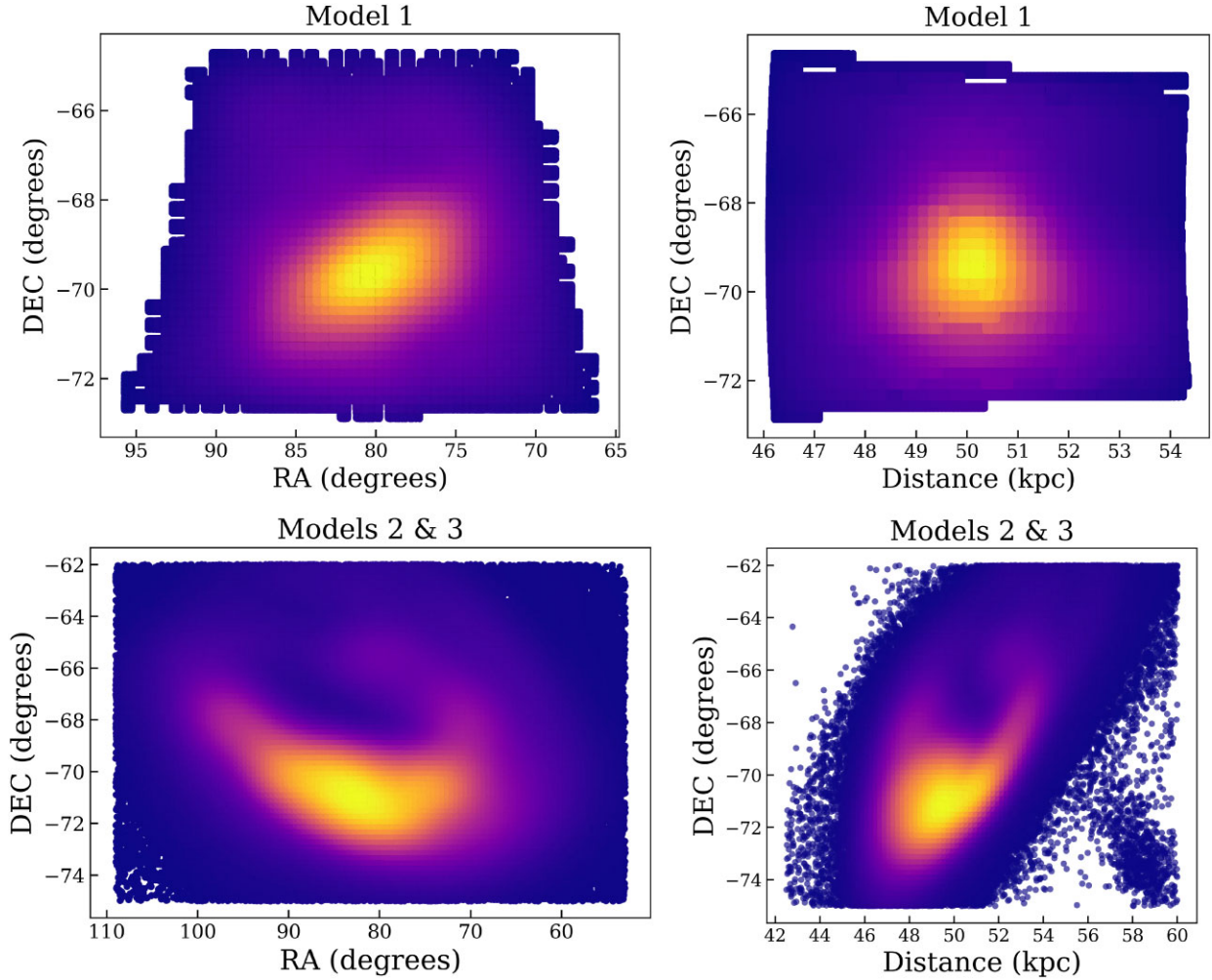


Figure 1. Top panels: Model 1 (the observation driven model)’s spatial distribution of binaries on the sky (RA, DEC; left) and along the line-of-sight (distance, DEC; right). Binaries are matched to the regions of Harris & Zaritsky (2009) based on each region’s observed star formation history, which are normalized so that integration gives the total mass of MSMS progenitors that should be assigned to each region. Binaries then adopt the sky location of their associated region. Distance is assigned based on a Hernquist distribution. Bottom panels: spatial distribution of binaries on the sky (RA, DEC; left) and along the line-of-sight (distance, DEC; right) for Models 2 and 3, the simulated and combined models. Distance and sky location are assigned based on stellar particles from the simulation of Lucchini et al. (2020). For all panels colour is indicative of local binary density with yellow corresponding to high density and blue corresponding to low density.

where $r_{\min} = 0.1 \text{ kpc}$, r_{hern} is taken as 2 kpc to reflect the full-width-at-half-maximum of the distribution of Cepheid variable stars (Jacyszyn-Dobrzniecka et al. 2016), and $r(z) = \sqrt{x^2 + y^2 + (z)^2}$. Finally, we invert equation (3) to find a binary’s final, deprojected coordinates.

The distribution of binaries on the sky based on 2D sky regions, as well as the distribution along the line of sight following equation (5), is given in top panels of Fig. 1. We note that this simplistic spatial distribution lacks small scale LMC features (cf. Section 2.1). In addition, the symmetric z -axis limits our examination of how LISA may utilize DWD distance to probe the LMC structure, especially crucial as lack of dust extinction makes GW ideal for such a purpose. It is therefore useful to consider a model with a more sophisticated spatial distribution as explored in the following section.

Model 2: simulation driven. To obtain a more realistic spatial distribution that includes characteristic LMC structural features and a realistic stellar distribution along z -axis, we utilize the N-body hydrodynamical simulation of Lucchini et al. (2020). It models

the entire magellanic system including both LMC and SMC, the magellanic bridge, the magellanic stream, and the leading arm. This simulation consists of $\sim 6.5 \times 10^5$ particles; their spatial distribution projected on the sky is shown in Fig. 1 (bottom panels). We highlight that the simulation clearly features a northern arm, and its central bar has the same characteristic warp as discovered by Subramaniam 2003. While this model does not perfectly replicate the LMC’s structure – e.g. having a northern arm that is closer to Earth than the central bar (Jacyszyn-Dobrzniecka et al. 2016) – it nevertheless provides a spatial distribution with detailed realistic features that potentially can be probed with GW observations.

We populate the model LMC by assigning DWD binaries the ages and locations of the simulated stellar particles. We note that the bulk of simulated particles are formed 6 Gyr ago, which represents the starting point of Lucchini et al. (2020)’s simulation. For DWD in these older particles we randomly assign a formation time between 0 and 4 Gyr. In this way we obtain an alternative star formation history from that of the observation driven model.

Model 3: combined. For a final, ‘combined’ model, we use the best features of the previous two models: the data-driven star formation history of Model 1 and the spatial distribution from Model 2. To achieve this in practice, we match each particle simulation particle of Lucchini et al. (2020)’s simulation to the nearest 2D sky region in Model 1 (cf. fig. 4 of Harris & Zaritsky 2009) and adopt its star formation history. In this way, we keep the detailed structure of the LMC seen in bottom panels of Fig. 1, but end up with a temporal distribution that is in agreement with electromagnetic observations.

2.3 Double white dwarfs synthetic population

We employ a synthetic collection of DWDs from Korol et al. (2020) generated using the SEBA stellar and binary evolution module (Portegies Zwart & Verbunt 1996; Nelemans et al. 2001; Toonen, Nelemans & Portegies Zwart 2012). Amongst their models we use the one with a fixed binary fraction of 50 per cent and a fixed metallicity $Z = 0.001$. Although LMC stars present a range of metallicities, Korol et al. (2020) showed that this assumption has only a moderate effect on the number of LISA DWD detections in Milky Way satellites. Note that by using a fixed binary fraction we neglect a potential correlation between the metallicity and primordial binary fraction that has been found in the Solar neighbourhood (Badenes et al. 2018; El-Badry & Rix 2019; Moe, Kratter & Badenes 2019).

