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Investigating the impact of reactions of C and CH with molecular hydrogen on a glycine gas-grain network

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ABSTRACT

The impact of including the reactions of C and CH with molecular hydrogen in a gas-grain network is assessed via a sensitivity analysis. To this end, we vary three parameters, namely, the efficiency for the reaction $C + H_2 \longrightarrow CH_2$, and the cosmic ray ionization rate, with the third parameter being the final density of the collapsing dark cloud. A grid of 12 models is run to investigate the effect of all parameters on the final molecular abundances of the chemical network. We find that including reactions with molecular hydrogen alters the hydrogen economy of the network; since some species are hydrogenated by molecular hydrogen, atomic hydrogen is freed up. The abundances of simple molecules produced from hydrogenation, such as CH₄, CH₃OH, and NH₃, increase, and at the same time, more complex species such as glycine and its precursors see a significant decrease in their final abundances. We find that the precursors of glycine are being preferentially hydrogenated, and therefore glycine itself is produced less efficiently.

Key words: astrochemistry – ISM: abundances – ISM: molecules.

1 INTRODUCTION

Interstellar dust plays a significant role in the rich chemistry that takes place in the interstellar medium. It is widely believed that complex-organic molecules (COMs) form on interstellar dust (Herbst & van Dishoeck 2009; Caselli & Ceccarelli 2012) since for certain molecules, grain-surface reactions are more efficient than gas-phase reactions. This is particularly important in cold astronomical environments where some gas-phase reactions may be highly inefficient, because a 'third body' is needed to take up the excess heat of an exothermic reaction. Dust grains thus act as an energy sink allowing the chemistry to thrive and this can lead to the formation of more complex organic molecules.

Both experimental work and modelling has shown that one such molecules, namely the amino acid glycine can be formed through energetic processing of the ices during the warm-up phase of star formation (Bernstein et al. 2002; Woon 2002; Bossa et al. 2009; Lee et al. 2009; Ciesla & Sandford 2012; Garrod 2013; Sato et al. 2018), although there is evidence to suggest that glycine would undergo destruction under increased irradiation (Pernet et al. 2013; Maté et al. 2015). In addition, in a joint experimental and modeling effort, Ioppolo et al. (2021) suggested that non-energetic mechanisms such as atom-addition reactions might be a promising route for glycine formation.

A new grain-surface reaction, inserting C atoms in H₂ to form CH_2 via $C + H_2 \longrightarrow CH_2$, was recently proposed to be barrierless by Simončič et al. (2020), based on earlier lab work by Krasnokutski

et al. (2016). They included this reaction in their network and found a far more rapid conversion of C-CH₄. Subsequently, Lamberts et al. (2022) performed a combined experimental and computational work to investigate the importance of reactions with molecular hydrogen for the formation of methane. It was found that while the former reaction might not be fully barrierless, and the barrier likely depends on the binding site, the reaction $CH + H_2 \longrightarrow CH_3$ does in fact proceed without a barrier. The reason these 'dihydrogenation' reactions might be of interest is that they make H₂ more chemically active, the importance of which was recognized already by Hasegawa & Herbst (1993) and by Meisner, Lamberts & Kästner (2017) in the context of water formation. Typically, H2 has one of the lowest binding energies of grain-surface species, lower than even atomic H (Al-Halabi & van Dishoeck 2007; Wakelam et al. 2017; Molpeceres & Kästner 2020), which allows the molecule to diffuse readily on the surface. Moreover, the molecular hydrogen abundance in molecular clouds and pre-stellar cores is much higher than that of atomic hydrogen (van Dishoeck & Black 1988; Goldsmith & Li 2005).

By including these reactions in chemical models, one might first of all expect changes in the CH₄ abundance, but it is equally interesting to consider the effect on downstream species such as complex organic molecules, whose typical abundances are far lower. Their sensitivity to new reactions should be considered, as their more abundant precursors might see changes in their abundances.

In this work, we look to build on the work by Simončič et al. (2020) and Lamberts et al. (2022) to investigate the impact of the dihydrogenation reactions of C and CH on our gas-grain chemical network. In particular, we are interested in observing the effect these reactions have on the production of glycine and its precursors. Our glycine network is based on the kinetic Monte Carlo network used

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Table 1. Table of the reactions added to the standard UCLCHEM network.

Reaction No.	Reaction	Reference
1	$CO + OH \longrightarrow HOCO$	Arasa et al. (2013)
2	$HOCO + H \longrightarrow H_2 + CO_2$	Goumans, Uppal & Brown (2008)
3	$HOCO + H \longrightarrow HCOOH$	Goumans et al. (2008), Ioppolo et al. (2011)
4	$CH_4 + OH \longrightarrow CH_3 + H_2O$	Lamberts et al. (2017)
5	$NH_2 + CH_3 \longrightarrow NH_2CH_3$	Ioppolo et al. (2021)
6	$NH_3 + CH \longrightarrow NH2CH_2$	Balucani et al. (2009)
7	$NH2CH_2 + H \longrightarrow NH_2CH_3$	Ioppolo et al. (2021)
8	$NH_2CH_3 + H \longrightarrow NH2CH_2 + H_2$	Oba et al. (2014)
9	$NH_2CH_3 + OH \longrightarrow NH2CH_2 + H_2O$	Ioppolo et al. (2021)
10	$NH2CH_2 + HOCO \longrightarrow NH_2CH_2COOH$	Woon (2002)
11	$H_2 + OH \longrightarrow H_2O + H$	Meisner et al. (2017)
12	$O_2 + H \longrightarrow HO_2$	Lamberts et al. (2013)
13	$HO_2 + H \longrightarrow OH + OH$	Lamberts et al. (2013)
14	$HO_2 + H \longrightarrow H_2 + O_2$	Lamberts et al. (2013)
15	$HO_2 + H \longrightarrow H_2O + O$	Lamberts et al. (2013)
16	$OH + OH \longrightarrow H_2O_2$	Lamberts et al. (2013)
17	$OH + OH \longrightarrow H_2O + O$	Lamberts et al. (2013)
18	$H_2O_2 + H \longrightarrow H_2O + OH$	Lamberts & Kästner (2017)
19	$N + O \longrightarrow NO$	Ioppolo et al. (2021)
20	$NO + H \longrightarrow HNO$	Fedoseev et al. (2012)
21	$HNO + H \longrightarrow H_2NO$	Fedoseev et al. (2012)
22	$HNO + H \longrightarrow NO + H_2$	Fedoseev et al. (2012)
23	$HNO + O \longrightarrow NO + OH$	Ioppolo et al. (2021)
24	$HN + O \longrightarrow HNO$	Ioppolo et al. (2021)
25	$N + NH \longrightarrow N_2$	Ioppolo et al. (2021)
26	$NH + NH \longrightarrow N_2 + H_2$	Ioppolo et al. (2021)
27	$C + O \longrightarrow CO$	Ioppolo et al. (2021)
28	$CH_3 + OH \longrightarrow CH_3OH$	Qasim et al. (2018)
29	$C + H_2 \longrightarrow CH_2$	Simončič et al. (2020), Lamberts et al. (2022)
30	$CH + H_2 \longrightarrow CH_3$	Lamberts et al. (2022)

