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# Vanadium oxide and a sharp onset of cold-trapping on a giant exoplanet

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The abundance of refractory elements in giant planets can provide key insights into their formation histories<sup>1</sup>. Owing to the low temperatures of the Solar System giants, refractory elements condense below the cloud deck, limiting sensing capabilities to only highly volatile elements<sup>2</sup>. Recently, ultra-hot giant exoplanets have allowed for some refractory elements to be measured, showing abundances broadly consistent with the solar nebula with titanium probably condensed out of the photosphere<sup>3,4</sup>. Here we report precise abundance constraints of 14 major refractory elements on the ultra-hot giant planet WASP-76b that show distinct deviations from proto-solar and a sharp onset in condensation temperature. In particular, we find nickel to be enriched, a possible sign of the accretion of the core of a differentiated object during the evolution of the planet. Elements with condensation temperatures below 1,550 K otherwise closely match those of the Sun<sup>5</sup> before sharply transitioning to being strongly depleted above 1,550 K, which is well explained by nightside cold-trapping. We further unambiguously detect vanadium oxide on WASP-76b, a molecule long suggested to drive atmospheric thermal inversions<sup>6</sup>, and also observe a global east–west asymmetry<sup>7</sup> in its absorption signals. Overall, our findings indicate that giant planets have a mostly stellar-like refractory elemental content and suggest that temperature sequences of hot Jupiter spectra can show abrupt transitions wherein a mineral species is either present or completely absent if a cold trap exists below its condensation temperature<sup>8</sup>.

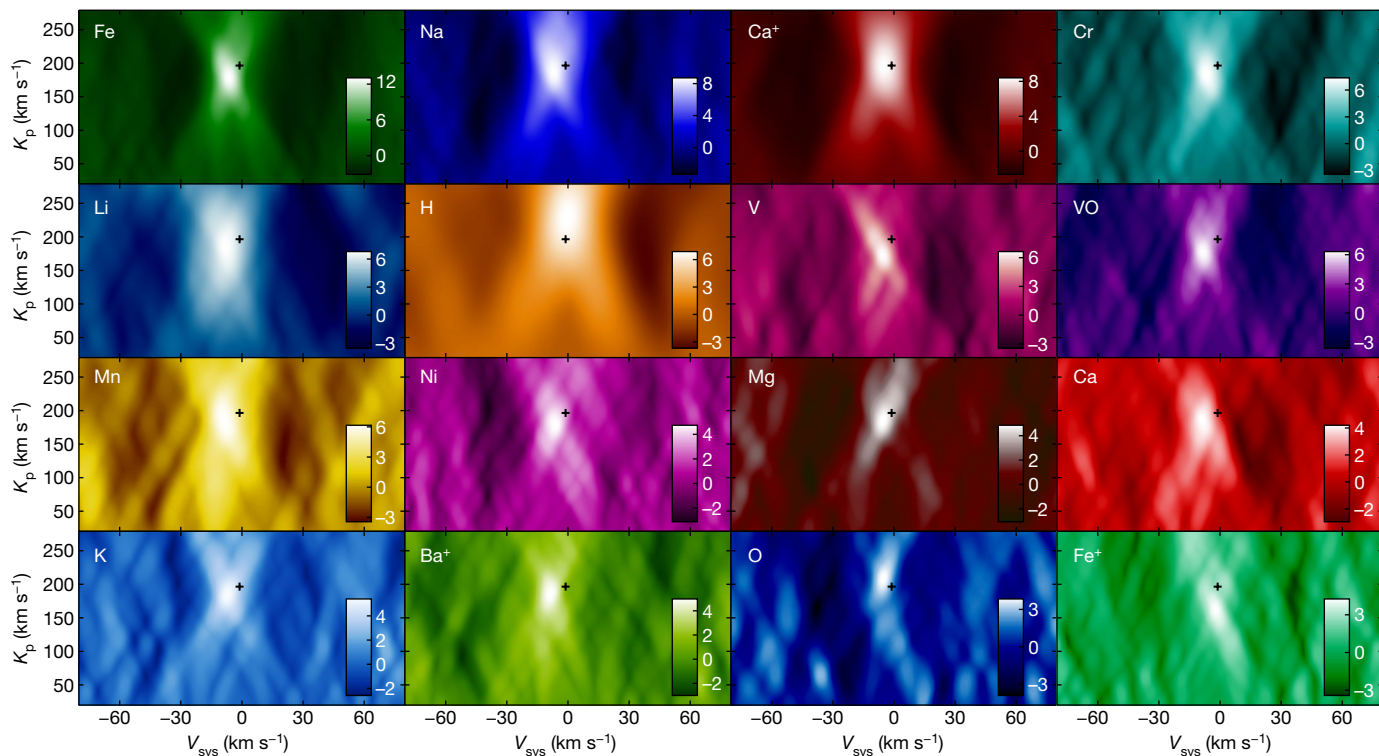
We observed three transits of the ultra-hot Jupiter WASP-76b (ref. 9) using the M dwarf Advanced Radial velocity Observer Of Neighboring eXoplanets (MAROON-X) high-resolution optical spectrograph<sup>10</sup> at the 8.1-m Gemini North Observatory in Hawaii. With a continuous wavelength coverage between 490 and 920 nm at a spectral resolution of  $\lambda/\Delta\lambda \approx 85,000$ , MAROON-X is well suited for investigating atomic species in planetary atmospheres. Two transits of WASP-76b were observed on 3 and 12 September 2020 and a third transit on 28 October 2021. Each transit observation consists of a time series of between 36 and 47 high-resolution spectra, each with integration times between 5 and 8 min.

The observed spectra contain contributions from WASP-76b (the planet), WASP-76 (the host star) and from the atmosphere of the Earth. The last two dominate the spectra and must be removed to uncover the comparatively much fainter planetary signal. We achieve this by using a principal component analysis (PCA)-based algorithm that removes the relatively stationary-in-velocity stellar ( $<1 \text{ km s}^{-1}$ ) and telluric ( $0 \text{ km s}^{-1}$ ) contributions from the data while leaving the rapidly Doppler shifting (about  $100 \text{ km s}^{-1}$ ) planetary signal largely unaffected

(see Methods and Extended Data Fig. 1). To uncover the atmospheric signature of WASP-76b, we cross-correlate the resulting cleaned-up spectra with transit-model templates as a function of radial velocity. If a model matches the data well, the cross-correlation function (CCF) of each in-transit spectrum should peak at the exact Doppler shift matching the orbital velocity of WASP-76b at that moment<sup>11</sup>. Viewing the cross-correlation as a function of time should then produce a trail that is centred at the velocity of the system ( $V_{\text{sys}}$ ) with a slope matching the radial velocity semi-amplitude of WASP-76b ( $K_p$ ). The CCFs at each phase can then be summed for different configurations of  $K_p$  and  $V_{\text{sys}}$  to produce two-dimensional signal-to-noise maps for a given atmospheric model. If a species is detected in the atmosphere of WASP-76b, a strong peak will be observed in the resulting map near the expected  $K_p$  and  $V_{\text{sys}}$ . We repeat this cross-correlation process using model templates containing absorption lines of only a single metal, ion or molecule (Extended Data Fig. 2) at a time to produce a chemical inventory of the atmosphere of WASP-76b.

We detect Fe, Na,  $\text{Ca}^+$ , Cr, Li, H, V, VO, Mn, Ni, Mg, Ca, K and  $\text{Ba}^+$  in the atmosphere of WASP-76b, as well as tentatively detect O and  $\text{Fe}^+$ ,

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**Fig. 1 | Cross-correlation results for species detected in the atmosphere of WASP-76b.** In each panel, the signal-to-noise velocity map of an element is shown, with the black cross indicating the expected location of the signal assuming a symmetric planet with a static atmosphere. Deviations from the

black cross in the  $K_p$ – $V_{\text{sys}}$  space can be indicative of chemical asymmetries and dynamics on WASP-76b. Clear signals, many of which are slightly offset, can be seen as bright-white blobs near the expected position for all 16 species shown.

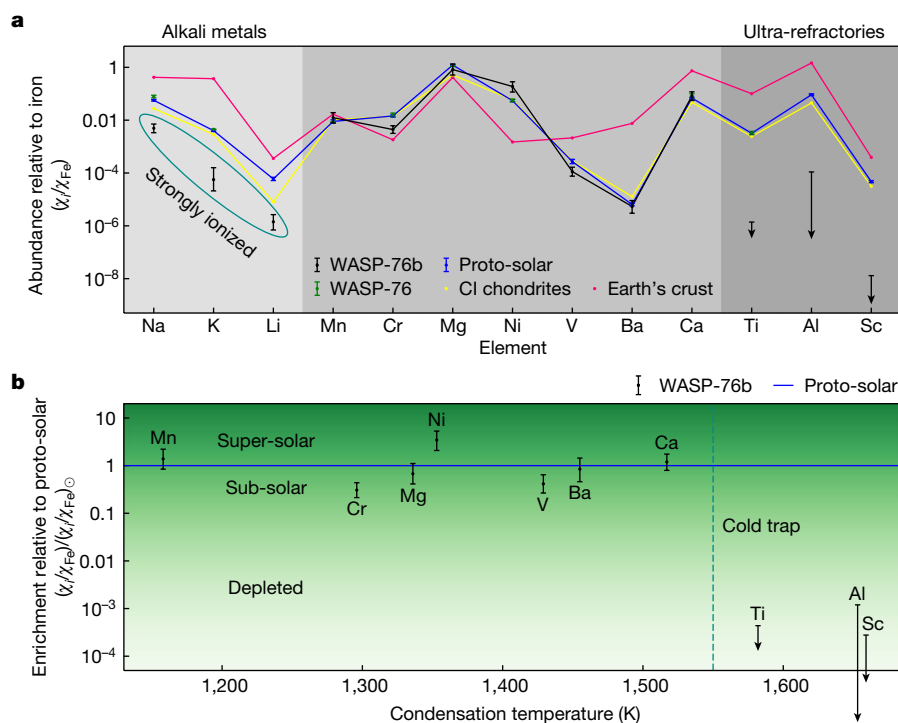
providing an unprecedentedly complete measurement of the chemical inventory of a gas giant planet (Fig. 1). We also notably do not find evidence of absorption from neutral or oxidized titanium (Extended Data Fig. 3), despite being highly sensitive to those species if they were in chemical-equilibrium abundances. The observed presence or absence of these species provides us with key insight into the conditions of the highly irradiated atmosphere of WASP-76b. Vanadium oxide (VO), in particular, is a strong optical absorber that has long been sought in ultra-hot Jupiter atmospheres owing to its theorized role as a driver of thermal inversions<sup>6</sup>. We detect the VO signal in all three individual MAROON-X transits and also with an independent analysis of two ESPRESSO transits of WASP-76b (see Methods and Extended Data Fig. 4). Our detection thus confirms that VO is present in hot Jupiter atmospheres as a source of atmospheric heating, alongside other shortwave absorbers, such as  $\text{H}^-$  and atomic metals<sup>12</sup>. Especially in the absence of titanium oxide, another highly potent optical absorber present in some even hotter exoplanets<sup>13</sup>, the VO molecules are directly exposed to the incoming short-wavelength stellar irradiation and become the dominant optical broadband absorbers above the  $\text{H}^-$  continuum, adding hundreds of Kelvins to the upper atmosphere of WASP-76b (ref. 14). The detection of ionized barium, with an atomic number  $Z = 56$ , also shows that, like Jupiter<sup>15</sup>, the atmosphere of WASP-76b is not notably fractionated by mass. Furthermore, the combined presence of V/VO, Ca/Ca<sup>+</sup> and Fe/Fe<sup>+</sup> provides a physical and chemical thermometer wherein the atmosphere of WASP-76b must have temperature regions in which these combinations of species can coexist, albeit potentially at different altitudes and/or longitudes. Most detections are slightly offset in velocity space from literature predictions for a symmetric and static atmosphere, indicative of dynamical<sup>17</sup> and chemical inhomogeneities<sup>13</sup> and/or three-dimensional effects<sup>16</sup> on WASP-76b.

