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# GOALS-JWST: Small Neutral Grains and Enhanced 3.3 $\mu\text{m}$ PAH Emission in the Seyfert Galaxy NGC 7469

Thomas S.-Y. Lai (賴劭愉)<sup>1</sup> , Lee Armus<sup>1</sup> , Marina Bianchin<sup>2</sup> , Tanio Díaz-Santos<sup>3,4</sup> , Sean T. Linden<sup>5</sup> , George C. Privon<sup>6,7,8</sup> , Hanae Inami<sup>9</sup> , Vivian U<sup>2</sup> , Thomas Bohn<sup>9</sup> , Aaron S. Evans<sup>6,7</sup> , Kirsten L. Larson<sup>10</sup> , Brandon S. Hensley<sup>11</sup> , J.-D. T. Smith<sup>12</sup> , Matthew A. Malkan<sup>13</sup> , Yiqing Song<sup>14,15</sup> , Sabrina Stierwalt<sup>16</sup> , Paul P. van der Werf<sup>17</sup> , Jed McKinney<sup>5</sup> , Susanne Aalto<sup>18</sup> , Victorine A. Buiten<sup>17</sup> , Jeff Rich<sup>19</sup> , Vassilis Charmandaris<sup>3,4,20</sup> , Philip Appleton<sup>1</sup> , Loreto Barcos-Muñoz<sup>6,7</sup> , Torsten Böker<sup>21</sup> , Luke Finnerty<sup>13</sup> , Justin A. Kader<sup>2</sup> , David R. Law<sup>22</sup> , Anne M. Medling<sup>23</sup> , Michael J. I. Brown<sup>24</sup> , Christopher C. Hayward<sup>25</sup> , Justin Howell<sup>1</sup> , Kazushi Iwasawa<sup>26,27</sup> , Francisca Kemper<sup>28,29,30</sup> , Jason Marshall<sup>31</sup> , Joseph M. Mazzarella<sup>1</sup> , Francisco Müller-Sánchez<sup>32</sup> , Eric J. Murphy<sup>6</sup> , David Sanders<sup>33</sup> , and Jason Surace<sup>1</sup>

<sup>1</sup> IPAC, California Institute of Technology, 1200 East California Boulevard, Pasadena, CA 91125, USA; [ThomasLai.astro@gmail.com](mailto:ThomasLai.astro@gmail.com)

<sup>2</sup> Department of Physics and Astronomy, 4129 Frederick Reines Hall, University of California, Irvine, CA 92697, USA

<sup>3</sup> Institute of Astrophysics, Foundation for Research and Technology-Hellas (FORTH), Heraklion, 70013, Greece

<sup>4</sup> School of Sciences, European University Cyprus, Diogenes Street, Engomi, 1516 Nicosia, Cyprus

<sup>5</sup> Department of Astronomy, University of Massachusetts at Amherst, Amherst, MA 01003, USA

<sup>6</sup> National Radio Astronomy Observatory, 520 Edgemont Road, Charlottesville, VA 22903, USA

<sup>7</sup> Department of Astronomy, University of Virginia, 530 McCormick Road, Charlottesville, VA 22903, USA

<sup>8</sup> Department of Astronomy, University of Florida, P.O. Box 112055, Gainesville, FL 32611, USA

<sup>9</sup> Hiroshima Astrophysical Science Center, Hiroshima University, 1-3-1 Kagamiyama, Higashi-Hiroshima, Hiroshima 739-8526, Japan

<sup>10</sup> AURA for the European Space Agency (ESA), Space Telescope Science Institute, 3700 San Martin Drive, Baltimore, MD 21218, USA

<sup>11</sup> Department of Astrophysical Sciences, Princeton University, Princeton, NJ 08544, USA

<sup>12</sup> Ritter Astrophysical Research Center, University of Toledo, Toledo, OH 43606, USA

<sup>13</sup> Department of Physics & Astronomy, 430 Portola Plaza, University of California, Los Angeles, CA 90095, USA

<sup>14</sup> European Southern Observatory, Alonso de Córdova, 3107, Vitacura, Santiago, 763-0355, Chile

<sup>15</sup> Joint ALMA Observatory, Alonso de Córdova, 3107, Vitacura, Santiago, 763-0355, Chile

<sup>16</sup> Physics Department, 1600 Campus Road, Occidental College, Los Angeles, CA 90041, USA

<sup>17</sup> Leiden Observatory, Leiden University, P.O. Box 9513, 2300 RA Leiden, The Netherlands

<sup>18</sup> Department of Space, Earth and Environment, Chalmers University of Technology, SE-412 96 Gothenburg, Sweden

<sup>19</sup> The Observatories of the Carnegie Institution for Science, 813 Santa Barbara Street, Pasadena, CA 91101, USA

<sup>20</sup> Department of Physics, University of Crete, Heraklion, 71003, Greece

<sup>21</sup> European Space Agency, Space Telescope Science Institute, Baltimore, MD 21218, USA

<sup>22</sup> Space Telescope Science Institute, 3700 San Martin Drive, Baltimore, MD 21218, USA

<sup>23</sup> Department of Physics & Astronomy and Ritter Astrophysical Research Center, University of Toledo, Toledo, OH 43606, USA

<sup>24</sup> School of Physics and Astronomy, Monash University, Clayton, VIC 3800, Australia

<sup>25</sup> Center for Computational Astrophysics, Flatiron Institute, 162 Fifth Avenue, New York, NY 10010, USA

<sup>26</sup> Institut de Ciències del Cosmos (ICCUB), Universitat de Barcelona (IEEC-UB), Martí i Franquès, 1, E-08028 Barcelona, Spain

<sup>27</sup> ICREA, Pg. Lluís Companys 23, E-08010 Barcelona, Spain

<sup>28</sup> Institut de Ciències de l'Espai (ICE, CSIC), Can Magrans, s/n, E-08193 Bellaterra, Barcelona, Spain

<sup>29</sup> ICREA, Pg. Lluís Companys 23, Barcelona, Spain

<sup>30</sup> Institut d'Estudis Espacials de Catalunya (IEEC), E-08034 Barcelona, Spain

<sup>31</sup> Glendale Community College, 1500 North Verdugo Road, Glendale, CA 91208, USA

<sup>32</sup> Department of Physics and Materials Science, The University of Memphis, 3720 Alumni Avenue, Memphis, TN 38152, USA

<sup>33</sup> Institute for Astronomy, University of Hawaii, 2680 Woodlawn Drive, Honolulu, HI 96822, USA

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## Abstract

We present James Webb Space Telescope (JWST) Near Infrared Spectrograph (NIRSpec) integral field spectroscopy of the nearby luminous infrared galaxy NGC 7469. We take advantage of the high spatial/spectral resolution and wavelength coverage of JWST/NIRSpec to study the 3.3  $\mu\text{m}$  neutral polycyclic aromatic hydrocarbon (PAH) grain emission on  $\sim 200$  pc scales. A clear change in the average grain properties between the star-forming ring and the central AGN is found. Regions in the vicinity of the AGN, with  $[\text{Ne III}]/[\text{Ne II}] > 0.25$ , tend to have larger grain sizes and lower aliphatic-to-aromatic (3.4/3.3) ratios, indicating that smaller grains are preferentially removed by photodestruction in the vicinity of the AGN. PAH emission at the nucleus is weak and shows a low 11.3/3.3 PAH ratio. We find an overall suppression of the total PAH emission relative to the ionized gas in the central 1 kpc region of the AGN in NGC 7469 compared to what has been observed with Spitzer on 3 kpc scales. However, the fractional 3.3  $\mu\text{m}$ -to-total PAH power is enhanced in the starburst ring, possibly due to a variety of physical effects on subkiloparsec scales, including recurrent fluorescence of small grains or multiple photon absorption by large grains. Finally, the IFU data show that while the 3.3  $\mu\text{m}$  PAH-derived star formation rate (SFR) in the ring is 27% higher than that inferred from the [Ne II] and [Ne III] emission lines, the integrated



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SFR derived from the 3.3  $\mu\text{m}$  feature would be underestimated by a factor of 2 due to the deficit of PAHs around the AGN, as might occur if a composite system like NGC 7469 were to be observed at high redshift.

*Unified Astronomy Thesaurus concepts:* Seyfert galaxies (1447); Active galactic nuclei (16); Polycyclic aromatic hydrocarbons (1280); Starburst galaxies (1570); Luminous infrared galaxies (946)

## 1. Introduction

The strong infrared emission bands at 3.3, 6.2, 7.7, 8.6, 11.3, and 17  $\mu\text{m}$  are detectable in the spectra of a diverse range of astrophysical sources, ranging from photodissociation regions (PDRs) in nebulae (Werner et al. 2004), to nearby star-forming (SF) galaxies (Peeters et al. 2004; Smith et al. 2007) and luminous infrared galaxies (LIRGs; see Armus et al. 2020, for a review), even to distant galaxies at high redshifts ( $z > 1$ ; Sajina et al. 2009; Pope et al. 2013; Riechers et al. 2014; Spilker et al. 2023). These prominent spectral features are often attributed to polycyclic aromatic hydrocarbon (PAH) molecules (Leger & Puget 1984; Allamandola et al. 1985; Tielens 2008; Li 2020), which can account for up to 15% of the cosmic carbon (Li & Draine 2001; Zubko et al. 2004; Draine & Li 2007; Jones et al. 2017) and contribute up to  $\sim 20\%$  of the total infrared luminosity from a galaxy (Smith et al. 2007).

