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Too dense to go through: The role of low-mass clusters in the pre-processing of satellite galaxies

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ABSTRACT

We study the evolution of satellite galaxies in clusters of the C-EAGLE simulations, a suite of 30 high-resolution cosmological hydrodynamical zoom-in simulations based on the EAGLE code. We find that the majority of galaxies that are quenched at $z = 0$ ($\gtrsim 80\%$) reached this state in a dense environment ($\log_{10} M_{200} [M_{\odot}] \geq 13.5$). At low redshift, regardless of the final cluster mass, galaxies appear to reach their quenching state in low mass clusters. Moreover, galaxies quenched inside the cluster that they reside in at $z = 0$ are the dominant population in low-mass clusters, while galaxies quenched in a different halo dominate in the most massive clusters. When looking at clusters at $z > 0.5$, their in-situ quenched population dominates at all cluster masses. This suggests that galaxies are quenched inside the first cluster they fall into. After galaxies cross the cluster's r_{200} they rapidly become quenched ($\lesssim 1$ Gyr). Just a small fraction of galaxies ($\lesssim 15\%$) is capable of retaining their gas for a longer period of time, but after 4 Gyr, almost all galaxies are quenched. This phenomenon is related to ram pressure stripping and is produced when the density of the intracluster medium reaches a threshold of $\rho_{\text{ICM}} \sim 3 \times 10^{-5} n_{\text{H}} \text{ (cm}^{-3}\text{)}$. These results suggest that galaxies start a rapid-quenching phase shortly after their first infall inside r_{200} and that, by the time they reach r_{500} , most of them are already quenched.

Key words: galaxies:clusters: general – galaxies: evolution – galaxies: formation – galaxies: star formation – galaxies: haloes

1 INTRODUCTION

Much effort has been devoted to understanding the physical processes that drive galaxy evolution. It has been well established that the environment plays a decisive role in shaping important properties of galaxies (Gunn & Gott 1972; Dressler 1980, 1984; Moore et al. 1996; Poggianti et al. 2001; Blanton & Moustakas 2009); while galaxies located in low-density environments typically show bluer colours and late-type morphology, galaxies located in dense environments, such as galaxy clusters, show redder colours and an early-type morphology (Gómez et al. 2003; Kauffmann et al. 2004; Poggianti et al. 2006; Moran et al. 2007; Blanton & Moustakas 2009). Even within clusters, there are differences between the population of galaxies located on the outskirts and those in the inner, denser regions

(Postman et al. 2005; Fasano et al. 2015; Haines et al. 2015; Brough et al. 2017; Cava et al. 2017).

This colour/morphology-density relation can be explained as a decrease in the star formation activity of galaxies located in denser environments and it is thought to be a consequence of the cold gas depletion of the satellite population (Haynes et al. 1984; Gavazzi et al. 2006; Fabello et al. 2012; Catinella et al. 2013; Hess & Wilcots 2013). Nevertheless, the main mechanism responsible for this gas depletion is still an open question. There are several processes that can cause a sharp decrease in the star formation activity of a galaxy. These can be broadly classified as events of ‘mass quenching’ and ‘environmental quenching’ (Peng et al. 2010). The former correspond to mechanisms related to internal galaxy processes influenced by the galaxy mass such as gas outflows driven by the presence of an active galaxy nucleus (AGN) (Croton et al. 2006; Fabian 2012; Ciccone et al. 2014), or supernova feedback and stel-

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lar winds (Larson 1974; Dekel & Silk 1986; Efstathiou 2000; Cantalupo 2010). Environmental quenching, on the other hand, refers to mechanisms related to the interaction between a galaxy and its local environment. Simulations and models have shown that in particular for the case of galaxy groups and clusters, the main mechanisms associated with environmental quenching can be separated into three broad categories (for details see Boselli & Gavazzi 2006; Jaffé et al. 2016): gravitational interactions between galaxies, gravitational interactions between cluster and galaxies and interactions between galaxies and the intra-cluster medium (ICM).

It is expected that at least some of these mechanisms act simultaneously, thus rendering the characterization of galaxy quenching in dense environments challenging. Nevertheless, over the last decade several studies have suggested that ram pressure and starvation are the main driven responsible for environmental quenching (De Lucia et al. 2012; Wetzel et al. 2013; Muzzini et al. 2014; Wetzel et al. 2015; Peng et al. 2015; Foltz et al. 2018). Wetzel et al. (2012) presented a model in which quenching is driven in a ‘*delayed-then-rapid*’ fashion in which galaxies, after its infall into a massive halo, keep forming stars as a central for several Gyr and, once the satellite star formation quenching begins, occurs rapidly.

In a recent observational study, Roberts et al. (2019, hereafter R19) showed that, inside clusters, galaxies experience what they call a ‘*slow-then-rapid*’ quenching scenario, especially for low mass galaxies. In the scenario proposed by R19, when a galaxy first enters the virial radius, R_{vir} , of a cluster it starts a slow quenching phase that can last between 1-2.5 Gyr until the galaxy reaches a region where the characteristic density of the ICM reaches $\rho_{\text{ICM}} \sim 10^{-28.3} \text{ gr cm}^{-3}$ ($\sim 3 \times 10^{-5} n_{\text{H}} \text{ cm}^{-3}$). At this point, the galaxy starts a rapid quenching phase, that can take between 0.5-1 Gyr. The authors suggest that this rapid quenching phase could be associated to ram pressure stripping events, but this is hard to conclude from observations alone.

Cosmological simulations are a powerful tool that allows us to disentangle the effects associated with the different processes that are typically coupled in a non-linear fashion. In the last years, analytic models (Fujita 2004; Mok et al. 2014; Contini et al. 2020), semi-analytic (De Lucia et al. 2012; Wetzel et al. 2013; Henriques et al. 2017; Stevens & Brown 2017; Cora et al. 2018; Contini et al. 2019) and hydrodynamical simulations (e.g. Bahé et al. 2013; Bahé & McCarthy 2015; Taylor et al. 2017; Wright et al. 2019; Pallero et al. 2019; Donnari et al. 2020) have allowed us to explore the relative importance of the different mechanisms that bring a galaxy to a final quenched state. In Pallero et al. (2019, hereafter P19), we studied the satellite population of galaxy clusters in the publicly available data from the EAGLE simulation. We found that most of these galaxies reach their quenching state inside galaxy clusters. However, our results suggested that the fraction of galaxies that reach clusters already quenched grows with the final cluster mass.

Quenching of satellites with stellar masses $\gtrsim 10^9 M_{\odot}$ typically took place in structures with masses in the range of low mass clusters ($\sim 10^{14} M_{\odot}$). Unfortunately, our study suffered from low number statistics as EAGLE only included 10 low and intermediate-mass clusters.

In this work, we will add new evidence to address these

open issues and to disentangle the mechanisms that finally lead to the quenching of the star formation in the satellite population of galaxy clusters.

To this end, we present a detailed study of the quenching history of the satellite population in galaxy clusters from the CLUSTER-EAGLE simulations (Bahé et al. 2017; Barnes et al. 2017b), a set of zoom-in high resolution hydrodynamical simulations of 30 galaxy clusters spanning a mass range of $14 \leq \log_{10} M_{200}^{z=0} [M_{\odot}] \leq 15.4$. Thanks to the wide mass range of the modelled clusters, the same mass resolution as EAGLE simulations, and an improved time resolution at low z , C-EAGLE is an ideal laboratory to study which is the dominant quenching mechanism associated with environmental quenching. This paper is organized as follows. Section 2 contains a brief summary of the simulations and a description of the method to identify structures in space and time. In Section 3 we characterize the properties of each cluster and their satellite populations at $z = 0$ and study how their properties evolve as a function of the time they spent inside clusters. In Section 4 we discuss how, in our models, ram pressure can account almost completely for the quenching of satellite galaxies. We also discuss the impact of numerical resolution on our results. Finally, in Section 5, we present a summary of our main conclusions.

2 SIMULATIONS

2.1 C-EAGLE simulations

Here we briefly summarize the main characteristics of the C-EAGLE simulations. For a more detailed description of the simulations and their subgrid model, see Barnes et al. (2017b) and Bahé et al. (2017). The C-EAGLE project comprises a suite of 30 cosmological hydrodynamical zoom-in simulations of massive galaxy clusters, spanning a mass range between $14 \leq \log_{10} M_{200}^{z=0} [M_{\odot}] \leq 15.4$. The clusters were selected from a parent N-body low resolution simulation with a box of size $(3.2 \text{ Gpc})^3$, first presented in Barnes et al. (2017a). Each zoom-in simulation was performed adopting the same flat Λ CDM cosmology that was used in the EAGLE simulation (Schaye et al. 2015; Crain et al. 2015) corresponding to $\Omega_{\Lambda} = 0.693$, $\Omega_m = 0.307$, $\Omega_b = 0.04825$, $\sigma_8 = 0.8288$, $Y = 0.248$ and $H_0 = 67.77 \text{ km s}^{-1}$ (Planck Collaboration et al. 2014).

All simulations were performed with the variant ‘AG-NdT9’ of the EAGLE simulations (Schaye et al. 2015), with a dark matter and baryonic mass resolution (gas and stars) of $m_{\text{DM}} = 9.7 \times 10^6 M_{\odot}$ and $m_{\text{gas}} = 1.8 \times 10^6 M_{\odot}$, respectively. The gravitational softening was set to $\epsilon = 2.66$ comoving kpc at $z \geq 2.8$, and set to $\epsilon = 0.7$ physical kpc at $z < 2.8$.

The zoom-in simulations were run with a high-resolution region extending at least to $5r_{200c}$, with the 24 galaxy clusters belonging to the HYDRANGEA sub-sample extending their high-resolution region to at least $10r_{200c}$. For this study, two of the C-EAGLE clusters were removed (CE-10 and CE-27). These clusters experienced a very dramatic numerical AGN outburst at high redshift that significantly and artificially affected the overall properties of the system.

The code used to run these models is an upgraded version of the N-Body Tree-PM SPH code GADGET 3 (described in Springel 2005). The modifications include updates to the hydrodynamics scheme, collectively known

as ‘ANARCHY’ (see Appendix A in Schaye et al. (2015) and Schaller et al. (2015) for details) and several subgrid physics models to simulate unresolved properties. Radiative cooling and photo-heating are implemented following Wiersma et al. (2009a), assuming an optically thin X-Ray/UV background (Haardt & Madau 2001). Star formation is implemented stochastically based on the Kennicutt-Schmidt law (Kennicutt 1998) in pressure law form as in Schaye & Dalla Vecchia (2008), using the metallicity-dependent density threshold as in Schaye (2004). Each particle is assumed to be a single-age stellar population, with a Chabrier initial mass function in the range $0.1M_{\odot} - 100M_{\odot}$ (Chabrier 2003).

Stellar evolution is modelled following Wiersma et al. (2009b), and chemical enrichment is followed for the 11 elements that most contribute to radiative cooling (i.e., H, He, C, N, O, Ne, Mg, Si, S, Ca, and Fe). The thermal energy released by stellar feedback is stochastically distributed among the gas particles surrounding the event without any preferential direction (Dalla Vecchia & Schaye 2012).

The galaxy formation model used in the C-EAGLE project has been widely tested for the EAGLE simulations. Schaye et al. (2015) showed that the model is successful at reproducing several observable quantities that were not directly calibrated in the simulations. As an example, galaxies in the EAGLE simulation have a realistic neutral gas content fraction Bahé et al. (2016); Crain et al. (2017), a key property when studying the evolution of the star formation in galaxies. Moreover, in Barnes et al. (2017b) it is shown that the model reproduces with reasonable good agreement important properties of the intracluster medium, such as the gas fraction-total mass relation and its metallicity. The predicted X-ray and Sunyaev-Zeldovich properties of clusters show a good match with the spectroscopic temperature, X-ray luminosity, and Y_{SZ} -mass relation. Regarding the stellar content in galaxies, the cluster satellite’s stellar mass function in the C-EAGLE project is in excellent agreement with observations up to $z \sim 1.5$ (Bahé et al. 2017; Ahad et al. 2021). Additionally, the satellite quenched fractions are also in reasonably good agreement with observations at $z = 0$ (Bahé et al. 2017) and subhalo masses are well matched to lensing observations (Bahé 2021).

2.2 Halo Identification

The main outputs from each simulation correspond to 30 snapshots spaced between $z = 14$ and $z = 0$, 28 of them equidistant in time with a $\Delta t = 500$ Myr. Two additional snapshots at $z = 0.366$ and $z = 0.101$ were added to facilitate comparisons with EAGLE. All these outputs were later post-processed to identify structures in each snapshot using the SUBFIND algorithm (Springel et al. 2001; Dolag et al. 2009) in a two-step procedure, as described below.