As the next step of our modelling, we adapt binaries’ parameters (as provided by SEBA) to resemble LMC’s star formation history and spatial distribution. Specifically, we use the DWD formation time, i.e. the time it takes to turn both main sequence stars in the binary into white dwarfs, the orbital separation at DWD formation, and the white dwarf masses. First, to account for the LMC stellar mass of $\sim \mathcal{O}(10^9) M_\odot$, we scale up our DWD population. We estimate the number of binaries injected in each LMC model based on the total stellar mass $M_\star$ by linearly re-scaling the synthetic collection by

$$N_{\text{DWD}} = \frac{M_\star}{M_{\text{SeBa}}} N_{\text{DWD,SeBa}}, \quad (6)$$

where $N_{\text{DWD,SeBa}}$ is the number of DWDs in the synthetic population and M_{SeBa} is the total simulated population mass. We consider several LMC mass estimates in the range between $1.5 \times 10^9 M_\odot$ and $3.0 \times 10^9 M_\odot$ (Kim et al. 1998; van der Marel et al. 2002; McConnachie 2012). Second, we employ our three alternative LMC models to assign sky locations and distances to synthetic DWDs. Third, we assign DWD ages according to the models’ star formation history. Lastly, we adjust binaries’ orbital separation to reflect energy loss from GW emission since the binaries first became DWD until the present day. This step allows us to determine which binaries are emitting in the LISA band today and to exclude those that have already merged. We compute DWD present-day GW frequencies as

$$f(t) = \frac{3^{3/8}}{2^{9/8}\pi} \left(\frac{96}{5}\right)^{-3/8} \left(\frac{G\mathcal{M}}{c^3}\right)^{-5/8} (\tau_{\text{merge}} - t)^{-3/8}, \quad (7)$$

where G is the gravitational constant, c is the speed of light, the chirp mass $\mathcal{M} = (m_1 m_2)^{3/5} / (m_1 + m_2)^{1/5}$ with m_1 and m_2 being the primary and the secondary white dwarf masses, and τ_{merge} is the merger time calculated from the initial frequency at the time of DWD formation (e.g. Maggiore 2008).

In Fig. 2 we summarize the obtained DWD formation histories in our three LMC models: observation driven (Model 1 in blue), simulation driven (Model 2 in grey), and combined (Model 3 in orange). We highlight a stark difference between the observation driven (Model 1) and simulation driven (Model 2) models. In the later, DWD formation is very steep up to ~ 4 Gyr and negligible since. In

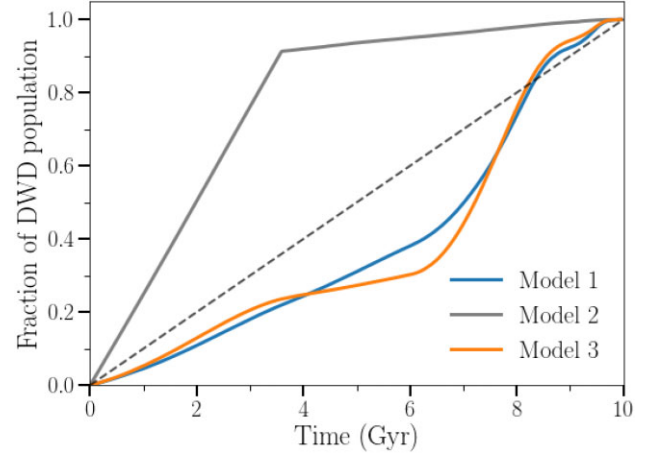


Figure 2. The DWD formation history for Model 1 (observation driven) in blue, Model 2 (simulation driven) in grey, and Model 3 (combined) in orange. The y-axis represents the cumulative fraction of the total DWD population emitting in the LISA frequency band today as they formed over time in the LMC (x-axis). As a reference, we also show the cumulative population for a constant star formation history as dashed black line (see Fig. A1 for the historic star formation rate).

the former, the star formation rate trend is reversed: DWD formation is shallow up to 6 Gyr followed by a significant grows in the last 3 Gyr. In particular, a burst around 500–800 Myr ago is responsible for a much larger young binary population. This is best seen with logarithmic spacing as given in Fig. A1, and, along with a more recent burst ~ 100 Myr ago is thought to correspond with interactions with the SMC (Gardiner & Noguchi 1996; Pietrzynski & Udalski 2000; Yoshizawa & Noguchi 2003; Harris & Zaritsky 2009; Indu & Subramaniam 2011; Besla et al. 2012). More quantitatively, we read from Fig. 2 that in Model 2 the DWD population reaches 90 per cent already at ~ 4 Gyr, whereas in Models 1 and 3, which share the same observation driven star formation history, it reaches 90 per cent only at 9 Gyr. This difference makes a significant impact on the size and the properties of the LISA detectable DWDs. As discussed in Korol et al. (2020), recent star formation activity adds newly more massive short period binaries to the LISA band that are also more easy to detect; in old stellar populations more massive short period binaries have already merged, whereas long period binaries have not yet evolved into the LISA band. We also note a slight difference in the DWD formation histories of the combined model and observation driven model. These arise because the number of binaries falling into the different LMC 2D sky regions is based on a different underlying mass distribution. Thus, the various star formation pixels in Models 1 and 3 have a different weight in the overall star formation history.

2.4 Double white dwarfs detectability with laser interferometer space antenna

As the result of the methodology described so far, we obtain a catalogue of $\sim \mathcal{O}(10^6)$ LMC DWDs emitting in the LISA frequency band. As the next step, we use the LISA data analysis pipeline presented in Karnesis et al. (2021), based on an signal-to-noise (S/N) evaluation using an iterative scheme for the estimate of the confusion foreground generated by Milky Way’s GW sources (see also Timpano, Rubbo & Cornish 2006; Crowder & Cornish 2007; Nissanke et al. 2012). We use LISA’s instrumental noise requirements defined in a technical note by LISA Science Study

Team (2018). We set a LISA mission duration of 4 yr and an S/N detection threshold of 7.

The scheme begins with the generation of the signal measured by LISA, by computing quasi-monochromatic waveform for each DWD and by projecting it on the LISA arms for a given duration of the mission. Such a waveform can be fully determined by eight parameters: amplitude $\mathcal{A}$, frequency f , frequency derivative or chirp $\dot{f}$, sky position (θ, ϕ) in ecliptic heliocentric coordinates, orbital inclination ι , polarization angle ψ , and initial orbital phase ϕ_0 (e.g. Cutler 1998). The first five parameters represent the outcome of our modelling procedure (cf. Section 2.3), whereas the remaining three are angular parameters that we assign randomly: ι is sampled from a uniform distribution in $\cos \iota$, whereas ψ and ϕ_0 are sampled from a flat distribution. We compute the gravitational amplitude as

$$\mathcal{A} = \frac{2(G\mathcal{M})^{5/3}(\pi f)^{2/3}}{c^4 d}, \quad (8)$$

and the frequency derivative as

$$\dot{f} = \frac{96}{5}\pi^{8/3} \left(\frac{G\mathcal{M}}{c^3} \right)^{5/3} f^{11/3}. \quad (9)$$

For each binary in the catalogue, the pipeline then computes S/N and performs an Fisher information matrix (FIM) analysis in order to estimate the accuracy of the parameter recovery. We recall that FIM for each recovered source can be evaluated as

$$F_{ij} = \left(\frac{\partial h(\vec{\theta})}{\partial \theta_i} \middle| \frac{\partial h(\vec{\theta})}{\partial \theta_j} \right) \bigg|_{\vec{\theta}=\vec{\theta}_{\text{true}}}, \quad (10)$$

where h is the template of each signal, and $\vec{\theta}$ represents the waveform parameter vector. The inverse of F_{ij} yields the covariance matrix with the diagonal elements being the mean square errors on each parameter θ_i , and the off-diagonal elements describing the correlations between the parameters $\rho_{\theta_i\theta_j}$. We remark that errors derived via FIM analysis are valid for relatively high S/N (e.g. Cutler 1998). Thus, we expect that in some cases the errors may be underestimated. A full Bayesian parameter estimation is required to derive more realistic uncertainties (e.g. Buscicchio et al. 2019; Finch et al. 2022; Katz et al. 2022).