Beyond identifying which species are present in the atmosphere of WASP-76b, we apply a high-resolution Bayesian atmospheric retrieval

framework<sup>17</sup> to the MAROON-X data and infer the presence of a stratosphere, bounded abundance constraints for 13 elements and molecules, as well as upper limits on several other species (see Extended Data Fig. 5 and Extended Data Table 1). We find that WASP-76b reaches temperatures substantially hotter than its equilibrium temperature of  $T_{\text{eq}} = 2,228$  K assuming zero albedo<sup>7</sup>. The hot stratosphere is also consistent with the presence of ionized species such as Ca<sup>+</sup> and Ba<sup>+</sup> at high altitude (Extended Data Fig. 6).

Comparing the inferred elemental abundances on WASP-76b to host star<sup>18</sup> and proto-solar<sup>5</sup> values, we find that the abundances of Mn, Cr, Mg, Ni, V, Ba and Ca all follow a remarkably similar trend, especially when taken relative to Fe (Fig. 2a). This agreement between chemically unfractured materials and planetary relative abundances spans several orders of magnitude and sharply contrasts compositions of highly differentiated bodies, such as the Earth's crust<sup>19</sup>. This indicates that the present-day atmosphere of WASP-76b, to the first order, has a similar refractory composition to the parent protoplanetary disk from which it was formed, similarly to what was found for a subset of these elements on another ultra-hot Jupiter, WASP-121b (refs. 3,4). We measure the abundance of neutral alkali metals Li, Na and K in the photosphere of WASP-76b to be markedly sub-solar, which is naturally explained by their relatively low ionization potentials causing these elements to be heavily ionized at the investigated temperatures and pressures. With their respective ions (Li<sup>+</sup>, Na<sup>+</sup>, K<sup>+</sup>) lacking spectral features because of not having any valence electrons, they cannot be remotely examined with MAROON-X. This naturally leads to their abundances being underpredicted from only considering their neutral forms and therefore not being representative of the bulk atmosphere.

The ultra-refractory elements Ti, Sc and Al, on the other hand, are severely depleted relative to proto-solar. Unlike alkali metals, these are not expected to be notably ionized. Instead, with their higher condensation temperatures<sup>20</sup> ( $T_{\text{cond}}$ ), we conclude that a cold trap on WASP-76b



**Fig. 2 | Retrieved elemental composition of the atmosphere of WASP-76b relative to iron.** **a**, Elemental-abundance ratios relative to iron on WASP-76b compared with proto-solar<sup>5</sup>, stellar<sup>18</sup> and CI chondrites<sup>20</sup> compositions. Other than alkali metals and ultra-refractories, elements in the atmosphere of WASP-76b follow a notably similar trend as these primitive, unprocessed materials. For comparison, we also show the crustal composition of Earth<sup>19</sup>, which is highly processed and—unsurprisingly—poorly represents the data. Alkali metals are measured to be underabundant owing to being strongly ionized at the low pressures and high temperatures examined and thus probably do not represent the true atmospheric abundance. **b**, Measured refractory abundance

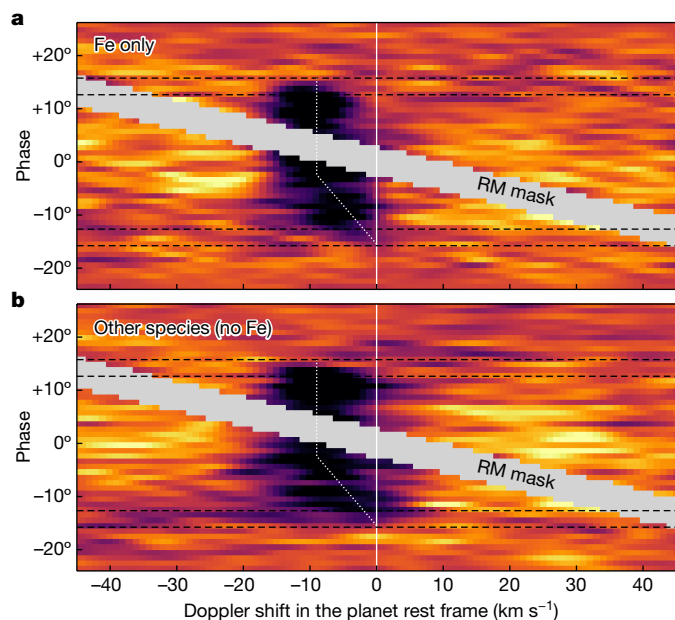
ratios in the atmosphere of WASP-76b relative to proto-solar. The atmospheric enrichment of WASP-76b is near proto-solar (blue line) for elements with condensation temperatures up to about 1,550 K, before sharply transitioning and showing orders of magnitude depletion levels. With their higher condensation temperatures, ultra-refractory elements (Ti, Sc, Al) probably seem depleted because of being cold-trapped on the colder nightside of the planet. From the near-proto-solar abundance of V/Ba/Ca and severe depletion of Ti/Al/Sc, we can constrain the cold-trap temperature to be between roughly 1,520 and 1,580 K. All error bars represent  $1\sigma$  uncertainties.

must cause these highly refractory elements to be removed from the gas phase of the upper atmosphere<sup>21</sup>. Indeed, the measured abundance ratios relative to solar show a steep transition, with elements having condensation temperatures below about 1,550 K being roughly in line with solar and elements with  $T_{\text{cond}} \geq 1,550$  K being substantially depleted (Fig. 2b). Such a cold-trap mechanism has also been proposed on the similar ultra-hot Jupiter WASP-121b ( $T_{\text{eq}} \approx 2,350$  K) to explain the lack of Ti and TiO (refs. 3,4). Measuring a wide range of refractory abundance ratios in other giant exoplanets progressively hotter and colder than WASP-76b will be necessary to better understand the condensation sequence of mineral species in exoplanet atmospheres and indirectly examine cloud compositions<sup>22,23</sup>. For example, if nucleation is efficient in hot Jupiter atmospheres, we may expect their transmission spectra to show sharp transitions as mineral species are depleted one by one as a function of their condensation temperature<sup>24</sup>. Similarly, hotter planets would progressively ‘unlock’ elements, as in the case of the similar but even warmer ultra-hot Jupiter WASP-189b ( $T_{\text{eq}} \approx 2,650$  K), which shows absorption from Ti and Sc species<sup>13</sup>. On the other hand, the condensation sequence of different mineral species is probably less straightforward, given that heterogeneous nucleation is strongly dependent on the availability of cloud seed particles, their nucleation rates<sup>25</sup> and their gravitational settling timescales<sup>26</sup>. Identifying and quantifying similar abundance transitions at which species become depleted in relation to their condensation temperatures will also serve as an indirect exploration of the nightside temperature profiles of hot Jupiters<sup>8</sup>.

Although most elements on WASP-76b are either consistent with proto-solar and stellar abundances or substantially depleted owing

to ionization/condensation, Cr, Ni and V do show differences at the approximately  $2-3\sigma$  level, even when using different model parameterizations (Extended Data Fig. 7). If representative of the bulk envelope, abundance ratios that deviate from proto-solar/stellar can shed light onto the composition of materials accreted during the formation and evolution of WASP-76b. For example, one possibility is that WASP-76b accreted a notable proportion of differentiated, non-solar-like material throughout its history. To quantify this hypothetical scenario, we use a toy model that calculates the final abundances in the atmosphere of WASP-76b post-accreting a body with a given composition and mass (see Methods). We find that measured refractory elemental ratios can be reasonably well matched if, for example, WASP-76b accreted Mercury-like material with a total mass half that of Earth (see Extended Data Fig. 8). If caused by a single accretion occurrence, such an event during the evolution of WASP-76b could resemble the giant collision that Jupiter has been considered to have undergone to explain its diluted core<sup>27</sup>. However, we notably are unable to perfectly match all measured abundance ratios on WASP-76b with the scenarios explored by our model (see Methods), highlighting the importance of precisely measuring a wide range of elements to help constrain the composition of any accreted material. We also cannot rule out that the inferred non-solar elemental-abundance ratios are instead the result of other physical or chemical atmospheric processes. For instance, if Ni spectral lines examine deeper, colder atmospheric layers on average, it may be that all other species are slightly more ionized, making Ni seem more abundant in comparison. Meanwhile, Cr and V could plausibly also seem underabundant because of being partially condensed<sup>28</sup> or partly bound in other compounds (for example, CrO, VO<sub>2</sub>; see Extended Data Fig. 6).





**Fig. 3 | Rest-frame absorption signals on WASP-76b.** **a**, Cross-correlation signature of iron absorption showing a leftward ‘kink’ (dotted white line) occurring during transit. The dark trail represents the absorption signature of the planet, which is only present during the transit event. The dashed black lines represent ingress and egress contact points, whereas the grey-shaded area is a region in radial velocity space not considered in the analysis owing to overlapping with the stellar Rossiter–McLaughlin effect (‘RM mask’ region). A uniform and static atmosphere on WASP-76b would produce an absorption trail centred at zero. Instead, the planetary signature is kinked, starting at a Doppler shift close to zero and becoming progressively more blueshifted during the first half of the transit before holding steady. **b**, Same as **a** but for other species included in the retrieval excluding iron (Na, K, Li, Mn, Cr, Mg, Ni, V, VO, Ba<sup>+</sup>, Ca and Ca<sup>+</sup>), also showing a similar asymmetric absorption trace (see Extended Data Fig. 9 for the individual contribution of each species). The asymmetry of the signal is not unique to iron, indicating that a global process affecting many elements similarly, such as an uneven temperature or distribution of clouds on the eastern and western limbs of WASP-76b, is causing this behaviour.