in Ioppolo et al. (2021), using in part updated rate constants from recent literature, as indicated in Table 1.

We start by describing the astrochemical model, our choice of parameters and how we evaluate the network sensitivity in Section 2. We then discuss the results as well as the astrochemical implications in Section 3 and summarize our conclusions in 4.

2 METHODOLOGY

2.1 The astrochemical model

In this work, the gas-grain chemical code UCLCHEM was used (Holdship et al. 2017). UCLCHEM makes use of a rate equation approach to modelling the gas and grain-surface and bulk abundances. The gas-phase reaction network is taken from the UMIST database (McElroy et al. 2013). The grain-surface network used was the default one as available on GITHUB.

Various reaction mechanisms are implemented in UCLCHEM. The grain-surface reaction mechanisms that exist in UCLCHEM include the Eley-Rideal mechanism as well as the Langmuir–Hinshelwood diffusion mechanism, which were implemented in Quénard et al. (2018), as was the competition formula from Chang, Cuppen & Herbst (2007) and Garrod & Pauly (2011). The binding energies that are used to calculate the diffusion reaction rate are taken from Wakelam et al. (2017). We also included an updated version of the glycine grain-surface network from Ioppolo et al. (2021), also including both the reactions $C + H_2 \longrightarrow CH_2$ and $CH + H_2 \longrightarrow CH_3$ as summarized

in Table 1. Note that the reaction $OH + H_2 \longrightarrow H_2O + H$ had been already included, based on previous work by, e.g. Meisner et al. (2017). The code also includes thermal and non-thermal desorption, such as due to H_2 formation, cosmic ray ionization as well as UV-induced desorption. Note that the astrochemical model used in Ioppolo et al. (2021) makes use of the non-diffusive grain-surface chemistry that is described in Garrod & Pauly (2011) and Jin & Garrod (2020). This is not used in UCLCHEM. The implications of this will be discussed later in this work.

UCLCHEM is used to model two distinct phases of the star formation process. Phase 1 is the free-fall collapse phase of a dark cloud for a default value of 5 million years, whereas Phase 2 models the warm-up phase immediately following Phase 1, with the initial density of Phase 2 equal to the final density of Phase 1. Phase 2 runs for 1 million years. Further details of the code can be found in Holdship et al. (2017).

2.2 Parameter selection

To assess the importance of the two proposed reactions to the network under various interstellar conditions, three parameters were varied, as listed in Table 2. The standard cosmic ray ionization rate in UCLCHEM is $\zeta=1.3\times10^{-17}~{\rm s}^{-1}$. This is in line with typical values that are of the order $10^{-17}~{\rm s}^{-1}$ in diffuse ISM conditions (O'Donnell & Watson 1974; Black, Hartquist & Dalgarno 1978; Hartquist, Doyle & Dalgarno 1978; Indriolo & McCall 2013). However, there exist observations of higher cosmic ray ionization rates (Indriolo et al. 2007; Indriolo & McCall 2012), which is why we also include analysis of a region with cosmic ray ionization rate of 10ζ . Cosmic

¹https://uclchem.github.io/

Table 2. The parameters that were varied in this work to assess the effect of the two reactions.

Parameter	Values	Comment
Final density of Phase 1 and initial density of Phase 2 Efficiency for barrierless $C + H_2 \longrightarrow CH_2$ Cosmic ray ionization rate	$10^5 \text{ cm}^{-3}, 10^6 \text{ cm}^{-3}, 10^7 \text{ cm}^{-3}$ 0, 0.05, 1 $\zeta, 10\zeta$	Final density of Phase 1 same as initial density of Phase 2 Efficiency of 0 is equivalent to reaction being excluded ζ is the standard cosmic ray ionization rate of 1.3 \times 10 ⁻¹⁷ s ⁻¹

ray ionization is typically expected to break larger molecules into smaller radicals. We did not consider lower values of the cosmic ray ionization rate, as these are typically not observed. The cosmic ray dependency on column density in O'Donoghue et al. (2022) covered a range of values that were, however, already covered by the factor of 10 we consider here. While they found differences for lower densities during the collapse phase, these were ironed out once the collapse reached larger final densities, which is why here we do not include this dependency on column density.

Three different astronomical regions were modelled:

- (i) a dark cloud with a final density of 10⁵ cm⁻³
- (ii) a low-mass protostar with a final density of 10⁶ cm⁻³
- (iii) a high-mass protostar, with a final density of 10⁷ cm⁻³

The heating profiles during Phase 2 for the last two cases are based on Viti et al. (2004) and differ for each astronomical object. The dark cloud simulation was only run for Phase 1, but was allowed to run for a further million years to allow the chemistry to settle.