The 3.3  $\mu\text{m}$  PAH band is the shortest-wavelength feature among all of the PAHs and primarily attributed to the C–H stretching mode. It is known to trace the neutral, smallest PAH population in the interstellar medium (ISM), typically with a radius of  $\sim 5 \text{ \AA}$  or  $N_C \sim 50$  carbon atoms (Schutte et al. 1993; Draine & Li 2007). The 3.3  $\mu\text{m}$  PAH has been the least-studied PAH feature due to its inaccessibility from ground-based observations and the lack of spectral coverage of Spitzer/IRS shortward of 5  $\mu\text{m}$ . Nonetheless, observations with the AKARI infrared satellite (Murakami et al. 2007) had previously demonstrated the power of using the 3.3  $\mu\text{m}$  PAH to diagnose star formation and active galactic nuclei (AGN) in (U)LIRGs (Imanishi et al. 2008, 2010; Ichikawa et al. 2014; Inami et al. 2018) and study the photoelectric efficiency that couples the stellar radiation field with gas temperatures (McKinney et al. 2021). Just as importantly, the 3.3  $\mu\text{m}$  band is the only PAH feature that will be accessible to JWST at high redshifts, making detailed studies of this feature at low redshift of paramount importance for understanding dust and metals at early epochs.

In a joint AKARI+Spitzer study, Lai et al. (2020) found that the fractional 3.3  $\mu\text{m}$  PAH power ( $L_{\text{PAH } 3.3}/L_{\Sigma \text{ PAH}}$ ) in nearby SF galaxies typically ranges from 1.5% to 3%, a value that is significantly smaller than that for the brightest 7.7  $\mu\text{m}$  PAH complex, which typically reaches  $L_{\text{PAH } 7.7}/L_{\Sigma \text{ PAH}} \sim 40\%$  in nearby SF galaxies (Smith et al. 2007). Despite its faintness, the 3.3  $\mu\text{m}$  PAH feature has been proposed as an accurate tool for measuring the average grain size when paired with the 11.3 and 17  $\mu\text{m}$  PAH bands, which also serve as tracers of neutral grains (e.g., Croiset et al. 2016; Lai et al. 2020; Maragkoudakis et al. 2020, 2022, 2023a, 2023b; Draine et al. 2021; Rigopoulou et al. 2021; Sidhu et al. 2022). The 3.3  $\mu\text{m}$  PAH holds the potential to be a sensitive diagnostic of ISM conditions, as well as an estimator of the star formation rate (SFR; e.g., Peeters et al. 2004; Shipley et al. 2016; Lai et al. 2020) accessible from nearby galaxies to high redshifts with JWST imaging and spectroscopy (Inami et al. 2018; Evans et al. 2022; Leroy et al. 2023; Sandstrom et al. 2023). These applications, however, strongly rely on our understanding of the properties of PAHs and how they depend on metallicity and

react to the ambient radiation field, as PAHs can be faint or even absent in low-metallicity environments (Wu et al. 2006; Hao et al. 2009), as well as in galaxies with powerful AGN (O’Dowd et al. 2009; Diamond-Stanic & Rieke 2010; Sajina et al. 2022). Surprisingly, some studies have shown that AGN photons can also be the source of PAH molecule excitation and emission (Howell et al. 2007; Jensen et al. 2017).

In this paper, we focus on tracing the physical conditions of the neutral dust grains probed by the 3.3  $\mu\text{m}$  PAH in NGC 7469 using the superb spatial and spectral resolution of the JWST Near Infrared Spectrograph (NIRSpec; Jakobsen et al. 2022). As a follow-up study to Lai et al. (2022, hereafter Paper I), the JWST/MIRI results are also included here in a joint analysis to provide a complete view of the aromatic bands in NGC 7469.

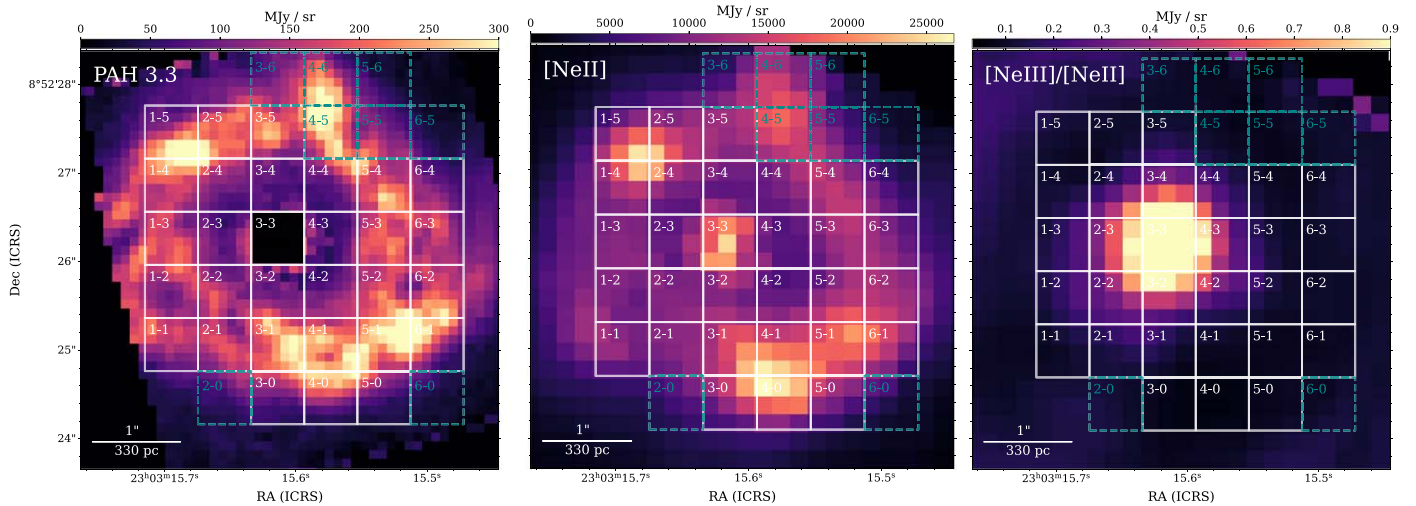
NGC 7469 (Arp 298, Mrk 1514, IRAS F23007+0836) is a galaxy located at  $D_L = 70.6 \text{ Mpc}$  and classified as a Seyfert 1.5 (Landt et al. 2008). It is an LIRG ( $L_{8-1000 \mu\text{m}} = 10^{11.6} L_{\odot}$ ) with a supermassive black hole mass of  $1.1 \times 10^7 M_{\odot}$  (Peterson et al. 2014; Lu et al. 2021) and X-ray luminosity of  $L_{2-10 \text{ keV}} = 10^{43.19} \text{ erg s}^{-1}$  (Asmus et al. 2015). NGC 7469 contains an accreting black hole and a starburst ring with a radius of  $\sim 500 \text{ pc}$  (e.g., Song et al. 2021). Recently, NIRCам imaging data revealed a factor of  $\sim 6$  more dusty, young ( $< 5 \text{ Myr}$ ) clusters compared to the previous identifications made by HST (Bohn et al. 2023). Using JWST integral field spectroscopy, emission lines with ionization potential up to 187 eV are detected within a distance of  $\sim 100 \text{ pc}$  from the AGN with an outflowing wind that can reach up to  $1700 \text{ km s}^{-1}$  traced by high-ionization lines (Armus et al. 2023), and the ionized coronal wind probed by [Mg V] (IP: 109 eV) extends out to  $\sim 400 \text{ pc}$  from the nucleus (Vivian et al. 2022). By studying the PAH ratios in the mid-infrared, Paper I found that feedback from the central black hole mostly impacts the inner ISM region, leading to a larger grain size distribution in the vicinity of the AGN than that in the ring (cf. García-Berete et al. 2022b).

Throughout this paper, a cosmology with  $H_0 = 70 \text{ km s}^{-1} \text{ Mpc}^{-1}$ ,  $\Omega_M = 0.30$ , and  $\Omega_{\Lambda} = 0.70$  is adopted. The redshift of NGC 7469 ( $z = 0.01627^{34}$ ) corresponds to a projected physical scale of  $330 \text{ pc arcsec}^{-1}$ .

## 2. Observations and Data Reduction

As part of Director’s Discretionary Time Early Release Science program 1328 (Co-PIs: L. Armus and A. Evans), the JWST near-infrared integral field unit (IFU) observations (Böker et al. 2022) of NGC 7469 were taken with NIRSpec on 2022 July 19. The NIRSpec IFU observations were carried out using a set of high-resolution gratings (with a nominal resolving power of  $R \sim 2700$ ), namely, G140H/100LP, G235H/170LP, and G395H/290LP, covering a wavelength range of 0.97–5.27  $\mu\text{m}$ . For each grating, the science exposure time was 817 s, and a four-point dither pattern was used to sample the extended SF ring. In this study, we focus on the G395H grating that contains the shortest PAH band at 3.3  $\mu\text{m}$ .

<sup>34</sup> NASA/IPAC Extragalactic Database.