Finally, the MAROON-X data also show a ‘kinked’ signal in the phase-resolved absorption of both iron and other species in the atmosphere of WASP-76b, with the absorption trails being progressively more blueshifted over the first half of transit before holding steady (Fig. 3). Previous observations of WASP-76b showing such a kink in iron absorption were interpreted as iron condensing out of the atmosphere from the dayside to the nightside as the cooler morning terminator rotated into view<sup>7</sup>. Given that we detect several species with differing condensation temperatures that all behave similarly (see also Extended Data Fig. 9), we conclude that, instead of rainout, a global process affecting most species systematically must be responsible. We thus argue that a substantial temperature asymmetry<sup>16</sup> and/or unevenly distributed high-altitude optically thick clouds<sup>29</sup> between the east and west terminators of the planet are more probable explanations. In these scenarios, most of the signal originates from the east terminator of WASP-76b, with contributions from the western limb damped resulting from a smaller scale height and/or clouds<sup>30</sup>. Although condensation of iron and other species may indeed be occurring from the dayside to the nightside, it is unlikely to be the sole cause of the observed asymmetric absorption trails of species in the atmosphere of WASP-76b.

**Online content**

Any methods, additional references, Nature Portfolio reporting summaries, source data, extended data, supplementary information,

acknowledgements, peer review information; details of author contributions and competing interests; and statements of data and code availability are available at <https://doi.org/10.1038/s41586-023-06134-0>.

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## Methods

### MAROON-X observations of WASP-76b

WASP-76b is a tidally locked ultra-hot Jupiter orbiting the F7-type star WASP-76 ( $V = 9.5$ ,  $T_{\text{eq}} = 6,329 \pm 65$  K). With a short orbital period of only 1.81 days, WASP-76b has a mass slightly less than that of Jupiter ( $0.894^{+0.014}_{-0.013} M_{\text{Jup}}$ ) but a radius that is much larger ( $1.854^{+0.077}_{-0.076} R_{\text{Jup}}$ )<sup>7</sup>. We used the MAROON-X spectrograph to obtain three continuous time series observations of WASP-76b as it transited its host star. The observations were part of Gemini North programmes GN-2020B-Q-122 (principal investigator: S.P.) and GN-2021B-Q-138 (principal investigator: F.D.). MAROON-X covers the full 491–678-nm wavelength range across 34 spectral orders in its ‘blue’ arm and the full 647–921-nm wavelength range across 28 spectral orders in its ‘red’ arm. As the detectors in the blue and red arms have different readout times, we integrated on the blue arm 40 s longer per exposure to maintain a constant cadence for the first two transits. The third transit maintained a fixed exposure time of 480 s for both channels. With the exception of the first transit, which misses the first 2–3 min of ingress, all three nights cover the entire transit, including start and end baseline, all under photometric sky conditions. For the analysis, we use the most up-to-date full-orbital solution<sup>7</sup>, revised with the latest ephemeris from TESS<sup>31</sup>. WASP-76 is part of a binary system<sup>32</sup>, with a companion near the limit of the 0.77'' field of view of MAROON-X. However, as the companion is effectively stationary in velocity space over the course of the observations, any starlight contamination is removed along with the primary star during the data detrending.

### Data reduction and detrending

We used the standard MAROON-X pipeline<sup>33</sup> to extract detector images into one-dimensional wavelength-calibrated spectra, order by order for each exposure. The pipeline outputs a  $N_{\text{exp}} \times N_{\text{order}} \times N_{\text{pixel}}$  high-resolution spectral time series for each channel of each transit sequence, with  $N_{\text{exp}}$ ,  $N_{\text{order}}$  and  $N_{\text{pixel}}$  being the number of exposures, number of spectral orders and number of pixels for each detector, respectively. The data contain the stellar spectrum, absorption from the transiting planet, Earth’s tellurics and noise. Our objective is to remove the stellar and telluric contributions that dominate the observed spectra, without removing the relatively much fainter planetary signal. This is possible to achieve because the absorption features of WASP-76b experience a change in radial velocity of about 100 km s<sup>-1</sup> over the duration of a full transit, whereas telluric and stellar spectral lines remain relatively stationary in wavelength. We apply a PCA-based reduction algorithm to remove all unwanted wavelength-stationary contributions and uncover the atmospheric signature of WASP-76b. The data-detrending procedure broadly follows ref. 17, with some minor alterations. Here we provide a summary of the steps applied to each order of the transit time series independently.

1. Outliers in the spectroscopic time series deviating by more than 5 $\sigma$  from the median of their spectral channel are flagged and corrected through interpolation of neighbouring wavelengths.
2. All observed spectra are aligned in velocity space to remove both the barycentric motion of Earth and the reflex motion of WASP-76. This procedure ensures that stellar lines are well wavelength-aligned and can be properly removed.
3. Blaze and throughput variations are removed by bringing each spectrum in the time series to the same continuum level. This consists of first dividing each pixel by its median in time, smoothing the residuals with a low-pass filter and then dividing the original spectra by this filter<sup>3,34,35</sup>. The smoothing is done by passing a median box filter with a width of 501 pixels that is subsequently convolved by a Gaussian with a standard deviation of 100 pixels. The large box kernel width and Gaussian standard deviation used relative to the 1.1 km s<sup>-1</sup> average velocity dispersion per pixel of MAROON-X are necessary to avoid altering the signal of the exoplanet<sup>3</sup>.

4. Each spectrum is divided by a second-order polynomial fit of the median of all out-of-transit exposures.
5. A PCA reconstruction of the first ten principal components summed is divided out of each spectrum. The choice to remove ten principal components is relatively arbitrary, as removing anywhere between approximately three and 20 components does not markedly change our results. At least a few must be removed to properly eliminate stellar and telluric contributions, but too many would remove the atmospheric signal of the planet.
6. Spectral channels that have a standard deviation greater than four times that of the median of their spectral order are masked and removed from the analysis.
7. An uncertainty of the form  $\sigma_n = (aF_n + b)^{0.5}$  is estimated for every flux value. Here  $F_n$  is the measured flux for each data point  $n$ , whereas  $a$  and  $b$  are fitted through a likelihood minimization of the data residuals after a PCA reduction. This is the same procedure used in refs. 3,35, to which the reader is referred for more details.

An overview of these steps applied on an example spectral order is shown in Extended Data Fig. 1. Although not all spectral orders are equal in terms of telluric and stellar line content, and may benefit from different levels of detrending, we refrain from optimizing the number of principal components removed in each order so as to not introduce further degrees of freedom that could potentially bias the analysis<sup>36,37</sup>. The 760–770-nm wavelength range dominated by deep telluric absorption is masked out and not considered in the analysis. The telluric-free and stellar-free cleaned data residuals and their associated uncertainties are then used to uncover the atmospheric trace of WASP-76b hidden in the noise by means of cross-correlation with atmosphere models.

### Atmospheric modelling

To properly study the atmosphere of WASP-76b, it is crucial to have a representative model. For our analysis, we generate synthetic high-resolution transmission spectra using the SCARLET framework<sup>17,38–42</sup> with cross-sections computed using HELIOS-K (<https://chaldene.unibe.ch/>)<sup>43,44</sup>. Opacities used for this work include AIO (ref. 45), CrH (refs. 46,47), K (ref. 48), Na (ref. 49), TiO (ref. 50), VO (ref. 51) and all other atoms and ions available in the VALD<sup>52</sup>, Kurucz<sup>53</sup> or NIST<sup>54</sup> databases. Our main results use the VALD database, but we use the Kurucz and NIST databases to verify the retrieved refractory abundance ratios (Extended Data Fig. 7a). The absorption cross-sections for a subset of these species that are of interest for this work are shown in Extended Data Fig. 5. Although the atomic cross-sections do not include pressure broadening (other than Na and K), this is unlikely to have a notable impact owing to low (sub-millibar) pressures at which metallic lines are investigated in this dataset. SCARLET models also include collision-induced absorption from H<sub>2</sub>–H<sub>2</sub> and H<sub>2</sub>–He interactions<sup>55</sup>, as well as Rayleigh scattering<sup>38</sup>. An optically thick grey cloud deck at a freely parameterized cloud-top pressure  $P_c$  is also included. Although H<sup>-</sup> is probably the dominant continuum opacity source on WASP-76b, because H<sup>-</sup> opacity<sup>56,57</sup> is nearly flat over the MAROON-X bandpass (Extended Data Fig. 2) and therefore affects modelled spectra similarly to a cloud deck, we opt to fit for  $P_c$  rather than H<sup>-</sup> because of its more easily interpretable effect on the resulting transmission spectrum. We use H<sub>2</sub>, He and H as filler gases such that the volume-mixing ratio of all species included in the model always sum to one in every atmospheric layer. The relative amounts of H<sub>2</sub>, He and H added depend on the temperature–pressure profile and is determined using FastChem<sup>58</sup>, with H<sub>2</sub> typically being most abundant deeper in the atmosphere and H dominating at higher altitudes (see Extended Data Fig. 6b). Accounting for the transition between atomic and molecular hydrogen is important, as it traces a notable change in mean molecular weight in the atmosphere. Models are generated at a resolution of  $R = 250,000$  over the full MAROON-X wavelength range and later broadened to the instrumental resolution of  $R = 85,000$ .

## Cross-correlation analysis

We perform a classical cross-correlation analysis to search for individual species in the atmosphere of WASP-76b. Given cleaned data residuals  $d_n$  with uncertainties  $\sigma_n$  (see Extended Data Fig. 1) and a model transmission spectrum  $m_n$ , the cross-correlation function is defined as

$$\text{CCF}(v) = \sum_n \frac{d_n m_n(v)}{\sigma_n^2}$$

and is calculated at a given radial velocity shift ( $v$ ) for each spectrum of the time series by summing over every wavelength bin ( $n$ ). Computed for a wide range of velocities and summing all spectral orders, this produces a CCF time series in which the absorption from the atmosphere of the planet should follow a distinct trail along the expected orbital path<sup>11</sup>. One can then integrate in time by phase-folding the data, summing contributions from all exposures for different combinations of  $K_p$  and  $V_{\text{sys}}$ , effectively generating a velocity–velocity map that should produce a peak near the known position of the planet if the modelled atmosphere matches the true signal (see Fig. 1). Deviations from the expected  $K_p$  and  $V_{\text{sys}}$  may then be indicative of dynamical or asymmetric chemical processes at play in the atmosphere of WASP-76b (ref. 13). Signal-to-noise values for a given map are computed by dividing summed CCF values by the standard deviation of all values away from the expected position of the signal of WASP-76b. As in previous work for this target<sup>30</sup>, values falling within  $\pm 10 \text{ km s}^{-1}$  of  $V_{\text{sys}}$  are masked out (see Fig. 3) and not included in the CCF or retrieval analyses to avoid contamination from the Rossiter–McLaughlin effect<sup>59,60</sup>.