Another parameter that was varied was the efficiency, α , of the extent to which the reaction $C + H_2 \longrightarrow CH_2$ is barrierless. While Simončič et al. (2020) considered the reaction to be fully barrierless, Lamberts et al. (2022) found that the reaction barrier likely depends on the binding site. As such, our grid of models considers efficiencies for the reaction of 0 (the reaction is not included), 0.05 (5 per cent of binding sites lead to a barrierless reaction and 95 per cent of the binding sites have an infinitely high barrier) and 1 (the reaction is fully barrierless). What this means practically is that the reaction rate is multiplied by the efficiency. The reaction $CH + H_2 \longrightarrow CH_3$ was included as only being barrierless, based on Lamberts et al. (2022).

2.3 Evaluating the network sensitivity

We quantify the effect of the new reactions on the model by considering the change in abundances of the species that are the most affected when taking the ratio of the abundances of the modified and original models. The modified model is the chemical network which has $\alpha=1$, whereas the original model was taken to be the network which had neither of the dihydrogenation reactions. These two scenarios were taken to be the extremes of the parameter range in terms of including these reactions. The ratio is most sensitive to strong deviations in the molecular abundances as a result of the dihydrogenation reactions.

This ratio is defined for each species i as

$$\delta_i(t) = \frac{x_i^M(t)}{x_i^O(t)},\tag{1}$$

where $x_i^M(t)$ is the abundance of species i in the modified model at time t and $x_i^O(t)$ is the abundance of the same species in the original model at time t.

We only considered species which had a value above a 'threshold of detectability'. This was to ensure that we did not look at species whose original and changed abundances were below what can be observed from an astronomical point-of-view. For grain-surface

species this threshold was set to 10^{-8} with respect to hydrogen whereas for gas-phase species this threshold equalled 10^{-12} with respect to hydrogen. We took 10^{-8} as a lower-limit threshold for grain-surface species, as this was the order of magnitude of the lowest reported abundances in Boogert, Gerakines & Whittet (2015). Similarly, the gas-phase threshold was taken based on the abundances of COMs typically observed in the gas-phase, such as in Jiménez-Serra et al. (2016, 2021).

We can also define a quantity that tracks the absolute change in the abundance of species,

$$\Delta_i(t) = x_i^M(t) - x_i^O(t) = x_i^O(t)[\delta_i(t) - 1], \tag{2}$$

This value indicates how species with relatively large abundances, such as elemental species or their hydrogenation products, are redistributed

3 RESULTS AND ASTROCHEMICAL IMPLICATIONS

We find that even though the amounts by which various species are affected differs for each stage of star formation, the general trends are broadly similar. As such, we group our analysis per phase. Tables 3 and 4 summarize the changes in terms of δ . The effect of the enhanced cosmic ray ionization rate is discussed in Section 3.1.3.

Our results differ from Ioppolo et al. (2021) in that, while glycine does form on the grains, it does not do so in Phase 1, as UCLCHEM does not utilize non-diffusive grain-surface mechanisms. Instead, glycine forms on the grains as the temperature increases in Phase 2.

3.1 Impact of the parameters

In this sub-section we consider the role that the physical and chemical parameters play. Tables 3 and 4 show the changes in abundance when we compare the original network without the dihydrogenation reactions with the $\alpha=1$ case. Figs 1 and 2 show the time series of the abundances for glycine and its precursors.

3.1.1 Final density

The final density of the collapsing cloud had a minor effect on the final abundances of the species in Phase 1. For all three astronomical objects modelled in Phase 1, we observe a significant decrease of grain-surface CH and C when the reactions are included and see an enhancement of grain-surface CH_2 , CH_3 , and CH_4 . However, the values of δ as well as their original abundances seem to be independent of the density, suggesting a saturation effect.

In Phase 2, we observe that the final density of the collapsing cloud does affect the extent to which the added reactions influence the final abundances. We notice that several hydrogenation-based species have greater abundances at lower densities, including species such as HOCO, H_2O_2 , CH_3CCH , and H_2CO .

506 J. Heyl et al.

Table 3. Summary of the species that experienced the greatest increases (top section) and decreases (bottom section) for each of the three astronomical objects in Phase 1.