**Figure 1.** Left: the  $3.3\ \mu\text{m}$  PAH map of NGC 7469. This image was created by subtracting a linear fit to the continuum in the unsmoothed NIRSspec cube (see Section 3), and it has a pixel size of  $0''.1$ , providing us with the highest spatial view of the dust distribution in NGC 7469 using JWST IFU. The spectra extracted in this study are based on the overlaid square grid, with each cell having a width of  $0''.6$ . Grids in white are cells with full NIRSspec (G395H)+MIRI coverage, while the ones in dashed blue have only partial coverage in MIRI. The central grid is masked due to noise in the linear continuum subtraction at the nucleus, where the PAH emission is very weak but has no effect on the extracted spectra analyzed in this paper. Middle: the [Ne II]  $12.81\ \mu\text{m}$  MIRI/MRS map for comparison, which has a pixel size of  $0''.2$  and traces the ionized atomic gas in the starburst ring. The same spectral grid is overlaid here. In general, there is good agreement between the  $3.3\ \mu\text{m}$  PAH and [Ne II] maps, but the  $3.3\ \mu\text{m}$  PAH map shows finer detail owing to its higher intrinsic spatial resolution, and differences are evident in the ratio of PAH to ionized gas emission around the ring. Right: emission line flux ratio map of [Ne III]/[Ne II], a measure of the radiation field hardness. Here the [Ne II] map was smoothed to match the PSF of [Ne III] before generating the line ratio map. As expected, the [Ne III]/[Ne II] ratio is highly peaked toward the center. These maps were created by applying a linear local continuum subtraction using the QFitsView tool (Ott 2012).

We downloaded the uncalibrated science observations through the MAST portal. The data reduction process was done using the JWST Science Calibration Pipeline version 1.8.3 (Bushouse et al. 2022), with a reference file of `jwst_12027.pmap`. Three stages of pipeline processing were applied, including `Detector1`, `Spec2`, and `Spec3`. While the dedicated background is not included in the observation strategy, “leakcal” was performed in the observation to mitigate the contamination due to a failed open microshutter assembly. We refer readers to our companion paper (Bianchin et al. 2023) for more details on the data processing.

To obtain a complete spectroscopic view of NGC 7469, we combined the NIRSspec IFU data with the MIRI IFU data presented in Paper I, allowing us to study all of the PAH features from 3 to  $20\ \mu\text{m}$ . To properly compare the NIRSspec and MIRI data, two modifications to the NIRSspec cube had to be made. First, a spatial offset of  $\sim 0''.24$  between the NIRSspec and MIRI IFUs was found, so we combined the two IFUs by shifting the WCS of NIRSspec to match with that of MIRI, whose astrometry has been verified with the MIRI imaging. This offset was measured by comparing the collapsed images from the spectral overlap region between  $5.0$  and  $5.1\ \mu\text{m}$  in both IFUs. Second, in order to combine the IFUs while matching their spatial resolution, we smoothed the NIRSspec G395H/290LP cube and MIRI channels 1 and 2 by convolving them with wavelength-dependent Gaussian convolution kernels to match with the FWHM measured at  $11.5\ \mu\text{m}$  (FWHM  $\sim 0''.46$ ), where the longest-wavelength PAH feature (PAH  $11.3\ \mu\text{m}$  in MIRI channel 2) studied in this paper is located. The convolution was performed as in Zhang & Ho (2023) and based on the `Gaussian2DKernel` function in `astropy.convolution`. The convolution kernel is calculated using the NIRSspec-observed standard star P330E (PID:

1538; PI: K. Gordon) and the MIRI observed star HD 159222 (PID: 1050; PI: B. Vandenbussche).

### 3. Analysis

Many studies have found that the hardness of the radiation field probed by [Ne III]/[Ne II] plays an important role in regulating PAH ratios in AGN and SF galaxies (e.g., Smith et al. 2007; Sales et al. 2010; García-Bernete et al. 2022a). Particularly, in Paper I, we have shown that the main driver of grain size variation in this galaxy is [Ne III]/[Ne II], which will also be used here to understand the properties of the  $3.3\ \mu\text{m}$  PAH. In Figure 1, we present the PAH and [Ne II] maps, along with a [Ne III]/[Ne II] emission line flux ratio map. The SF ring is dominated by bright PAH emission, which provides a dust map with the highest available spatial resolution of FWHM  $\sim 0''.19$  using the NIRSspec IFU data compared to the MIRI channel 3 [Ne II] map, which has an FWHM  $\sim 0''.48$ . In general, the PAH  $3.3$  and [Ne II] emission show similar morphology, having four clumps along the ring, but the relative intensities differ. For example, the neon emission is brighter in the southern clump (4-1), whereas PAH  $3.3$  is brighter in the northern clump (4-5).

In Paper I, the study of the MIR dust and gas properties in NGC 7469 was carried out on an aperture-based analysis of the MIRI IFU data. In the study presented here, we maximize the full potential of an IFU observation by applying a grid extraction to both the NIRSspec and MIRI IFUs. This allows us to obtain spectra across a wide wavelength range from 3 to  $28\ \mu\text{m}$  in a systematic approach. The grid extraction was performed on the smoothed NIRSspec and MIRI cubes using the spectral extraction tool included in the most recent version of the Continuum And Feature Extraction (CAFE) software. This software was initially developed by Marshall et al. (2007) for Spitzer/IRS and has recently been updated for JWST (see T. Díaz-Santos et al. 2023, in preparation), which includes the



capability of handling both spectral extraction of the IFU cube and spectral decomposition of the extracted spectra.

For the grid extraction, we initially set the center of the grid at the nucleus (23:03:15.6139, +8:52:26.266 J2000) and used  $7 \times 7$  grid cells, each  $0''.6$  wide, to extract spectra across the full field of view (FoV) from both the smoothed NIRSpec and MIRI cubes. The grid presented in Figure 1 shows the positions of our extracted spectra. No extraction is done if the grid cell exceeds the coverage of the IFU FoV, so the total number of extractions is 38 instead of 49, including the 30 cells (white) having full G395H and MIRI coverage and eight cells (dark blue) having complete NIRSpec coverage but incomplete MIRI coverage (missing some PAH emission). Of the eight cells, five (2-0, 3-6, 4-5, 5-5, and 6-0) exhibit the  $3.3 \mu\text{m}$  PAH, the  $3.4 \mu\text{m}$  aliphatic feature, and the [Ne II] and [Ne III] lines, even in the absence of MIRI coverage. These cells will be discussed in the analysis focusing on the aliphatic-to-aromatic ratio (see Figure 3, left panel). This partial coverage of MIRI is due to the slight spatial mismatch between each MIRI subchannel. To obtain the full spectral coverage of  $3\text{--}28 \mu\text{m}$ , we stitched the spectrum of NIRSpec to MIRI.

The 35 cells with  $3.3$ ,  $3.4$ , [Ne II], and [Ne III] are divided into three groups: SF, gap, and nucleus. Most of the extracted spectra belong to the SF ring, except for the five cells 2-3, 2-4, 3-4, 4-3, and 4-4, which are designated as gap region cells, and the nuclear extraction (3-3). Note that the nuclear spectrum presented here differs from the AGN spectrum analyzed in Armus et al. (2023), which was extracted using an expanding cone with a small inner diameter of  $0''.3$  at  $5.5 \mu\text{m}$  to isolate the central source and the base of the outflow.

#### 4. Results

The FoV of the dithered JWST NIRSpec IFU observations ( $4''.2 \times 4''.8$ ) provides complete coverage of the NGC 7469 circumnuclear ring and inner ISM, allowing for high spatial and spectral resolution in the study of dust and star formation properties in the near-infrared. To demonstrate the variation of the spectra across the IFU, we summed the individual grid spectra to create three NGC 7469 spectra: nuclear, ring, and gap (see Figure 2(a)). The low-resolution ( $R \sim 100$ ) AKARI + Spitzer spectrum taken from the AKARI–Spitzer Extragalactic Spectral Survey (ASESS; Lai et al. 2020) is also presented for comparison. Clearly, PAHs are heavily suppressed in the nuclear spectrum, while the high-ionization fine-structure lines are noticeably strong. For example, the nuclear spectrum has the highest [Ne III]/[Ne II] and [S IV]/[Ne II] line flux ratios among the region spectra.

To fit and analyze the NIRSpec spectra, we use CAFE, a spectral decomposition tool that is able to simultaneously fit the PAH features, dust continuum, absorbed ice and gas features, and narrow fine-structure atomic and molecular gas emission lines. The stacked SF ring spectrum, together with its spectral fit, is presented in Figure 2(b), and a zoom-in view at the  $3 \mu\text{m}$  region is presented in Figure 2(c). In this regime, the spectrum is dominated by the  $3.3 \mu\text{m}$  PAH band and  $3.4 \mu\text{m}$  aliphatic feature, sitting on top of a broad plateau that can extend up to  $\sim 3.8 \mu\text{m}$ . In addition, absorption features due to water ice at  $3.05 \mu\text{m}$ ,  $\text{CO}_2$  ice at  $4.27 \mu\text{m}$ , and CO ice at  $4.67 \mu\text{m}$  can also be seen. We report the central wavelengths and widths of the three dust components seen in the stacked ring spectrum in Table 1. We find the  $3.3$  and  $3.4 \mu\text{m}$  features to be narrower, while the plateau is broader than previous measurements made

with AKARI (Lai et al. 2020). This is undoubtedly due to the substantially higher spectral resolution of NIRSpec, enabling an accurate and detailed deblending of the dust components in the  $3 \mu\text{m}$  complex for the first time. The measurements of the extinction-corrected PAH and neon fluxes in each cell, as derived from the  $9.7 \mu\text{m}$  silicate absorption depth and fit with CAFE, are presented in Table 2.