3 RESULTS

3.1 Number of detections

Out of $\sim \mathcal{O}(10^6)$ emitting in the LISA frequency band, we find from a few to several $\sim \mathcal{O}(10^2)$ to be detectable with S/N > 7 assuming 4 yr for the mission duration with 100 per cent duty cycle (however see Amaro Seoane et al. 2022b). Our results are reported in Table B1. All detected LMC DWDs have $f > 1.7$ mHz (orbital period < 20 min), which reflects LISA's selection effects with frequency and distance (e.g. see Fig. 1.1 of Amaro-Seoane et al. 2022a). Fig. 3 illustrates our results. In the top panel it shows the trend for the total number of DWDs emitting in the LISA band and the number of detectable ones with the LMC mass; in the bottom panel it shows the break down of detectable DWDs in different categories (detached, interacting, and supernova type Ia progenitors, SNIa) for a fixed LMC mass of $2.7 \times 10^9 M_\odot$. We find that the number of DWDs emitting in the LISA band primarily scales with the LMC stellar mass, with a little difference between our LMC models. The number of detected DWDs across considered LMC models ranges between 119 and 543 (see Table B1). As we keep DWD synthetic population fixed, differences in the number of detections can be attributed to differences in the LMC stellar

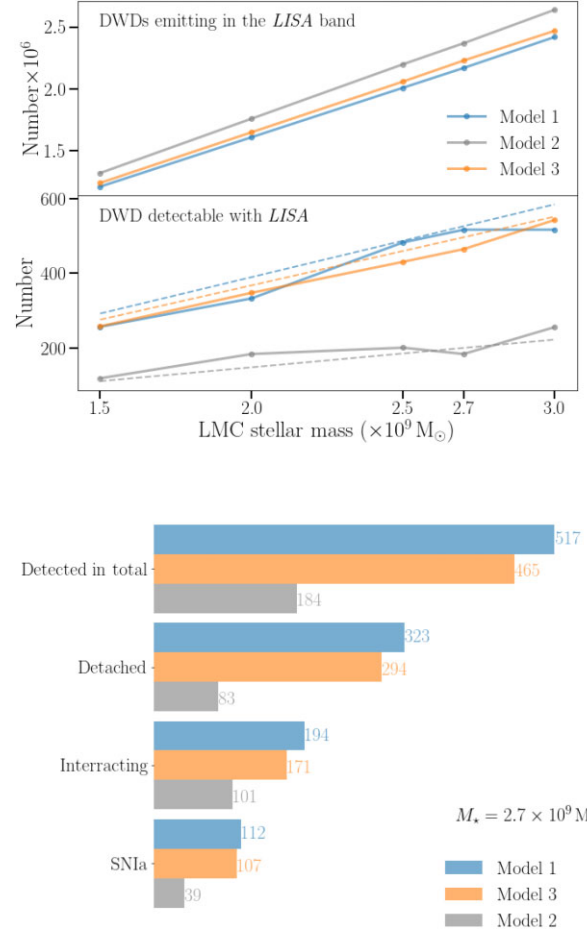


Figure 3. Summary plots for the number of detected LMC DWDs. Top panels: solid lines represent the trend for the total DWD number emitting in the LISA band and those detectable with LISA within 4 yr of mission; dashed lines show linear regression fit [cf. equation (12)]. Bottom panel: the break down of detected DWDs in different categories for a fixed LMC mass of $2.7 \times 10^9 M_\odot$.

mass and/or star formation history and/or spatial distribution within the LMC. We next discuss the relative importance of these model ingredients.

It is immediately evident from Fig. 3 that the number of DWDs (total in the LISA band and detected) scale linearly with the assumed LMC stellar mass. For a fixed LMC mass, by comparing Models 1 and 3 (i.e. models with the same star formation history but different spatial distributions), we can deduce that the spatial distribution alone has a minor impact of less than 10 per cent. This is because the LMC's size is still fairly small compared to its distance so that re-distribution does not cause substantial differences. The largest variation (up to ~ 60 per cent) in total number of detected binaries is due to different underlying star formation history. This is evident by comparing Model 2 and Model 3.

Our results reveal that up to 50–200 LISA detectable DWDs may be in an interacting phase, with possible electromagnetic counterpart signals. Here we flag as ‘interacting’ those binaries in which one of the two white dwarfs overfill its Roche Lobe as defined by (Eggleton 1983)

$$r_{\text{RL}} = \frac{0.49q^{2/3}a}{0.6q^{2/3} + \log(1 + q^{1/3})}, \quad (11)$$

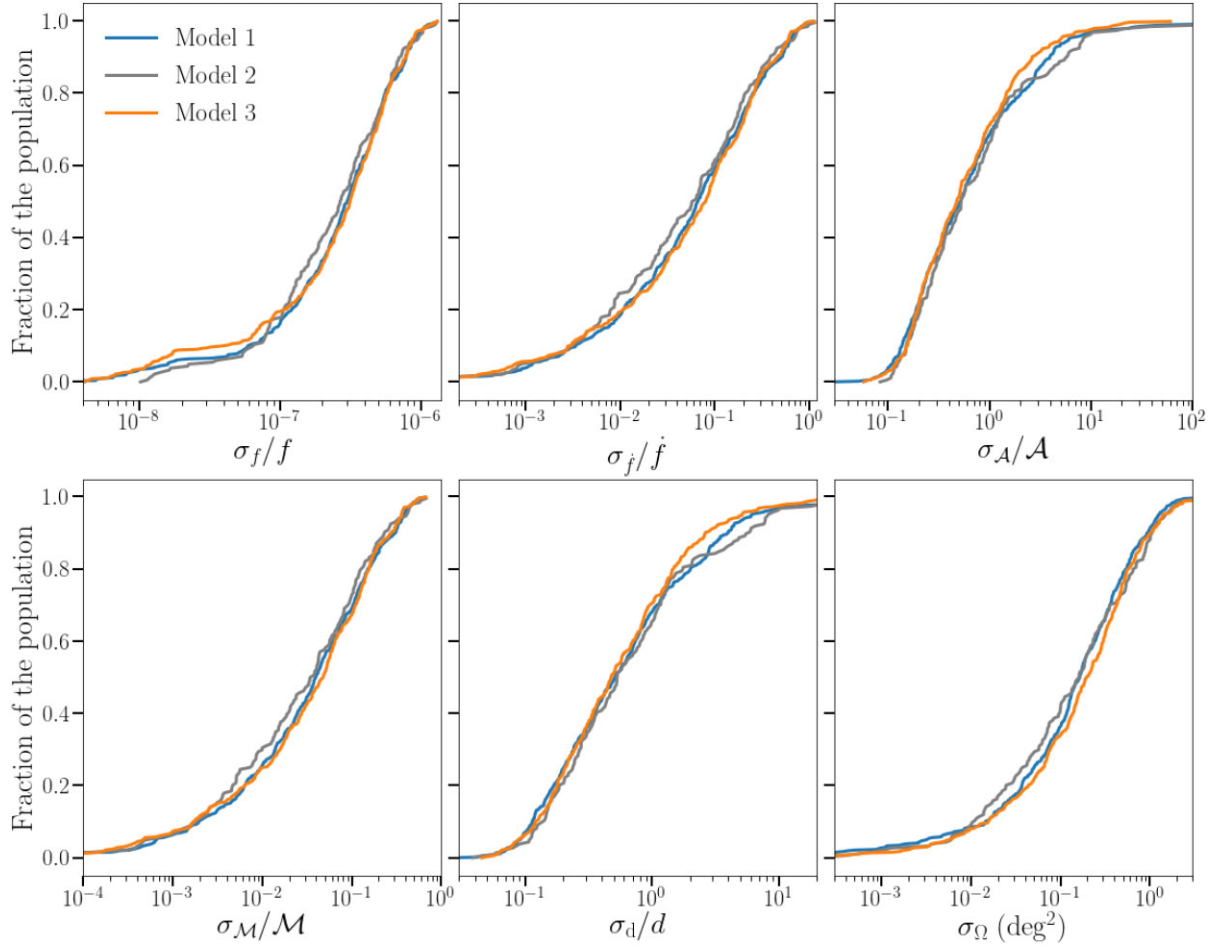


Figure 4. Cumulative distributions of errors for the LISA mission after 4 yr. These errors are estimated with the Fisher information matrix technique as part of the LISA data analysis pipeline of Karnesis et al. (2021).