Dark cloud		Low-mass star				High-mass star		
		Original			Original			Original
Species	δ	abundances	Species	δ	abundances	Species	δ	abundances
#CH2	2.8	4.1×10^{-7}	#CH2	2.8	4.1×10^{-7}	#CH2	2.8	4.1×10^{-7}
#CH3	2.3	2.6×10^{-7}	#CH3	2.3	2.6×10^{-7}	#CH3	2.3	2.6×10^{-7}
#CH4	1.3	4.0×10^{-6}	#CH4	1.3	3.8×10^{-6}	#CH4	1.3	3.8×10^{-6}
#NH3	1.1	3.8×10^{-6}	#NH3	1.1	3.7×10^{-6}	#NH3	1.1	3.7×10^{-6}
#H2CS	1.1	2.4×10^{-8}	#H2CS	1.1	2.4×10^{-8}	#H2CS	1.1	2.4×10^{-8}
#CH3OH	1.04	1.5×10^{-5}	#CH3OH	1.04	1.3×10^{-5}	#СНЗОН	1.04	1.3×10^{-5}
#HNC	1.03	2.3×10^{-8}	#HNC	1.04	2.3×10^{-8}	#HNC	1.04	2.3×10^{-8}
#H2SiO	1.03	3.3×10^{-7}	#H2SiO	1.03	1.1×10^{-8}	#H2SiO	1.03	3.4×10^{-7}
#HCN	1.02	1.7×10^{-7}	#HO2	1.03	2.3×10^{-7}	#HO2	1.03	2.3×10^{-7}
#O2	1.02	1.8×10^{-6}	NO	1.03	1.0×10^{-10}	#HCN	1.02	1.6×10^{-7}
#CH	1.1×10^{-15}	7.2×10^{-7}	#CH	2.0×10^{-15}	7.2×10^{-7}	#CH	2.1×10^{-15}	7.2×10^{-7}
#C	2.4×10^{-13}	1.4×10^{-6}	#C	2.5×10^{-13}	1.4×10^{-6}	#C	2.5×10^{-13}	1.4×10^{-6}
#NCH4	3.7×10^{-13}	1.5×10^{-7}	#NCH4	3.4×10^{-13}	1.5×10^{-7}	#NCH4	3.4×10^{-13}	1.5×10^{-7}
#NH2CH3	8.0×10^{-13}	1.9×10^{-7}	#NH2CH3	8.3×10^{-13}	2.0×10^{-7}	#NH2CH3	8.3×10^{-13}	2.0×10^{-7}
NH2CH3	1.5×10^{-12}	8.7×10^{-10}	#Si	0.98	5.6×10^{-8}	#Si	0.98	5.6×10^{-8}
CH	0.96	9.3×10^{-10}	#SiH	0.99	2.5×10^{-8}	#SiH	0.99	2.5×10^{-8}
CH3	0.98	1.5×10^{-9}	#SiH2	0.99	1.3×10^{-8}	#SiH2	0.99	1.3×10^{-8}
#Si	0.98	5.7×10^{-8}	#O	0.99	7.8×10^{-5}	#SI	0.99	6.7×10^{-5}
#SiH	0.99	2.6×10^{-8}	#H3CO	0.99	1.7×10^{-6}	#H3CO	0.99	1.7×10^{-6}
#SiH2	0.99	1.4×10^{-8}	#HNO	0.99	1.2×10^{-5}	#HNO	0.99	1.2×10^{-5}

Species with a '#' are grain-surface species. All other species are gas-phase.

Table 4. Summary of the species that experienced the greatest increases (top section) and decreases (bottom section) for each of the three astronomical objects in Phase 2.

Low-mass star					
Species	δ	Original abundances	Species	δ	Original abundances
HOCO	3.7	9.3×10^{-10}	НОСО	2.1	4.3×10^{-8}
H_2O_2	2.6	4.3×10^{-9}	СНЗОН	2.0	1.8×10^{-9}
CH ₃ CHO	2.2	1.0×10^{-7}	СН3СНО	2.0	1.5×10^{-7}
CH ₃ OH	2.1	3.7×10^{-9}	C2H4	2.0	2.5×10^{-9}
CH3CN	1.7	1.0×10^{-9}	CH2CO	1.9	1.8×10^{-10}
C4H	1.6	3.2×10^{-10}	H2CO	1.7	9.3×10^{-9}
C3H2	1.5	5.6×10^{-9}	CH3	1.7	1.1×10^{-10}
CH3CCH	1.5	2.4×10^{-8}	NH3	1.6	1.3×10^{-8}
NH3	1.5	2.7×10^{-7}	CH3CN	1.5	7.1×10^{-10}
NH2CHO	1.4	2.7×10^{-7}	C2H2	1.5	1.1×10^{-8}
NCH4	3.8×10^{-5}	9.2×10^{-7}	NCH4	4.7×10^{-5}	8.3×10^{-7}
NH2CH3	2.4×10^{-3}	1.6×10^{-7}	NH2CH3	2.5×10^{-3}	1.7×10^{-7}
	6.0×10^{-2}	6.3×10^{-9}		6.3×10^{-3}	7.2×10^{-8}
NH2CH2COOH			NH2CH2COOH		
H2S	0.88	2.0×10^{-9}	NO	0.82	4.0×10^{-6}
SO2	0.92	4.4×10^{-8}	NCCN	0.96	3.9×10^{-7}
MG+	0.93	8.0×10^{-8}	O2	0.96	7.1×10^{-6}
O	0.95	1.3×10^{-5}	HCOO	0.96	1.9×10^{-10}
CH2OH	0.95	6.4×10^{-8}	C2N	0.97	3.5×10^{-8}
O2	0.96	4.2×10^{-5}	O	0.97	3.6×10^{-8}
SO	0.97	1.9×10^{-6}	CO2	0.97	7.6×10^{-6}

All species listed are gas-phase.

3.1.2 Efficiency

For more abundant species, such as H_2O and CH_3OH , we find that the results obtained from using a branching fraction of 0.05 for the barrierless dihydrogenation of C are essentially the same as using a efficiency of 1 (the reaction is fully barrierless).

We do find that the efficiency parameter plays a role in the final abundances of glycine and its precursors during the warmup phase of low and high-mass stars. This can be seen in Figs 1 and 2. For Phase 1, the species are not detectable except for the original configuration. However, we still observe that for the other three configurations an increasing value of α corresponds to an increased level of depletion. In Phase 2, the configurations are all detectable and this same hierarchy remains in the gas-phase.