##### 4.1. Aliphatic versus Aromatic Emission in NGC 7469

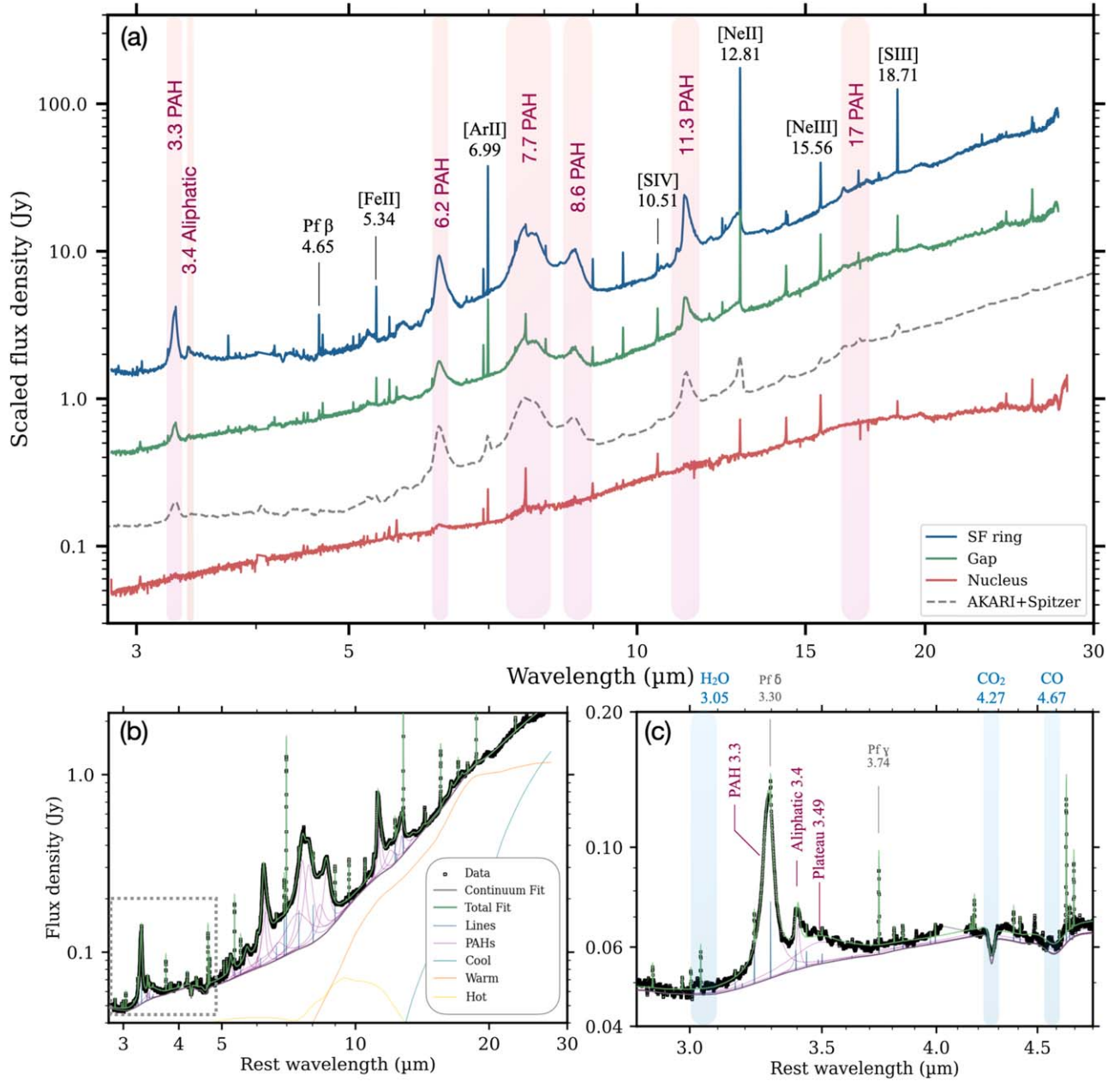
The abundance ratios of different types of carbonaceous dust can provide insights into their formation mechanisms and histories. The  $3.4 \mu\text{m}$  emission feature is often attributed to the vibrations of  $-\text{CH}_3$  (methyl) and  $-\text{CH}_2$  (methylene) groups in aliphatic hydrocarbons (Joblin et al. 1996; Yang et al. 2016), and it can also be seen in absorption along sight lines with large extinction (Pendleton & Allamandola 2002; Chiar et al. 2013; Hensley & Draine 2020). The variation of the ratio between this  $3.4 \mu\text{m}$  aliphatic feature and the  $3.3 \mu\text{m}$  aromatic feature is indicative of the processing of dust particles in the ISM, as the chainlike aliphatic bonds are more fragile compared to the ringlike aromatic bonds.

The  $3.4 \mu\text{m}$  aliphatic feature was detected in all of the grid cells of NGC 7469 except for the nucleus. The  $3.4/3.3$  ratio varies by about a factor of 2 across the IFU FoV. Regions with a harder radiation field, as measured by the [Ne III]/[Ne II] line flux ratio, tend to have a lower average  $3.4/3.3$  ratio and a larger scatter. The  $3.4/3.3$  ratio is lower in cells between the ring and the nucleus, decreasing by about 20% compared to those in the SF ring, suggesting that the aliphatic bonds are more susceptible to photodestruction than the aromatic bonds (Figure 3, left panel). There is no detection of the  $3.4 \mu\text{m}$  emission at the nucleus, and the range in  $3.4/3.3$  ratio values is largest in the cells with [Ne III]/[Ne II]  $> 0.5$ . While visible, the variation of the  $3.4/3.3$  ratio in the SF ring of NGC 7469 is relatively small, ranging only from 5% to 8%, compared to Galactic regions that can range widely from  $\sim 1\%$  to 17% in PDRs (Pillari et al. 2015, see Figure 7(c)), suggesting that the photoenergy density in the ring is high, with  $\log(G_0) > 3$  in units of the Habing field (Habing 1968).

##### 4.2. PAH Size Distribution

Small particles like PAHs tend to have larger cross sections per unit mass than larger grains when assuming a continuous grain size distribution (Mathis et al. 1977), rendering them sensitive probes for local physical conditions such as size and charge state within the ISM. Smaller PAHs tend to emit at short wavelengths because upon absorbing a UV photon, each bond in a smaller PAH molecule gets excited to a higher vibrational level, which will subsequently emit in short wavelengths. Typically, PAH size estimates have relied on the ratio of  $6.2/7.7$  PAH, while the  $11.3/7.7$  PAH ratio has been used for determining charge states (Draine & Li 2001). The ratio of  $6.2/7.7$ , however, is not an ideal tracer for the size because the difference between the two wavelengths is relatively small, so the ratio has only limited diagnostic power, as shown in Figure 7 of Draine & Li (2007).

One of the key advantages of NIRSpec is accessing the  $3.3 \mu\text{m}$  PAH band, which has been shown to be a more effective proxy for the size of the small dust populations, particularly when used together with the  $11.3 \mu\text{m}$  band, since both trace neutral PAHs (e.g., Croiset et al. 2016;



**Figure 2.** (a) The stacked (summed over all corresponding grid cells) spectra in the three regions of NGC 7469, including the SF ring (blue), the gap between the ring and the nucleus (green), and the nucleus (red), with the key diagnostic emission features labeled. The PAHs are heavily suppressed in the nucleus, where the  $[\text{Ne III}]/[\text{Ne II}]$  is highest, indicative of a harder radiation field dominated by the central AGN. The low-resolution ( $R \sim 100$ ) AKARI+Spitzer spectrum, which covers the nucleus, gap, and SF ring taken from the ALESS catalog (Lai et al. 2020), is also presented for comparison (dashed gray line). Panel (b) shows the multicomponent CAFE fit of the stacked SF ring spectrum, and a zoom-in view of the  $3 \mu\text{m}$  regime is shown in panel (c), delineated by the dashed box. As in panel (a), key features are labeled here. The aromatic feature at  $3.3 \mu\text{m}$  and the aliphatic feature at  $3.4 \mu\text{m}$  can be readily differentiated with the high spectral resolution of NIRSpect, allowing us for the first time to confidently separate these features in galaxies and use their flux ratio as an independent diagnostic of the dusty ISM.

**Table 1**  
The  $3 \mu\text{m}$  Dust Features

Component (1)	$\lambda$ [ $\mu\text{m}$ ] (2)	$\Gamma$ (3)
3.3 PAH	3.288	0.012
3.4 aliphatic	3.404	0.005
Plateau	3.484	0.070

**Note.** Central wavelengths and widths of the dust feature as measured in the stacked NGC 7469 SF ring spectrum. Column (1): dust feature name; column (2): central wavelength of the feature; column (3): fractional FWHM, defined by the ratio of FWHM and central wavelength.