## Bayesian inference framework

Although a cross-correlation analysis is the method of choice for investigating whether a given element or molecule is present in the atmosphere being examined, it provides only limited quantitative information on their abundances. We detect iron on WASP-76b for example, but how much iron is there? To further characterize the atmosphere of WASP-76b, we perform a Bayesian high-resolution retrieval analysis using the likelihood prescription of refs. 3, 35 given by

$$\ln(L) = -\frac{N}{2} \ln \left( \frac{1}{N} \sum_n \frac{(d_n - \alpha m_n)^2}{\sigma_n^2} \right),$$

in which  $N$  is the total number of data points and  $\alpha$  is a model-scaling parameter. We do not include a noise-scaling term  $\beta$ , which has a negligible impact on large, well-behaved datasets<sup>35</sup>.

To accurately infer atmospheric parameters, we mimic the effects that both the rotation of WASP-76b and the reduction steps applied to the data have on the underlying signal to the model. For this, the model  $m_n$  is convolved with a rotational broadening kernel assuming tidally locked rotation, projected in time for a given combination of  $K_p$  and  $V_{\text{sys}}$  and multiplied by a stellar quadratic limb-darkening model. We follow the approach outlined in ref. 61 to compute the limb-darkening model (see their equations 16–22) but assume the planet to be uniform. We use an impact parameter  $b = 0.027$  and quadratic limb-darkening coefficients  $\mu_1 = 0.393$  and  $\mu_2 = 0.219$  (ref. 7). The model is then injected in the removed PCA reconstruction of the data (step 5 of the data reduction) and the PCA algorithm is run once again on this mock dataset including the injected model<sup>17</sup>. The output is a modified  $m_n$  better representative of what we would expect the true signal to be after the application of the detrending algorithm. This modified  $m_n$  is used as an input to the equation and ensures that accurate parameters can be inferred<sup>62,63</sup>.

For the retrieval, we fit the log volume-mixing ratio ( $\log_{10}\chi$ , prior between  $-14$  and  $0$ ) of all 20 species simultaneously, making no a priori assumptions on the chemistry and instead setting abundances to be

uniform in pressure. We also simultaneously fit for the temperature structure ( $T_0 - T_9$ , prior between  $0$  and  $7,000 \text{ K}$ ), log continuum pressure ( $\log_{10}P_c$ , prior between  $-8$  and  $2 \text{ bar}$ ), log scale factor ( $\log_{10}\alpha$ , prior between  $-2$  and  $2$ ) and orbital and systemic velocities  $K_p$  (prior between  $166.5$  and  $226.5 \text{ km s}^{-1}$ ) and  $V_{\text{sys}}$  (prior between  $-21.1$  and  $18.9 \text{ km s}^{-1}$ ). The temperature–pressure profile is freely fit as parameterized in ref. 17 (see their equation 11) using ten temperature points uniformly distributed in log pressure between  $10^{-10}$  ( $T_0$ ) and  $1$  ( $T_9$ ) bar and a smoothing prior  $\sigma_s = 160 \text{ K dex}^{-2}$ . All parameters are fit simultaneously to ensure that all degeneracies are marginalized in the parameter inference<sup>62</sup>. The Bayesian inference and parameter exploration is done using the emcee<sup>64</sup> Markov chain Monte Carlo Python package. We run the Markov chain Monte Carlo for 30,000 steps with 200 walkers and use the last 10,000 steps (2,000,000 samples) to compute posterior distributions. The retrievals are computationally expensive, needing approximately three months to run when using 18 CPU cores in parallel. We test convergence by using the Gelman–Rubin diagnostic for each chain. We also supplement our retrieval with the white-light-curve transit depth from existing HST/STIS data<sup>65</sup> to anchor models at the correct planetary radius.

To estimate total elemental abundances in the atmosphere of WASP-76b, vanadium is calculated from the sum of V and VO, titanium from Ti and TiO, calcium from Ca and Ca<sup>+</sup>, chromium from Cr and CrH, aluminium from Al and AlO, barium from Ba<sup>+</sup> and other species from their neutral forms (for example, magnesium from Mg). In most cases, the species included in the retrieval make up the dominant portion of the expected chemical inventory at the pressures studied (Extended Data Fig. 6). We note that a limitation of our atmospheric model is that it does not account for abundance variations in pressure owing to ionization, which can potentially bias inferred volume-mixing ratios, as is evidently the case for the alkali metals. For example, iron may seem less abundant than in reality because of partial ionization, especially given the tentative Fe<sup>+</sup> signal observed in the data—although this signal probably originates from higher altitudes than those examined by neutral Fe. Similarly, retrieved abundances might be underestimated owing to partial condensation or from not considering further species. For example, V and Ti may be partly bound in VO<sub>2</sub> and TiO<sub>2</sub> molecules that we cannot measure because of the lack of suitable line lists<sup>66</sup>.

We note that the retrieved velocity parameters are driven by the strongest absorber (that is, Fe) and that other species can have small relative offsets in  $K_p$  and  $V_{\text{sys}}$ . However, we refrain from fitting  $K_p$  and  $V_{\text{sys}}$  freely for each species to obtain meaningful abundance ratios. The extreme examples, H, O and Fe<sup>+</sup>, are excluded from the retrieval as their peaks are substantially offset from the other detected species (Fig. 1), potentially because of strong vertical winds or outflows<sup>30,59</sup>. Similarly, some absorption features investigate up to non-hydrostatic layers of the atmosphere<sup>4,18</sup> (for example, Ca<sup>+</sup>) and may not be well represented by our model, which could bias the abundances we retrieve. We also note that our retrieval is based on a one-dimensional atmospheric model and only captures average properties of the inherently three-dimensional planet, which could result in further biases of inferred parameters<sup>67</sup>. However, ref. 61 shows that the retrieved iron abundance on WASP-76b from one-dimensional retrievals is roughly consistent with results from two-dimensional retrievals. Nevertheless, we chose not to overinterpret retrieved absolute abundances and rather focus on relative abundance measurements, which can be constrained much more precisely and are less sensitive to systematic biases<sup>3</sup>.

We recover orbital velocity parameters  $K_p = 180.7 \pm 0.6 \text{ km s}^{-1}$  and  $V_{\text{sys}} = -7.0 \pm 0.1 \text{ km s}^{-1}$ . The discrepancy between these and the known literature values results from the planetary rotation, day-to-night atmospheric winds and an artefact of fitting a kinked absorption trail with a traditional planetary orbit and should be interpreted as an indicator of an asymmetry in the eastern and western limbs of WASP-76b (ref. 7). Similarly, the retrieved continuum pressure level (in bar) of  $\log_{10}P_c = -2.21_{-0.28}^{+0.27}$  and temperature profile are an average

of both hemispheres, when—in reality—this value is probably different on each terminator<sup>61</sup>. We also recover a scaling factor of  $\alpha = 0.52^{+0.08}_{-0.09}$ , which may be decreasing the amplitude of the models to compensate for an evening terminator signal that is damped owing to a colder temperature and/or clouds<sup>16,29,30</sup>. The retrieved temperature profile is consistent with previous findings of the upper atmosphere of WASP-76b being hotter than expected from model predictions<sup>59,61,68,69</sup>. We estimate the total refractory metallicity of WASP-76b (that is, the sum of all retrieved absolute abundances) to be  $\log_{10}(\sum X_i) = -3.45^{+0.30}_{-0.32}$  slightly enriched relative to both solar<sup>5</sup> (+0.62 dex) and stellar<sup>18</sup> (+0.28 dex) compositions (Extended Data Fig. 5b). We further test the robustness of our results by performing separate retrievals using both different line lists and temperature-structure parameterizations. We find that, regardless of the atomic line lists used (VALD<sup>52</sup>, Kurucz<sup>53</sup> or NIST<sup>54</sup>) and assumed temperature–pressure profile form (free<sup>17</sup>, Guillot<sup>70</sup> or isothermal), the inferred elemental-abundance ratios are all relatively consistent (Extended Data Fig. 7a). However, we warn that fitted temperature profiles with less flexibility can lead to the shape of the temperature structure being overconstrained (Extended Data Fig. 7b).

### Notable non-detections

We notably do not detect Ti, TiO, Ti<sup>+</sup>, V<sup>+</sup>, Sc, CrH, AlO or Al (Extended Data Fig. 3). A previous analysis of ESPRESSO data<sup>30</sup> has also shown Sr<sup>+</sup> and Co detections on WASP-76b; however, MAROON-X does not cover the bluer wavelengths at which these species have most of their absorption features. We therefore cannot confirm or exclude their presence in the atmosphere of WASP-76b owing to lack of sensitivity. On the other hand, MAROON-X can very effectively investigate Ti, TiO, Sc and AlO. The absence of detections and derived strict upper limits for these species indicate that they are indeed depleted from the upper atmosphere of WASP-76b. Given that Ti, Sc and Al compounds all seem to be heavily underabundant, and that these elements have a higher condensation temperature than other metals that are readily detected, we concluded that these species must be condensed out of the upper atmosphere and be cold-trapped on the nightside.

### VO detection

The detection of VO on WASP-76b is, to our knowledge, the first time this molecule has been unambiguously observed using high-resolution spectroscopy on an exoplanet. VO, along with TiO, have long been theorized to be drivers for thermal inversions in ultra-hot Jupiter atmospheres owing to the very strong opacity of these molecules at optical wavelengths<sup>6,71</sup> (see also Extended Data Fig. 2). Despite this interest, VO has proved to be notoriously difficult to robustly detect in exoplanet atmospheres. For example, VO was tentatively observed on a different ultra-hot Jupiter, WASP-121b (refs. 72–74), but further observations later refuted this claim<sup>75–77</sup>. The ambiguity arises because, at low resolution, the broad absorption bands of VO can be hard to distinguish from other opacity sources or from systematic noise in the data<sup>34,75</sup>. High-resolution spectroscopy can mitigate these difficulties, but is also more sensitive to the accuracy of the theoretical opacities used to generate the modelled spectra<sup>78</sup>. Owing to the most up-to-date high-temperature VO line list being known to have inaccuracies in some wavelength ranges<sup>79,80</sup>, we confirm our results using an independent analysis method applied on a different dataset. We analysed existing WASP-76b ESPRESSO<sup>81</sup> transit data<sup>7</sup> (programme 1102.C-744) using the same methodology outlined in ref. 13 and also recover a clear VO detection (Extended Data Fig. 4b). This means that the VO signal on WASP-76b is especially strong such that we are able to detect it, even with an imperfect line list. We warn that our inferred abundance of VO is only as good as the cross-sections used, which may bias our results<sup>62</sup>. We note that a new VO line list is, at present, being computed by the ExoMol group (J. Tennyson, private communication), which—once released—may

increase our capabilities of detecting VO on other less favourable exoplanets.