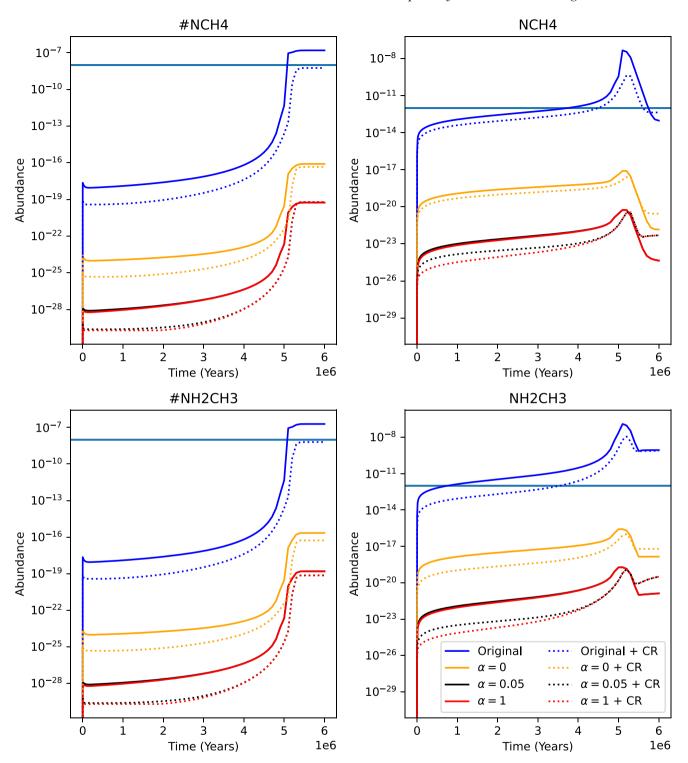


Figure 1. Time series of the abundances of grain-surface and gas-phase NH_2CH_2 and NH_2CH_3 in Phase 1 of a dark cloud. Furthermore, we observe that the inclusion of the dihydrogenation reactions, regardless of efficiency α severely depletes the abundances of the glycine precursors in both phases relative to the original model which did not include either of the dihydrogenation reactions. Also plotted are the limits of detectability we have used for gas and grain-surface species. We do not plot glycine, as it is not formed at all in Phase 1. We observe that only the original model is capable of producing 'detectable' levels of methylamine and the methylamine radical. For the other configurations, an increase in α results in increased depletion of the species relative to the original model. We also observe that enhanced cosmic ray ionization depletes the abundances on the grains but not in the gas.

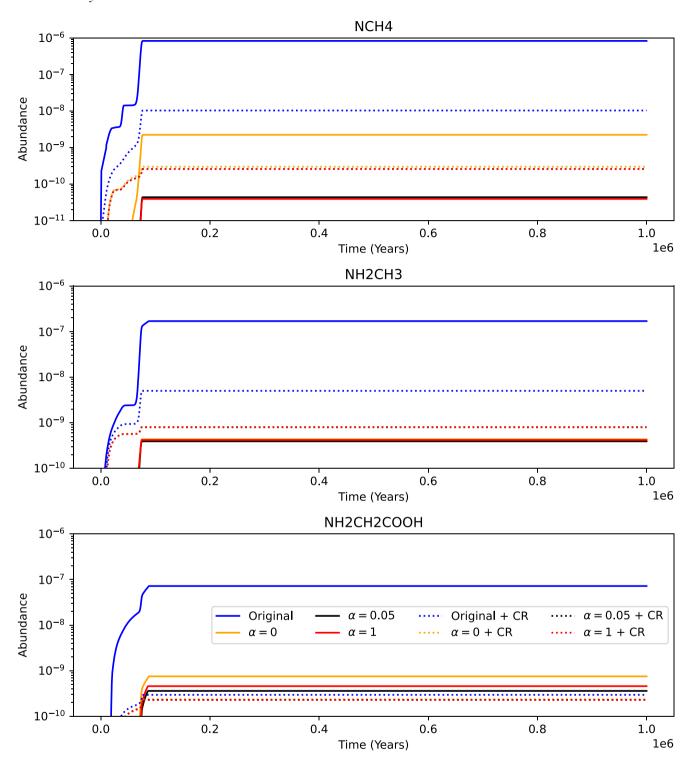


Figure 2. Time series of the abundances of gas-phase NH₂CH₂, NH₂CH₃, and NH₂CH₂COOH in Phase 2 of a high-mass star. We observe that glycine is produced in the warm-up phase. The enhanced cosmic ray ionization rate is found to significantly deplete all three species in the gas-phase for the original model. For NH₂CH₂ and NH₂CH₃, when $\alpha = 0$, $\alpha = 0.05$ or $\alpha = 1$, the enhanced cosmic ray ionization rate results in an increase of their abundances. For glycine, the enhanced cosmic ray ionization rate seems to decrease its gas-phase abundance.

3.1.3 Cosmic ray ionization rate

The degree of cosmic ray ionization is found to play an important role in enhancing or counteracting the role of the dihydrogenation reactions. The cosmic ray destruction routes we include in our standard network are from Garrod, Widicus Weaver & Herbst (2008).

These consist of hydrogen abstraction reactions and reactions that produce radical-radical pairs of products. An enhanced cosmic ray ionization rate leads to the destruction of many of hydrogenated species, such as CH₄, NH₃, H₂O, and CH₃OH, as well as their precursors. This leads to further hydrogen reservoirs being released

and radicals being formed which can go on to form glycine and its precursors. Because no cosmic ray destruction mechanisms for these complex, larger, species are included, we find that these are more abundantly produced.

This is important to consider in the context of glycine. In Figs 1 and 2, we plot the time dependence of the abundance of glycine precursors for eight different parameter sets, including the enhanced cosmic ray ionization rate. In Phase 1, we find that on the grains, the enhanced cosmic ray ionization rate depletes the species. In Phase 2, the effect varies by configuration and species. The original configuration consistently leads to a decrease of all plotted species in the presence of enhanced cosmic ray ionization. The $\alpha=0$ configuration is depleted for the methylamine radical and glycine, but enhanced for methylamine. The $\alpha=0.05$ and $\alpha=1$ configurations are depleted for methylamine and glycine, but enhanced for the methylamine radical.

3.2 General implications

As can be seen in Tables 3 and 4, the inclusion of reactions with molecular hydrogen affects the hydrogen economy of the reaction network. Previously, the reaction network had a significant amount of H₂ being adsorbed or produced on the surface with no chemical destruction mechanisms. The H₂ molecules are a previously untapped hydrogen reservoir that is now being utilized (Hasegawa & Herbst 1993). Because one H₂ frees up two H atoms on the surface, other atomic hydrogenation reactions can take place more easily. Therefore, we observe the increase in the abundances of species in Phases 1 and 2 that are the products of hydrogenation. While for many of the more common species, the relative increase, i.e. δ is small, the abundance increases in absolute terms. There are large relative and absolute changes in the network of less abundant species, such as NH₂CH₂, NH₂CH₃, and NH₂CH₂COOH and there are fairly large absolute changes in the network of highly abundant species, such as C and its hydrogenation products.