Maragkoudakis et al. 2020; Sidhu et al. 2022). Figure 3 (right) shows the  $11.3/3.3$  PAH ratio as a function of the hardness of the radiation field. In the SF spectra of NGC 7469, the PAH  $11.3/3.3$  ratios cluster around a value of  $2.5 \pm 0.3$ , with a relatively smaller scatter. The gap spectra, on the other hand, exhibit, on average, a  $2\times$  increase in the  $11.3/3.3$  PAH ratios with a much broader distribution ( $5.2 \pm 1.1$ ). This agrees with the findings presented in Paper I, which suggested that the average grain size becomes larger closer to the nucleus, consistent with the destruction of small grains in the harder UV fields closer to the AGN (Puget & Leger 1989; Voit 1992). The nucleus appears to display an anomalously low  $11.3/3.3$  ratio

**Table 2**  
Extinction-corrected PAH-band and Emission Line Fluxes in NGC 7469

Cell	Group	PAH 3.3 $\mu\text{m}$ ( $\times 10^{-17}$ W m $^{-2}$ )	Aliphatic 3.4 $\mu\text{m}$ ( $\times 10^{-17}$ W m $^{-2}$ )	PAH 11.3 $\mu\text{m}$ ( $\times 10^{-17}$ W m $^{-2}$ )	$\Sigma\text{PAH}$ ( $\times 10^{-17}$ W m $^{-2}$ )	[Ne II] 12.81 $\mu\text{m}$ ( $\times 10^{-17}$ W m $^{-2}$ )	[Ne III] 15.56 $\mu\text{m}$ ( $\times 10^{-17}$ W m $^{-2}$ )
1-1	SF	6.17 $\pm$ 0.04	0.41 $\pm$ 0.02	14.00 $\pm$ 0.22	101.02 $\pm$ 1.03	3.39 $\pm$ 0.14	0.38 $\pm$ 0.02
1-2	SF	5.53 $\pm$ 0.04	0.34 $\pm$ 0.01	15.76 $\pm$ 0.41	121.54 $\pm$ 1.34	3.76 $\pm$ 0.19	0.69 $\pm$ 0.04
1-3	SF	6.39 $\pm$ 0.05	0.38 $\pm$ 0.01	15.26 $\pm$ 0.28	125.88 $\pm$ 1.29	3.77 $\pm$ 0.23	0.78 $\pm$ 0.05
1-4	SF	7.07 $\pm$ 0.04	0.47 $\pm$ 0.02	17.82 $\pm$ 0.02	146.49 $\pm$ 0.04	4.57 $\pm$ 0.21	0.57 $\pm$ 0.04
1-5	SF	4.73 $\pm$ 0.05	0.30 $\pm$ 0.01	12.13 $\pm$ 0.33	95.55 $\pm$ 1.35	5.20 $\pm$ 0.25	0.50 $\pm$ 0.04
2-1	SF	6.93 $\pm$ 0.04	0.47 $\pm$ 0.02	13.81 $\pm$ 0.26	107.66 $\pm$ 1.08	3.50 $\pm$ 0.26	0.62 $\pm$ 0.03
2-2	SF	8.36 $\pm$ 0.08	0.54 $\pm$ 0.02	21.86 $\pm$ 0.24	157.91 $\pm$ 2.17	4.87 $\pm$ 0.33	2.16 $\pm$ 0.20
2-3	G	2.22 $\pm$ 0.06	0.12 $\pm$ 0.01	8.90 $\pm$ 0.24	88.83 $\pm$ 2.63	3.83 $\pm$ 0.26	3.40 $\pm$ 0.53
2-4	G	4.72 $\pm$ 0.11	0.25 $\pm$ 0.01	17.97 $\pm$ 0.26	132.54 $\pm$ 2.02	5.67 $\pm$ 0.23	1.88 $\pm$ 0.10
2-5	SF	6.14 $\pm$ 0.11	0.36 $\pm$ 0.01	16.09 $\pm$ 0.30	128.06 $\pm$ 1.55	5.77 $\pm$ 0.39	0.64 $\pm$ 0.03
3-0	SF	5.83 $\pm$ 0.10	0.38 $\pm$ 0.02	14.46 $\pm$ 0.33	102.47 $\pm$ 1.44	4.12 $\pm$ 0.22	0.31 $\pm$ 0.03
3-1	SF	8.55 $\pm$ 0.08	0.48 $\pm$ 0.01	20.44 $\pm$ 0.28	173.06 $\pm$ 2.27	5.66 $\pm$ 0.38	0.78 $\pm$ 0.03
3-2	SF	10.11 $\pm$ 0.18	0.60 $\pm$ 0.04	25.29 $\pm$ 0.44	149.96 $\pm$ 4.62	4.68 $\pm$ 0.25	2.70 $\pm$ 0.17
3-3	N	2.97 $\pm$ 1.24	0.01 $\pm$ 0.01	5.28 $\pm$ 3.96	98.24 $\pm$ 6.49	4.52 $\pm$ 0.58	6.72 $\pm$ 0.27
3-4	G	1.80 $\pm$ 0.04	0.05 $\pm$ 0.01	11.61 $\pm$ 0.18	80.51 $\pm$ 2.81	3.64 $\pm$ 0.33	2.81 $\pm$ 0.13
3-5	SF	6.00 $\pm$ 0.10	0.30 $\pm$ 0.01	15.56 $\pm$ 0.33	134.35 $\pm$ 2.00	4.34 $\pm$ 0.21	0.65 $\pm$ 0.05
4-0	SF	7.92 $\pm$ 0.05	0.40 $\pm$ 0.02	17.19 $\pm$ 0.25	134.94 $\pm$ 1.47	5.91 $\pm$ 0.29	0.37 $\pm$ 0.03
4-1	SF	7.91 $\pm$ 0.08	0.41 $\pm$ 0.01	20.90 $\pm$ 0.02	169.17 $\pm$ 0.08	7.25 $\pm$ 0.24	0.67 $\pm$ 0.02
4-2	SF	10.86 $\pm$ 0.21	0.55 $\pm$ 0.03	38.13 $\pm$ 0.49	278.49 $\pm$ 3.59	6.53 $\pm$ 0.15	1.52 $\pm$ 0.05
4-3	G	3.35 $\pm$ 0.06	0.16 $\pm$ 0.02	21.54 $\pm$ 0.31	150.41 $\pm$ 2.67	4.32 $\pm$ 0.43	2.91 $\pm$ 0.09
4-4	G	2.53 $\pm$ 0.06	0.13 $\pm$ 0.01	13.69 $\pm$ 0.15	89.54 $\pm$ 1.13	3.97 $\pm$ 0.26	1.39 $\pm$ 0.10
5-0	SF	6.78 $\pm$ 0.03	0.42 $\pm$ 0.01	17.30 $\pm$ 0.24	110.82 $\pm$ 1.44	3.47 $\pm$ 0.17	0.29 $\pm$ 0.02
5-1	SF	9.01 $\pm$ 0.06	0.47 $\pm$ 0.02	22.35 $\pm$ 0.21	184.15 $\pm$ 1.77	6.76 $\pm$ 0.31	0.53 $\pm$ 0.03
5-2	SF	7.32 $\pm$ 0.20	0.40 $\pm$ 0.02	18.72 $\pm$ 0.37	157.29 $\pm$ 1.47	5.09 $\pm$ 0.27	0.66 $\pm$ 0.04
5-3	SF	6.96 $\pm$ 0.17	0.39 $\pm$ 0.02	18.19 $\pm$ 0.41	151.53 $\pm$ 3.57	4.37 $\pm$ 0.31	0.66 $\pm$ 0.04
5-4	SF	5.60 $\pm$ 0.12	0.30 $\pm$ 0.02	14.14 $\pm$ 0.60	120.60 $\pm$ 2.26	4.61 $\pm$ 0.27	0.59 $\pm$ 0.04
6-1	SF	8.01 $\pm$ 0.19	0.44 $\pm$ 0.03	20.88 $\pm$ 0.38	168.17 $\pm$ 1.71	4.47 $\pm$ 0.18	0.39 $\pm$ 0.02
6-2	SF	6.77 $\pm$ 0.24	0.36 $\pm$ 0.02	16.76 $\pm$ 0.29	142.90 $\pm$ 1.48	4.96 $\pm$ 0.20	0.48 $\pm$ 0.03
6-3	SF	6.49 $\pm$ 0.14	0.33 $\pm$ 0.02	15.42 $\pm$ 0.39	118.36 $\pm$ 1.56	4.23 $\pm$ 0.16	0.42 $\pm$ 0.02
6-4	SF	4.94 $\pm$ 0.13	0.24 $\pm$ 0.02	11.36 $\pm$ 0.32	78.83 $\pm$ 1.43	2.86 $\pm$ 0.14	0.34 $\pm$ 0.02

**Note.** In the second column, “SF” refers to the SF ring cell, “G” refers to the gap cell, and “N” refers to the nucleus cell.

but with a relatively large uncertainty due to very weak PAH emission (see the nucleus spectrum in Figure 2(a)). This will be further discussed in Section 5.1.