### Accretion model

We explore theoretical scenarios in which the atmospheric composition of WASP-76b may have been enriched by the accretion of non-solar, differentiated materials. For this, we build a chemical composition toy model that takes the following parameters as input: the envelope mass of WASP-76b, its pre-merger volatile and refractory abundances and the accreted body's core and mantle masses and their chemical composition. The model then outputs the predicted post-merger atmospheric composition, assuming the accreted body is well mixed into the envelope of WASP-76b and to the low pressures examined by these data (for example, through strong winds both predicted by general circulation models<sup>82</sup> and observationally inferred<sup>59</sup> to exist on WASP-76b). The chemistry of the core of the accreted object is sampled from the range of abundances found in iron meteorites<sup>83</sup>, which are thought to be representative of the core composition of Mercury<sup>84</sup>. For the mantle, we explore different chemical compositions, namely those of CI chondrites<sup>20</sup>, CB meteorites<sup>85</sup> and the surface of Mercury<sup>84</sup>. On the basis of a thorough exploration of parameter space, we find no single model that perfectly matches all of the observational data. Models accreting at least about 0.1–0.5  $M_{\oplus}$  ( $M_{\oplus}$  being the mass of the Earth) of purely iron-meteorite-like material (resembling the core of Mercury) can fit well Cr, Mg, Ni, V and Ba but not Mn or Ca. Adding material representative of the surface of Mercury on top allows us to fit Ca as well without losing the other good fits but not Mn, whose abundance in the accreted materials is always too low and thus is never fitted. Keeping the iron core but accreting a solar-composition mantle instead is practically identical to the pure-core case. The assumed initial C/O ratio and metallicity of the parent-body atmosphere only affect the mass needed to be accreted to match observed abundance ratios and otherwise have no impact on the final composition. The accreted mass is inversely correlated with the abundance of Ni in the accreted material, in which assuming a high Ni abundance decreases the total mass necessary to obtain the same enrichment levels. We note that, although our model assumes the enrichment of WASP-76b to have come from a single accreted body, other scenarios in which the equivalent material instead originated from upwards dredging of a dilute core<sup>86,87</sup> or an influx of smaller comets and meteorites are similarly plausible.

### Data availability

The MAROON-X data used in this work are available at [https://udemontreal-my.sharepoint.com/:f/g/personal/stefan\\_pelletier\\_umontreal\\_ca/EkYThK-JMKFHlclxy7RnlIABySi6V60HuZCOc\\_9m6lfe6Q?e=vGerbT](https://udemontreal-my.sharepoint.com/:f/g/personal/stefan_pelletier_umontreal_ca/EkYThK-JMKFHlclxy7RnlIABySi6V60HuZCOc_9m6lfe6Q?e=vGerbT). The ESPRESSO data used to confirm the VO detection are publicly available on Dace (<https://dace.unige.ch/dashboard/>). Source data are provided with this paper.

### Code availability

The MAROON-X reduction pipeline<sup>33</sup> used by the instrument team to perform the data extraction is public software available from Gemini at <https://github.com/GeminiDRSoftware/MAROONXDR>. The atmospheric modelling and retrievals use SCARLET<sup>17,40</sup>, HELIOS-K<sup>44</sup> (<https://helios-k.readthedocs.io>), FastChem<sup>58</sup> (<https://github.com/exoclimate/FastChem>), emcee<sup>64</sup> (<https://emcee.readthedocs.io/en/stable/>) and corner.py<sup>88</sup> (<https://corner.readthedocs.io/en/latest/>). The ESPRESSO data analysis was performed using Tayph<sup>66</sup> (<https://github.com/Hoeijmakers/tayph>). The main analysis routines written for this work and using the astropy<sup>89,90</sup>, matplotlib<sup>91</sup>, numpy<sup>92</sup>, scipy<sup>93</sup> and scikit-learn<sup>94</sup> Python libraries are available at [https://udemontreal-my.sharepoint.com/:f/g/personal/stefan\\_pelletier\\_umontreal\\_ca/EmXMwsPp2JFCnckKJNWkf7ABrEomi5EqmadxK4HofD7ltQ?e=73Gblp](https://udemontreal-my.sharepoint.com/:f/g/personal/stefan_pelletier_umontreal_ca/EmXMwsPp2JFCnckKJNWkf7ABrEomi5EqmadxK4HofD7ltQ?e=73Gblp).



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**Author contributions** S.P., B.B. and L.P. conceived the project. S.P. wrote the original MAROON-X observing proposal and the manuscript and carried out the analysis of the MAROON-X data, with B.B. and H.J.H. providing guidance. M.A.-D. performed the accretion modelling portion of the analysis. B.P. independently analysed the ESPRESSO data to confirm the VO detection. D.K., A.S., J.L.B. and J.S. assisted with the observational setup, carried out the observations and performed the MAROON-X data extraction. F.D., B.K., T.H. and A.C. acquired and contributed extra data for the project. L.B. implemented the FastChem equilibrium chemistry code in the modelling framework. A.Y.K., O.L. and N.C.-B. contributed to the stellar

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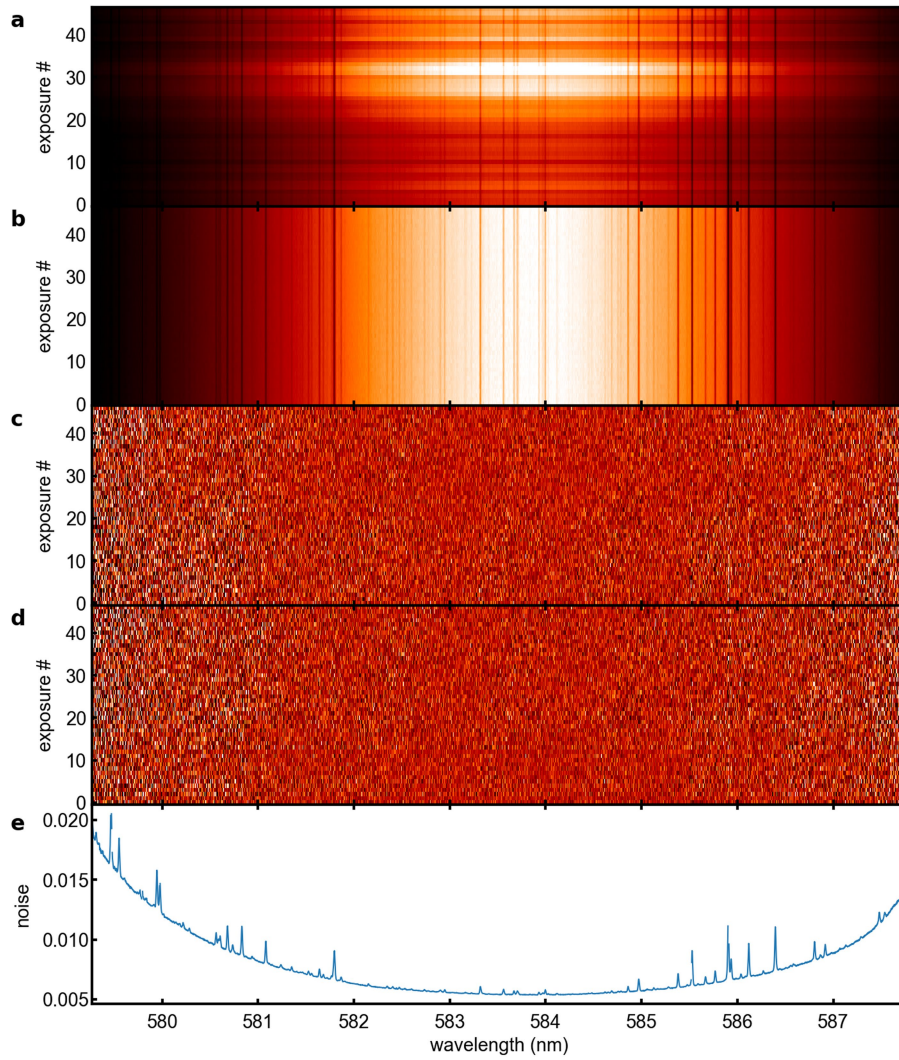
**Additional information**

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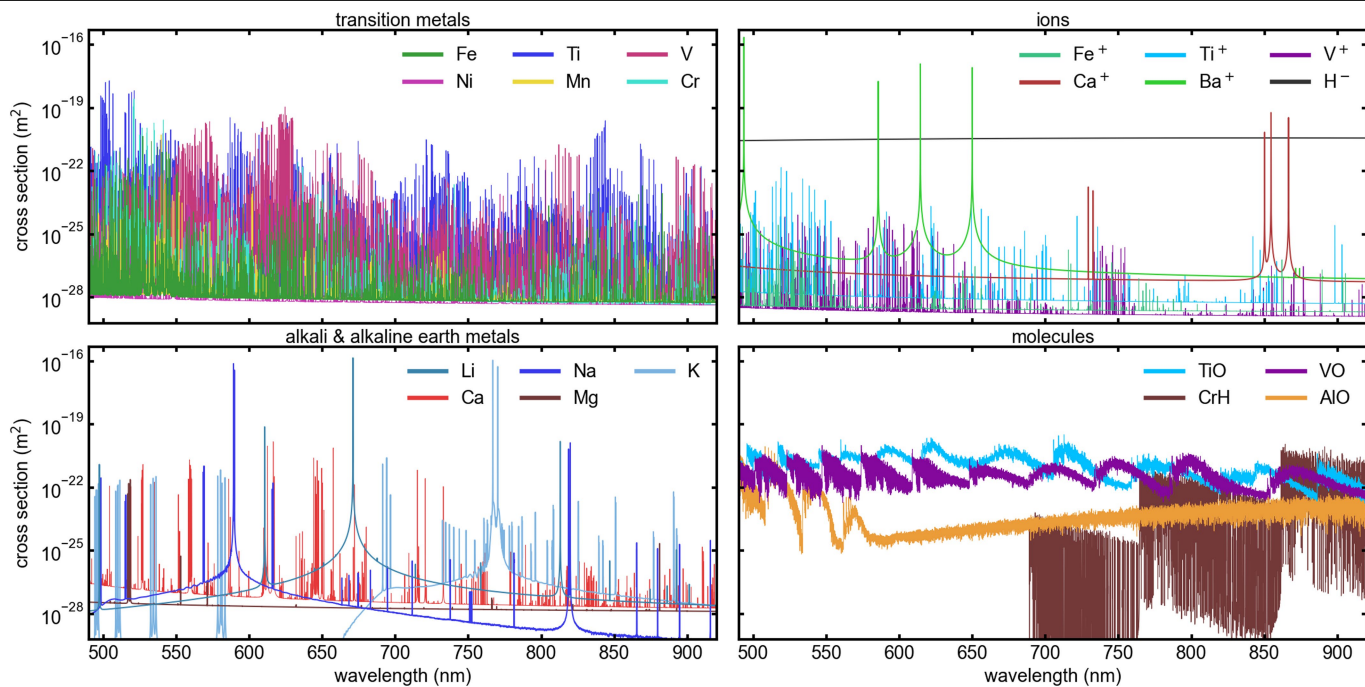
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**Extended Data Fig. 1 | Data detrending steps and noise model.** **a**, A single MAROON-X spectral order for the 5.3-h transit time series obtained on 12 September 2020. **b**, After continuum alignment. **c**, With the median out-of-transit spectrum divided out. **d**, Post-removal of stellar and telluric contaminants using PCA, with ten principal components removed in this case. **e**, The noise model used for an example exposure. The reduction steps take care of