We can also comment on the carbon budget. The previously defined Δ parameter allows us to consider how carbon is redistributed as a result of the new reactions being included. For instance, for the dark cloud during Phase 1, the total Δ for the main carbon-based grain-surface species that increase

$$\Delta_{\text{total}}(\text{\#CH}_2 + \text{\#CH}_3 + \text{\#CH}_4 + \text{\#H}_2\text{CS} + \text{\#CH}_3\text{OH}) = 2.9 \times 10^{-6}.$$

is nearly equal to that of the total decrease $\boldsymbol{\Delta}$ of main grain-surface species:

$$\Delta_{\text{total}}$$
 (#C + #CH + #NCH₄ + #NH₂CH₃ = 2.5 × 10⁻⁶.

From this we can see that the dihydrogenation reactions redistribute the carbon between the aforementioned species. The remaining carbon is redistributed to other species in the network in smaller amounts. We also observe that besides the methyl radical, also species that contain the CH_3 group, such as CH_3OH and CH_3CN see increases in their abundances, via the reactions $CH_3 + OH \longrightarrow CH_3OH$ and $CH_3 + CN \longrightarrow CH_3CN$.

In a similar fashion, nitrogen is redistributed throughout the network. The grain-surface ammonia abundance increases by 10 per cent, i.e. 3.8×10^{-7} . The decrease in #NCH₄ and #NH₂CH₃ accounts for 3.4×10^{-7} or ~ 90 per cent.

3.3 Implications for simple grain-surface species

In the light of the recent ice observations with the *James Webb Space Telescope*, both published (Yang et al. 2022) and upcoming (McClure

et al. 2017), it is important to consider the effect on the main ice constituents. Fig. 3 shows the time-evolution of the abundances of grain-surface H₂O, CO, CO₂, CH₃OH, H₂CO, NH₃, and CH₄ in Phase 1 of a dark cloud. These are species that have been securely or likely identified in the ices (Boogert et al. 2015). The shaded areas in the plots indicate the 68 per cent confidence interval for the measured abundances, taken from Boogert et al. (2015). In Boogert et al. (2015), the abundances were given in terms of the median value as well as the upper and lower quartiles. It was assumed that the spread in the measurements was Gaussian, which meant that the interquartile range represented 1.36σ . This spread in measurements is due to both observational error and source-to-source variation. We observe that we recover the measured abundances for most of the species within the uncertainty, with the exception of grain-surface CO₂. The inclusion of the dihydrogenation reactions does not change how well the models agree with the abundance measurements; however, for all hydrogenation products we observe that the inclusion of reactions with molecular hydrogen increases their abundance, as a result of the additional atomic hydrogen on the surface. In short, despite uncertainties surrounding activation energies, networks, and binding energies, we are able to recover observational abundances reasonably well when we include the reactions with molecular hydrogen and this gives us confidence that the predictions we make for glycine and its precursors are accurate.

3.4 Implications for glycine and its precursors

In Tables 3 and 4, we observe that the abundances of glycine and its precursors decreases if molecular hydrogen is part of the reaction network. We can also explain why the abundance of precursors of glycine, gas and grain NH_2CH_3 and NH_2CH_2 decrease. The former is formed through the reaction $NH_2 + CH_3$, but since more atomic H is present on the grains, both radical species are preferentially hydrogenated. The inclusion of H_2 as a reacting species, not just in the context of the two reactions we consider in this work, introduces greater competition for radicals that are needed for the formation of complex organic molecules. This results in the lower abundances of NH_2CH_3 and NH_2CH_2 .

We can also use this to justify the impact of the efficiency. Figs 1 and 2 plot the time series for the various efficiencies as well as with enhanced cosmic ray ionization in Phase 1 and 2, respectively. We previously remarked that the original configuration produced the most of glycine and its precursors. For the other configurations, the greater the value of α , the greater the depletion of these species. This makes sense when one considers that an increasing value of α results in more H_2 being consumed and therefore more atomic H becoming available to hydrogenate precursors.

We now look to compare our results with observations. We do this separately for glycine and its precursors. We also discuss the implications of not using non-diffusive grain-surface mechanisms in our code, such as the ones discussed in Garrod & Pauly (2011) and Jin & Garrod (2020).

3.4.1 Methylamine and the methylamine radical

Methylamine (NH₂CH₃) and the methylamine radical (NH₂CH₂) are important precursors of glycine. The hydrogen abstraction of methylamine to form the methylamine radical is crucial, as there is growing evidence to suggest that the reaction NH₂CH₂ + HOCO \longrightarrow NH₂CH₂COOH is a feasible glycine