In order to define bound halos, first, a friends-of-friends algorithm (FoF) is applied to all dark matter particles using a linking length $b = 0.2$ times the mean interparticle distance. Baryons are then assigned to the FoF (if any) associated with their nearest dark matter particle. If a FoF halo possesses fewer than 32 dark matter particles, it is considered unresolved and discarded. As a second step, SUBFIND identifies any gravitationally self-bound substructures (or ‘subhaloes’) within a FoF halo taking dark matter and

baryons into consideration. These subhaloes are identified as local overdensities using a binding energy criterion. For a more detailed description of the method, we refer to Springel et al. (2001) and Dolag et al. (2009).

SUBFIND identifies structures at a single point in time. Since we need to follow galaxies both in space and time an additional procedure is required to create merger trees. The trees analyzed in this work were obtained with the SPIDERWEB algorithm (see Bahé et al. 2019, Appendix A). Differently from other ‘merger tree’ algorithms, SPIDERWEB was designed with the purpose to identify descendants for as long as possible in crowded environments. For that purpose, SPIDERWEB consider as descendants all subhaloes that share particles between consecutive snapshots. This treatment becomes especially relevant in groups and cluster of galaxies, where galaxy-galaxy encounters become more common. The satellite galaxies in our sample correspond to those satellite galaxies that are inside the r_{200} of the most massive structures of each zoom-in simulation. Throughout this paper, we will define as galaxies all those subhalos that posses an stellar mass $M_{\text{star}} > 10^8 M_{\odot}$, and a total mass $M_{\text{galaxy}} > 10^9 M_{\odot}$. As a result, we ensure that each of our galaxies is resolved with at least, 1000 dark matter particles at any snapshot. Additionally, as we also want to understand the reasons behind the satellite quenching in our simulations, we choose to keep the mass threshold in which the star formation evolution, stellar mass function, and mass-size relation are reasonably well-calibrated for the galaxy formation model (Schaye et al. 2015; Furlong et al. 2015, 2017). To avoid detection of galaxies with no primordial origin, we only followed the evolutionary history of those galaxies whose merger trees can be continuously followed at least up to $z = 2$. It is worth noting that, as discussed by Bahé et al. (2017), the C-EAGLE simulation show an excess of quenched galaxies in the low-mass end ($M_{\star} < 10^{9.5} M_{\odot}$). This result is believed to be a combination of environmental quenching effects and limitations related to the resolution of the simulation (Schaye et al. 2015). Note, however, that in our work we focus on galaxies that undergo their quenching process purely due to environmental mechanisms. This means that most of our selected galaxies would have reached a higher stellar mass at $z = 0$ if not affected by the physical processes explored in this article. As a result, our conclusions are not significantly affected by the stellar mass threshold chosen to select satellite galaxies. This is further discussed in Appendix A, where we explored the effects of choosing a higher stellar mass threshold ($M_{\star} > 10^{9.5} M_{\odot}$) for the satellite selection. We discuss this over-quenching in more detail in Section 4.2 and the impact that could have on our results, in order to avoid spurious results, we will neglect the population of low mass galaxies that reach their quenching state as centrals, and only take into account those galaxies that reach their quenched state as satellites. Nonetheless, regardless of the cluster mass, galaxies that reach their quenching state while centrals correspond to less than 10% of all the galaxies in our cluster sample.

2.3 Ram Pressure and Restoring Force models

To measure the instantaneous ram pressure experienced by cluster satellites galaxies, we follow the methodology described in Vega-Martínez et al. (2021). They introduce a

ID Cluster	$M_{500}^{z=0.1}$	$R_{500}^{z=0.1}$	$E_{\text{kin}}/E_{\text{thermal}}$	Virialization State
CE-0	13.905	0.65	0.11	Unrelaxed
CE-1	13.982	0.69	0.24	Unrelaxed
CE-2	13.920	0.66	0.05	Relaxed
CE-3	13.926	0.66	0.06	Relaxed
CE-4	13.853	0.62	0.11	Unrelaxed
CE-5	13.892	0.64	0.09	Relaxed
CE-6	14.128	0.77	0.09	Relaxed
CE-7	14.176	0.80	0.09	Relaxed
CE-8	14.070	0.74	0.09	Relaxed
CE-9	14.244	0.84	0.16	Unrelaxed
CE-11	14.290	0.87	0.21	Unrelaxed
CE-12	14.422	0.96	0.08	Relaxed
CE-13	14.341	0.91	0.06	Relaxed
CE-14	14.563	1.08	0.31	Unrelaxed
CE-15	14.407	0.95	0.22	Unrelaxed
CE-16	14.311	0.89	0.12	Unrelaxed
CE-17	14.537	1.05	0.25	Unrelaxed
CE-18	14.639	1.14	0.09	Relaxed
CE-19	14.586	1.09	0.30	Unrelaxed
CE-20	14.482	1.01	0.12	Unrelaxed
CE-21	14.800	1.29	0.29	Unrelaxed
CE-22	14.837	1.33	0.12	Unrelaxed
CE-23	14.426	0.97	0.26	Unrelaxed
CE-24	14.821	1.31	0.13	Unrelaxed
CE-25	15.045	1.56	0.31	Unrelaxed
CE-26	14.899	1.39	0.08	Relaxed
CE-28	14.902	1.39	0.13	Unrelaxed
CE-29	15.077	1.60	0.30	Unrelaxed

Table 1. Main properties for the 30 clusters composing the Cluster-EAGLE suite. In this table are specified the M_{500} , R_{500} and virialization state of each cluster. The $E_{\text{kin}}/E_{\text{thermal}}$ correspond the ration between the kinetic and thermal energy of the gas particles inside one R_{500} as detailed in [Barnes et al. \(2017b\)](#). Clusters with $E_{\text{kin}}/E_{\text{thermal}} < 0.1$ are defined as relaxed.

general analytic profile for the instantaneous ram pressure experienced by satellite galaxies that is a function of their host mass and redshift. The profile is described by a damped power law as:

$$P_{\text{ram}}(M, z) = P_0(z) \left[\frac{1}{\xi(z)} \left(\frac{r}{R_{200}} \right) \right]^{-\frac{3}{2}\alpha(M, z)} \quad (1)$$

This depends directly on r/r_{200} , the relative distance of the satellite galaxy from the halo center. $P_0(z)$ sets the normalization of the profile; $\xi(z)$ determines the radial scaling, and the exponent $\alpha(M, z)$ regulates the dependence on the host halo mass according to

$$\alpha(M, z) = \alpha_M(z) \log \left(M_{200} h^{-1} [\text{M}_{\odot}] \right) - 5.5. \quad (2)$$

Thereby, the shape of the profile is fully described by $P_0(z)$, $\xi(z)$, and $\alpha_M(z)$. Their dependence with the redshift is expressed in terms of the scale factor $a = 1/(1+z)$, as follows:

$$\log \left(\frac{P_0(z)}{10^{-12} h^2 \text{dyn cm}^{-2}} \right) = 7.01 a^{-0.122} - 9.1, \quad (3)$$

$$\xi(z) = -3.4 a^{-0.42} + 10.2, \quad (4)$$

$$\alpha_M(z) = 3.3 \times 10^{-3} a^{1.33} + 0.512. \quad (5)$$

The numerical values in these relations were obtained from a χ^2 minimization to fit ram pressure measurements from hydrodynamical resimulations of groups and clusters of galaxies.

Given the low scatter found in the relation by the authors, and given the great agreement between the values found with this method and direct measurement inside the clusters in our $z = 0$ sample, throughout this work, we will use this expression to estimate the ram pressure acting on our satellite galaxies at any time.

A galaxy will suffer ram pressure stripping once the ram pressure overcomes the restoring force exerted by the galaxy itself. To measure the restoring pressure, we follow the methodology presented in [Simpson et al. \(2018\)](#). The restoring force per area on the satellite's gas can be expressed by:

$$P_{\text{rest}} = \left| \frac{\partial \Phi}{\partial z_h} \right|_{\text{max}} \Sigma_{\text{gas}}, \quad (6)$$

where Φ corresponds to the gravitational potential, z_h correspond to the direction motion of the gas (and in opposite direction the gas displacement), $|\partial \Phi / \partial z_h|_{\text{max}}$ is the maximum of the derivative of Φ along z_h , and Σ_{gas} is the surface gas density of the satellite ([Roediger & Hensler 2005](#)). We adopt a simple estimate for both $|\partial \Phi / \partial z_h|_{\text{max}}$ and Σ_{gas} that can be applied to all the galaxies in our sample regardless of their morphology. Additionally, we will measure $|\partial \Phi / \partial z_h|_{\text{max}}$ as:

$$|\partial \Phi / \partial z_h|_{\text{max}} = v_{\text{max}}^2 / r_{\text{max}} \quad (7)$$

where v_{max} is the maximum circular velocity of the galaxy and r_{max} is the distance to the center of the galaxy where the maximum circular velocity is reached. For the gas surface density, we will measure it using:

$$\Sigma_{\text{gas}} = M_{\text{gas}} / \pi (2r_{1/2}^{\text{gas}})^2, \quad (8)$$

where M_{gas} is the total gas mass of the galaxy and $(r_{1/2}^{\text{gas}})$ is the 3D radius enclosing half of the gas mass.

As shown by [Vega-Martínez et al. \(2021\)](#), the profile presented has the strength of providing accurate and quick ram pressure measurements for galaxies in different environments, from the outskirts ($> 2r_{200}$) of low mass groups to the inner parts of massive galaxy clusters. Additionally, the profile works in a wide range of redshift, from $0 < z < 2$. This profile allows us to study the ram pressure efficiency of the different structures. Even though this profile does not take into account the asymmetric behaviour of the intracluster medium, the agreement between our simulations and the profile is remarkable and allows us to characterize the pressure exerted by the different haloes in which galaxies reside throughout its life (for details, see Appendix B).

3 PROPERTIES OF CLUSTERS: QUENCHED POPULATION AND GAS DENSITY PROFILES

As previously discussed, it is well established that the environment can play an important role in quenching the star formation activity of galaxies, especially within the inner and denser environments of clusters. However, how the mechanisms associated with environmental quenching lead

to the final quenching state is still an open question. In this work, we want to study the relation between the local density of the ICM and the moment of quenching for the satellite population. Several criteria have been used in the past to select quenched galaxies. Following previous studies, we use a threshold in specific star formation rate (sSFR), defined as the instantaneous star formation rate (SFR) divided by their total stellar mass (M_\star) (Weinmann et al. 2010; De Lucia et al. 2012; Wetzel et al. 2012, 2013) such that if a galaxy satisfies *i*) that its $\text{sSFR} \leq 10^{-11} \text{yr}^{-1}$ and *ii*) it never reaches a sSFR greater than that threshold value again, it is classified as quenched. It is worth mentioning that we tried a different threshold sensitive to redshift to split our sample between star-forming and quenched galaxies, following Davé et al. (2019). Using this criterion, a galaxy is classified as quenched if its $\text{sSFR} \leq 10^{-10.8+0.3z} \text{yr}^{-1}$. We found none or little differences in our results. In order to make our results directly comparable with previous works, we choose to keep the fixed sSFR threshold.

P19 studied the quenching history of the satellite population of galaxy clusters in the EAGLE simulation. They used the same quenching classification criteria that we will use throughout this paper. They found that cluster satellite galaxies reach their quenching state preferentially in dense environments such as massive galaxy groups or low-mass galaxy clusters. This holds even for galaxies that become quenched before their infall into the $z = 0$ host cluster and those that suffer strong drops in their SFR (but not get completely quenched) prior to their infall to any cluster. They also found that the quenched fraction in $z = 0$ cluster satellites grows significantly after the first r_{200} crossing. However, this trend depends strongly on the cluster mass, with low mass clusters showing the largest increase of the quenched fraction during satellite infall. Here we aim to determine the physical mechanisms behind satellite quenching during the r_{200} crossing and to understand why low mass clusters seem to play a larger role in this phenomenon.

3.1 Quenching of the star formation

Because of their high resolution and wide range of halo masses, spanning from groups to clusters, the C-EAGLE re-simulations are an excellent laboratory to understand the different environmental processes that affect the quenching of star formation and its dependence on cluster mass. Under the assumption that more massive clusters are associated with denser environments, we first search for correlations between the properties of galaxies at their quenching time and the mass of the halo where they reach their quenching state.

Following P19, and to improve the number statistics of our results, we stack galaxy in four bins of host cluster mass $M_{200}(z = 0)$. P19 split the 10 most massive halos in the EAGLE simulations in three bins of mass, ranging from $14 < \log_{10} M_{200}/M_\odot < 14.8$, and found significant differences between the properties of satellites residing in the lowest and highest cluster masses at quenching.