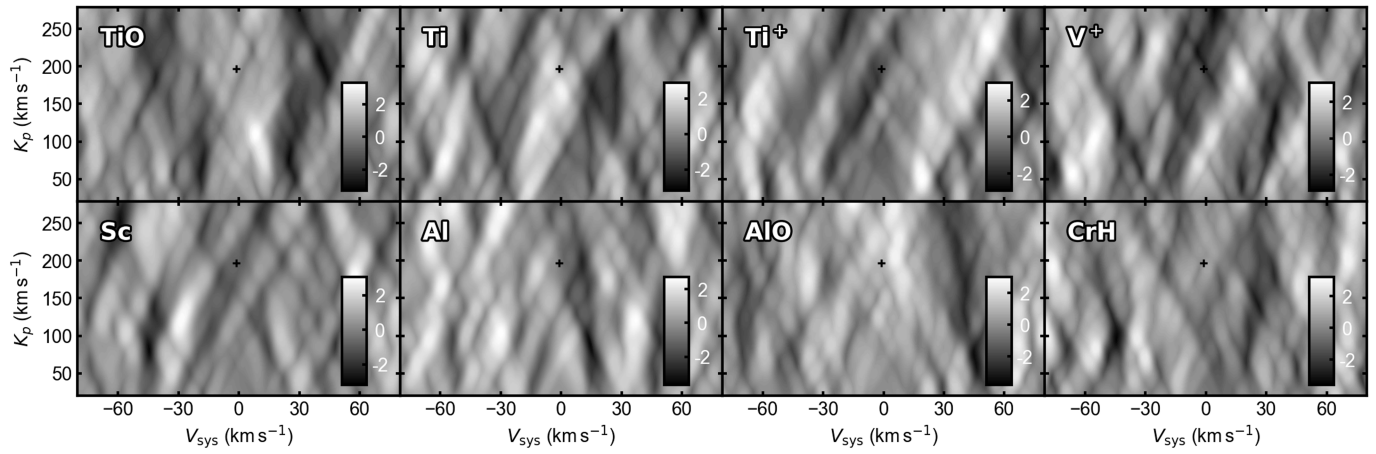
removing all telluric and stellar lines, as well as the continuum, whereas the noise model serves to downweigh regions of the spectrum that are noisier. The end product of the data reduction (**d**) contains the planetary trace buried in noise and serves as the input for the later cross-correlation and retrieval analyses.



**Extended Data Fig. 2 | Cross-sections of prominent metals, ions and molecules.** Values are computed for a temperature of 2,500 K and include  $10^{-4}$  bar pressure broadening where applicable. Atoms with only a single valence electron (Li, Na, K,  $\text{Ca}^+$  and  $\text{Ba}^+$ ) have few very strong lines, whereas most other metals tend to be composed of 'line forests' on top of a low

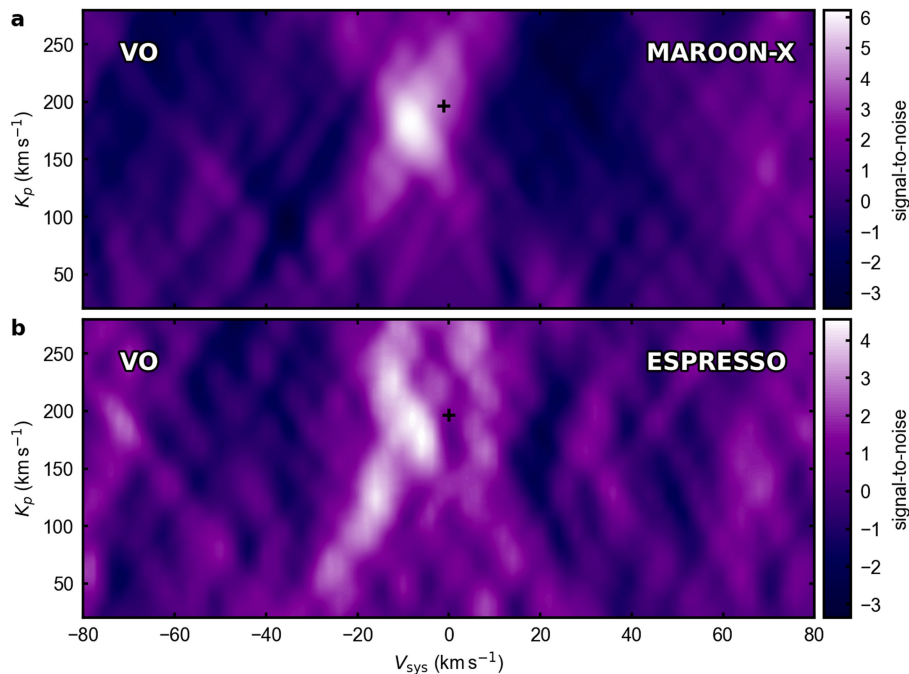
continuum. Molecules also have forests of spectral lines, but on top of distinct absorption bands. Meanwhile,  $\text{H}^-$  acts as a flat source of continuum opacity in this wavelength range, similar to a grey cloud deck. The product of the cross-section and volume-mixing ratio of a species in the atmosphere of WASP-76b determines how detectable it is.





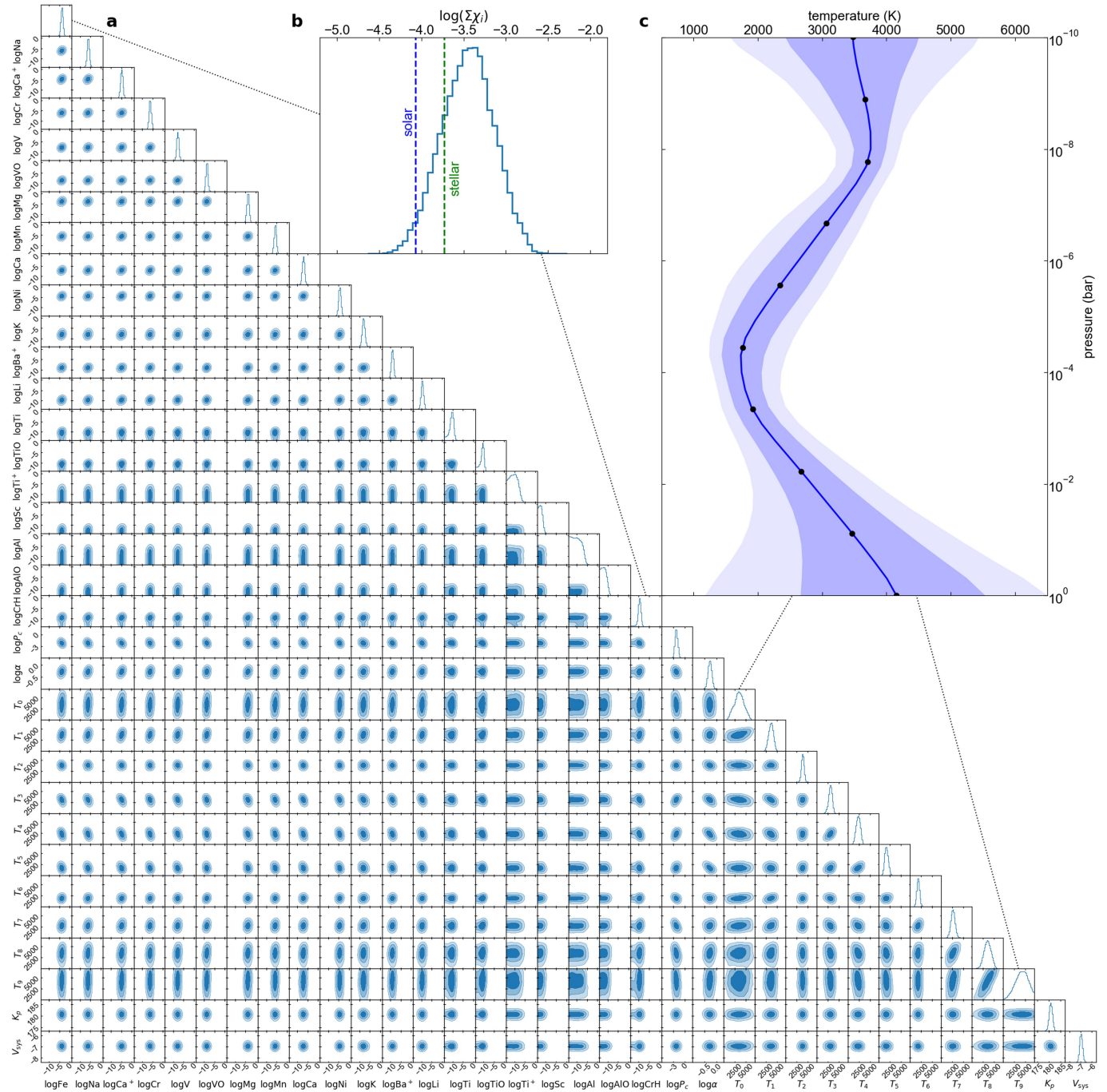
**Extended Data Fig. 3 | Cross-correlation results for species not detected in the atmosphere of WASP-76b.** Same as Fig. 1 but for species of interest that were notably not detected. Particularly noticeable is the absence of Ti, TiO, Sc and AlO, which MAROON-X should be able to easily detect. Other non-detections

are less surprising, as they either have few strong lines in the MAROON-X bandpass (Al) or are not expected to be particularly abundant from chemical-equilibrium predictions ( $\text{Ti}^+$ ,  $\text{V}^+$ , CrH).