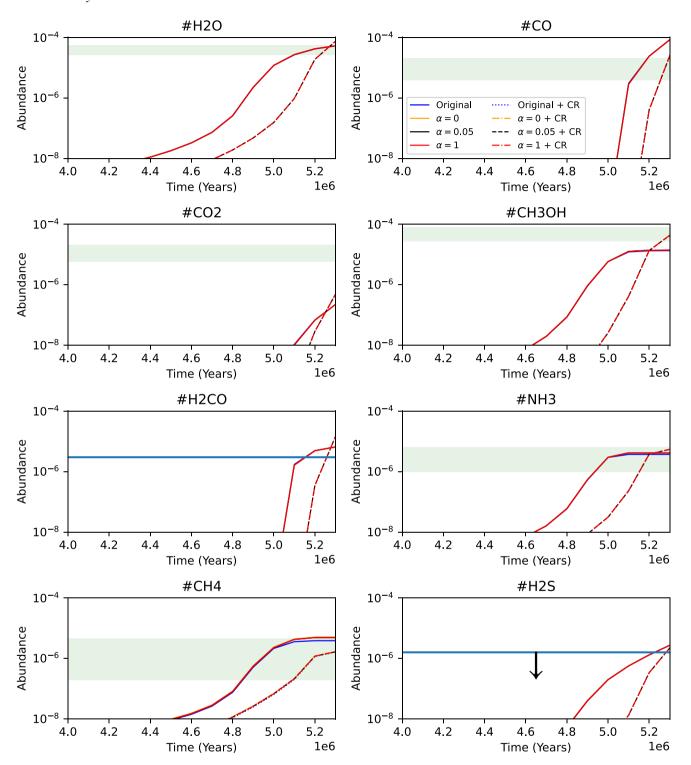


Figure 3. Time series of the abundances of grain-surface H_2O , CO, CO_2 , CH_3OH , H_2CO , NH_3 , CH_4 , and H_2S in Phase 1 of a dark cloud. We include the species that have securely identified or likely identified. The abundances were adapted from Boogert et al. (2015). The shaded areas include the 1σ region of abundances. In the case of H_2CO , no uncertainty was provided in the original source, so there is no shaded area. Grain-surface H_2S only has an upper limit on its abundance. For both normal and enhanced cosmic ray ionization rates, the time-series differ very little, which is why it is difficult to distinguish them visually.

formation route (Ramesh & Yuan-Pern 2022). Confirmed detections of methylamine in high-mass star forming regions are summarized in Table 5. We observe improved level of agreement between our model outputs and observations when the reactions are included with $\alpha=1$. We observed significant enhancement when the cosmic ray ionization rate was increased. This suggests that if dihydrogen is

chemically active on the grains, one would need to consider regions of high cosmic ray ionization rate to detect these precursors of glycine, as these reactions reduce the abundance of methylamine. In the case of the Bøgelund et al. (2019) observation, we have confidence in the value of our ratio, as the chemical network for methanol is well-established.

Table 5. Table of methylamine abundance measurements relative to reference molecules for high-mass stars.

Abundance measurements (relative to					
Reference molecule	Reference	reference molecule)	Original model ratio	New model ratio	
CH ₃ OH	Bøgelund et al. (2019)	$8 \times 10^{-3} - 0.1$	37	0.02	
H_2	Ohishi et al. (2019)	$1.5 \pm 1.1 \times 10^{-8}$	3.5×10^{-7}	3.9×10^{-10}	

Also included are the corresponding ratios obtained in this work for high-mass stars with the standard cosmic ray ionization rate.

However, the entirety of the above discussion regarding the agreement of our results with observations is incomplete without discussing the effect of the non-diffusive reaction mechanisms being absent in our modelling. These mechanisms are of particular use when considering reactions between reactants which are likely to react very slowly via the Langmuir–Hinshelwood diffusion mechanism, such as the reaction between CO and OH to form CO₂. Methylamine and the methylamine radical are formed via reactions 6 and 7, which involve species with high binding energies, thereby making their formation at 10 K inefficient via diffusion. As a result, the fact that we do not include the non-diffusive mechanisms means that methylamine and its radical are under-produced.

3.4.2 Glycine

While there may be no confirmed detection for glycine in the literature, various estimates exist. In Gibb et al. (2004), an upper limit of 0.3 per cent with respect to water was determined, whereas in Jiménez-Serra et al. (2014), this was estimated to be around 0.1 per cent. In this work, we find that when the dihydrogenation reactions are not included this value is 0.07 per cent and when we include both reactions then it is 2×10^{-4} per cent. We should note that in the absence of experimentally-motivated gas-phase glycine destruction reactions the values derived in this work are only upper limits, if one neglects non-diffusive mechanisms. In the previous sub-section, we discussed that methylamine and its radical are underproduced. This will result in glycine being underproduced as well, not just due to the underproduction of its precursors, but also because reaction 10 is less efficient if assumed to be diffusion-only.

4 CONCLUSION

In this work, we considered the effect of including the reactions of H_2 with C and CH in our grain-surface network. We ran a grid of 12 models that vary the final density of the collapsing cloud, the efficiency for the 'barrier' of $C + H_2 \longrightarrow CH_2$ as well as the cosmic ray ionization rate.

Making molecular hydrogen chemically active unlocks a previously untapped reservoir of hydrogen, and therefore freeing up the use of atomic hydrogen for hydrogenation reactions. A particularly interesting consequence of this is that making H₂ more chemically active decreased the abundances of glycine and its precursors. This may aid in explaining why methylamine, the methylamine radical as well as glycine have remained undetected so far.

We note that we do not have a comprehensive gas-phase network for glycine and its precursors. That is likely to be a limitation. While it is still likely that glycine and its precursors form on the grains and then evaporate into the gas-phase, it is possible that there would be gas-phase destruction routes as well. Additionally, cosmic-ray ionization destruction routes on the grains and in the gas-phase are likely also needed, as these typically break large molecules down into smaller radicals which are then recycled for further gas-phase

reactions. As such, the abundances we obtain for glycine and its precursors are likely to only be upper limits.

An additional limitation is the absence of the non-diffusive reaction mechanisms discussed in Garrod & Pauly (2011) and Jin & Garrod (2020). The consequence is that glycine and its precursors do not form efficiently on the grains at 10 K, which is different to what was found in Ioppolo et al. (2021). As such, they are under-produced in our models, whereas diffusion-efficient reactions overproduce certain species. However, without implementing this formalism in the code, it is difficult to assess the relative impacts of these mechanisms on the final abundances.

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DATA AVAILABILITY

The data underlying this article are available in the article and in its online supplementary material.