Taking advantage of the larger C-EAGLE sample, we will be able to see if the trends found in P19 prevail at higher masses. To this end we split our $z = 0$ cluster sample as follows:

- High mass: $14.9 < \log_{10} M_{200}/M_\odot < 15.4$

- Intermediate-high mass: $14.6 < \log_{10} M_{200}/M_\odot < 14.9$
- Intermediate-low mass: $14.3 < \log_{10} M_{200}/M_\odot < 14.6$
- Low mass: $14.0 < \log_{10} M_{200}/M_\odot < 14.3$.

The number of clusters in each mass bin from high to low mass is 8, 8, 6 and 7, respectively.

To characterize the relation between the environment in which galaxies reside and the quenching of their star formation, we will further split our galaxy sample as follows:

- Pre-quenched as satellites: quenched as satellites of another structure, i.e., outside the cluster.
- Quenched on cluster outskirts: satellites quenched outside clusters r_{200} but while part of the cluster FoF.
- Quenched in-situ: satellites quenched inside the r_{200} of the final cluster.

We note that the term pre-processing is commonly used in the literature to refer to all environmental effects that could affect galaxies prior to their accretion into clusters (eg. Boselli et al. 2005; Wetzel et al. 2012; De Lucia et al. 2012; Einasto et al. 2020). This is, regardless of whether the result is the cease of their star formation or not. In this work, we are focusing on galaxies that fully cease their star formation activity before being accreted onto clusters. For this particular “pre-processed” population we adopted the term “pre-quenched”. This definition was taken from Pallero et al. (2019), in which “pre-processing” accounted for those moments when galaxies experienced strong drops in their star formation rate, while “pre-quenching” accounted for those galaxies that cease their star formation prior to their accretion into the studied cluster.

Figure 1 shows the distribution of $z = 0$ satellite properties at their quenching time, t_q , for the four cluster mass bins. From top to bottom, we show the results for *high*, *intermediate-high*, *intermediate-low* and *low* cluster mass bins respectively. In each panel, blue, orange, and red histograms represent the population of satellites pre-quenched as satellites, quenched on cluster outskirts, and quenched in-situ, respectively. Each histogram is normalized by the total number of satellite galaxies per cluster mass bin. The dashed lines correspond to the median of each population.

In the left column, we show the distribution of satellite host mass at their t_q . We can see that, regardless of the mass of the cluster, quenching tends to happen in high-mass haloes. We find that the typical host mass at the time of quenching increases with $z = 0$ cluster mass; this change is more appreciable in the pre-quenched population. However, note that the bins in these histograms are 0.5 dex wide; i.e. the peak of the distribution is in-between $13 \leq \log_{10} M_{200}/M_\odot \leq 14.5$, and correspond to the typical mass range for low mass clusters. We will further explore this in detail in Section 3.2. From the pre-quenched satellite population, we see that the peak in the host mass distribution at t_q moves from galaxy groups at the low-mass bin to low-mass clusters at the high-mass bin. This is a consequence of the hierarchical nature in which the clusters in our simulations are built. This result also shows that galaxies that arrive already quenched to the present-day clusters reached their quenched state earlier when they entered another massive halo ($\log_{10} M_{200}/M_\odot \geq 12.0$).

The middle and right columns show the total stellar mass (M_\star) and the stellar mass fraction for the galaxies at

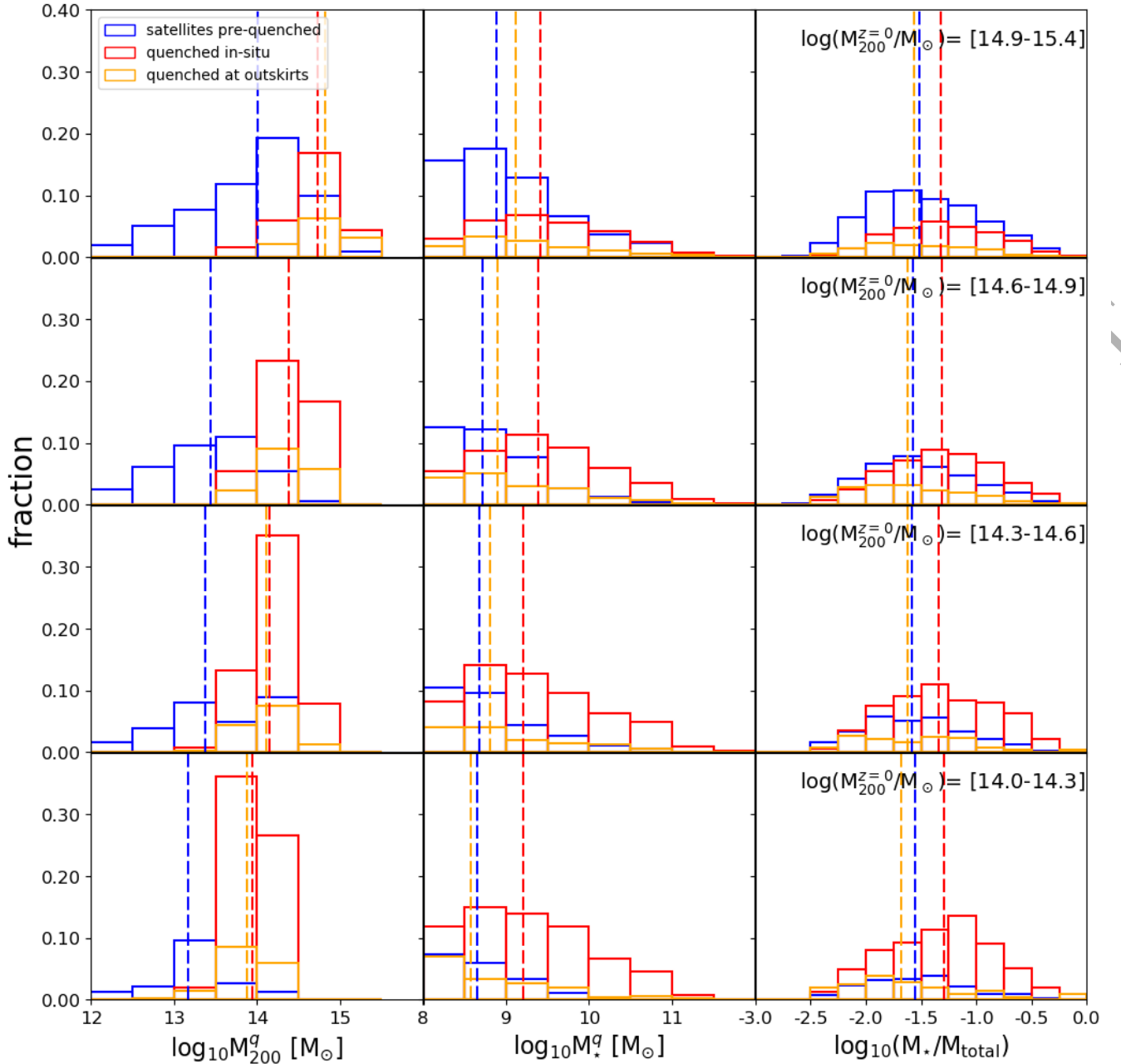


Figure 1. Mass distribution of galaxies and their hosts at the moment when they quench. Each row corresponds to the results obtained after stacking the distribution of galaxies associated with clusters within (from upper to lower row) high, intermediate-high, intermediate-low and low mass bin respectively. The left column shows the mass distribution of the host of each galaxy at the moment of quenching. The middle column shows the distribution of galaxies’ stellar mass at quenching time and the right hand column shows the distribution of stellar mass fractions. The galaxies are separated into ‘pre-quenched as satellites’ (blue), ‘quenched on cluster outskirts’ (orange), and galaxies ‘quenched in-situ’ (red). We can see that there is a trend between the $z = 0$ cluster mass and the halo mass at the quenching time for the satellite population. Nonetheless, independent of the M_{200} , more than 80% of cluster galaxies get quenched inside a halo with $M_{200} \geq 10^{13.5} M_{\odot}$.

their t_q . The stellar mass fraction is defined as the stellar mass divided by their total mass, $M_{\star}/M_{\text{galaxy}}$. Here the total mass includes the stellar, gas, black-holes and dark matter mass. Galaxies quenched in-situ show a higher stellar mass and stellar ratio compared to the pre-quenched and quenched at the outskirts populations. Moreover, galaxies that quenched on the outskirts of their $z = 0$ clusters fol-

low the same distribution as galaxies quenched as satellites of other clusters than their $z = 0$ host, in both their stellar mass and stellar mass fraction. This suggests that the mechanisms affecting these two populations are similar. On the right column, we can also see that pre-quenched galaxies tend to be more dark matter dominated than their “in-situ” quenched counterparts. We can also see that there is no re-

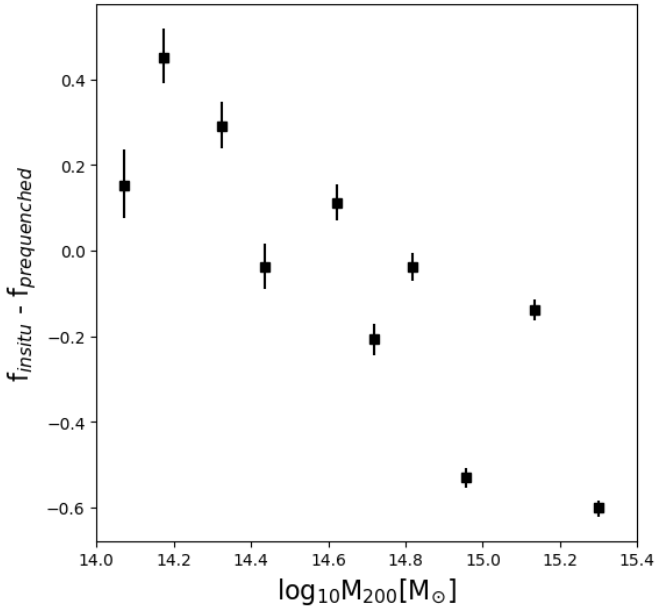


Figure 2. Variation of the fraction between the population pre-quenched as satellites and the population of galaxies quenched inside the cluster's r_{200} ($f_{\text{insitu}} - f_{\text{pre-quenched}}$) as a function of their halo cluster mass. Each dot represents a triplet of clusters ranked by mass. The halo mass plotted corresponds to the average of each group of three clusters, with the exception of the most massive bin, in which only two clusters were used. The error bar correspond to the binomial error associated with each measurement. In this plot we can see that the predominant population changes as the cluster mass grows. This result supports the scenario where low mass clusters are responsible for the high pre-quenching fraction found in high mass clusters.

lation between the median of the stellar to total mass ratio and the cluster mass in which galaxies reside at $z = 0$, for all populations.

As previously discussed, more than 60% of the galaxies quenched in our sample reach this state in groups or low mass clusters; i.e. with $13.5 \leq \log_{10} M_{200} [M_{\odot}] \leq 14.5$. This suggests that satellites get quenched after their first interaction with the ICM of structures within this mass range, independently of any subsequent mergers into more massive clusters. As massive clusters grow by the accretion of lower mass structures, most galaxies arrive in the final cluster already quenched in less massive progenitors. In the following sections, we will study and discuss the possible physical processes that can lead to the quenching of galaxies in these structures.

Figure 2 quantifies the relative abundances of the pre-quenched and in-situ quenched satellite populations. To increase the satellite number statistics, and to obtain more reliable trends, we bin our cluster sample in 10 sets based on their M_{200c} at $z = 0$. We divide the 30 available clusters into ten sets of three, except for the most massive set that contains only two clusters. The figure shows the difference between the fraction of in-situ quenched and pre-quenched satellites,

$$\frac{N_q^{\text{insitu}} - N_q^{\text{preq}}}{N_{\text{total}}} = f_{\text{insitu}} - f_{\text{preq}},$$

as a function of the mean M_{200c} of each bin. The error bar on each dot corresponds to the binomial error associated with each measurement. This figure clearly shows that for low mass clusters, galaxies quenched in-situ are the predominant population. On the other hand, for high mass clusters, the population of pre-quenched satellites becomes highly dominant. At $M_{200}^{z=0} \sim 10^{14.6} M_{\odot}$ the dominant population changes from insitu-quenched dominated to pre-quenched dominated. It should be noted that more than 60% of cluster galaxies that are quenched at present-day reached this state in structures with $13.5 \leq \log_{10} M_{200} [M_{\odot}] \leq 14.6$.

3.2 Clusters at different times

In this Section, we explore the evolution of the cluster satellite properties over a range of redshifts $0 < z < 1$, typical for current photometric surveys. As in the previous sections, we study the satellite population of the most massive structures in each zoom-in simulation, at each redshift. It should be noted that their corresponding satellite galaxies can be different from the $z = 0$ population, as many of them will be disrupted before the present day (but see Bahé et al. 2019), while others will be accreted after the selected redshift. We choose to compare the satellite population of clusters at four different times: $z = 0.90, 0.54, 0.30$ and $z = 0$. These correspond to lookback times of 7.5, 5.5, 3.5, and 0.0 Gyr, respectively. As before, for each satellite galaxy, we record the values of their total stellar mass and the host halo mass at their corresponding t_q .

