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Insight into the chromophore of rhodopsin and its Meta-II photointermediate by ^{19}F solid-state NMR and chemical shift tensor calculations†

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^{19}F nuclei are useful labels in solid-state NMR studies, since their chemical shift and tensor elements are very sensitive to the electrostatic and space-filling properties of their local environment. In this study we have exploited a fluorine substituent, strategically placed at the C-12-position of 11-*cis* retinal, the chromophore of visual rhodopsins. This label was used to explore the local environment of the chromophore in the ground state of bovine rhodopsin and its active photo-intermediate Meta II. In addition, the chemical shift and tensor elements of the chromophore in the free state in a membrane environment and the bound state in the protein were determined. Upon binding of the chromophore into rhodopsin and Meta II, the isotropic chemical shift changes in the opposite direction by +9.7 and −8.4 ppm, respectively. An unusually large isotropic shift difference of 35.9 ppm was observed between rhodopsin and Meta II. This partly originates in the light-triggered 11-*cis* to all-*trans* isomerization of the chromophore. The other part reflects the local conformational rearrangements in the chromophore and the binding pocket. These NMR data were correlated with the available X-ray structures of rhodopsin and Meta II using bond polarization theory. For this purpose hydrogen atoms have to be inserted and hereto a family of structures were derived that best correlated with the well-established ^{13}C chemical shifts. Based upon these structures, a 12-F derivative was obtained that best corresponded with the experimentally determined ^{19}F chemical shifts and tensor elements. The combined data indicate strong changes in the local environment of the C-12 position and a substantially different interaction pattern with the protein in Meta II as compared to rhodopsin.

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Introduction

Rhodopsin is the visual pigment of the rod photoreceptor cells in the vertebrate retina, which mediate dim-light vision.^{1–3} As a member of the membrane protein superfamily of the heptahelical G protein-coupled receptors, rhodopsin is specialized for activation by photons in such a way that its ligand is the

photosensory element (chromophore) that has become covalently bound to the apoprotein, opsin. The ligand is a vitamin A derivative, 11-*cis* retinal (Fig. 1), which is attached to Lys-296 in the binding site through formation of a protonated Schiff base.^{1,4–7} This local charge is stabilized by a nearby located soft counterion formed by a H-bonded network concentrated around a glutamate residue of opsin (Glu-113).^{8–13} The 11-*cis* isomer acts as a potent inverse agonist of opsin, nearly completely eliminating its basal activity, but light absorption can convert it into the all-*trans* configuration that acts as a full agonist.^{2,14} This photo-isomerization of the ligand induces conformational changes in the protein, which are driven by about 35 kcal of photon energy stored in the photoproduct Batho.^{15–17} Within several milliseconds these structural transitions culminate in generation of the active state, Meta II, which binds and activates its cognate G protein transducin.^{14,18} The active state is characterized by deprotonation of the Schiff base^{14,19,20} and by subtle rearrangements in several structural elements like helices V, VI and VII, intracellular loops and the binding site with the

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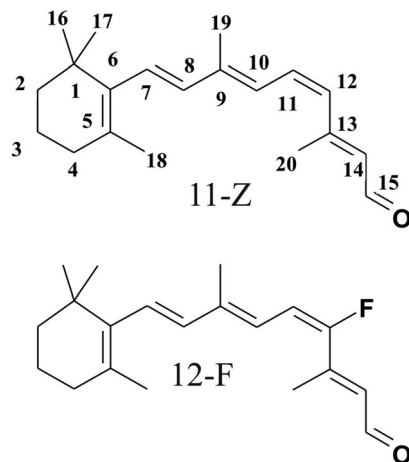


Fig. 1 Chemical structure of 11-*cis* retinal (11-Z) and the 12-fluoro derivative used in this study. The structures are shown in the 12-*s trans* conformation, which is dictated by the opsin binding pocket.