REFERENCES

Al-Halabi A., van Dishoeck E. F., 2007, MNRAS, 382, 1648

Arasa C., van Hemert M. C., van Dishoeck E. F., Kroes G. J., 2013, J. Phys. Chem., 117, 7064

Balucani N., Bergeat A., Cartechini L., Volpi G. G., Casavecchia P., Skouteris D., Rosi M., 2009, J. Phys. Chem., 113, 11138

Bernstein M. P., Dworkin J. P., Sandford S. A., Cooper G. W., Allamandola L. J., 2002, Nature, 416, 401

Black J. H., Hartquist T. W., Dalgarno A., 1978, ApJ, 224, 448

Bøgelund E. G., McGuire B. A., Hogerheijde M. R., van Dishoeck E. F., Ligterink N. F. W., 2019, A&A, 624, A82

Boogert A. A., Gerakines P. A., Whittet D. C., 2015, ARA&A, 53, 541

Bossa J. B., Duvernay F., Theulé P., Borget F., D'Hendecourt L., Chiavassa T., 2009, A&A, 506, 601

Caselli P., Ceccarelli C., 2012, A&AR, 20, 56

Chang Q., Cuppen H. M., Herbst E., 2007, A&A, 469, 973

Ciesla F. J., Sandford S. A., 2012, Science, 336, 452

Fedoseev G., Ioppolo S., Lamberts T., Zhen J. F., Cuppen H. M., Linnartz H., 2012, J. Chem. Phys., 137, 054714

Garrod R. T., 2013, ApJ, 765, 60

Garrod R. T., Pauly T., 2011, ApJ, 735, 15

Garrod R. T., Widicus Weaver S. L., Herbst E., 2008, ApJ, 682, 283

Gibb E. L., Whittet D. C. B., Boogert A. C. A., Tielens A. G. G. M., 2004, ApJS, 151, 35

512 *J. Heyl et al.*

Goldsmith P. F., Li D., 2005, ApJ, 622, 938

Goumans T. P. M., Uppal M. A., Brown W. A., 2008, MNRAS, 384, 1158

Hartquist T. W., Doyle H. T., Dalgarno A., 1978, A&A, 68, 65

Hasegawa T. I., Herbst E., 1993, MNRAS, 261, 83

Herbst E., van Dishoeck E. F., 2009, ARA&A, 47, 427

Holdship J., Viti S., Jiménez-Serra I., Makrymallis A., Priestley F., 2017, AJ, 154, 38

Indriolo N., Geballe T. R., Oka T., McCall B. J., 2007, ApJ, 671, 1736

Indriolo N., McCall B. J., 2012, ApJ, 745, 91

Indriolo N., McCall B. J., 2013, Chem. Soc. Rev., 42, 7763

Ioppolo S. et al., 2021, Nature Astron., 5, 197

Ioppolo S., Cuppen H. M., van Dishoeck E. F., Linnartz H., 2011, MNRAS, 410, 1089

Jiménez-Serra I. et al., 2016, ApJ, 830, L6

Jiménez-Serra I., testi L., Caselli P., Viti S., 2014, ApJ, 787, L33

Jiménez-Serra I., Vasyunin A. I., Spezzano S., Caselli P., Cosentino G., Viti S., 2021, ApJ, 917, 44

Jin M., Garrod R. T., 2020, ApJS, 249, 26

Krasnokutski S. A., Kuhn M., Renzler M., Jäger C., Henning T., Scheier P., 2016, ApJ, 818, L31

Lamberts T., Cuppen H. M., Ioppolo S., Linnartz H., 2013, Phys. Chem. Chem. Phys., 15, 8287

Lamberts T., Fedoseev G., Kästner J., Ioppolo S., Linnartz H., 2017, A&A, 599, A132

Lamberts T., Fedoseev G., van Hemert M. C., Qasim D., Chuang K.-J., Santos J. C., Linnartz H., 2022, ApJ, 928, 48

Lamberts T., Kästner J., 2017, ApJ, 846, 43

Lee C.-W., Kim J.-K., Moon E.-S., Minh Y. C., Kang H., 2009, ApJ, 697, 428
 Maté B., Tanarro I., Escribano R., Moreno M. A., Herrero V. J., 2015, ApJ, 806, 151

McClure M. et al., 2017, IceAge: Chemical Evolution of Ices during Star Formation, JWST Proposal ID 1309. Cycle 0 Early Release Science

McElroy D., Walsh C., Markwick A. J., Cordiner M. A., Smith K., Millar T. J., 2013, A&A, 550, A36

Meisner J., Lamberts T., Kästner J., 2017, ACS Earth Space Chem., 1, 399 Molpeceres G., Kästner J., 2020, Phys. Chem. Chem. Phys., 22, 7552

O'Donnell E. J., Watson W. D., 1974, ApJ, 191, 89

O'Donoghue R., Viti S., Padovani M., James T., 2022, ApJ, 934, 63

Oba Y., Chigai T., Osamura Y., Watanabe N., Kouchi A., 2014, Meteorit. Planet. Sci., 49, 117

Ohishi M., Suzuki T., Hirota T., Saito M., Kaifu N., 2019, PASJ, 71, 86 Pernet A., Pilmé J., Pauzat F., Ellinger Y., Sirotti F., Silly M., Parent P., Laffon C., 2013, A&A, 552, A100

Qasim D., Chuang K. J., Fedoseev G., Ioppolo S., Boogert A. C. A., Linnartz H., 2018, A&A, 612, A83

Quénard D., Jiménez-Serra I., Viti S., Holdship J., Coutens A., 2018, MNRAS, 474, 2796

Ramesh J. P., Yuan-Pern L., 2022, Commun. Chem., 5, 1

Sato A. et al., 2018, Mol. Astrophys., 10, 11

Simončič M., Semenov D., Krasnokutski S., Henning T., Jäger C., 2020, A&A, 637, A72

van Dishoeck E. F., Black J. H., 1988, ApJ, 334, 771

Viti S., Collings M. P., Dever J. W., McCoustra M. R. S., Williams D. A., 2004, MNRAS, 354, 1141

Wakelam V., Loison J. C., Mereau R., Ruaud M., 2017, Mol. Astrophys., 6,

Woon D. E., 2002, ApJ, 571, L177

Yang Y.-L., 2022, The Astrophysical Journal Letters, 941, L13

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