The top panels of Figure 3 show the relation between the median host mass when satellites quench and the mass of the cluster they belong to at the corresponding z . The different columns show the results at the four redshifts considered. Each symbol represents the median of the distribution, \bar{M}_{200}^q ; squares and circles represent the population of pre-quenched and in-situ quenched satellites, respectively. As before, clusters have been combined into 10 mass bins. Thus, each symbol represents the population of satellites of 3 clusters (2 in the last bin). The colour coding indicates the fraction that each population represents with respect to the total, at the given redshift. For clarity, galaxies quenched as centrals are not shown in this figure. Nevertheless, they are considered when estimating the percentage that each population represents.

In the top left panel we show the relation between \bar{M}_{200}^q and the cluster's M_{200} at $z = 0$. The panel shows a mild growth of \bar{M}_{200}^q with $z = 0$ cluster mass. The dashed lines show the result of a linear regression fit, obtained combining all satellites inside each cluster-mass bin, regardless of whether they were quenched in-situ or pre-quenched. Not surprisingly, this relation is stronger for the in-situ quenched than for the pre-quenched satellites. From the colour coding, we can also see that, at $z = 0$, the dominant population changes with cluster mass (see also Fig. 2). For low $z = 0$ cluster masses, most satellites were quenched in-situ, while for their high mass counterparts satellites arrive into the cluster pre-quenched. However, the values of \bar{M}_{200}^q are all within the mass range of low mass clusters, i.e. $\log_{10} M_{200} [M_{\odot}] \sim 14.0$.

As we move towards a higher redshift, we can see that the slope of the fit increases, reaching a value ≈ 1 at $z = 0.9$. In section 4.1 we will study in detail the mechanisms that

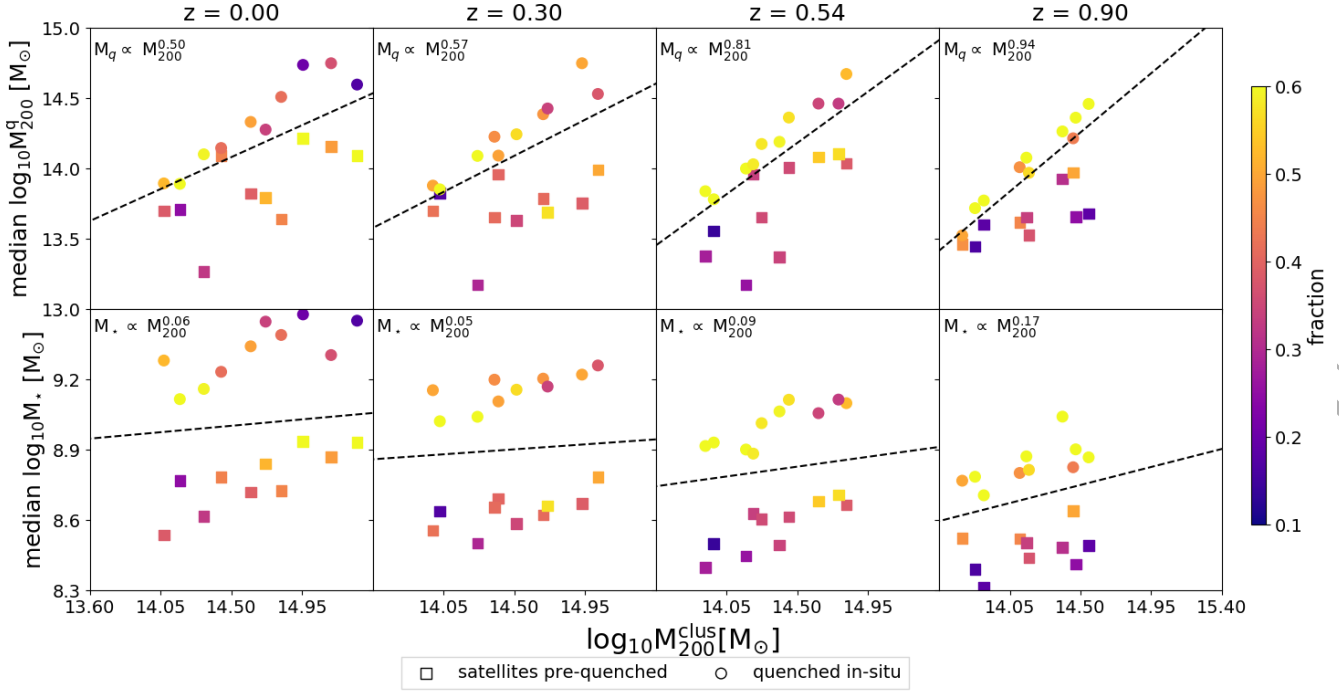


Figure 3. Median of the halo and stellar mass for the satellite galaxies at their quenching time as a function of their halo mass at $z = 0.0; 0.30; 0.54; 0.90$ (from left to right). Squares represent the median of the pre-quenched population while the circles represent the population quenched in-situ. Each dot represents a triplet of clusters ranked by mass, color-coded by the fraction that the pre-quenched or in-situ quenched population represents in a given bin of mass and redshift from the whole population for those clusters. Additionally, a fit to the whole population is plotted in dashed black lines. As we look at higher redshift clusters, in-situ quenching becomes dominant, and the median of the halo mass where galaxies are quenched is similar to the mass of the cluster itself. As we go down in redshift pre-quenching becomes dominant. This highlights the fact that galaxies reach their quenching state in the first transition from central to satellite galaxy. Regarding the stellar mass, as we look at galaxies at higher redshift, the stellar content is reduced systematically, while the slope of the fit shows a mild systematic growth.

produce this trend. Note that the most massive structures at this z have $M_{200} \sim 10^{14.6} M_{\odot}$. Interestingly, according to Fig. 2, this mass corresponds to the transition where the pre-quenched population at $z = 0$ becomes dominant over the in-situ quenched satellites. A comparison between the left and right top panels ($z = 0$ and $z = 0.9$, respectively) shows that, even at this transition mass, the in-situ quenched population is more dominant at high than low redshifts. This could be produced by the different assembly histories that clusters at $z = 0.90$ and $z = 0.00$ have experienced. High redshift clusters have mainly accreted star-forming galaxies residing in low mass structures, while low redshift clusters were able to grow by the accretion of more massive, pre-quenched substructures.

In the bottom panels, we show the median of the galaxies' stellar mass distribution, \bar{M}_* , at their quenching time. Here we find a clear correlation (but a very mild one) between \bar{M}_* and the cluster mass. Instead, we find a correlation between \bar{M}_* and redshift, independently of the cluster mass. Similar to what was found in the previous section, a difference of 0.5 dex in \bar{M}_* between the pre-quenched and the in-situ quenched population is found at all z .

In order to highlight the role played by the environment in the quenching of the star formation activity of satellites, in Figure 4 we show, for each cluster, the satellite quenched fraction as a function of satellite accretion time onto the corresponding cluster. The panels show the results obtained at

different z . Each line is associated with one of the 30 simulated clusters, and the colour coding indicates the cluster M_{200} at the given z . Following P19, to generate this figure we computed, for each galaxy within the cluster R_{200} (at the corresponding z), the time (t_{infall}) when it first crossed the cluster's R_{200} . We then define, for each galaxy, the variable $t - t_{\text{infall}}$, and identify the moment when it became quenched on this new time scale. Finally, we compute the cumulative quenched galaxy fraction as a function of $t - t_{\text{infall}}$.

The bottom right panel of Fig 4 shows the result obtained at $z = 0$. For all clusters, we find that the satellite quenched fraction rises rapidly after infall ($t - t_{\text{infall}} > 0$), regardless of the cluster mass. The jump is, however, significantly more abrupt for low mass clusters. High mass clusters show a larger fraction of quenched satellites for $t - t_{\text{infall}} < 0$ and rapidly reach a fraction of 90% after satellite infall. In general, we find that the fraction of pre-quenching varies between 20-80%, with a strong dependence on the cluster mass. Almost every satellite galaxy reaches its quenched state 4 Gyr after their accretion time (see also P19). As previously discussed, most $z = 0$ satellites get quenched while inhabiting low mass clusters and arrive at the final massive structure without any star formation activity. As a result, the ρ_{ICM} of these $z = 0$ massive clusters ($M_{200} \gtrsim 10^{14.6} M_{\odot}$) are not playing a significant role in the overall quenching of (massive) cluster galaxies in the local Universe. Instead, their main role is to end the star formation activity of that small

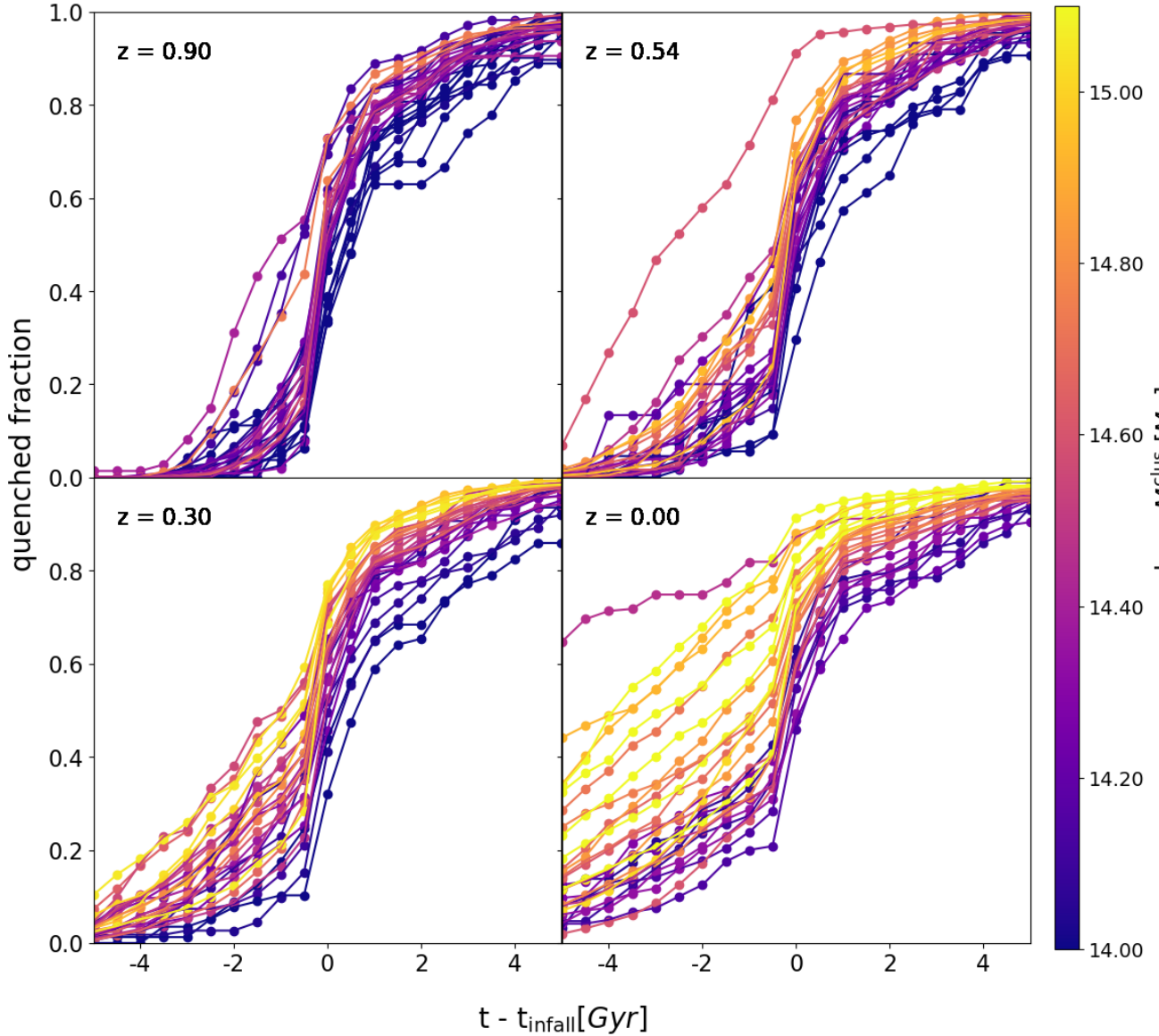


Figure 4. Growth of the quenched fraction of the satellite population of galaxy clusters as a function of the normalized timescale $t - t_{\text{infall}}$ for clusters at $z = 0.00; 0.30; 0.54; 0.90$. Each line represents one cluster in the simulation, color-coded by their M_{200} at the given redshift. Negative (positive) values correspond to times before (after) the first crossing of the R_{200} of the final cluster. As we look at clusters at higher redshift, we can see that the fraction of pre-quenching is reduced significantly ($t - t_{\text{infall}} < 0$). At the time of their infall ($t - t_{\text{infall}} = 0$) we can see a sharp increase in quenching and an abrupt change in the slope of the fraction of quenched galaxies. This result is even clearer when looking at higher redshifts, where even the most massive clusters present a rise in the quenched fraction $\gtrsim 60\%$ when $-1 \lesssim t - t_{\text{infall}} \lesssim 1$. As we decrease the redshift, the fraction of pre-quenching increases, especially for the most massive clusters, which at $z = 0.30$ presents $> 50\%$ of their galaxies already quenched before accretion. And by $z = 0$, clusters with $M_{200} \geq 10^{15} M_{\odot}$ possess a pre-quenching fraction $> 80\%$.

amount of galaxies that is not accreted in a subgroup into the cluster.