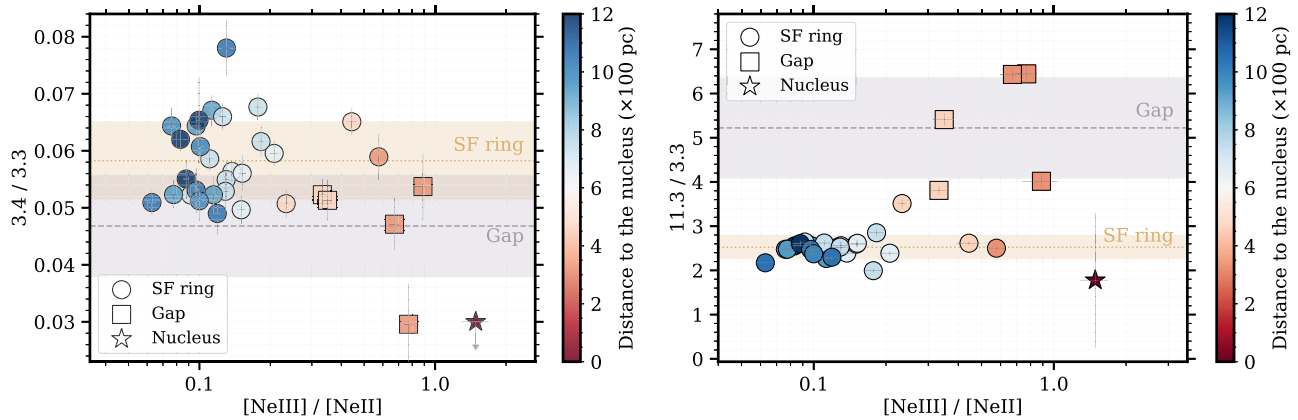
#### 4.3. Enhancement of the 3.3 $\mu\text{m}$ PAH

In Paper I, it was shown that the hardness of the radiation field was correlated with the PAH size variations seen in NGC 7469. Here we expand that analysis using the 3.3  $\mu\text{m}$  PAH feature. Figure 4(a) shows the variation of the total PAH-to-[Ne II]+[Ne III] ratio as a function of the [Ne III]/[Ne II] line flux ratio, color-coded by the total ionized PAH surface density, which traces the amount of photoelectric heating in the ISM. Here the total PAH ( $\Sigma\text{PAH}$ ) refers to the sum of the 3.3, 6.2, 7.7, 8.6, 11.3, and 17  $\mu\text{m}$  PAH features, while the total ionized PAH surface density is derived from only the 6.2, 7.7, and 8.6  $\mu\text{m}$  PAH features. In cells with large [Ne III]/[Ne II] line flux ratios ( $>0.25$ ), the  $\Sigma\text{PAH}$ -to-Ne ratio drops, resulting in a downward trend of  $\Sigma\text{PAH}$  to Ne with increasing radiation hardness. Their low  $\Sigma\text{PAH}$ -to-Ne ratios may be due to excess emission from AGN, together with a possible deficit of total PAH emission in highly ionized regions, similar to what has recently been observed in four nearby SF galaxies with JWST and MUSE (Egorov et al. 2023).

Besides the grid cell spectral extractions, we include two additional points to Figure 4. The diamond represents the total extraction of the IFU using a radius of 1''.6 that covers the AGN and the entire ring, while the triangle represents the integrated

light obtained from the combined AKARI+Spitzer NGC 7469 spectrum in the ASESS sample (Lai et al. 2020). The AKARI+Spitzer measurement probes an area that is three times wider than the coverage provided by the IFU, with a measured  $\Sigma\text{PAH}$ -to-Ne ratio that is larger than all of the individual JWST points. This is consistent with Díaz-Santos et al. (2011), who showed that PAHs are more extended, on average, than the [Ne II] emission in the GOALS galaxies. The two sets of contours indicate the distributions based on 244 galaxies within GOALS (red; Stierwalt et al. 2014) and 62 galaxies that are in both GOALS and the bright PAH samples in ASESS (blue; Lai et al. 2020) but refitted with CAFE to rule out the potential discrepancy caused by the use of different spectral decomposition tools. The GOALS sample includes LIRGs powered by starbursts and/or AGN, while the bright PAH sample consists only of SF galaxies. The NGC 7469 data points overlap with the two sets of contours but lie at the lower envelope of both distributions, suggesting slightly weaker overall PAH emission (with respect to the fine-structure line emission) in NGC 7469 compared to global measurements of other LIRGs and SF galaxies. We note that only PAHs at wavelengths  $>5 \mu\text{m}$  are considered in Stierwalt et al. (2014), so there is no inclusion of the 3.3  $\mu\text{m}$  PAH, which typically contributes less than 3% of the total PAH emission in SF galaxies (Lai et al. 2020) but would nevertheless raise the GOALS points even further.

In Figure 4(b), the PAH 3.3-to-Ne ratios do not show a trend when [Ne III]/[Ne II]  $<0.25$ , but as the [Ne III]/[Ne II] ratio becomes larger ( $>0.25$ ), the PAH 3.3/neon ratio drops by



**Figure 3.** Left: the aliphatic ( $3.4 \mu\text{m}$ ) to aromatic ( $3.3 \mu\text{m}$ ) ratio as a function of the hardness of the radiation field probed by  $[\text{Ne III}]/[\text{Ne II}]$  in NGC 7469. The circles indicate cells in the SF ring, the squares indicate cells in the gap region, and the star indicates the nucleus. The color indicates the projected linear distance from the AGN. The two dashed lines and the shaded regions indicate the averages and  $1\sigma$  of  $3.4/3.3$  in the gap region cells (squares) and the SF ring (circles). On average, the  $3.4/3.3$  ratio in the gap decreases by 20% compared to that of the SF ring. The  $3.4/3.3$  ratio shows a large scatter when  $[\text{Ne III}]/[\text{Ne II}] > 0.5$ . The  $3.4 \mu\text{m}$  feature at the nucleus is not detected; a  $3\sigma$  upper limit is shown here. Right: the  $11.3/3.3$  PAH ratio (a diagnostic of the average sizes of the PAH grains) as a function of the hardness of the radiation field. The two dashed lines and the shaded regions indicate the averages and  $1\sigma$  of PAH  $11.3/3.3$  in the gap region cells (squares) and the SF ring (circles). Cells in the gap region show, on average, a factor of 2 higher  $11.3/3.3$  compared to those in the SF ring, suggesting a larger average PAH grain size in regions with a harder radiation field. This may be due to the preferential destruction of small grains by the AGN. The nucleus appears to show an anomalously low  $11.3/3.3$  ratio but with a relatively large uncertainty.

at least 50%, most clearly seen in the gap cells between the nucleus and the ring. The relative enhancement of the  $3.3 \mu\text{m}$  PAH emission compared to other PAH features is evident in Figure 4(c). Here a striking picture emerges when we zoom in on the scale of the SF ring, with each cell having a width of  $\sim 200$  pc, and find PAH  $3.3/\Sigma\text{PAH}$  ratios that can be up to  $4\times$  higher than the global measurements seen in other starburst galaxies. This excess  $3.3 \mu\text{m}$  PAH persists compared to most starburst galaxies even in cells close to the nucleus, suggesting that the regions viewed by the IFU in NGC 7469 have a smaller grain size distribution than those measured in starburst galaxies on kiloparsec scales. A detailed discussion of this excess can be found in Section 5.2. Another PAH band exhibiting a noteworthy enhancement is the  $6.2 \mu\text{m}$  feature; however, its maximum increase is limited to 30%, a magnitude that falls significantly short of the enhancement observed in the  $3.3 \mu\text{m}$  PAH band.

## 5. Discussion

### 5.1. Photodestruction of PAH

In Section 4, we showed that the hardness of the radiation field plays a pivotal role in shaping the properties of the small dust grains responsible for PAH emission in NGC 7469, particularly the aliphatic-to-aromatic ratio and the grain sizes. The typical ratio of  $3.4/3.3$  in NGC 7469 is 3%–8% (Figure 3, left panel), with the lowest values seen where  $0.5 < [\text{Ne III}]/[\text{Ne II}] < 1$ . The  $3.4/3.3$  has been shown to increase by up to  $\sim 20\%$  in some PDR environments, such as the reflection nebula NGC 7023 (Pilleri et al. 2015) and normal SF galaxies (Lai et al. 2020). The decrease of the  $3.4/3.3$  seen in the gap between the nucleus and the SF ring of NGC 7469 indicates further grain processing within 500 pc of the AGN. The nucleus itself ( $r < 100$  pc) shows no detectable aliphatic emission.