where $q = m_2/m_1$ is the mass ratio and a is the DWD orbital separation derived from equation (7) using the Kepler’s law. Material extending beyond this tear drop shaped gravitational equipotential is not bound to its white dwarf and mass transfer with the binary partner may occur, potentially impacting the GW signal. From this point onwards, it is yet unclear if the system onsets a stable mass transfer or merges (e.g. Marsh, Nelemans & Steeghs 2004; Shen 2015; Tauris 2018, and Amaro-Seoane et al. 2022a for a review). Thus, we do not attempt to model further the evolution of these systems in the present work.

Importantly, amongst detected DWDs 25–125 (including both detached and interacting) have chirp masses $>0.6 M_\odot$. This chirp mass threshold corresponds to an equal mass binary with a total mass exceeding the Chandrasekhar limit. When such binaries merge, they may result in SNIa explosions (e.g. Webbink 1984), although an alternative outcome may be an accretion-induced collapse resulting in a neutron star (e.g. Nomoto & Iben 1985; Shen et al. 2012).

3.2 Estimated laser interferometer space antenna measurement errors and association with the large magellanic cloud

Fig. 4 gives the cumulative distribution of estimated LISA errors using the FIM technique for frequency, frequency derivative, GW amplitude, chirp mass, distance, and sky position. Since at the LMC distance LISA can only access DWDs with $f > \text{a few mHz}$, frequency, frequency derivative, and GW amplitude across the whole sample

can be measured, which makes it possible to constrain binaries chirp masses and distances [cf. equations (8) and (9)]. We also note that sky location is remarkably well constrained, with half of the sample having localization errors of $\sigma_\Omega < 0.2 \text{ deg}^2$. Using Bayesian techniques for the binary parameter estimation Roebber et al. (2020) found that high frequency DWDs in Milky Way satellites can have small sky localization uncertainty, which improves exponentially with increasing GW frequency (or equivalently period). However, uncertainties in distance measurements have no such dependence. Indeed, half of the sample in Fig. 4 have completely unconstrained distances because of the large errors on GW amplitude. This means that for the majority of LMC DWDs, error bars on the distance will be larger than the LMC physical size.

Detected DWD binaries can be associated with the LMC because of the small errors on the sky position ($< \text{a few deg}^2$) compared to the angular size of the galaxy ($\sim 10 \text{ deg}$) and because of its location in a sparsely populated part of the sky away from the Galactic plane. We quantify the probability of contamination by the (foreground) Milky Way DWDs by following Roebber et al. (2020, see their Section 4). They make use of the Milky Way (disc +bulge) DWD population from Korol et al. (2019) and include a stellar halo generated with a single burst star formation history (SFH), a power law density distribution according to Iorio et al. (2018), and a total mass of $1.4 \times 10^9 M_\odot$ (e.g. Mackereth et al. 2019). They find that this Milky Way foreground model yields $\sim 1 \text{ DWD/deg}^2$ in the sky region

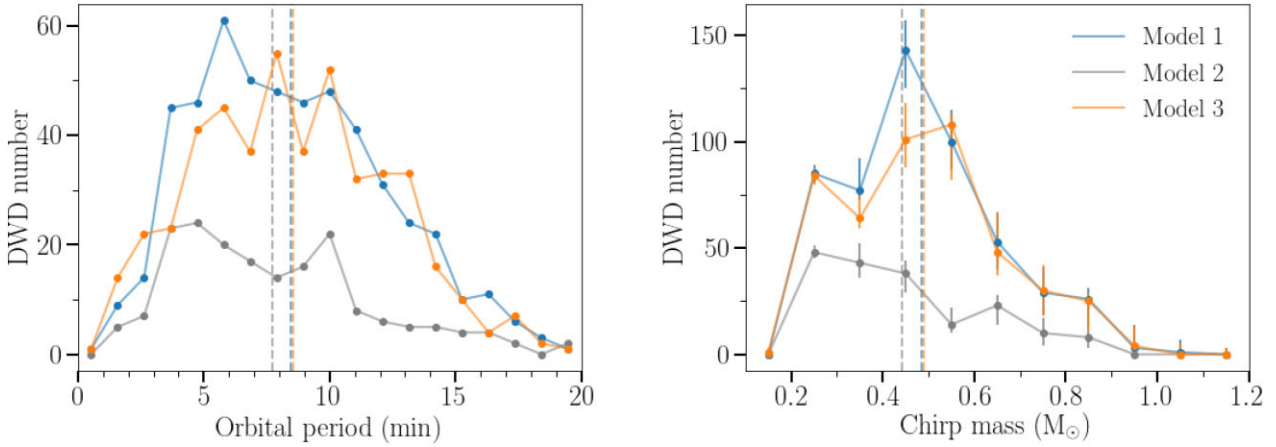


Figure 5. Distributions of orbital periods (left) and chirp masses (right) of LMC DWD with $S/N > 7$. As before we use blue colour for Model 1, grey for Model 2, and orange for Model 3. Error bars represent the minimum and maximum of each histogram bin in 100 realizations wherein periods (chirp masses) are drawn from a normal distribution based on $\sigma_{P,i}$ ($\sigma_{M,i}$) computed with Fisher information matrix. Vertical dashed lines indicate the mean of the respective distributions.

occupied by the LMC. Thus, with this foreground density and an average sky localization error of 0.2 deg^2 , as estimated in this study, we obtain 0.2 contamination foreground sources or a typical false alarm probability of ~ 0.01 . We note that this result is frequency dependent: at lower (higher) frequencies, the sky localization gets worse (better) and as a consequence the false alarm rate increases (decreases). In addition, Roebber et al. (2020) point out that the situation may be further complicated by the fact that the LMC is located at the ecliptic pole, where LISA suffers by a partial degeneracy in the LISA response. This may result in a larger foreground due to the presence of Milky Way DWDs at the other ecliptic pole. Naturally, for binaries with precise distance estimates the associations will be robust.

3.3 Parameter distributions of detectable double white dwarfs

In Fig. 5 we compare the distribution of detectable binary orbital periods and chirp masses for our LMC models. Error bars in Fig. 5 are estimated by taking the minimum and maximum of each histogram bin in 100 realizations, wherein each binary period/chirp mass are drawn from a normal distribution with a mean at the true period and a standard deviation set by the period/chirp mass uncertainty (cf. Section 3.2). We note that the error bars associated with the LMC DWD orbital period distribution (left panel) are very small because GW frequencies are very well constraint with $\sigma_{f/f} < 10^{-6}$ for all binaries. From the comparison, it is clear that distributions of DWD orbital periods and chirp masses are similar in Model 1 and Model 3, whereas Model 2 differ from the other two. The similarities/differences can be attributed to similarity/differences in the star formation histories. Recall that Models 1 and 3 have similar star formation histories but different spatial distributions, whereas Model 2 has a different star formation history and shares the same spatial distribution as Model 3. In Model 2 the majority of DWD is formed within the first 4 Gyr, whereas the majority of DWDs in Models 1 and 3 formed is the last 4 Gyr (cf. Fig. 2). In addition, ~ 100 detectable binaries in Models 1 and 3 became DWD just a few Myr ago – a truly ‘fresh’ stock. This result reflects how period shortening through GW emission is a poor source of LISA-detectable binaries compared to the direct formation at low orbital period. In other words, it is more efficient to freshly populate the LISA band

with newly born milli-Hz binaries than with old binaries that slowly shorten their period enough to be bright source in LISA.