As we move towards a higher redshift, the fraction of satellites arriving into the clusters as star-forming galaxies increases. Note that at $z = 0.9$ (top left panel) the typical pre-quenched satellite fraction is $\lesssim 20\%$, and that just after t_{infall} these fractions quickly jump to values of $\sim 70\%$. Most cluster satellites are suffering strong environmental effects for the first time in structures that are within the mass range $10^{13.5} \lesssim M_{200} \lesssim 10^{14.5} M_{\odot}$, i.e. low mass clusters (see also Fig. 1)

It should be noted that there is a clear outlier at $z = 0$, that can also be well distinguished at $z = 0.54$. This line corresponds to the cluster CE-9 which underwent a major merger at $z \sim 0.6$. This structure collides with another cluster of similar mass which already had a large fraction of quenched galaxies. The result is the accretion onto the main cluster of a large number of pre-quenched satellite in a very short timescale. As a result, this cluster possesses $\sim 40\%$ larger pre-quenched fraction with respect to other clusters with similar mass.

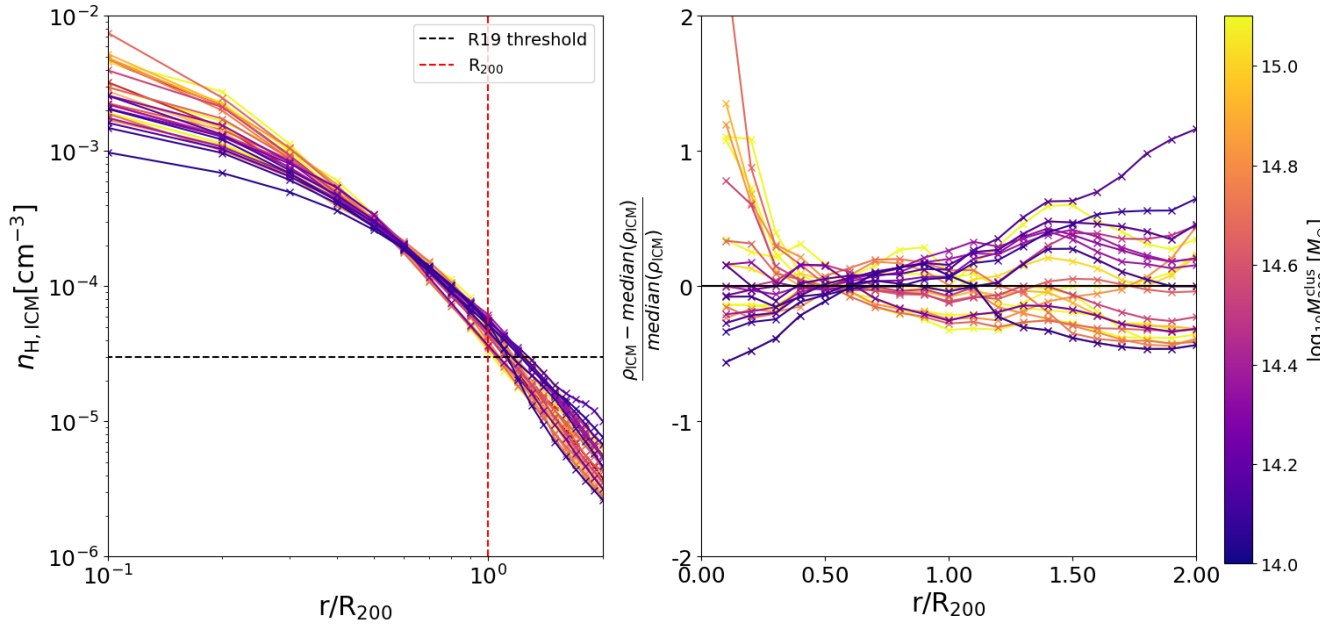


Figure 5. *Left panel:* Azimuthally averaged cluster gas density profiles for the $z = 0$ sample. *Right panel:* Fractional deviation from the median cluster gas density profiles. Each line is color-coded by cluster M_{200} . We can see that low mass clusters possess a more extended gas component in comparison with high mass clusters that are more concentrated. This result is more evident when looking at the normalized distribution in the right panel. Regardless of the cluster mass, the threshold in ρ_{ICM} found by Roberts et al. (2019) is reached near the r_{200} of our clusters.

3.3 Gas density of the intracluster medium

Our analysis points towards a preferential cluster mass range where environmental quenching is more efficient. Interestingly, using a sample of cluster galaxies selected from the Sloan Digital Sky Survey with high-quality Chandra X-ray data, R19 found a threshold in ρ_{ICM} for quenching, likely related to the effects of ram pressure stripping. Their study shows that the fraction of quenched galaxies in dense environments grows with ρ_{ICM} . For galaxies with a stellar-mass $\log_{10} M_{\star} [M_{\odot}] > 9.9$ (intermediate and high mass galaxies), the relation between the quenched fraction and ρ_{ICM} can be described with a single power law. However, for galaxies in the mass range of $9.0 < \log_{10} M_{\star} [M_{\odot}] < 9.9$ (low mass galaxies) a broken power law is needed. The ‘knee’ in this double power law is located at densities $\rho_{\text{ICM}} \approx 10^{-28.3} \text{ gr cm}^{-3}$ ($n_{\text{H}} = 3 \times 10^{-5} [\text{cm}^{-3}]$). According to the model presented in R19, when a galaxy reaches this density threshold it starts to experience a ‘rapid-quenching’ mode. It is worth noting that, as discussed by Simpson et al. (2018), the ICM of hosts with total masses in the range $12 < \log_{10} M_{200} [M_{\odot}] < 13$ are inefficient at quenching this type of satellites. Thus, they are bound to reach the final quenching state within the environments of low mass clusters and massive groups.

To explore this scenario with our models, in Figure 5 we show azimuthally averaged gas density profiles of the 28 $z = 0$ main clusters in the C-EAGLE simulation. Each line on the left panel corresponds to the gas density distribution of one cluster in our simulation, colour-coded by their M_{200} at $z = 0$. The cluster-centric distances have been normalized by the corresponding r_{200} . Note that low mass clusters typically have relatively more extended gas distributions than their more massive counterparts. These differences are bet-

ter highlighted in the right panel, where we show

$$\Delta \rho_{\text{ICM}} = \frac{\rho_{\text{ICM}} - \bar{\rho}_{\text{ICM}}}{\bar{\rho}_{\text{ICM}}}.$$

Here $\bar{\rho}_{\text{ICM}}$ represents, at each radius, the median of the distribution of ρ_{ICM} for all clusters. These panels clearly show that the distribution of gas in high mass clusters is more concentrated. These clusters also possess lower amounts of cold gas than lower mass objects.

The black dashed line on the left panel indicates the ρ_{ICM} threshold introduced by R19 and the red dashed line shows the density of each cluster at their r_{200} . In general, we find that low mass clusters reach this threshold a bit further out, $\sim 1.3 R_{200}$, than high mass clusters, at $\sim 1.0 R_{200}$. Note that, as shown in Figure 4 (see also P19), a drastic change in the slope of the cumulative satellite quenched fraction as a function of time takes place at $(t - t_{\text{infall}} = 0)$ Gyr; this is, during the first crossing of the clusters r_{200} .

Using the density profiles shown in Fig 5 we estimate, for each cluster, the local value of ρ_{ICM} at the time where in-situ quenched satellites ceased their star formation activity. These values are obtained by interpolating the density profiles shown in Fig 5 at the locations where satellites are identified as quenched. We find that the peak of the local density experienced by each galaxy at their quenching time is $n_{\text{H, ICM}} \sim 1.8 \times 10^{-4} \text{ gr cm}^{-3}$ briefly after $0.7 R_{200}$ which roughly corresponds to r_{500} (Ettori & Balestra 2009).

To extend these results to the pre-quenched population, we estimate the local densities where galaxies reach their quenching state, either in the final host or in an accreted substructure. This time we use the ram pressure profile shown in Section 2.3. As discussed, ram pressure can be written as

$$P_{\text{ram}} = \rho_{\text{ICM}} v^2,$$

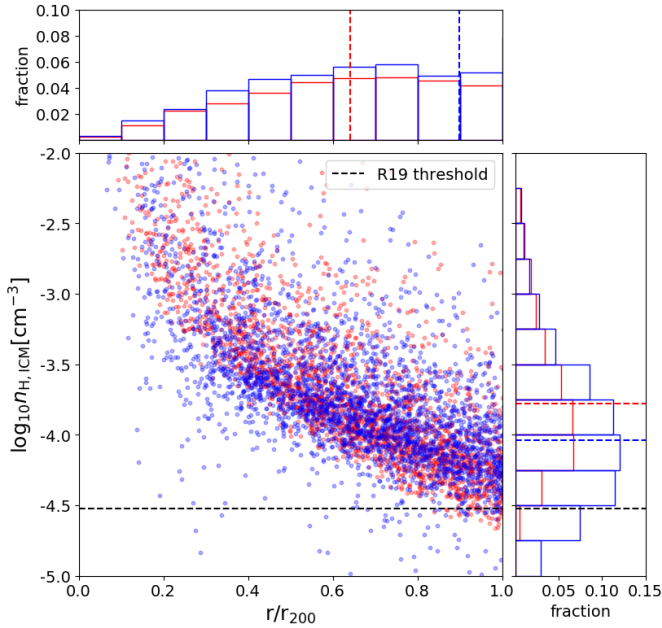


Figure 6. Density of the local environment in which galaxies reside at their quenching time as a function of the distance from the halo center in units of r_{200} . Blue dots correspond to galaxies that are quenched in a different halo than the final cluster (pre-quenched) and red dots correspond to galaxies quenched in the final cluster (quenched in-situ). The histograms show the distribution of density and distance separately. The black dashed lines show the threshold found in R19 and the blue and red dashed lines in the histograms correspond to the median of each distribution. We can see that the majority of the population in our sample, regardless of their status as pre-quenched or in-situ quenched, suffer quenching in high density environments ($\log_{10} n_{H, ICM} \geq 3 \times 10^{-5} [\text{cm}^{-3}]$) with the majority of the population reaching this state in the outskirts of clusters ($r > 0.5 r_{200}$).

where ρ_{ICM} corresponds to the density of the local environment of the galaxy, and v^2 corresponds to the square of the relative velocity between the galaxy and its environment. Based on this, we estimate the local environment density using the analytic ram pressure profile and the velocity of the galaxies directly measured from the simulations. The velocity is measured as the relative velocity between the galaxy and the centre of the host in which each galaxy resides at their t_q . In Fig 6 we show the values of the local density at which galaxies reach their quenching state as a function of the normalized distance to the host-centre at their t_q . The red dots correspond to the in-situ quenched population, while blue dots stand for the pre-quenched population. The histograms show both populations separately, and the red and blue dashed lines correspond to the median of each distribution. The dashed black line corresponds to the density threshold proposed by R19. We can see that in general, regardless of their in-situ or pre-quenched condition, galaxies reach their quenching state in a local environment denser than the threshold proposed by R19. We can also see that the majority of galaxies reach their quenched state at the outskirts of their host ($\sim r_{200}$), and just a small percentage of the whole sample the galaxies reach their quenched state in the inner parts of their hosts ($r \lesssim 0.5 r_{200}$). It is worth to emphasise that we define the accretion time of a galaxy as

the first time when a galaxy crosses the virial radii of its host cluster. This is done by individually tracing the histories of each galaxy throughout the simulation. On this way, we are sure that we are looking at galaxies that are undergoing their first pericentral passage. We note that, given the time resolution of the simulation, we are not able to recover the exact moment and place where the galaxy gets quenched. Furthermore, for this particular result, we are not taking into account the time spent within the host in which they eventually get quenched. Nonetheless, as we see in Figure 4, galaxies get quenched shortly after their first accretion event (but see Oman et al. 2020, for observational evidence to the contrary).