ligand in the all-*trans* configuration.^{21–31} Meta II decays with a half time of minutes, whereby the all-*trans* retinal is released from the orthosteric binding site and the opsin structure partially reverts to the original rhodopsin conformation.^{1,32–35} The availability of crystal structures of rhodopsin, Meta II and opsin presents a major step forward in understanding important elements of the activation mechanism of G protein-coupled receptors.²⁶ On the other hand, the resolution varies between 2.2 and 3.3 Å,^{11,12,24,25,36–38} which implies that atomic details cannot be exactly derived from the electron density. For instance, modeling is used to derive the position of the ligand, and this position is not equivalent to the one obtained by quantum mechanics/molecular dynamics.¹² In addition, the crystal lattice imposes limitations on the 3D conformation of the protein, in particular the more flexible intracellular surface, and may select one that is not the dominant one in the natural aqueous membrane dispersion.^{11,12,31,39,40}

Hence it is essential to confront crystal structures with experimental evidence derived from alternative conformational analyses allowing a natural membrane environment. While biochemical and biophysical (ESR, NMR, wide-angle X-ray scattering) data generally qualitatively agree with the structural alterations in the transition from rhodopsin to Meta II as indicated by the crystal structures,^{3,10,13,21–23,27,28,35,41,42} quantitative agreement has proven to be difficult to achieve, as yet.^{31,43}

Here we present a solid-state NMR approach with fluorine-labeling in combination with computational chemistry. The advantage of the fluorine label as a NMR probe is its high gyromagnetic ratio, large chemical shift range and high sensitivity to its local environment.^{44–46} In addition, from the spinning side bands in solid state NMR the principal values of the chemical shift (CS) tensors can be easily derived, which provide additional geometrical information. Prerequisite for drawing sensitive conclusions from the tensor elements is that the orientation of the fluorine CS tensor is known. Fluorine labels have been used before in solution NMR studies on rhodopsin solubilized in detergent micelles to probe its photoactivation. The label was

introduced either by chemical modification of cysteine mutants or by phosphorylation^{35,47,48} or by incorporation of fluoro-labeled retinals.^{49–53} We have opted for an 11-*cis* retinal analog, since the microenvironment of the rhodopsin ligand is relatively well defined in the crystal structures. A fluorine label at C-12 was selected (Fig. 1) since the C-10–C-13 segment in the chromophore is the trigger point for the photoactivation of rhodopsin^{54–61} and reorientation of the chromophore proceeds up to the Meta II state.^{13,24,62,63} Furthermore, the ¹⁹F-chemical shift at the C-12 position was shown not to be very sensitive to the group at C-15 (free aldehyde or (un)protonated Schiff base),^{49,50,52} and in a recent study we showed that photo-activation triggers changes in the vibrational properties of a fluorine label at C-12.⁶⁴

In this report we demonstrate that incorporation of 11-*cis* 12-F-retinal results in a significant positive ligation shift of the chemical shift of the fluorine label, primarily reflected in the δ_{11} and δ_{22} tensor elements, while Meta II exhibits a significant negative ligation shift, primarily reflected in the δ_{33} tensor element. Our subsequent computational studies demonstrate that the isotropic ligation shift and changes of the tensor elements originate in changes of the F–C bond length and the polarizing influence of the microenvironment of the retinal.

Experimental section

Materials

All chemicals were of analytical grade. Detergents were obtained from Anatrace (Maumee, OH, USA). Unmodified 11-*cis* retinal was provided by Dr Rosalie Crouch (Medical University of South Carolina, Charleston, USA) through financial support from NEI.