In Fig. 5 we also show the mean values of the respective distributions as dashed vertical lines, which fall at longer DWD orbital periods for Model 1 and Model 3 ($\sim 8.5 \text{ min}$) and larger chirp masses ($\sim 0.49 M_{\odot}$) compared to Model 2 ($\sim 7.7 \text{ min}$ and $\sim 0.44 M_{\odot}$, respectively). In particular, we highlight that even a small difference of $\sim 0.7 \text{ min}$ is well constrained by LISA considering that the largest relative error on the DWD orbital period across all three models is of 0.025 ms . To give a more robust evidence that this lower relative period distribution is statistically meaningful, we apply a two sample Kolmogorov–Smirnov (KS) test which seeks to quantify the differences between two populations and to ascertain if they originated from the same probability distribution. We find that the difference between Model 1/Model 3 and Model 2 distribution is significant at the 0.005 level. Although, the chirp masses for the LMC DWDs will be overall well constrained (with the largest fractional error of 0.3), the average error is comparable to the difference between the means of the chirp mass distributions. Still, the KS test confirms that the chirp mass distribution are different at 0.001 significance level.

3.4 Investigating constraints on the large magellanic cloud stellar mass

Fig. 3 clearly shows that the number of detectable DWDs is linearly dependent on the LMC’s total stellar mass. Indeed, by performing a linear regression we find that our three models are well fit by

$$\begin{aligned} N_{M1} &\approx 195 \left(\frac{M_{\star}}{10^9 M_{\odot}} \right) \\ N_{M2} &\approx 74.2 \left(\frac{M_{\star}}{10^9 M_{\odot}} \right) \\ N_{M3} &\approx 184 \left(\frac{M_{\star}}{10^9 M_{\odot}} \right) \end{aligned} \quad (12)$$

where N_{Mi} is the number of detectable DWD as a function of the LMC stellar mass $M_{\star}$. Note that the intercepts are approximately zero because zero LMC mass yields no DWDs. The coefficient of determination for each model takes on values that are reasonably close to 1 ($\sim 0.97, 0.75$, and 0.97 for Models 1, 2, and 3, respectively)

suggesting the relationship is indeed linear. We overplot these results as dashed lines in Fig. 3.

We can now exploit this linear dependence between the number of LISA detectable DWDs and the LMC mass, and reverse-engineer our modelling to infer what stellar mass is required to generate the ‘observed’ number of DWD with LISA. Using as the study case of Milky Way satellite galaxies, Korol et al. (2021) showed that this can be done in analogy with simple stellar populating models widely used for estimating the stellar mass of galaxies based on their total light (e.g. Tinsley & Gunn 1976; Bruzual 1983; Maraston 1998). Using the same fiducial binary population synthesis model as in our work, they constructed a Bayesian inference scheme considering the number of detected (detached) DWDs associated with the satellite and the measured distance to the satellite as the only two input parameters. The authors demonstrated that counting the number of LISA sources in a satellite provides a measurement of its stellar mass. Even when the number of detections in a (known) satellite is zero, it is still possible to place an upper limit on the satellites mass using the same argument. We highlight that such an inference scheme recovers the ‘true’ LMC mass and gives smaller errors when the star formation history used for the inferences coincides with the ‘true’ one. Here, we employ the same inference method as in Korol et al. 2021 using a constant star formation history, which is different from star formation histories considered here to model the LMC. Korol et al. 2020 predict that a constant star formation history for a total stellar mass of $2.7 \times 10^9 M_\odot$ would yield ~ 235 LISA detectable DWDs. When comparing DWD formation histories in Fig. 2 and our results in Table B1, not surprisingly this number falls in between that based on Model 2 (with no recent star formation) and that of Models 1 and 3 (with active star formation at recent times). We apply a Gaussian prior on DWD distances ($\mu = 49.21$ kpc and $\sigma = 35$ kpc), a flat prior for the LMC mass in the range $[10^6 - 10^{10}] M_\odot$ and a flat prior for the LMC age in the range $[1 - 10]$ Gyr. We plug in the number of detected binaries estimated with Models 1, 2, and 3 for the true LMC mass of $2.7 \times 10^9 M_\odot$ (cf. Table B1).

We recover the LMC mass of $3.01^{+0.4}_{-1.1} \times 10^9 M_\odot$ based on the number of detections generated by Model 1, $0.7^{+0.19}_{-0.25} \times 10^9 M_\odot$ based on Model 2, and $2.7^{+0.39}_{-1.0} \times 10^9 M_\odot$ based on Model 3. In Fig. 6 we show the obtained posterior distribution for the LMC stellar mass, whereas the LMC age remains an unconstrained parameter. Fig. 6 reveals that even using ‘incorrect’ constant star formation history assumption, which is close enough to the ‘true’ one as for the case of Model 1 and Model 3, we can recover the LMC mass within 1σ .

4 DISCUSSION AND CONCLUSIONS

In this work we assembled a realistic population of DWD binaries for a range of LMC stellar masses. We used a fiducial DWD synthetic model from Korol et al. (2020) based on binary population synthesis simulations of Toonen et al. (2012) computed with the SEBA code. These binaries were injected in three different LMC models conceived to test the effect of different star formation models and spatial distributions of stars on the size and properties of LISA detectable DWD population. In Model 1 we assigned synthetic DWD binaries ages and positions based on the completeness-corrected regional star formation histories of Harris & Zaritsky (2009). This model provides an observation based star formation history, but because it uses sky-projected 2D stellar density distribution it lacks small-scale LMC’s distinct features. To account for this later deficiency and to probe whether LISA can spatially resolve the LMC structure, in Model 2 we employ the N-body hydrodynamical simulation of Lucchini et al. (2020, see bottom panels of Fig. 1). In addition to the spatial

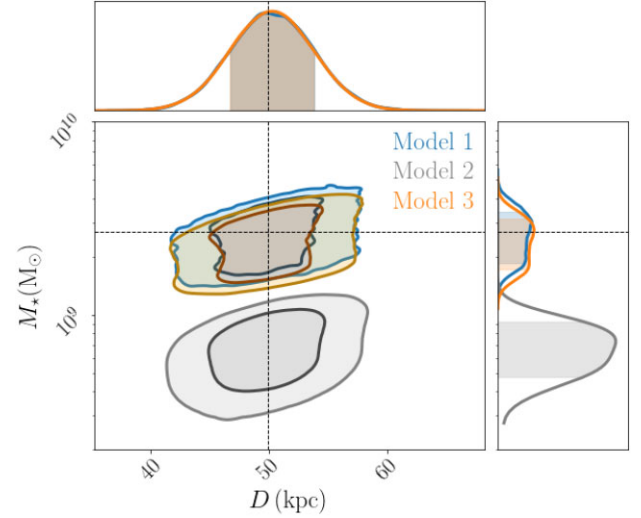


Figure 6. Posteriors for the LMC total stellar mass and distance for our three models: Model 1 (blue), Model 2 (grey), and Model 3 (orange). The assumed ‘true’ LMC mass and distance are $2 \times 10^9 M_\odot$ and 49.21 kpc (dashed black lines). We note that our Model 1 yields essentially the same contours as Model 3 – given that these models share the same star formation history.

distribution of the simulation particles we also make use of their mass and ages, which provides us with an alternative (simulation driven) star formation history to that of Model 1. In Model 3 we combined the most realistic features of the previous two models: observation based star formation history from Model 1 and spatial distribution of stars from Model 2. For all three models we then compute how many DWD could be detected by LISA over a 4 yr mission and how well we could constrain their parameters. Our main findings can be summarized as follows.