4 DISCUSSION

4.1 Ram pressure as the main culprit

The R19 ‘slow-then-rapid’ quenching scenario proposes that galaxies first experience a slow quenching phase as they approach the cluster due to what is known as starvation (Larson et al. 1980). During this phase galaxies are not able to replenish their gas reservoir and, as a result, slowly deplete their gas available to form stars. In the absence of other mechanisms, starvation would slowly lead to the final quenching of galaxies, within a typical timescale $\gtrsim 3$ Gyr (Peng et al. 2010, 2015). However, and in particular for low mass galaxies, once they reach the inner cluster region, where the intracluster-medium (ICM) reaches $\rho_{ICM} \sim 10^{-28.3} \text{ gr cm}^{-3}$ ($n_{H, ICM} \sim 3 \times 10^{-5} [\text{cm}^{-3}]$), the process switches to a rapid quenching phase due to the effect of ram pressure. The ram pressure experienced by satellites in these inner regions becomes large enough to overcome their restoring force, triggering the rapid depletion of their gas component. The typical quenching timescale in this phase is of the order of $\lesssim 1$ Gyr. As we have just shown in the previous section, this density is reached in these simulations at the outskirts of low mass clusters ($\sim R_{200}$).

In Figure 7, we further explore whether the rapid quenching phase we observe in our simulations is associated with the R19 ICM density threshold, $n_{H, ICM} \sim 3 \times 10^{-5} [\text{cm}^{-3}]$. In the left panel of this figure, we show the time evolution of the balance between the ram pressure and the restoring force per area felt by the satellites’ gas component. As an example, we focus on a subset of randomly selected galaxies. As before, the timescale is measured as $(t - t_{\text{infall}})$. This panel shows that, after crossing the virial radius of the cluster, the influence of ram pressure rapidly grows until it completely dominates over the restoring force. The middle and right panels show the time evolution of the total gas and HI content, respectively. We can see that the gas content of these galaxies starts to slowly decrease prior to accretion (slow phase) and then rapidly drops after infall (rapid phase). To generalize these results, we show in the left panel of Figure 8 the cumulative fraction of galaxies dominated by ram pressure ($P_{\text{ram}} > P_{\text{rest}}$) for each cluster, as a function of their infall time. Note that ram pressure is estimated for each satellite based on the density of their environment at the corresponding time, using the methodology shown in Section 2.3. In other words, if galaxies are arriving into the cluster as satellites of another substructure, the ICM considered is

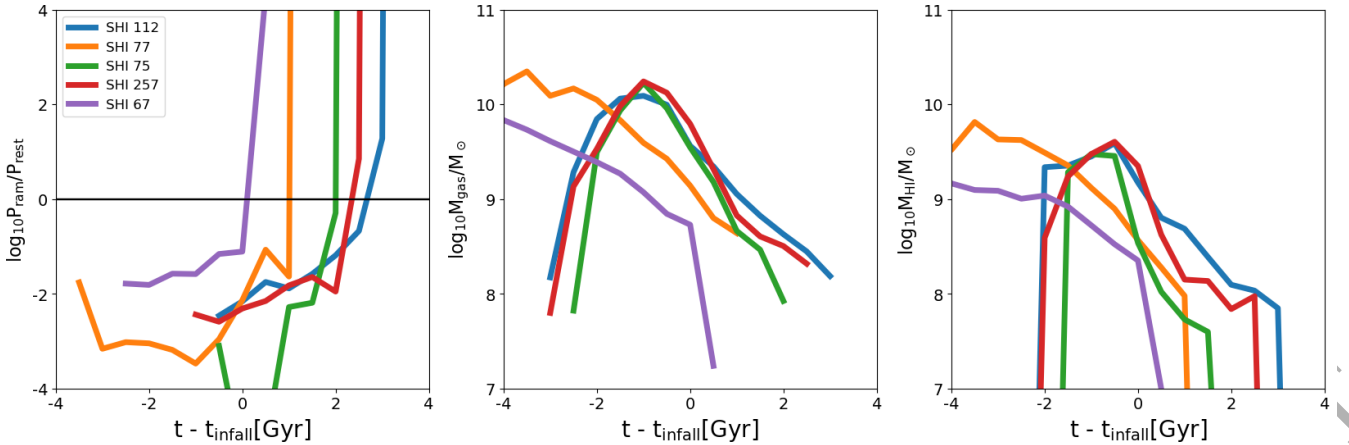


Figure 7. Evolution of properties for a subset of galaxies in our sample. SHI correspond to the SubHaloIndex associated with each galaxy in the zoom-in simulation CE-0. *Left:* Ratio between ram pressure and the restoring force per area, *Middle:* evolution of the total gas mass and *Right:* HI content in the galaxy, as a function of their time since accretion. The restoring force per area dominates over the ram pressure until they arrive to the cluster. At the time that galaxies become ram-pressure dominated, they lose most of their total gas content.

the one associated with the corresponding substructure. For comparison, in the right panel of this figure, we also show the evolution of the quenched galaxy fraction. Each line is color-coded by the cluster M_{200} at $z=0$. As we saw in the previous section, present-day satellites with masses $M_{\star} \gtrsim 10^9 M_{\odot}$ (see Fig. 1), in low mass clusters, tend to arrive as star-forming objects but the quenched fraction grows rapidly after infall (right panel). This trend becomes progressively less pronounced as the mass of the cluster increases. The left panel of this figure shows a very similar trend when considering the fraction of satellites that are ram pressure dominated (with respect to their restoring force). We can clearly see that, for low mass clusters, the vast majority of satellites reach the cluster's r_{200} for the first time, dominated by their restoring force, but this abruptly changes after infall. Instead, in more massive clusters, a significant fraction of satellites are arriving already dominated by ram pressure (and already quenched). As previously discussed, and shown in P19, this is due to the accretion of larger substructures into more massive clusters, bringing a significant fraction of the present-day quenched satellite population. Nonetheless, and regardless of the mass of the cluster, we see immediately after the t_{infall} a sharp rise in the fraction of ram pressure dominated fraction ($\sim 40\%$ in the most extreme cases). As the cluster mass grows, the rise becomes less pronounced.

4.2 Over-quenching in C-EAGLE

It was reported in Bahé et al. (2017) that the fraction of low mass ($\log_{10} M_{\star} [M_{\odot}] \leq 9.5$) quenched galaxies in the C-EAGLE simulations is higher than observed in data. In fact, while observations show that the fraction of quenched galaxies decreases steadily as we look at lower stellar masses, in C-EAGLE this fraction rises towards low masses (See figure 6 of Bahé et al. 2017). In a similar way, the EAGLE simulations also possess higher fraction of quenched galaxies towards lower masses ($\log_{10} M_{\star} [M_{\odot}] \lesssim 9.5$), mainly due to numerical resolution artifacts (Schaye et al. 2015). Since C-EAGLE uses the same numerical code and resolution as

EAGLE, Bahé et al. (2017) argue that the excess of passive galaxies found in these simulations is also mainly due to resolution effects, as over-quenching is measured only in low mass satellite galaxies. They propose that this resolution issue can be connected to the formation of large holes ($\sim 20\text{kpc}$) in the atomic hydrogen discs of many dwarf galaxies, similar to what was found in the EAGLE simulation (see Bahé et al. 2016). The limited resolution and high energetic feedback model implemented in EAGLE develop these large holes in the discs of dwarfs, making them more susceptible to ram pressure stripping events.

As we showed in the previous section, in general, C-EAGLE clusters possess an extended gas envelope that makes them relatively denser than the observational clusters shown in R19, reaching the density needed for ‘rapid-quenching’ even at the outskirts of low mass clusters. Given all these considerations it is straightforward to conclude that the excess of dwarf galaxies quenched inside clusters is mainly due to a combination of the extreme environment in which galaxies reside and the feedback phenomenology implemented in the code (for a more detailed explication regarding the feedback model we refer the reader to Schaye et al. 2015; Crain et al. 2015; Bahé et al. 2017). To address this issue, simulations containing galaxy clusters with higher resolution are needed. However no suitable simulation suites are available at present; TNG-50 (Nelson et al. 2019; Pillepich et al. 2019) and ROMULUSC (Tremmel et al. 2019) only partly fulfill this requirement. These simulations lack a statistical sample of clusters to carry out our study, with just one halo with $\log_{10} M_{\star} [M_{\odot}] \geq 14$. On the other hand, simulations as the ones presented in THE THREE HUNDRED PROJECT (Cui et al. 2018), or in the TNG-300 (Marinacci et al. 2018; Naiman et al. 2018; Nelson et al. 2018; Pillepich et al. 2018; Springel et al. 2018) can be useful to statically study the properties of galaxies inside clusters. However, due to the low mass and spatial resolution, it is impossible to study in detail the mechanisms that lead to the quenching of star formation in these simulations.

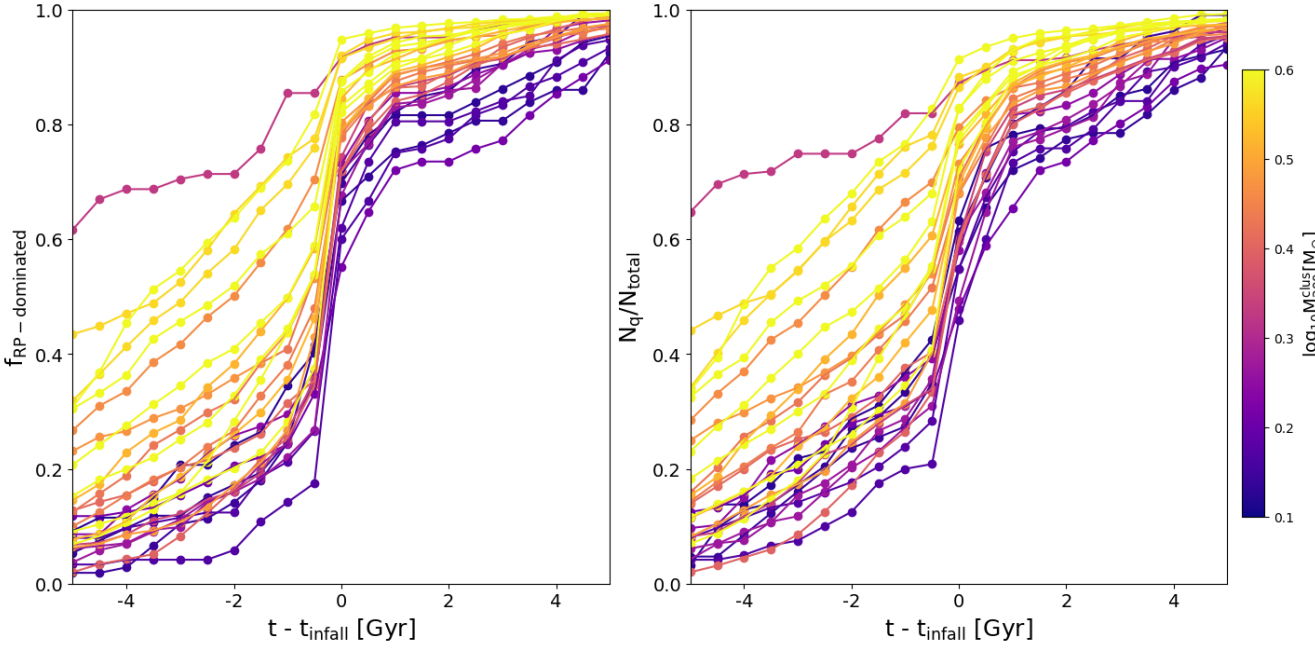


Figure 8. *Left:* Evolution of the ram-pressure dominated population of galaxies as a function of time since infall. *Right:* Growth of the quenched fraction of the satellite population in clusters as a function of the time since infall. Each line represents one cluster in the simulation, color-coded by their M_{200} at $z=0$. These results are coincident with the evolution of the ram pressure dominated population; as galaxies become more ram pressure dominated the quenched fraction also increases. This suggests that most galaxies quenched in our sample were quenched due to ram pressure inside clusters of galaxies.