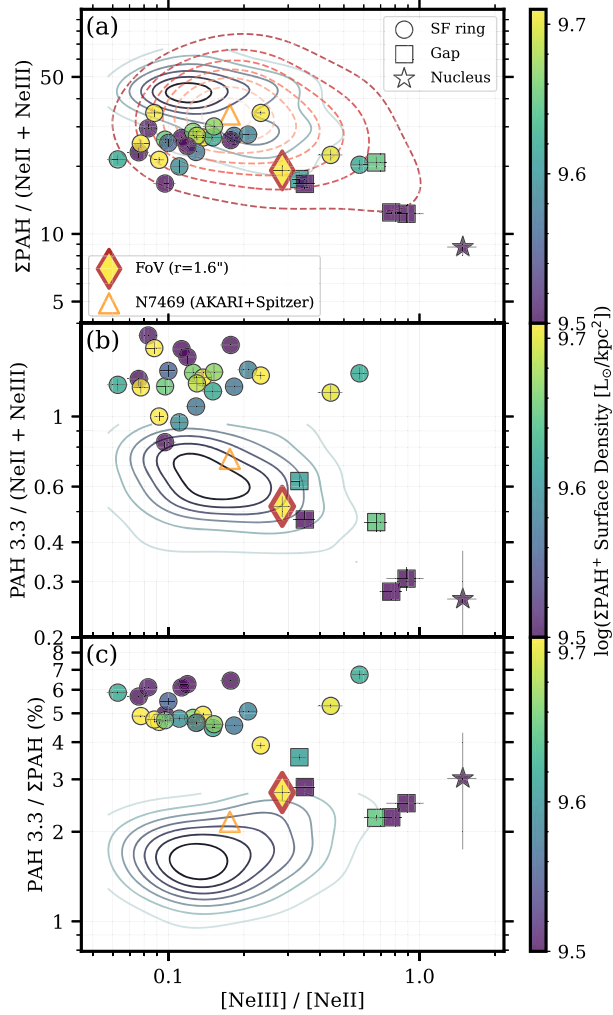
Likewise, in Figure 3 (right), the grain sizes in the gap cells are larger compared to those in the SF ring, suggesting the absence of smaller grains closer to the AGN. Nevertheless, a relatively low  $11.3/3.3$  ratio is apparently seen on the nucleus,

albeit with a large uncertainty and weak overall PAH emission. The relatively low  $11.3/3.3$  ratio on the NGC 7469 nucleus is similar to the low  $7.7/6.2$  nuclear ratio presented in Lai et al. (2022). Recently, Donnelly et al. (2023) successfully modeled suppressed long-to-short PAH ratios in the nucleus of NGC 4138 (seen in the Spitzer data) by combining AGN, stellar bulge, and disk emission. That AGN can directly excite PAHs has also been suggested by Jensen et al. (2017). To explain the low  $11.3/3.3$  in the nucleus of NGC 7469, careful modeling of the balance between the destruction and illumination of PAHs in the vicinity of a powerful AGN, including ambient stellar light from old and young stars, would be necessary.

### 5.2. The Excess of PAH 3.3

The apparent relative excess of the  $3.3 \mu\text{m}$  PAH in NGC 7469 compared to SF galaxies is surprising, since the  $3.3 \mu\text{m}$  PAH emission traces the smallest PAH population in the ISM ( $N_C \sim 50$ ; Draine et al. 2021), which is thought to be most vulnerable to photodestruction. In fact, all of the PAHs are suppressed in NGC 7469 (Figure 4(a)) compared to galaxies powered by starbursts, while the  $3.3 \mu\text{m}$  PAH is not suppressed as much as the other PAHs in the SF ring (Figure 4(b)). This leads to the relative enhancement in the  $3.3 \mu\text{m}$  PAH flux. The  $3.3 \mu\text{m}$  PAH does exhibit a slight positive correlation with radiation hardness, dropping at the highest values of the  $[\text{Ne III}]/[\text{Ne II}]$  ratio. The enhancement of the  $3.3 \mu\text{m}$  PAH emission in NGC 7469 can be seen even more clearly when studying the fractional PAH 3.3 in the ring, which can be  $\sim 3\text{--}4\times$  higher compared to starburst galaxies observed with AKARI+Spitzer (Figure 4(c)). Even in the gap regions, where the PAHs are clearly affected by the AGN, the fractional  $3.3 \mu\text{m}$  PAH power still lies at the upper end of the galaxy distribution. This enhancement of the  $3.3 \mu\text{m}$  PAH emission is consistent with other recent findings based on the JWST NIRCcam and MIRI imaging of normal SF galaxies that seem to suggest an increase of the  $3.3 \mu\text{m}$  PAH contribution in the F335M filter in regions with harder and more intense radiation fields (Chastenot et al. 2023). Our result also suggests that





**Figure 4.** Fractional PAH emission as a function of radiation field hardness in NGC 7469. The data are color-coded by the surface density of the total ionized PAH, consisting of the 6.2, 7.7, and 8.6  $\mu\text{m}$  PAH bands, which traces the amount of photoelectric heating in the ISM. Symbols are the same as in Figure 3, with the addition of the red diamond representing values measured from the total IFU spectrum and the orange triangle representing values measured from the combined AKARI+Spitzer spectrum. For comparison, the red dashed contours indicate the distribution of LIRGs taken from Stierwalt et al. (2014), and the blue solid contours indicate the distribution of the 62 PAH-bright galaxies (mostly LIRGs that are also in GOALS) taken from Lai et al. (2020) but refitted with CAFE. While the  $\Sigma\text{PAH}$ -to-Ne ratio of the SF ring in NGC 7469 (top) is consistent with the lower envelope of LIRG values, the PAH 3.3-to-Ne ratio (middle) is high, consistent with a relative excess of 3.3 PAH emission compared to other PAH features (bottom), where the values in the ring of NGC 7469 are up to  $4\times$  higher than in most starbursting LIRGs (see Sections 4.3 and 5.2 for more detailed discussions).

those cells in the SF ring having higher fractional 3.3  $\mu\text{m}$  PAH emission also have a lower total ionized PAH surface density, which is not surprising, as 3.3  $\mu\text{m}$  PAH emission mainly traces neutral PAHs.

When comparing our JWST IFU to Spitzer results, it is important to realize that the physical areas covered by our single cell in the JWST IFU and the Spitzer slit in NGC 7469 differ substantially, with the cell probing a region of 200 pc versus the IRS slit, which summed the light over  $\sim 3$  kpc scales. Based on Figures 4(a) and (c), the total PAH is more suppressed in the central 1 kpc region (diamond; extracted by an aperture with  $r = 1''.6$ ) compared to the coarser 3 kpc area (triangle), whereas the fractional 3.3  $\mu\text{m}$  PAH is more

enhanced. We have ruled out the possibility of systematic differences in the flux estimates by CAFE when analyzing spectra with diverse spectral resolutions ( $R = 100$  versus  $R = 2700$ ) being the cause of such a PAH 3.3 enhancement. A test has been done to smooth the high-resolution JWST spectrum of the total IFU extraction to match with Spitzer’s, and the ratios of  $3.3/\Sigma\text{PAH}$  show only a negligible difference ( $< 2\%$ ).

Understanding the physics behind the enhanced durability of the smallest dust particles, which give rise to the excess emission from PAH 3.3 emitters, carries significant implications for simulating the dust mass buildup and chemical enrichment throughout cosmic time (e.g., Narayanan et al. 2023). Theoretical studies have shown that molecules consisting of fewer carbon atoms ( $N_c \lesssim 40$ ) are able to efficiently dissipate the absorbed UV energy via recurrent fluorescence (RF), with a relaxation timescale of  $\sim$ milliseconds as opposed to IR emission with timescales of  $\sim$ seconds (Leger et al. 1988). Recently, laboratory studies have reproduced RF as a cooling mechanism for PAH molecules (e.g., Bernard et al. 2017, 2023; Navarro Navarrete et al. 2023; Stockett et al. 2023), indicating that RF can be a powerful mechanism to significantly enhance the survival rate of small particles under high UV radiation densities. Another potential explanation for excess short-wavelength PAH emission is the fact that larger PAHs can “act” like smaller PAHs in high radiation field environments. This arises due to the shorter mean time between photon absorptions in a high UV radiation density environment, leading PAH molecules to have higher vibrational energy, which in turn will emit at shorter wavelengths (Draine et al. 2021). The nature of the small grains and their energy-loss channels, along with their resilience under harsh conditions, will greatly impact the evolution of the ISM, the heating of gas via the photoelectric effect, and the UV/optical extinction law. JWST now allows us to probe significantly smaller physical scales than was previously possible and trace the true variation in the 3.3  $\mu\text{m}$  emission within galaxies over a wide range of environments.

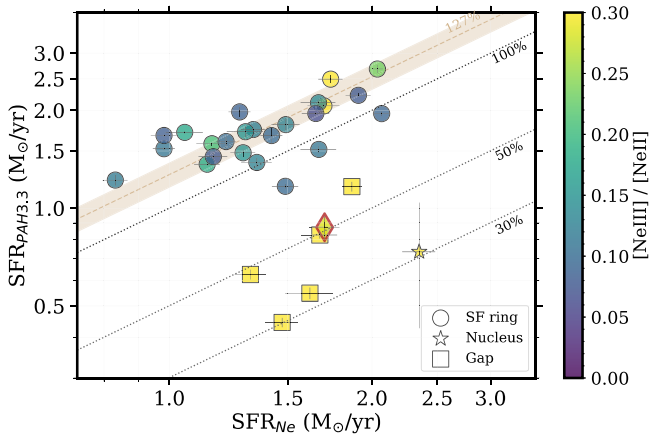
### 5.3. PAH 3.3 as an SFR Indicator

JWST has opened up the possibility of spectroscopically detecting PAH emission in large numbers of high- $z$  galaxies (Spilker et al. 2023) and, in particular, using this feature as an SFR indicator. The 3.3  $\mu\text{m}$  PAH is the last band to shift beyond the long-wavelength coverage of MIRI, making it a prime tool to probe star formation out to redshifts of  $z \sim 6$ . At low redshifts, it has been shown that there is a tight correlation between the SFR probed by PAH 3.3 and  $[\text{Ne II}] + [\text{Ne III}]$  emission in metal-rich SF galaxies; metal-poor galaxies, on the other hand, exhibit peculiar behavior that may lead to an underestimation of the SFR via the use of PAHs by nearly an order of magnitude in some cases (Lai et al. 2020).