LISA will be able to detect from a hundred to several hundred DWDs in the LMC. These include both detached and accreting systems. In this study we do not provide a detailed modelling of the accreting DWD systems; thus, we cannot confidently state if they establish a stable of mass transfer or merge. Even when we consider detached DWDs only, our most realistic LMC model (Model 3) yields from 159 to 357 LISA detections. This means that the LMC will stand out as an extended and spatially resolved structure on the LISA’s GW sky.

Importantly, LISA will provide scores of double degenerate SNIa progenitors, extremely challenging to find with electromagnetic telescopes. Our estimates range from 25 to 125 detectable SNIa progenitors. To quote these numbers we arbitrarily define as SNIa progenitor a DWD with $\mathcal{M} > 0.6 M_\odot$, which corresponds to an equal-mass binary exceeding the Chandrasekhar mass limit. Although our SNIa progenitor definition is oversimplified, it suggests that LISA will play a crucial role in understanding SNIa double-degenerate formation channel. SNIas are useful as standard candles to measure cosmological distances (Riess et al. 1998; Perlmutter et al. 1999), and therefore a number of works have explored the role of DWDs in total SNIa merger rates (Badenes & Maoz 2012; Maoz, Badenes & Bickerton 2012; Toonen et al. 2012; Claeys et al. 2014; Maoz, Hallakoun & Badenes 2018; Toonen, Perets & Hamers 2018) as opposed to other SNIa formation channels such as those reviewed in Livio & Mazzali (2018). While identifying double degenerate SNIa progenitors can help in determining their contribution to the SNIa merger rate, no previously observed stellar system have been robustly confirmed to be such a progenitor (Rebassa-Mansergas et al. 2019).

Our results show that LISA will be fundamental in discovering SNIa progenitors not only within our own Galaxy, but also in a completely different lower metallicity and actively star-forming environment as that of the LMC.

We demonstrate that while the number of LISA detections is linearly dependent on the LMC total stellar mass, it is also sensitive to recent star formation activity (cf. Fig. 3). This can be deduced when comparing Model 1 and Model 3 (with recent star formation) to Model 2 (with most of DWDs being formed at early times). We find that Model 1 and 3 generate 1.8–2.8 times more LISA detectable DWDs than Model 2. In particular, a large fraction of this difference comes from binaries formed $\sim \mathcal{O}(10^2)$ Myr ago, which are ~ 10 times more numerous in Models 1 and 3 compared to Model 2 and became DWDs just in the last few Myr. To sum up, these results show that the total stellar mass sets the total number of DWDs in the LMC, whereas the star formation history (in combination with the age) determines how many detectable DWDs emit in the LISA frequency band at the present time.

Besides the total number of binaries, the difference is also noticeable when comparing DWD period and chirp mass distributions (cf. Fig. 5). We show that DWDs in Model 1 and Model 3 on average have ~ 1 min shorter orbital periods and $\sim 0.05 M_\odot$ higher chirp masses than DWDs in Model 2. This highlights how period shortening through GW emission is a less efficient way of forming LISA detectable binaries compared to the formation of binaries directly in the LISA band. The model comparison also reveals that the numbers of LISA detectable DWDs do not significantly change with the LMC spatial distribution. Further analysis may narrow down what can be learned about the LMC’s morphology and characterize the GW confusion noise from unresolvable LMC DWDs. We note that the impacts of both stellar mass and the current star formation rate on the number of detectable binaries could be disentangled by considering the period distribution of detected binaries. One could imagine quantitatively relating the number of detectable binaries, the median period of detectable binaries, the stellar mass of the LMC, and the current star formation rate. However, such a relation is out of the scope of this work since our models only considered high star formation activity and near complete inactivity, and our approach includes only period distributions from a limited number of simulations. Nevertheless, our findings suggests that the period and chirp mass distributions and number of LISA detectable DWD binaries may probe both LMC stellar mass and recent star formation, motivating more careful study.

While we explore different masses and star formation histories, our methodology relies on a fiducial DWD population model. Though since the DWD formation we only consider energy loss through GWs, other processes such as tidal interactions may alter the size and the properties of the LISA sample (e.g. Biscoveanu, Kremer & Thrane 2022). As we have already mentioned, we do not provide detailed modelling for the accreting DWD, but only flag as accreting binaries that overflow their Roche Lobe at the time they are observed with LISA. It is yet unclear what is the most likely fate for these binaries, and so more development is needed from both theory and observations to make reliable forecasts for LISA (e.g. Amaro-Seoane et al. 2022a). However, we chose to not remove these binaries from our simulations so that these can be followed up in the future studies. We also keep some of the key quantities entering in the binary evolution such as the metallicity, common envelope model and its efficiency, and binary fraction. Earlier work of Korol et al. (2020) showed these have minor effect compared to the assumptions on the properties of the satellite galaxy, which we extensively explore here.

Given the number and well constrained sky locations of detectable binaries, future structural analysis of the LMC using our population is highly promising. Additionally, while this work focuses on detectable binaries, far more will be undetectable and contribute to a background confusion noise. Future efforts may try to constrain this noise in case it is relevant for the LISA sensitivity curve (as in Cornish & Robson 2017) or useful for probing structure (as in Benacquista & Holley-Bockelmann 2006). Similarly, the past observed and simulated studies our research depended on have also been conducted for the SMC (Harris & Zaritsky 2004; Lucchini et al. 2020) making similar analysis straightforward. Even so, the SMC is estimated to contain a factor of $\lesssim 5$ fewer detectable DWD binaries compared to the LMC (Korol et al. 2020) making the utility of gravitational probes less optimistic. Finally, even though the LMC is predicted to host only ~ 2 – 5 detectable double neutron star binaries (Seto 2019; Lau et al. 2020), their inclusion together with other types of compact object binaries in our work could only improve LISA’s ability to extract LMC features.

ACKNOWLEDGEMENTS

This work would not have been possible without Silvia Toonen, who provided us with a mock DWD population, and Scott Lucchini, who provided us with the LMC simulation. The authors thank Nikolaos Karnesis for allowing us to use the data analysis pipeline outlined in Section 2.4. The code is publicly available in the following GitHub repository: https://gitlab.in2p3.fr/Nikos/galactic_properties_from_cgbs.

The authors also thank Jeffrey Pflueger, Fraser Evans, Stella Reino, and Pieter van Dokkum for useful conversations on implementing our methods and reporting results.

VK acknowledges support from the Netherlands Research Council NWO (Rubicon 019.183EN.015 grant). EMR acknowledges that this project has received funding from the European Research Council (ERC) under the European Union’s Horizon 2020 research and innovation programme (Grant agreement No. 101002511—VEGA P).

This research made use of the open source python package PyGaia developer by Anthony Brown, publicly available at <https://github.com/agabrown/PyGaia>, and tools provided by the LISA Data Processing Group (LDPG) and the LISA Consortium LISA Data Challenges (LDC) working group <https://lisa-ldc.lal.in2p3.fr/>.

DATA AVAILABILITY

Our fiducial LMC DWD population catalogue, which specifies the frequency, chirp, ecliptic latitude, ecliptic longitude, amplitude, inclination, polarization, orbital phase, age, and white dwarf masses of present day LISA band LMC DWD binaries are available on Zenodo and can be accessed via 10.5281/zenodo.6918083. This public release represents a $2.7 \times 10^9 M_\odot$ stellar mass LMC since we find the estimate of van der Marel et al. (2002) to be the most reliable of those considered here (being based on Zaritsky 1999’s analysis of Zaritsky, Harris & Thompson 1997 which is a more recent observational study than the work referenced in de Vaucouleurs et al. 1991 as was considered by McConnachie 2012). This public release excludes Model 2 since its star formation history does not reflect observation. Other population models are currently available upon request. We request that researchers utilising any of these populations cite this work.