4.3 Comparison with other studies

Previous studies based on numerical simulations have also studied the physical process intervening in the quenching of cluster satellite galaxies. Jung et al. (2018) used the Yonsei Zoom-in Cluster Simulations (YZiCS Choi & Yi 2017), a suite of 16 galaxy clusters within a mass range of $5 \times 10^{13} < M_{200}/M_{\odot} < 10^{15}$ ran with the adaptive mesh refinement (AMR) code RAMSES (Teyssier 2002), to investigate the environmental mechanisms behind the origin of gas-poor galaxies inside galaxy clusters. They define their sample of gas-poor galaxies on the $M_{\text{gas}}/M_{\star}-M_{\star}$ plane as the population of galaxies that deviates from the gas-rich galaxy sequence towards the low gas mass side by 2 sigmas. This sample is composed of galaxies that generally possess less than 10% of the gas mass found on their gas-rich counterparts, at the galaxy same mass. Similar to our results, the authors find that most galaxies get their gas component stripped within haloes with typical masses of massive groups or low mass galaxy clusters. In particular, they split their sample between three populations; 1) galaxies that experienced their gas stripping before their accretion into the studied cluster (pre-processed population), 2) galaxies that experience a fast stripping inside the studied cluster (fast cluster processing) and 3), those galaxies that experience a slow stripping inside the final cluster (slow cluster processing). They found that 33.7% of the galaxies in their sample suffered pre-processing. This population is dominated by satellite galaxies that underwent their processing inside haloes with typical masses of massive groups. An additional 42.7% experienced “fast cluster processing”, losing most of their gas mass before the first pericentric passage. These galaxies usually (but not exclusively) have low gas mass at their accretion time in

comparison to field galaxies, and where satellites of other haloes before accretion. Finally, 23.7% of the galaxies retain their gaseous components for several Gyrs after their accretion (after their first pericentric passage). This population is mainly composed of galaxies that were central and gas-rich before being accreted into the final cluster. The authors show that galaxies can lose their gaseous component slowly if they are accreted into low mass clusters ($\leq 10^{14} M_{\odot}$) on circular orbits, i.e. that do not approach the cluster centre close enough. This slow cluster processing is particularly relevant for massive groups ($M_{200} < 10^{14} M_{\odot}$) and rapidly lose relevance with cluster mass. On the other hand, both pre-processing and fast cluster processing rise in importance with cluster mass, in a similar fashion to what was shown in our work.

By combining the star formation histories of observed galaxies with the evolutionary tracks of simulated galaxies, some studies have tried to constrain the accretion history of individual cluster galaxies to better understand their quenching process. Based on satellite orbital model extracted from the HYDRANGEA and the YZiCS simulations Oman et al. (2020) and Rhee et al. (2020), respectively, were able to characterize the different stages of the quenching process proposed by Wetzel et al. (2013); i.e. the “delayed-then-rapid” quenching model. The quenching timescale obtained by both authors, for the observed sample of galaxies, is larger than what is found in our work and in Jung et al. (2018), even though the same simulation were considered. In particular, Rhee et al. (2020) found that the total timescale for quenching to take place is around ~ 5 Gyr. In the other hand, Oman et al. (2020) finds that the quenching timescales for satellites in clusters could be as

large as > 5.5 Gyr after their first pericentral passage. This is much larger than the timescales observationally measured by other authors (eg. [Wetzel et al. 2013](#); [Rhee et al. 2020](#)).

Within this context, the quenching timescale estimated from observations is generally larger than what estimated from hydrodynamical simulations ([Jung et al. 2018](#); [Pallero et al. 2019](#); [Lotz et al. 2019](#)). As we show in Figure 4, galaxies in our simulation reach their quenching state rapidly after their accretion into the cluster. Regardless of the cluster mass, $\geq 70\%$ of galaxies in our sample reach their quenching state within ≤ 1 Gyr after their accretion into the cluster. Similarly, in [Lotz et al. \(2019\)](#) using the hydrodynamical simulations *Magneticum Pathfinder*¹ the authors found that within 1 Gyr of their accretion into clusters, the large majority of galaxies reach their quenching time. This usually can be related to the time when galaxies reach their first pericentral passage, and they argue that this is probably related to a ram pressure stripping event, as was probed within our work. Additionally, when taking into consideration more massive galaxies ($M_\star \geq 1.5 \times 10^{10} M_\odot$) the quenching timescale becomes larger ($t_q \sim 2 - 3$ Gyr). This latter result is similar to the quenching timescale that we found when using a larger threshold in stellar mass, as shown in Appendix A.

Is important to note however that, even though both [Rhee et al. \(2020\)](#) and [Oman et al. \(2020\)](#) consider pre-processing as an important aspect in galaxy evolution, they cannot directly account for the pre-processing effects as we do in our work. Given the nature of combining data from observations and orbits from simulations, in their work it is not possible to measure if the observed satellites inhabited other massive halo prior to their accretion into their final cluster. Furthermore, both [Rhee et al. \(2020\)](#) and [Oman et al. \(2020\)](#) measure the quenching timescale as the moment when galaxies start to decrease their star formation rate, up to the moment when their star formation rate is completely depleted. Instead, in our work we consider the timescale between the moment when galaxies crosses the virial radii of a cluster for the first time, and their last star formation rate event. It should be noted that, as stated in [Pallero et al. \(2019\)](#), galaxies suffer strong drops in their star formation several Gyr before their accretion into the present-day clusters. However, despite this strong drops experienced, they remain forming stars at slow pace up to their accretion into a massive halo ($\sim 10^{14} M_\odot$).

These discrepancies in the quenching timescales highlights the importance to include the pre-processing in models to constrain the accretion history of galaxies. Accounting for episodes in which galaxies underwent the extreme environmental conditions provided by galaxy clusters, prior to their accretion into the present-day host, can improve the agreement between theoretical and observational works. Finally, most simulations nowadays agree that the pre-processed, or pre-quenched galaxies as referred in this work, represent the largest satellite population of massive clusters at the present day. As examined in the IllustrisTNG simulations, the pre-processing is a key aspect needed to explain the large number of passive galaxies in clusters nowadays [Gu et al. \(2020\)](#); [Donnari et al. \(2020\)](#).

5 SUMMARY AND CONCLUSIONS

In this paper, we have studied the evolution of the quenched fraction of satellites in the galaxy clusters of the C-EAGLE simulation. C-EAGLE is a suite of 30 cosmological hydrodynamic zoom-in simulations of massive galaxy clusters in the mass range of $14.0 \leq \log_{10} M_{200}/M_\odot \leq 15.4$ at $z = 0$, using the state-of-the-art EAGLE galaxy formation code. We select satellite galaxies with stellar mass $\log_{10} M_\star/M_\odot \geq 8$ that at $z = 0$ are bound to a main cluster and are located inside the cluster's r_{200} . First, using the merger trees obtained with the SPIDERWEB post-processing software (see details in [Bahé et al. 2019](#)), we study the time evolution of the quenched population ($\text{sSFR} < 10^{-11} \text{yr}^{-1}$) as they traverse different environments. To study the quenching of star formation of galaxies inside and outside the main cluster, we define three types of populations:

- pre-quenched as satellites: quenched as satellites in a structure outside the final $z = 0$ cluster.
- quenched on cluster outskirts: quenched outside the cluster R_{200} while in cluster FoF
- quenched in-situ: quenched inside the main clusters R_{200} .

We found that regardless of the final cluster mass, most galaxies ($> 80\%$) reach their final quenching state inside dense structures ($\log_{10} M_{200}/M_\odot \geq 13.5$); the remaining galaxies ($\lesssim 20\%$) are quenched in a medium sized halo ($12.0 < \log_{10} M_{200}/M_\odot < 13.5$) as satellites. We study the relation between the median of the distribution of galaxy properties at their quenching time and the cluster final mass. The first thing to notice is that, regardless of the cluster mass, quenching occurs mainly in halos of $\log_{10} M_{200}/M_\odot \sim 14$; in fact, as the cluster mass increases the predominant population in our sample changes from galaxies quenched in-situ, to galaxies pre-quenched. This highlights the fact that the quenching of the star formation happens in low mass clusters rather than in the most massive structures, even though they have the environmental conditions to strip the gas reservoirs of these satellites. Most of the galaxies that these massive clusters accrete are already quenched due to their earlier accretion onto another low mass cluster.

Regarding their stellar content, galaxies quenched inside the cluster contain 0.5 dex higher stellar masses than the pre-quenched population and show 0.3 dex higher stellar mass fraction.

We also studied the gas density profiles of clusters and the relation between the local density and the quenching suffered by satellite galaxies. Our clusters possess similar gas profiles. There are no cool-core clusters in the sample and all are widely hot-gas dominated. In general, high mass clusters show a more concentrated hot gas component, while low mass clusters exhibit a more extended envelope. In the case of low mass clusters, they reach the critical threshold in ρ_{ICM} , discussed by [Roberts et al. \(2019\)](#) at $\sim 1.3 r_{200}$. This suggests that galaxies in our sample experience the 'rapid-quenching' scenario starting at the outskirts of these clusters. When comparing to the time evolution of the quenched fraction we found that, at the moment of the first r_{200} crossing, the quenched fraction increases rapidly, with low mass clusters being the ones showing the most extreme change in this fraction; for these, the rise in their quenched fraction

¹ www.magneticum.org

($\sim 60\%$) is produced in a very short timescale ($\lesssim 1\text{Gyr}$). After this period, the fraction grows more slowly ($\sim 15\%$ within $1 \lesssim t - t_{\text{infall}} \lesssim 4\text{Gyr}$). High mass clusters contain higher fractions of quenched galaxies regardless of their accretion time. In this case, the effect of crossing the cluster's virial radius is not as significant as in the case of low mass clusters. This is due to the earlier passage of the galaxies through the high-density structure of previous substructures. Nevertheless, we still found a measurable rise in the fraction of quenched galaxies after infall into their final massive clusters. 1Gyr after infall, the fraction of quenched galaxies rise from 10% to 40%, and between [1-4]Gyrs after their infall, the fraction of quenched galaxies grow at a much slower rate than in low mass clusters ($\lesssim 10\%$). This result is slightly different at high redshift, where most of our simulated clusters are significantly less massive. At higher z , we find that most galaxies reach their quenching state in-situ, regardless of the cluster mass. This is a consequence of the different assembly history of clusters at high and low redshift. In other words, our results show that galaxies in the mass range of $9.0 < \log_{10} M_{\star} [\text{M}_{\odot}] < 9.9$ reach their quenching state inside the first dense structure they fall into. Regardless of z and cluster mass, 4Gyr after their infall, almost all galaxies ($\gtrsim 90\%$) are quenched.

Even though the excess of quenching found in our results, especially at the outskirts of clusters for low mass galaxies, can be related to the limited resolution of the simulation, we show that the ram pressure experienced by intermediate and high mass galaxies at the moment when they reach this threshold in gas density is high enough to strip their gas content, shortly after their first infall.

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DATA AVAILABILITY

The data presented in the figures are available upon request from the corresponding author. The raw simulation data can be requested from the C-EAGLE team (Bahé et al. 2017; Barnes et al. 2017b).

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APPENDIX A: A HIGHER STELLAR MASS THRESHOLD FOR GALAXY SELECTION

Here we discuss the variations that our main results have when selecting only galaxies with a stellar-mass greater $M_{\star} > 10^{9.5} M_{\odot}$.

Figure A1 shows the variation in the predominant population between pre- and insitu-quenched as a function of the cluster mass. The error bar corresponds to the binomial error associated with each value. As shown in this figure, when selecting galaxies using a higher mass threshold the same results discussed throughout the article are found. The importance of pre-quenching grows towards higher cluster mass, becoming dominant for massive clusters. Nevertheless, it should be noted that the mass at which the predominant population changes shifts towards higher cluster masses, as expected.

Figure A2 shows the relation between the median host mass when satellites quench and the mass of the cluster they belong to at the corresponding z . The different columns show the results at the four considered redshifts. Each symbol represents the median of the distribution, \bar{M}_{200}^q ; squares and circles represent the population of pre-quenched and in-situ quenched satellites, respectively. As before, clusters have been combined into 10 mass bins. Thus, each symbol represents the population of satellites of 3 clusters (2 in the last bin). The colour coding indicates the fraction that each population represents with respect to the total, at the given redshift. The dashed lines show the result of a linear regression fit, obtained combining all satellites inside each cluster-mass bin, regardless of whether they were quenched in-situ or pre-quenched. In this case, in the top panel, as we move towards a higher redshift, the slope of the fit reach a value ≈ 1 at $z = 0.54$, earlier than what was found when selecting galaxies with stellar mass $M_{\star} > 10^8 M_{\odot}$. In the bottom panels, we show the median of the galaxies’ stellar mass distribution, \bar{M}_{\star} , at their quenching time. In this case, the results do not show any correlation with cluster mass.