Synthesis of 12-F retinal

The 12-F derivative was prepared in the 11-*cis* configuration as described before.⁶⁵ In this case the C15 precursor was synthesized as a mixture of the all-*trans* and 11-*cis* isomers. The 11-*cis* isomer was isolated by chromatography and extended into the corresponding retinal. The spectral and ¹H NMR characteristics of the obtained 11-*cis* 12-F-retinal agreed with published data.^{64,65}

Isolation of bovine opsin and regeneration with 11-*cis* retinals

Bovine rod outer segment membranes in the opsin form (opsin membranes) were prepared from fresh, light-adapted cattle eyes as described.^{66,67} The regeneration capacity of these preparations was estimated from the A_{280}/A_{500} ratio measured after subsequent incubation with a 3-fold excess of 11-*cis* retinal, whereby a ratio of 2.1 ± 0.1 was taken to represent membranes with a maximal rhodopsin content. Rhodopsin and the 12-fluoro-rhodopsin analog were prepared with opsin membranes showing a regeneration capacity in the range 90–100%. Regeneration and all further manipulations with the pigments were done under dim red light (> 620 nm, RG645 longpass filter).

The 12-F analog pigment was generated by incubating a suspension of opsin membranes (50–100 μM opsin in buffer A: 20 mM PIPES, 130 mM NaCl, 5 mM KCl, 2 mM CaCl₂, 0.1 mM EDTA, 1 mM DTE, pH 6.5) with a 2–3 fold molar excess of the

11-*cis* fluororetinol at room temperature. After two hours a small aliquot was assayed for the extent of regeneration by addition of 11-*cis* retinal in a 1/1 molar ratio to the original opsin. Regeneration appeared to be nearly complete and to achieve full regeneration an additional aliquot of the fluororetinol was added and the incubation continued overnight at 4 °C. Excess retinal was then converted into the corresponding oxime by addition of an 1 M hydroxylamine solution (pH 6.5) to a final concentration of 10 mM. After cooling in ice and 30 min incubation, the oxime was largely removed by two extractions with 50 mM heptakis-(2,6-di-*O*-methyl)- β -cyclodextrin,⁶⁶ which, however, also removes some lipids thereby perturbing the Meta I to Meta II transition. To restore a native lipid/protein ratio, the membrane pellet was dissolved in 20 mM nonylglucoside in buffer A (to *ca.* 50 μ M of pigment) by incubation for 1 h on ice. Undissolved material was removed by centrifugation (30 min, $80\,000 \times g$, 4 °C), and the supernatant mixed with a solution of asolectin (100 mg ml⁻¹ in 50 mM nonylglucoside in buffer A) to achieve a 50-fold molar excess of asolectin with respect to pigment. After 15 min incubation on ice, detergent was extracted by addition of solid β -cyclodextrin to a slight excess over nonylglucose and the resulting proteoliposomes isolated by overnight sucrose step-gradient centrifugation at $200\,000 \times g$ and 4 °C as described before.⁶⁸ The proteoliposome band was recovered from the 20%/45% sucrose interface, diluted with 2 volumes of doubly distilled water, pelleted by centrifugation (60 min, $200\,000 \times g$, 4 °C) and stored in aliquots under Argon in a light-tight container at -80 °C.

UV/vis spectroscopy

The spectral properties of the pigments were determined in micellar solution, by solubilization to about 2.5 μ M in 20 mM dodecylmaltoside in buffer A containing 10 mM hydroxylamine. The wavelength of maximal absorbance in the visible region (λ_{max}) was determined as the peak position in the difference spectrum obtained after subtraction of the spectrum after illumination (300 s; 150 W halogen light through a Schott OG530 longpass filter and fiber optics) from the dark state spectrum.

The photoequilibrium obtained after illumination was analysed using the end-on photomultiplier set-up of a Perkin-Elmer Lambda 18 spectrophotometer. Pigment proteoliposomes were suspended to about 30–40 μ M of pigment in buffer B (similar to buffer A, but with 20 mM MES as a buffering compound and a pH of 5.0). The sample was maintained at 0 °C in a melting ice bath and illuminated (conditions as above). Aliquots were taken before and after 30 s, 2 min and 4 min of illumination, and spectra were recorded in a cuvette with a 2 mm light path. After 4 min of illumination the proteoliposomes were spun down (20 min, $100\,000 \times g$, 0 °C) and the pellet was frozen in a solid CO₂/alcohol mixture.