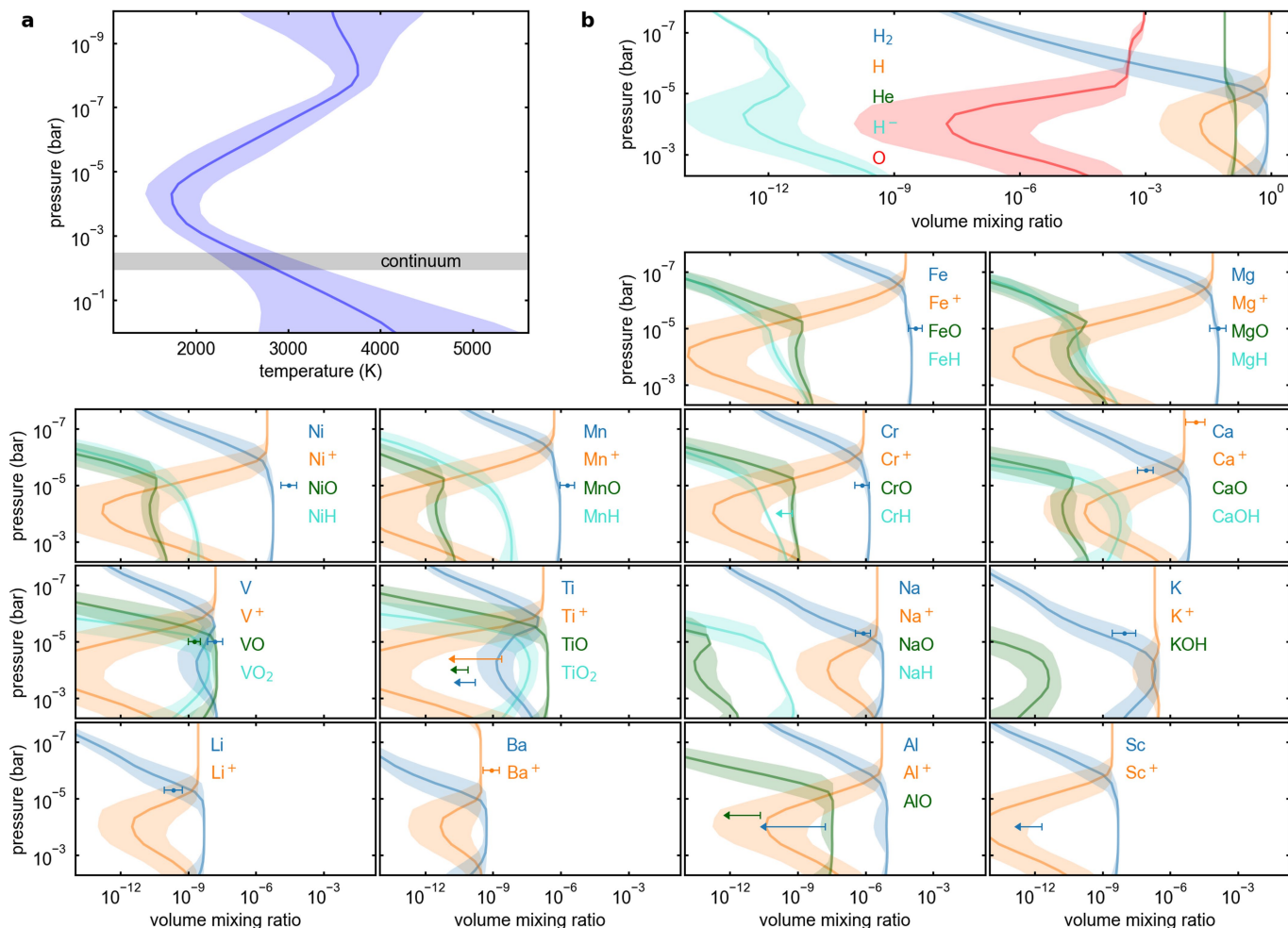
**Extended Data Fig. 4 | Independent confirmation of the VO detection on WASP-76b.** **a**, The VO detection with MAROON-X (490–920 nm,  $R = 85,000$ , three transits). **b**, The VO detection with ESPRESSO (380–790 nm,  $R = 140,000$ , two transits<sup>7</sup>). The ESPRESSO analysis was done with an independent framework<sup>13</sup>, using an independent model template, on a separate dataset, and also shows a clear VO signal slightly offset to the expected location (black

cross) that is consistent with the MAROON-X data. Differences in the overall detection strength and shape are because of MAROON-X having a lower resolution but a redder wavelength coverage, and three instead of two transits. The presence of a VO signal in both the MAROON-X and ESPRESSO transit datasets provides a high degree of confidence as to the robustness of the signal.



**Extended Data Fig. 5 | Retrieved constraints on the atmospheric and orbital properties of WASP-76b obtained from three MAROON-X transits. a,** Corner plot of the marginalized posterior distributions for the abundance of included species, cloud-top pressure  $P_c$  (in bar), scaling parameter  $\alpha$ , temperature of different atmospheric layers, Keplerian velocity  $K_p$  and systemic velocity  $V_{sys}$ . The shaded regions respectively depict the 39.3%, 86.5%, and 98.9% confidence

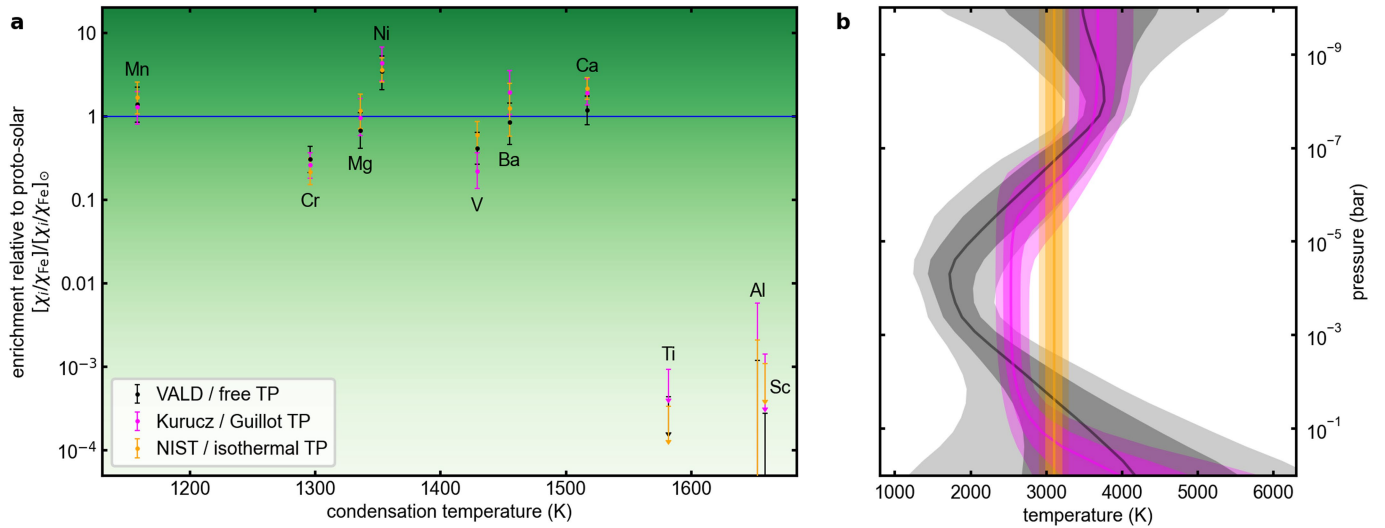
intervals. **b,** The sum of the volume-mixing ratio of individual metals, ions and molecules included in the model. The equivalent sums for solar and stellar compositions (dashed lines) are also shown for comparison, with WASP-76b being consistent with slightly (+0.28 dex) super-stellar. **c,** The resulting vertical temperature structure from the ten temperature points ( $T_0$ - $T_9$ , black dots), showing the presence of a stratosphere.



**Extended Data Fig. 6 | Chemical-equilibrium predictions of the atmospheric composition of WASP-76b.** **a**, The retrieved temperature–pressure profile and cloud-top continuum pressure. **b**, Chemical-equilibrium-abundance predictions<sup>58</sup> for a wide range of elements given the retrieved temperature–pressure structure (**a**) and assuming a stellar atmospheric composition<sup>18</sup>. Measured abundances for WASP-76b at the estimated investigated altitudes are also shown for comparison. Most refractory species (for example, Fe, Mg,

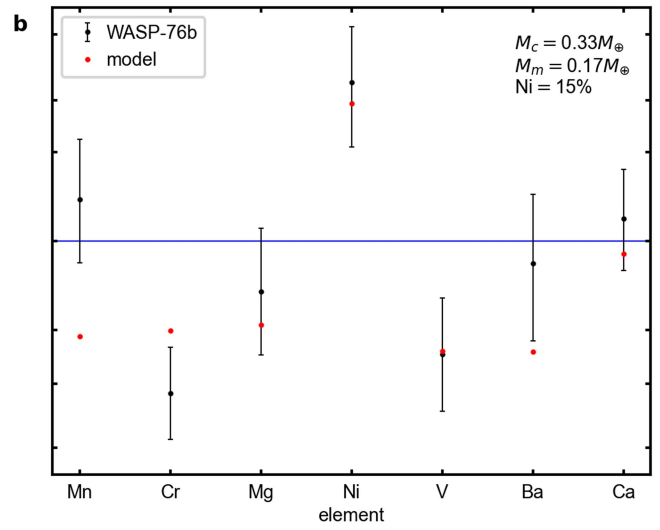
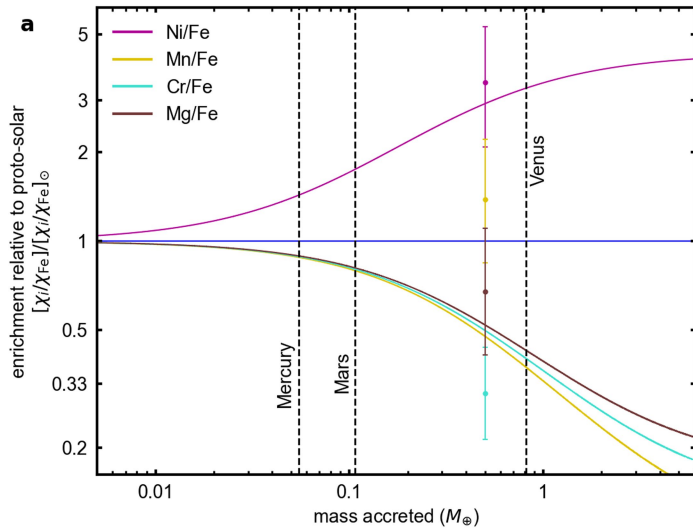
Ni, Mn and Cr) are not expected to be substantially ionized below the microbar level and are relatively well approximated by a constant-with-altitude volume-mixing-ratio model. With the exception of V and Ti, most elements are only expected to be bound in molecular form in trace amounts. Alkali metals, calcium and barium all ionize more readily and have deep spectral features and, thus, are expected to be substantially ionized at the lower pressures analysed. Error bars and shaded regions represent  $1\sigma$  uncertainties.