REFERENCES

- Adams M. R., Cornish N. J., Littenberg T. B., 2012, *Phys. Rev. D*, 86, 124032
- Amaro-Seoane P. et al., 2017, preprint (arXiv:1702.00786)
- Amaro-Seoane P. et al., 2022a, preprint (arXiv:2203.06016)
- Amaro Seoane P. et al., 2022b, *Gen. Relativ. Gravitation*, 54, 335221342
- Andrews J. J., Breivik K., Pankow C., D’Orazio D. J., Safarzadeh M., 2020, *ApJ*, 892, L9
- Badenes C., Maoz D., 2012, *ApJ*, 749, L11
- Badenes C. et al., 2018, *ApJ*, 854, 147
- Belokurov V. A., Erkal D., 2019, *MNRAS*, 482, L9
- Benacquista M., Holley-Bockelmann K., 2006, *ApJ*, 645, 589
- Besla G., Kallivayalil N., Hernquist L., van der Marel R. P., Cox T. J., Kereš D., 2012, *MNRAS*, 421, 2109
- Besla G., Martínez-Delgado D., van der Marel R. P., Beletsky Y., Seibert M., Schlafly E. F., Grebel E. K., Neyer F., 2016, *ApJ*, 825, 20
- Biscoveanu S., Kremer K., Thrane E., 2022, preprint (arXiv:2206.15390)
- Brüns C. et al., 2005, *A&A*, 432, 45
- Bruzual A. G., 1983, *ApJ*, 273, 105
- Buscicchio R., Roebber E., Goldstein J. M., Moore C. J., 2019, *Phys. Rev. D*, 100, 084041
- Butcher H., 1977, *ApJ*, 216, 372
- Choi Y. et al., 2018, *ApJ*, 866, 90
- Claeys J. S. W., Pols O. R., Izzard R. G., Vink J., Verbunt F. W. M., 2014, *A&A*, 563, A83
- Cornish N., Robson T., 2017, in *J. Phys. Conf. Ser.* p. 012024, preprint (arXiv:1703.09858)
- Crowder J., Cornish N. J., 2007, *Phys. Rev. D*, 75, 043008
- Cullinan L. R. et al., 2020, *MNRAS*, 497, 3055
- Cutler C., 1998, *Phys. Rev. D*, 57, 7089
- De Leo M., Carrera R., Noël N. E. D., Read J. I., Erkal D., Gallart C., 2020, *MNRAS*, 495, 98
- Eggleton P. P., 1983, *ApJ*, 268, 368
- El-Badry K., Rix H.-W., 2019, *MNRAS*, 482, L139
- Erkal D., Belokurov V. A., 2020, *MNRAS*, 495, 2554
- Erkal D. et al., 2019, *MNRAS*, 487, 2685
- Finch E. et al., 2022, *MNRAS*, preprint (arXiv:2210.10812)
- Gardiner L. T., Noguchi M., 1996, *MNRAS*, 278, 191
- Georgousi M., Karnesis N., Korol V., Pieroni M., Stergioulas N., 2022, *MNRAS*, 519, 2552
- Harris J., Zaritsky D., 2004, *AJ*, 127, 1531
- Harris J., Zaritsky D., 2009, *AJ*, 138, 1243
- Haschke R., Grebel E. K., Duffau S., 2012, *AJ*, 144, 106
- Hernquist L., 1990, *ApJ*, 356, 359
- Indu G., Subramaniam A., 2011, *A&A*, 535, A115
- Iorio G., Belokurov V., Erkal D., Koposov S. E., Nipoti C., Fraternali F., 2018, *MNRAS*, 474, 2142
- Jacyszyn-Dobrzaniecka A. M. et al., 2016, *AcA*, 66, 149
- Kallivayalil N., van der Marel R. P., Besla G., Anderson J., Alcock C., 2013, *ApJ*, 764, 161
- Karnesis N., Babak S., Pieroni M., Cornish N., Littenberg T., 2021, *Phys. Rev. D*, 104, 043019
- Katz M. L., Danielski C., Karnesis N., Korol V., Tamanini N., Cornish N. J., Littenberg T. B., 2022, *MNRAS*, 517, 697
- Kim S., Staveley-Smith L., Dopita M. A., Freeman K. C., Sault R. J., Kesteven M. J., McConnell D., 1998, *ApJ*, 503, 674
- Korol V., Koop O., Rossi E. M., 2018, *ApJ*, 866, L20
- Korol V., Rossi E. M., Barausse E., 2019, *MNRAS*, 483, 5518
- Korol V. et al., 2020, *A&A*, 638, A153
- Korol V., Belokurov V., Moore C. J., Toonen S., 2021, *MNRAS*, 502, L55
- LISA Science Study Team, 2018, Technical Report ESA-L3-EST-SCI-RS-001, LISA Science Requirements Document, www.cosmos.esa.int/web/lisa/lisa-documents/, ESA, www.cosmos.esa.int/web/lisa/lisa-documents/
- Laporte C. F. P., Gómez F. A., Besla G., Johnston K. V., Garavito-Camargo N., 2018, *MNRAS*, 473, 1218
- Lau M. Y. M., Mandel I., Vigna-Gómez A., Neijssel C. J., Stevenson S., Sesana A., 2020, *MNRAS*, 492, 3061
- Livio M., Mazzali P., 2018, *Phys. Rep.*, 736, 1
- Lucchini S., D’Onghia E., Fox A. J., Bustard C., Bland-Hawthorn J., Zweibel E., 2020, *Nature*, 585, 203
- Luo J. et al., 2016, *Classical Quantum Gravity*, 33, 035010
- Mackereth J. T. et al., 2019, *MNRAS*, 489, 176
- Mackey D., Koposov S., Da Costa G., Belokurov V., Erkal D., Kuzma P., 2018, *ApJ*, 858, L21
- Maggiore M., 2008, *Gravitational Waves. Volume 1: Theory and Experiments*. Oxford University Press, New York
- Maoz D., Badenes C., Bickerton S. J., 2012, *ApJ*, 751, 143
- Maoz D., Hallakoun N., Badenes C., 2018, *MNRAS*, 476, 2584
- Maraston C., 1998, *MNRAS*, 300, 872
- Marsh T. R., Nelemans G., Steeghs D., 2004, *MNRAS*, 350, 113
- McConnachie A. W., 2012, *AJ*, 144, 4
- Moe M., Kratter K. M., Badenes C., 2019, *ApJ*, 875, 61
- Nelemans G., Yungelson L. R., Portegies Zwart S. F., Verbunt F., 2001, *A&A*, 365, 491
- Nissanke S., Vallisneri M., Nelemans G., Prince T. A., 2012, *ApJ*, 758, 131
- Nomoto K., Iben I. J., 1985, *ApJ*, 297, 531
- Olsen K. A. G., Zaritsky D., Blum R. D., Boyer M. L., Gordon K. D., 2011, *ApJ*, 737, 29
- Peñarrubia J., Gómez F. A., Besla G., Erkal D., Ma Y.-Z., 2016, *MNRAS*, 456, L54
- Perlmutter S. et al., 1999, *ApJ*, 517, 565
- Pietrzyński G., Udalski A., 2000, *AcA*, 50, 355
- Pietrzyński G. et al., 2013, *Nature*, 495, 76
- Portegies Zwart S. F., Verbunt F., 1996, *A&A*, 309, 179
- Rebassa-Mansergas A., Toonen S., Korol V., Torres S., 2019, *MNRAS*, 482, 3656
- Riess A. G. et al., 1998, *AJ*, 116, 1009
- Roebber E. et al., 2020, *ApJ*, 894, L15
- Ruan W.-H., Guo Z.-K., Cai R.-G., Zhang Y.-Z., 2018, *Int. J. Mod. Phys. A*, 35, 17
- Ruiz-Lara T., 2019, in *A Synoptic View of the Magellanic Clouds: VMC, Gaia and Beyond*. p. 20
- Schommer R. A., Suntzeff N. B., Olszewski E. W., Harris H. C., 1992, *AJ*, 103, 447
- Seto N., 2019, *MNRAS*, 489, 4513
- Shao S., Cautun M., Deason A. J., Frenk C. S., Theuns T., 2018, *MNRAS*, 479, 284
- Shen K. J., 2015, *ApJ*, 805, L6
- Shen K. J., Bildsten L., Kasen D., Quataert E., 2012, *ApJ*, 748, 35
- Storck A., Church R., 2022, *MNRAS*, preprint (arXiv:2205.01507)
- Subramaniam A., 2003, *ApJ*, 598, L19
- Tauris T. M., 2018, *Phys. Rev. Lett.*, 121, 131105
- Timpano S. E., Rubbo L. J., Cornish N. J., 2006, *Phys. Rev. D*, 73, 122001
- Tinsley B. M., Gunn J. E., 1976, *ApJ*, 203, 52
- Toonen S., Nelemans G., Portegies Zwart S., 2012, *A&A*, 546, A70
- Toonen S., Perets H. B., Hamers A. S., 2018, *A&A*, 610, A22
- Wan Z., Guglielmo M., Lewis G. F., Mackey D., Ibata R. A., 2020, *MNRAS*, 492, 782
- Webbink R. F., 1984, *ApJ*, 277, 355
- Weinberg M. D., Nikolaev S., 2001, *ApJ*, 548, 712
- Wilhelm M. J. C., Korol V., Rossi E. M., D’Onghia E., 2021, *MNRAS*, 500, 4958
- Yoshizawa A. M., Noguchi M., 2003, *MNRAS*, 339, 1135
- Zaritsky D., 1999, *AJ*, 118, 2824
- Zaritsky D., Harris J., Thompson I., 1997, *AJ*, 114, 1002
- Zaritsky D., Harris J., Thompson I. B., Grebel E. K., 2004, *AJ*, 128, 1606
- de Grijs R., Wicker J. E., Bono G., 2014, *AJ*, 147, 122
- de Vaucouleurs G., de Vaucouleurs A., Corwin Herold G. J., Buta R. J., Paturel G., Fouque P., 1991, *Third Reference Catalogue of Bright Galaxies*
- van der Marel R. P., Kallivayalil N., 2014, *ApJ*, 781, 121
- van der Marel R. P., Alves D. R., Hardy E., Suntzeff N. B., 2002, *AJ*, 124, 2639