Preparation of 12-F rhodopsin and 12-F metarhodopsin II (Meta II)

The 11-*cis* 12-F-retinal derivative (Fig. 1) was prepared and purified as described before.^{64,65} All manipulations with rhodopsin were performed under dim red light (Schott RG645 longpass filter)

except for the preparation of Meta II. The rhodopsin 12-F analogue was prepared from opsin and 11-*cis* 12-F-retinal, also as described before,⁶⁴ yielding a pigment with its maximal absorbance at 510 ± 2 nm (Fig. 2) in agreement with earlier data.^{64,65} In comparison to rhodopsin the Meta I to Meta II transition of 12-F rhodopsin is strongly shifted to a lower pH range.⁶⁴ Therefore, 12-F Meta II was generated by illumination of 12-F rhodopsin in proteoliposomes at a pH of 5.0, where close to maximal conversion to Meta II can be realized.⁶⁴ Since the decay of Meta II is strongly reduced at lower temperature, illumination was performed at 0 °C, where the half-life of Meta II is several hours. Under the selected conditions maximal conversion to 12-F Meta II could be achieved after 4 min of illumination (Fig. 2). Based upon the pH dependence of the 12-F Meta I to Meta II transition and the absorbance shift^{64,69} it was estimated that the photoproduct consisted of about 10% remaining rhodopsin, 20% Meta I and 70% Meta II. Meta II can be stabilized by freezing below -13 °C (260 K)⁷⁰ and, after filling the rotor with the Meta II proteoliposomes by centrifugation at 0 °C, the rotor was stored at -80 °C.

Solid-state NMR spectroscopy

Varian 4 mm rotors were used, where the standard Vespel spacer and endcap had been replaced by ones machined from Kel-F. All subsequent manipulations were performed in the dark or under dim red light (Schott RG645 longpass filter). The rotors were filled with pigment proteoliposomes by centrifugal force (120 min, $30\,000 \times g$, 0 °C).⁹ Per rotor about 300 nmoles (12 mg) of pigment was applied. The frozen pellet of the Meta II sample, generated as described above, was directly applied into the centrifuge tube insert, and allowed to melt during centrifugation, thereby pelleting the released membranes. The filled rotors were stored at -80 °C for maximally two days before the NMR analysis was started.

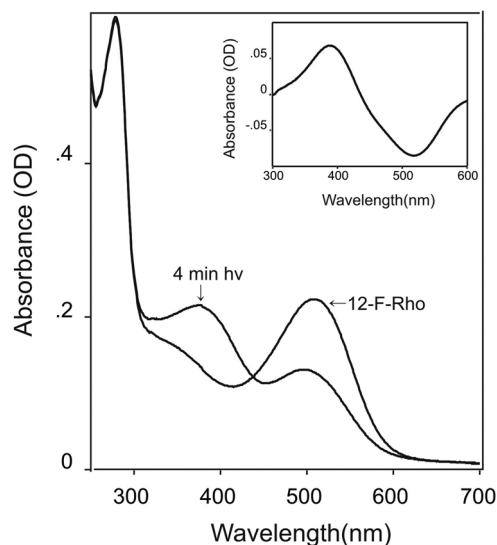


Fig. 2 UV/vis spectra of 12-F rhodopsin in proteoliposome suspension before (12-F-Rho) and after four min of illumination (4 min *hν*). The inset presents the difference between these spectra, demonstrating that the major photoproduct is Meta II (385 nm absorbance band).