Figure 5 shows the comparison between the SFR derived from the 3.3  $\mu\text{m}$  PAH (Lai et al. 2020), which can be written as

$$\log\left(\frac{\text{SFR}}{M_{\odot} \text{ yr}^{-1}}\right) = -(6.80 \pm 0.18) + \log\left(\frac{L_{\text{PAH } 3.3}}{L_{\odot}}\right), \quad (1)$$

and that from the neon lines, which is based on Equation (13) in Ho & Keto (2007), with the fraction of ionized neon emission derived from Equation (5) in Zhuang et al. (2019). An independent test has also been conducted by comparing our



**Figure 5.** The PAH 3.3–derived SFR vs. the neon–derived SFR in NGC 7469. The two SFR indicators follow each other within the uncertainties in SF cells. On average,  $\text{SFR}_{\text{PAH}3.3}$  is slightly (27%) higher compared to  $\text{SFR}_{\text{Ne}}$ , as indicated by the 1:1 line (100%). The shaded area indicates the  $1\sigma$  range of the SF ring cells. However, the gap and nucleus show significant lower PAH–derived SFRs, as expected, due to both weaker PAH emission and enhanced neon emission from the AGN. The red diamond represents the total extraction from the FoV of the IFU (arbitrarily divided by 20 on both axes to allow it to be plotted with the starburst ring points) and shows an at least  $\sim 50\%$  drop in the  $3.3\ \mu\text{m}$  PAH–inferred SFR, as might be measured for NGC 7469 at high redshift and lower spatial resolution.

PAH–derived SFR to the PAH 6.2–to–IR luminosity relation reported in Table 4 of Gruppioni et al. (2016). By assuming the median 6.2/3.3 PAH ratio of 3.08 found in our SF cells, the SFRs provided in Gruppioni et al. (2016) and Lai et al. (2020) are consistent at the 10% level.

In Figure 5, on average, the measurements in the SF ring follow the 1:1 line, showing slightly higher (27%)  $\text{SFR}_{\text{PAH}3.3}$  than  $\text{SFR}_{\text{Ne}}$  with a small spread of 0.09 dex. However, the gap region cells and the nucleus show substantial deviations from the correlation due to the suppressed PAH emission in and around the AGN. An observation of a galaxy like NGC 7469 at high redshift, where the AGN and ring were spatially unresolved, would result in an underestimate of the SFR by at least 50% if measured from PAH 3.3 alone due to the mixing of the AGN and starburst emission in the beam.

## 6. Summary

With the advent of JWST, we can now probe the ISM in the vicinity of a powerful AGN on scales of a few hundred parsecs. In this paper, we present observations of the ISM in the nearby Seyfert galaxy NGC 7469 using NIRSpc IFU, supported by MIRI observations first reported in Lai et al. (2022). In particular, we focus on the properties of small and neutral aromatic grains traced by the  $3.3\ \mu\text{m}$  PAH band, together with a nearby aliphatic feature at  $3.4\ \mu\text{m}$ . Our findings can be summarized as follows.

1. With the NIRSpc IFU, we are able for the first time to confidently separate the dust features at  $3\ \mu\text{m}$  in a resolved extragalactic source. In NGC 7469, the aliphatic ( $3.4\ \mu\text{m}$ ) to aromatic ( $3.3\ \mu\text{m}$ ) ratio varies by a factor of 2 across the IFU FoV, and this ratio is slightly lower near the AGN. This low 3.4/3.3 flux ratio, although with scatter, is most visible in the region between the Seyfert nucleus and the starburst ring, where the  $[\text{Ne III}]/[\text{Ne II}]$  line flux ratio reaches values of 0.5–1.0. The 3.4/3.3 ratio

at the nucleus is very low compared to all other regions, suggesting extreme photodestruction of the grains by the AGN.

2. The 11.3/3.3 PAH ratio, which is a measure of the average grain size, is significantly larger ( $>2\times$ ) in the vicinity of the AGN than it is in the starburst ring, suggesting a destruction of small grains within 300 pc of the nucleus, consistent with the results of Paper I. The nuclear spectrum displays a seemingly low 11.3/3.3 ratio, although with a large uncertainty, as all of the PAH fluxes are suppressed in this region.
3. The total PAH–to–Ne ratio throughout the NIRSpc field is low in NGC 7469 compared to most local LIRGs and star-forming galaxies. However, the 3.3–to–Ne ratio and the fractional  $3.3\ \mu\text{m}$  PAH power ( $3.3/\Sigma\text{PAH}$ ) in the starburst ring of NGC 7469 are high compared to most LIRGs, which may be due to recurrent fluorescence of small grains or multiple photon absorption by large grains on 100 pc scales.
4. An unresolved measurement of the PAH and Ne line emission in the central region of NGC 7469 that includes both the AGN and the star-forming ring, such as might be observed with JWST at high redshift, would lead to an underestimation of its PAH–derived star formation rate (SFR) by at least 50%. Thus, caution should be exercised when using PAHs to measure SFR in high- $z$  AGN hosts or galaxies where AGN cannot be safely ruled out.

Here we have shown that the  $3.3\ \mu\text{m}$  PAH feature varies on subkiloparsec scales in NGC 7469, and that these variations can reveal changes in the grain populations as they react to changes in the ISM properties around an AGN and an intense starburst in this composite galaxy. Further observations of other nearby sources like NGC 7469 will be critical to fully understand, from a statistical standpoint, the expected variations in the PAH–band fluxes and ratios that we can expect to see at high redshift with JWST in the coming years.

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*Facilities:* JWST (NIRSpec & MIRI), MAST, NED.

*Software:* Astropy (Astropy Collaboration et al. 2013, 2018, 2022), CAFE (Marshall et al. 2007), JWST Science Calibration Pipeline (Bushouse et al. 2022), Imfit (Newville et al. 2014), Matplotlib (Hunter 2007), Numpy (van der Walt et al. 2011; Harris et al. 2020), QFitsView (Ott 2012), SciPy (Virtanen et al. 2020).

### ORCID iDs

Thomas S.-Y. Lai (賴劭愉) <https://orcid.org/0000-0001-8490-6632>  
 Lee Armus <https://orcid.org/0000-0003-3498-2973>  
 Marina Bianchin <https://orcid.org/0000-0002-6570-9446>  
 Tanio Díaz-Santos <https://orcid.org/0000-0003-0699-6083>  
 Sean T. Linden <https://orcid.org/0000-0002-1000-6081>  
 George C. Privon <https://orcid.org/0000-0003-3474-1125>  
 Hanae Inami <https://orcid.org/0000-0003-4268-0393>  
 Vivian U <https://orcid.org/0000-0002-1912-0024>  
 Thomas Bohn <https://orcid.org/0000-0002-4375-254X>  
 Aaron S. Evans <https://orcid.org/0000-0003-2638-1334>  
 Kirsten L. Larson <https://orcid.org/0000-0003-3917-6460>  
 Brandon S. Hensley <https://orcid.org/0000-0001-7449-4638>  
 J.-D. T. Smith <https://orcid.org/0000-0003-1545-5078>  
 Matthew A. Malkan <https://orcid.org/0000-0001-6919-1237>  
 Yiqing Song <https://orcid.org/0000-0002-3139-3041>  
 Sabrina Stierwalt <https://orcid.org/0000-0002-2596-8531>  
 Paul P. van der Werf <https://orcid.org/0000-0001-5434-5942>  
 Jed McKinney <https://orcid.org/0000-0002-6149-8178>  
 Susanne Aalto <https://orcid.org/0000-0002-5828-7660>  
 Jeff Rich <https://orcid.org/0000-0002-5807-5078>  
 Vassilis Charmandaris <https://orcid.org/0000-0002-2688-1956>  
 Philip Appleton <https://orcid.org/0000-0002-7607-8766>  
 Loreto Barcos-Muñoz <https://orcid.org/0000-0003-0057-8892>  
 Torsten Böker <https://orcid.org/0000-0002-5666-7782>  
 Luke Finnerty <https://orcid.org/0000-0002-1392-0768>  
 Justin A. Kader <https://orcid.org/0000-0002-6650-3757>  
 David R. Law <https://orcid.org/0000-0002-9402-186X>  
 Anne M. Medling <https://orcid.org/0000-0001-7421-2944>  
 Michael J. I. Brown <https://orcid.org/0000-0002-1207-9137>  
 Christopher C. Hayward <https://orcid.org/0000-0003-4073-3236>  
 Justin Howell <https://orcid.org/0000-0001-6028-8059>  
 Kazushi Iwasawa <https://orcid.org/0000-0002-4923-3281>  
 Francisca Kemper <https://orcid.org/0000-0003-2743-8240>  
 Jason Marshall <https://orcid.org/0000-0001-7712-8465>

Joseph M. Mazzarella <https://orcid.org/0000-0002-8204-8619>  
 Francisco Müller-Sánchez <https://orcid.org/0000-0002-2713-0628>  
 Eric J. Murphy <https://orcid.org/0000-0001-7089-7325>  
 David Sanders <https://orcid.org/0000-0002-1233-9998>  
 Jason Surace <https://orcid.org/0000-0001-7291-0087>

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