Figure A3, shows the growth of the quenched fraction as a function of galaxy’s accretion time, for all clusters in

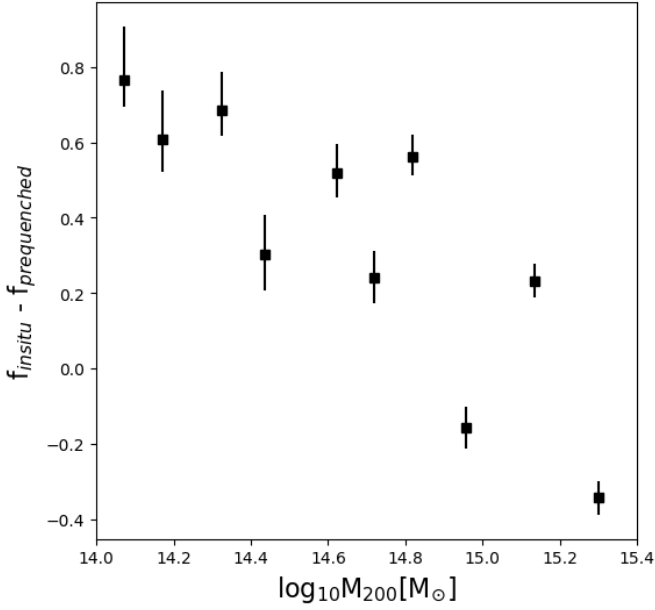


Figure A1. Variation of the fraction between the population pre-quenched as satellites and the population of galaxies quenched inside the cluster's r_{200} ($f_{\text{insitu}} - f_{\text{pre-quenched}}$) as a function of their halo cluster mass, when selecting satellites with stellar mass greater than $M_{\star} \geq 10^{9.5} M_{\odot}$. Each dot represents a triplet of clusters ranked by mass. The halo mass plotted corresponds to the average of each group of three clusters, with the exception of the most massive bin, in which only two clusters were used. The error bar correspond to the binomial error associated with each measurement. In this plot we can see that the predominant population changes as the cluster mass grows. This result supports the scenario where low mass clusters are responsible for the high pre-quenching fraction found in high mass clusters.

our sample at redshift $z = 0.00; 0.30; 0.54; 0.90$. Each line corresponds to one of our clusters and are colour-coded by their M_{200} at given z . We can see a strong variation in the quenching fraction when galaxies cross the virial radii of a cluster for the first time, and that the fraction of pre-quenched galaxies grow with cluster mass, as we show throughout the manuscript using the lower mass threshold in stellar mass ($M_{\star} > 10^8 M_{\odot}$).

APPENDIX B: AGREEMENT BETWEEN THE ANALYTIC PROFILE AND THE SIMULATIONS

As mentioned in section 2.3, instead of directly using the simulated gas distributions, we used an analytic profile to measure the ram pressure experience by the satellites in our sample. The implementation of this analytic method has the important advantage of significantly reducing computational costs. As we show in Figure B1, and further discussed in Vega-Martínez et al. (2021), both methods to compute ram pressure profiles show, statistically, very good agreement. The figure compares the result of measuring ram pressure profiles using an spherically averaged gas density profile extracted from the 30 C-Eagle simulations, and that obtained from the analytic profile. Results are shown up to R_{200} , distance that we used in this work to define membership. As we

can see from this Figure, the ram pressure measured with the spherically averaged density profiles and the analytic formulae are in good agreement, specially for the regions where the ram pressure becomes relevant ($\log_{10} P_{\text{ram}} < 10 [\text{dyn cm}^{-2}]$).

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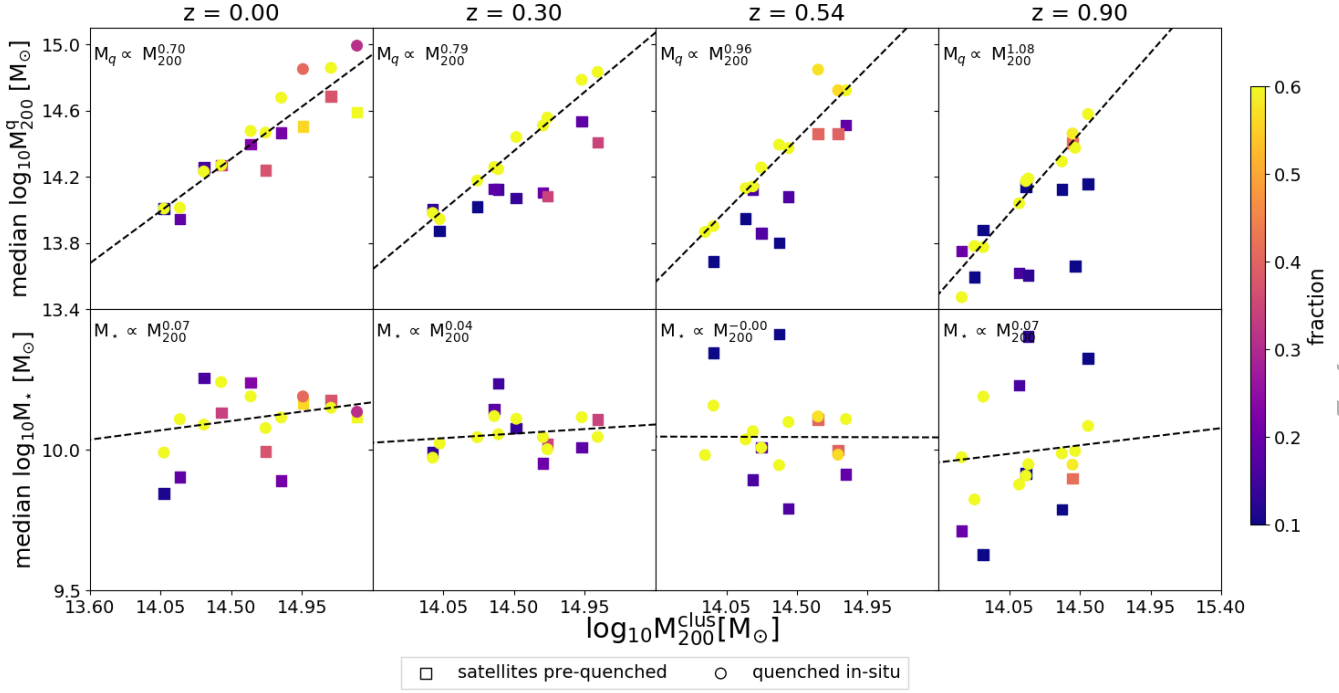


Figure A2. Median of the halo and stellar mass for the satellite galaxies at their quenching time as a function of their halo mass at $z = 0.0; 0.30; 0.54; 0.90$ (from left to right). Satellites were selected using a threshold in stellar mass $M_\star \geq 10^{9.5} M_\odot$. Squares represent the median of the pre-quenched population while the circles represent the population quenched in-situ. Each dot represents a triplet of clusters ranked by mass, colour-coded by the fraction that the pre-quenched or in-situ quenched population represents in a given bin of mass and redshift from the whole population for those clusters. Additionally, a fit to the whole population is plotted in dashed black lines. As we look at higher redshift clusters, in-situ quenching becomes dominant, and the median of the halo mass where galaxies are quenched is similar to the mass of the cluster itself. No relation can be found between the stellar mass at quenching time and the cluster mass can be found.

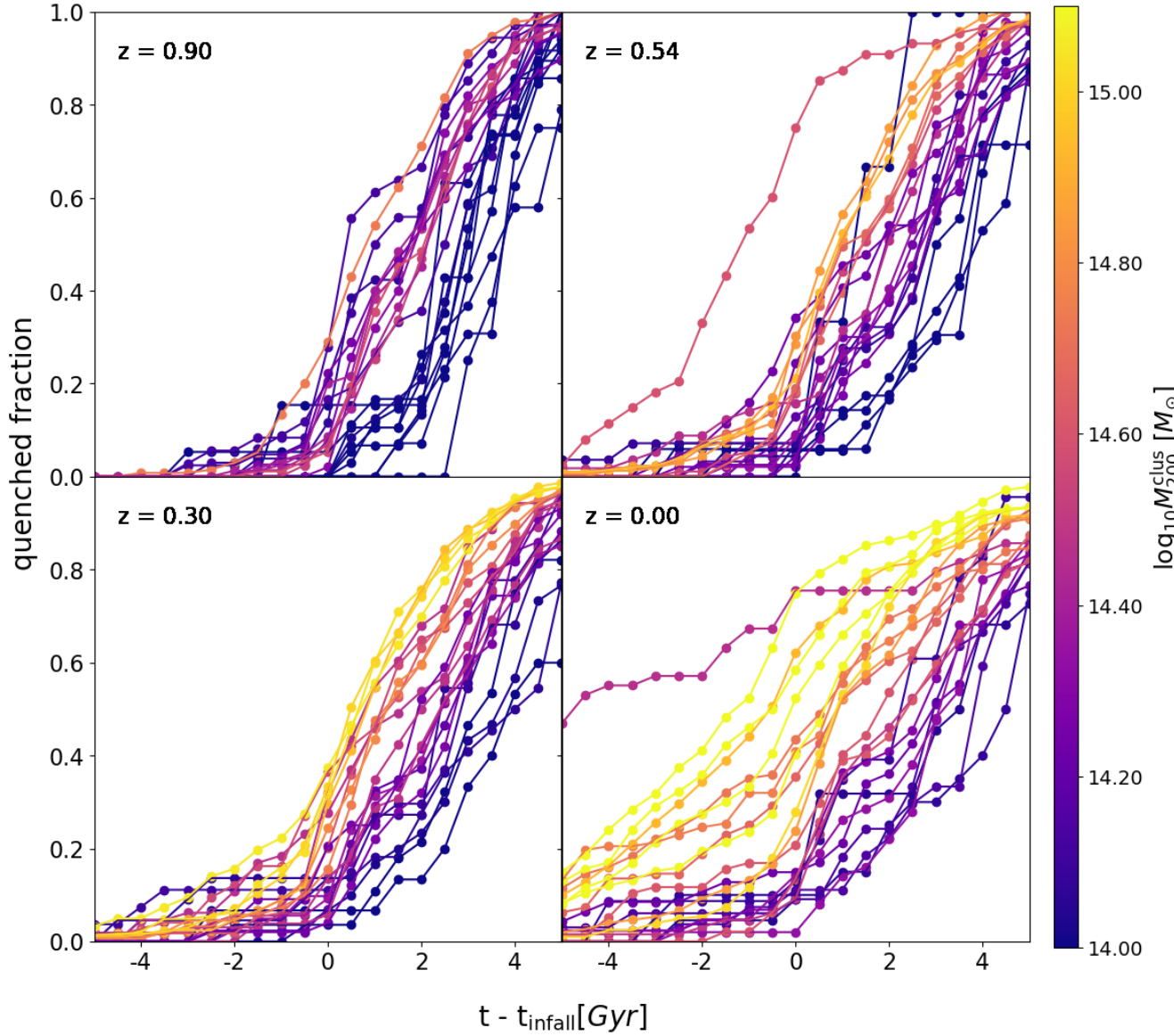


Figure A3. Growth of the quenched fraction of the satellite population of galaxy clusters as a function of the normalized timescale $t - t_{\text{infall}}$ for clusters at $z = 0.00; 0.30; 0.54; 0.90$. Each line represents one cluster in the simulation, color-coded by their M_{200} at the given redshift. Negative (positive) values correspond to times before (after) the first crossing of the R_{200} of the final cluster. Galaxies were selected using a threshold in stellar mass $M_{\star} \geq 10^{9.5} M_{\odot}$. As we look at clusters at higher redshift, we can see that the fraction of pre-quenching is reduced significantly ($t - t_{\text{infall}} < 0$). At the time of their infall ($t - t_{\text{infall}} = 0$) we can see a sharp increase in quenching and an abrupt change in the slope of the fraction of quenched galaxies.

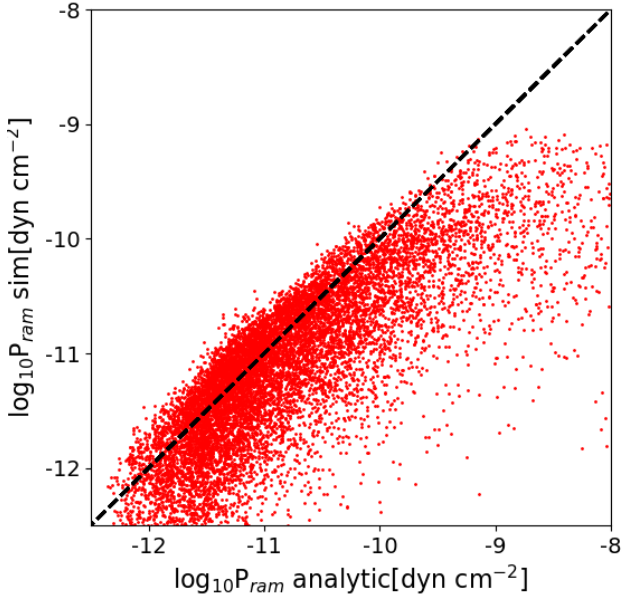


Figure B1. A comparison between the ram pressure measured from the simulations using the gas density profiles of all main clusters and the relative velocity of the galaxies with respect to the BCG and the ram pressure measured with the profile provided in [Vega-Martínez et al. \(2021\)](#). The black dashed line shows the value if the relation were one to one. As we can see both values are in great agreement, specially for the low ram pressure region.