All experiments were performed on either a Varian Infinity or VNMRs console at a static magnetic field of 14.1 T, employing a Chemagnetics 4 mm HFXQ quadrupole resonance probehead. The top and the bottom of the magnet bore were covered by blackout curtains to prevent any light from reaching the sample during the NMR experiments. The ^{19}F cross-polarization (from ^1H) magic-angle-spinning (CP-MAS) experiments⁷¹ were performed at ^1H and ^{19}F Larmor frequencies of 600 and 564.5 MHz, respectively, using spinning frequencies of 10 or 12 kHz and a ^{19}F Hahn echo⁷² of two rotational periods before ^{19}F signal detection under proton decoupling. The ^{19}F spectra were referenced to CCl_3F via the ^{13}C chemical shift of the Adamantane resonance following the IUPAC recommendations.^{73,74} All ^{19}F experiments were setup using a powdered sample of Fmoc-(4-F)-Phe-OH at room temperature. The calibration of the variable temperature (VT) controller using a sample of lead nitrate ($\text{Pb}(\text{NO}_3)_2$)^{75,76} showed that a readout temperature of -80°C used in our low-temperature experiments on rhodopsin corresponds to sample temperatures of -71.5°C at 12 kHz MAS frequency and -74.8°C at 10 kHz MAS frequency, hence well below the temperature where Meta II starts to decay.

The ^{19}F CP-MAS spin-echo experiments on the sample of 11-*cis* 12-F-retinaloxime in liposomes were performed at 10 and 12 kHz MAS frequency, at a sample temperature of -74.8 and -71.5°C , respectively, utilizing a 90° proton pulse of 3.1 μs , and a ^{19}F 180° pulse of 8 μs . The nutation frequency of the proton rf field during CP was given by 40 kHz and the rf field on the ^{19}F was adjusted to fulfill the (+1) Hartman-Hahn sideband matching condition.⁷¹ XiX proton decoupling⁷⁷ with a nutation frequency of 58 kHz was applied during ^{19}F signal detection. In total 5120 transients were accumulated with a relaxation delay of 10 s between each transient. An exponential apodization function with a width of 500 Hz was applied before Fourier transformation of the ^{19}F time domain signal.

The experiments with 12-F rhodopsin were also performed at 10 and 12 kHz MAS frequency, at a sample temperature of -74.8 and -71.5°C , respectively, and the same experimental parameters were used as for the 11-*cis* 12-F-retinaloxime sample. The number of total accumulated transients was 56 432 and 15 984 for the ^{19}F CP-MAS spin echo experiments at 10 and 12 kHz MAS frequency, respectively, using a relaxation delay of 5 s. Exponential apodization with a width of 500 Hz was again applied to the time domain signal.

The CP-MAS experiments with illuminated 12-F rhodopsin were solely performed at a MAS frequency of 12 kHz and a sample temperature of -71.5°C . The ^1H 90° and ^{19}F 180° pulses were given by 3 and 8 μs , respectively. The proton rf field during the cross-polarization was set to 53 kHz and the ^{19}F rf field was adjusted to match the (−1) Hartman-Hahn sideband condition. The nutation frequency of the XiX decoupling during ^{19}F signal detection was given by 53 kHz. 14 048 transients with a relaxation delay of 5 s were recorded, and an exponential apodization function with a width of 500 Hz was again applied before Fourier transformation.

Analysis of solid-state NMR spinning-sideband patterns

The experimental MAS sideband patterns of the ^{19}F chemical shift interaction were fitted to calculations to obtain the 3 principal