APPENDIX A: MODEL STAR FORMATION RATES

In Fig. A1, we show the total effective star formation history of each model assuming the LMC stellar mass of van der Marel et al. (2002). This is calculated by breaking up our simulation into logarithmically spaced time steps and dividing the mass of MS+MS binaries formed in each time step, totalled across the LMC, by the length of each time step. We see that Model 1's total star formation history agrees with the measurements of Harris & Zaritsky (2009) once they are totalled across the LMC and re-normalized to reflect only the star formation rate of our DWD progenitors. Model 2 has little star formation in the last 6 Gyr, corresponding to the simulated stellar particle's large ages. Model 3's bursty star formation is highly reminiscent of Model 1, though slight differences occur since the matching procedure makes different LMC regions play a larger, or smaller, role in the total star formation history compared to Model 1.

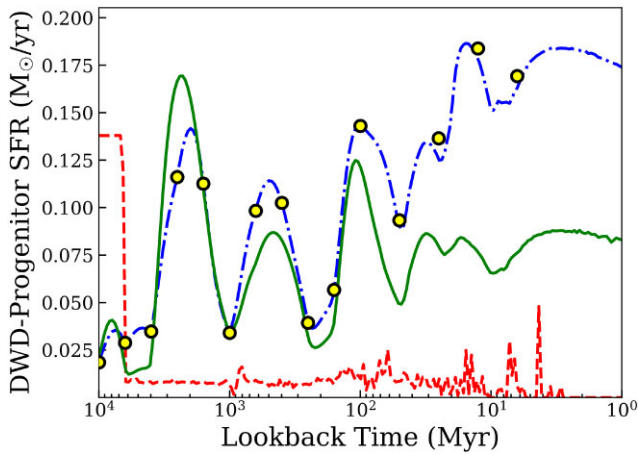


Figure A1. The total effective star formation history of MSMS binaries destined to become DWD binaries for Model 1 (the observation driven model; blue, dash-dotted), Model 2 (the simulated model; red, dashed), and Model 3 (the combined model; green, solid), along with the measurements of from Harris & Zaritsky (2009, yellow, circles) totalled across all regions and re-normalized by the global constant used in Model 1. While the x -axis is given in log scale, integrating over each star formation history gives the same MSMS population total mass. Note that this graph is given in logarithmic look-back time as is probed by observation, rather than linear LMC time as given in Fig. 2.

APPENDIX B: NUMBER OF LISA DETECTIONS

Table B1 shows our results for the total number of binaries in the LISA band ($N_{\text{DWD}}^{\text{LISA}}$) and binaries detected with $S/N > 7$ (N_{DWD}^d). N_{DWD}^d further breaks down into detached systems (N_{DDWD}^d), whereas the number of detectable accreting DWDs can be derived by subtracting these from the total. In the last column we highlight a sub-sample of LISA-detectable DWD that are potential double degenerate SNIa progenitors with $\mathcal{M} > 0.6 M_{\odot}$. We predict a few to several hundreds of individually resolved DWD binaries in the LMC, that correspond to a detection efficiency in the LISA band of $N_{\text{DWD}}^d / N_{\text{DWD}}^{\text{LISA}} \sim 10^{-4}$. Therefore, for each resolved binary there are 10^4 low signal-to-noise ones that can contribute to a –spatially localised– background signal. The detected ones are between 5 percent and 10 percent of the $N_{\text{DWD}}^{<20\text{min}}$ populations in the LISA band.

Table B1. The impact of stellar mass and our different models for age and position assignment on the population of DWD binaries with a LISA $S/N > 7$ after the nominal 4 yr mission length.

Mass ($M_{\odot}$)	Model	$N_{\text{DWD}}^{\text{LISA}}$	N_{DWD}^d	N_{DDWD}^d	N_{SNIa}^d
1.5×10^9	1	1.21×10^6	257	161	71
1.5×10^9	2	1.32×10^6	119	70	25
1.5×10^9	3	1.24×10^6	258	159	68
2.0×10^9	1	1.61×10^6	333	213	90
2.0×10^9	2	1.76×10^6	184	92	39
2.0×10^9	3	1.65×10^6	348	218	77
2.5×10^9	1	2.01×10^6	483	303	93
2.5×10^9	2	2.20×10^6	201	105	39
2.5×10^9	3	2.06×10^6	431	282	111
2.7×10^9	1	2.17×10^6	517	323	112
2.7×10^9	2	2.37×10^6	184	83	41
2.7×10^9	3	2.23×10^6	465	294	107
3.0×10^9	1	2.42×10^6	517	318	111
3.0×10^9	2	2.64×10^6	256	128	39
3.0×10^9	3	2.47×10^6	543	357	125

Note. Column 1: total LMC stellar mass. 2: distribution model (1 $\equiv$ observation driven, 2 $\equiv$ simulation driven, 3 $\equiv$ combined). 3: number of present day DWD binaries in the LISA-band with $f > 10^{-4}$ Hz. 4: number of LISA-detectable DWD binaries as calculated from the MLDC pipeline. 5: number of LISA-detectable DWD binaries that are detached, engaging in negligible mass transfer. 6: number of LISA detectable DWDs that are potential double degenerate SNIa progenitors with $\mathcal{M} > 0.6 M_{\odot}$.

This paper has been typeset from a $\text{\LaTeX}$ file prepared by the author.