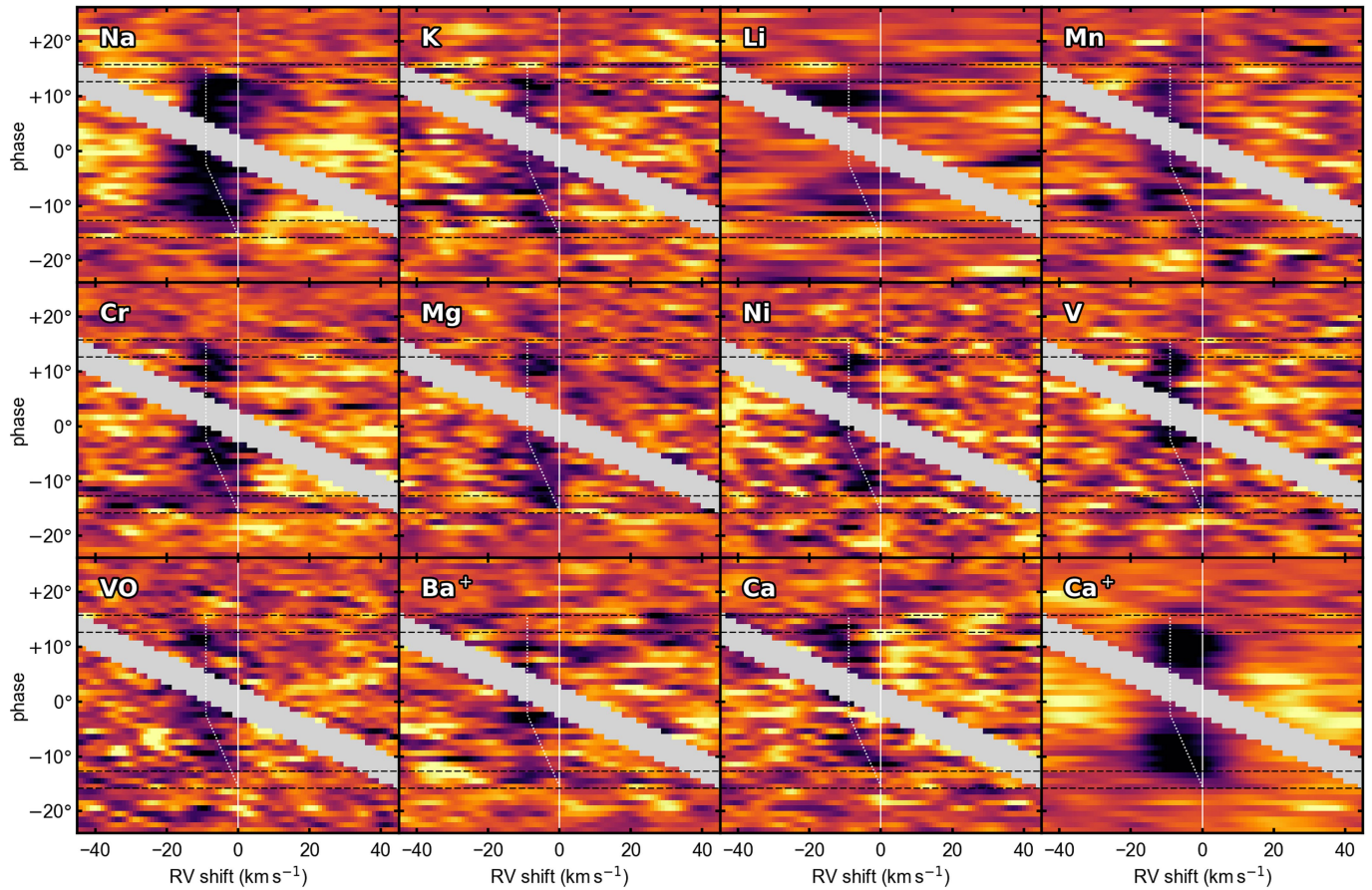
**Extended Data Fig. 7 | Comparison of results from different retrieval prescriptions.** **a**, Retrieved abundance ratios given different model parameterizations. **b**, Inferred temperature profile for each associated retrieval prescription. The three retrievals use separate combinations of the VALD<sup>52</sup>, Kurucz<sup>53</sup> and NIST<sup>54</sup> opacities and fitted free<sup>17</sup>, Guillot<sup>70</sup> and isothermal temperature–pressure profiles. Despite using differing temperature–structure

prescriptions and opacities, the recovered abundance ratios in all three retrievals are consistent within uncertainties. However, assuming an isothermal or Guillot profile can overconstrain the temperature structure. Error bars represent  $1\sigma$  uncertainties. Shaded regions represent  $1\sigma$  and  $2\sigma$  contours.



**Extended Data Fig. 8 | Accretion toy model exploring the scenario of WASP-76b accreting a body with a Mercury-like composition.** **a**, The change in enrichment of elemental-abundance ratios relative to proto-solar as a function of the accreted mass (V/Fe and Ba/Fe behave similarly to Mn/Fe and are not shown for clarity). In this example, the accreted body has a core-mass fraction of 1.95, a mantle composition matching that of the surface of Mercury<sup>84</sup> and a core composition of 15% Ni and predominantly Fe as the rest<sup>83</sup>. The horizontal blue line denotes a proto-solar composition and the vertical dashed lines show the masses of Mercury, Mars and Venus for reference. If too small a body is

accreted ( $\leq 0.1 M_{\oplus}$ ) onto the initial  $284 M_{\oplus}$  of WASP-76b, the composition does not change greatly from proto-solar. If too large a body is accreted ( $\geq 3 M_{\oplus}$ ), the overall composition begins to change too exceptionally. Although all abundance ratios require different masses to be perfectly matched under this assumed enrichment material composition, the overall best fit occurs if a large object between roughly the size of Mars and Venus is added to WASP-76b. **b**, A comparison between the data and the toy model assuming an accreted mass of  $0.5 M_{\oplus}$ , in which the data points except Mn/Fe are reasonably well matched. All error bars denote  $1\sigma$  uncertainties.



**Extended Data Fig. 9 | Rest-frame absorption signals of individual species on WASP-76b.** Shown in each panel are the cross-correlation trails for the species combined in Fig. 3b. Despite their wide range in condensation temperature, most species have a similar 'kinked' absorption trail as Fe (dotted

white line), probably indicating that condensation is not the sole culprit for the asymmetric signature. One notable exception is ionized calcium, which does not show such an asymmetry, probably because of Ca<sup>+</sup> triplet having an absorption depth consistent with analysing an escaping atmosphere<sup>18</sup>.

**Extended Data Table 1 | Retrieved abundance constraints for species on WASP-76b**

Element	$\log(\chi_i)$	$\log(\chi_i/\chi_{\text{Fe}})$	$\log\left(\frac{[\chi_i/\chi_{\text{Fe}}]}{[\chi_i/\chi_{\text{Fe}}]_{\odot}}\right)$	$\log\left(\frac{[\chi_i/\chi_{\text{Fe}}]}{[\chi_i/\chi_{\text{Fe}}]_{*}}\right)$	$T_{\text{cond}}^{20}$ [K]
Na	$-6.10^{+0.33}_{-0.34}$	$-2.31^{+0.16}_{-0.17}$	$-1.06^{+0.17}_{-0.18}$	$-1.20^{+0.18}_{-0.19}$	958
K	$-8.04^{+0.51}_{-0.52}$	$-4.25^{+0.45}_{-0.43}$	$-1.85^{+0.45}_{-0.43}$	$-1.90^{+0.46}_{-0.43}$	1006
Li	$-9.65^{+0.38}_{-0.41}$	$-5.85^{+0.27}_{-0.31}$	$-1.65^{+0.27}_{-0.32}$	–	1142
Mn	$-5.70^{+0.33}_{-0.33}$	$-1.91^{+0.19}_{-0.20}$	$0.14^{+0.20}_{-0.21}$	$0.16^{+0.21}_{-0.22}$	1158
Cr	$-6.15^{+0.32}_{-0.32}$	$-2.35^{+0.14}_{-0.15}$	$-0.51^{+0.16}_{-0.16}$	$-0.55^{+0.18}_{-0.18}$	1296
CrH	$< -9.24$	$< -5.52$	$< -3.68$	$< -3.71$	1296
Fe	$-3.79^{+0.29}_{-0.31}$	0	0	0	1334
Mg	$-3.87^{+0.35}_{-0.36}$	$-0.08^{+0.21}_{-0.21}$	$-0.17^{+0.21}_{-0.21}$	$-0.15^{+0.22}_{-0.22}$	1336
Ni	$-4.54^{+0.35}_{-0.35}$	$-0.73^{+0.18}_{-0.21}$	$0.54^{+0.19}_{-0.22}$	$0.53^{+0.20}_{-0.23}$	1353
V	$-7.80^{+0.34}_{-0.34}$	$-4.00^{+0.18}_{-0.18}$	$-0.44^{+0.20}_{-0.21}$	–	1429
VO	$-8.72^{+0.26}_{-0.26}$	$-4.94^{+0.19}_{-0.17}$	$-1.37^{+0.21}_{-0.19}$	–	1429
Ba	$-9.07^{+0.34}_{-0.36}$	$-5.27^{+0.22}_{-0.26}$	$-0.07^{+0.23}_{-0.26}$	–	1455
Ca	$-7.07^{+0.33}_{-0.37}$	$-3.28^{+0.15}_{-0.17}$	$-2.11^{+0.16}_{-0.18}$	$-2.18^{+0.19}_{-0.21}$	1517
Ca <sup>+</sup>	$-4.87^{+0.40}_{-0.44}$	$-1.09^{+0.16}_{-0.17}$	$0.07^{+0.17}_{-0.18}$	$-0.00^{+0.20}_{-0.21}$	1517
Ti	$< -9.78$	$< -5.97$	$< -3.47$	$< -3.48$	1582
TiO	$< -10.10$	$< -6.32$	$< -3.83$	$< -3.83$	1582
Ti <sup>+</sup>	$< -8.61$	$< -4.79$	$< -2.30$	$< -2.31$	1582
Al	$< -7.77$	$< -3.96$	$< -2.92$	–	1653
AlO	$< -10.67$	$< -6.88$	$< -5.84$	–	1653
Sc	$< -11.67$	$< -7.88$	$< -3.56$	–	1659

⊙ Proto-solar

\* Stellar