components of the chemical shift tensor. We chose to present our results using both the Haeberlen and Maryland conventions^{73,78,79} since both have their merits when discussing results from experiments and theoretical calculations. In the Haeberlen convention^{73,78} the three principal components in the deshielding convention are labeled δ_{xx} , δ_{yy} and δ_{zz} and ordered according to $|\delta_{zz} - \delta_{iso}| \geq |\delta_{xx} - \delta_{iso}| \geq |\delta_{yy} - \delta_{iso}|$, where the isotropic chemical shift is given by $\delta_{iso} = (\delta_{xx} + \delta_{yy} + \delta_{zz})/3$. In addition, the anisotropic chemical shift δ_{aniso} and asymmetry parameter η of the chemical shift anisotropy (CSA) are defined as $\delta_{aniso} = \delta_{zz} - \delta_{iso}$ and $\eta = (\delta_{yy} - \delta_{xx})/\delta_{aniso}$, respectively. On the other hand, in the Maryland convention^{73,79} the three principal components are labeled δ_{11} , δ_{22} and δ_{33} and are ordered according to $\delta_{11} \geq \delta_{22} \geq \delta_{33}$. The isotropic chemical shift δ_{iso} , the span Ω and skew κ are given by $\delta_{iso} = (\delta_{11} + \delta_{22} + \delta_{33})/3$, $\Omega = \delta_{11} - \delta_{33}$ and $\kappa = 3(\delta_{22} - \delta_{iso})/\Omega$. The Maryland convention is useful when discussing the chemical shift tensor as a whole, while the Haeberlen convention is more applicable when discussing the CSA principal tensor values separately from the isotropic chemical shift. Below we will report the three principal components δ_{11} , δ_{22} and δ_{33} as well as δ_{iso} , δ_{aniso} , η , Ω and κ . The spinning-sideband amplitudes for a single crystallite orientation are calculated by discrete Fourier transformation of the signal amplitudes obtained by evolution under the chemical shift interaction over one rotational period.⁸⁰ The final spinning-sideband manifold is obtained by powder averaging using the octant Lebedev set containing 109 pairs of Euler angles.⁸¹ The experimental ^{19}F spectra were fitted in a single step to the calculated sideband manifold, hence fitting parameters included not only the parameters (δ_{iso} , δ_{aniso} , η), but also the parameters describing the shape of the centerband line according to a linear combination of a Gaussian and Lorentzian line shape (overall scaling, width, fraction Lorentzian). Hence, a total of 6 fitting parameters were used for each sample component. Experimental results obtained at different spinning frequencies were simultaneously fitted. We obtained the 68.27% confidence interval, corresponding to an error of plus/minus one standard deviation for a normal distribution, for δ_{aniso} and η , in a separate step by calculating the sum-squared deviation S between the experimental and calculated integrated sideband amplitudes for a series of different values for either δ_{aniso} or η , where the other fitting parameters were optimized as to minimize S in each case. The confidence interval is determined by the set of values for either δ_{aniso} or η for which $S \leq S_{\min}\{1 + F_{1,n-p}^\alpha/(n - p)\}$, where $S_{\min}$ is the sum-squared deviation minimized by optimizing all fitting parameters, n is the number of integrated sideband amplitudes used in the fit, p is the total number of fitting parameters, $\alpha = 0.3173$ and F_{p_1, p_2}^α is the upper α probability point of the F distribution with p_1 and p_2 degrees of freedom.^{82,83} The error of δ_{aniso} is determined in an analogous manner from the direct line shape fitting of the spinning sideband manifold. The errors for Ω , κ , δ_{11} , δ_{22} and δ_{33} were estimated by error propagation.

Computational modeling and NMR tensor calculations by bond polarization theory (BPT)

The bond polarization theory (BPT) works with localized bond contributions and the influence of changes of the structure and

charge distribution on molecular properties can be readily separated and discussed. In this investigation the BPT is used for atomic partial charges⁸⁴ and ¹³C and ¹⁹F chemical shifts.^{46,85} For the ¹⁹F chemical shift the BPT equation has a simple form since only the polarization of the F–C bond has to be discussed. Eqn (1) gives the BPT expression for the chemical shift tensor δ in the bond coordinate system of the F–C bond:

$$\delta_{\alpha\beta} = n\delta_{\alpha\beta}^{\text{F-C}} + n^2 A_{\alpha\beta}^{\text{F-C}} (\langle \chi_{\text{F}} | \hat{V} | \chi_{\text{F}} \rangle - \langle \chi_{\text{C}} | \hat{V} | \chi_{\text{C}} \rangle) \quad (1)$$

In the bond coordinate system (the z-axes are pointing into bond direction) the chemical shift tensor is regarded as diagonal with the principal components $\{\delta_{xx}, \delta_{yy}, \delta_{zz}\}$ (the discussion of eqn (1) is given in terms of chemical shift and in this case we have to add -93.68 ppm as sum of the ¹⁹F inner shell contribution of -306.3 ppm⁸⁶ and a reference shielding of 212.62 ppm, see ref. 46). The principal components are calculated as a sum of two terms: a bond tensor $\delta^{\text{F-C}}$ and a second term describing the change of the nuclear shielding by polarization. The bond tensor stands for a bond with the configuration weight $n = 1.0$ with no polarizing charges in its surrounding.

The principal components of this CS bond tensor are given in Table 1 and in CS calculations this bond tensor is scaled by the configuration weight of the bond n . The introduction of n means that the bonds are not treated as strictly localized but bond electrons can occupy anti-bond states in other bond configurations. An empirical relation connects the configuration weight to valence of the bond and to the bond length.^{87,88} In this BPT description the bond tensor contains also the shift contributions of the three fluorine lone pair orbitals.

The term in brackets in eqn (1) is called polarization energy and is calculated from the potentials V of the charges q in the surrounding of the F–C bond and from the fluorine hybrid orbitals χ_{F} and the carbon hybrid orbitals χ_{C} forming the bond. Since the influence of potentials q/r get weaker with increasing distance r only charges in the vicinity of the F–C bond have a large influence. To get the polarizing influence on the chemical shift the polarization energy is multiplied by the tensor $A^{\text{F-C}}$ with the dimension ppm/energy (Table 1). These principal components can be (in good approximation) assigned to the bond coordinate system in the following way: component yy perpendicular to the carbon π system, component xx within the π plane and zz in F–C bond direction. The polarization energy is a scalar but the tensor $A^{\text{F-C}}$ can have very different elements even with different signs and therefore the polarizing influence in different directions is highly anisotropic. Since both tensors $\delta^{\text{F-C}}$ and $A^{\text{F-C}}$ are diagonal with respect to the bond coordinate system the final result of eqn (1) will be diagonal in this coordinate system as well. Only through-space shieldings of other bonds could contribute off-diagonal elements but these contributions are neglected in the framework of the BPT.

The problem of quantum chemical CS calculations is the need for precise structures. Minute bond length changes can cause large changes in chemical shifts and even X-ray structures are not precise enough. The next problem of the X-ray investigations is that in the case of the rhodopsin structures no proton positions are refined. Starting with the X-ray structure of the dark adapted rhodopsin (PDB 1U19) by Okada *et al.*¹² and the structure of the light activated form (PDB 4A4M) by Deupi *et al.*,³⁶ protons are substituted and possible van der Waals contacts were healed out by short MD simulations (for the complete modeling protocol see the ESI†).

Starting from the BPT equations derivatives of the chemical shifts with respect to Cartesian coordinates have been developed⁸⁹ and these CS derivatives allow geometry optimizations with chemical shifts as constraints. It is possible to derive pseudo-energies from the differences of calculated and experimental chemical shifts and the derivatives of these pseudo-energies are CS pseudo-forces that are needed to drive the geometry optimizations. The CS pseudo-energies are added to the regular energies of the COSMOS-NMR force field.^{90,91} The COSMOS-NMR force field has the possibility to work with several forms of NMR constraints and it has been applied to protein structure investigations using ¹³C and ¹⁵N chemical shifts.⁸⁵ This opens the possibility to refine the structure of the retinal molecule and its binding pocket within the retinal protein. We used literature data for the ¹³C chemical shifts of the retinal as collected by Fröhmknecht *et al.*⁹² as constraints for a global minimum search of both retinal forms. Since gradient procedures for the minimization always run into the next local minimum we refine a large library of more than 1000 MD snapshots and selected the 20 best fitting structures as final result. As criterion we used the CS pseudo energy of the retinal (for details see ESI†). The proton of C12 of retinal was converted into a fluorine and the geometry of the fluorinated compounds was geometry optimized to adjust the C–F bond lengths. Assuming that the fluorination has only a minor influence on the structure these models were used for fluorine tensor calculation. To adjust the retinal structure and its binding surrounding to the fluorination we geometry optimized the models using the experimental isotropic ¹⁹F chemical shifts as constraints.

A better way to refine the fluorinated structure would be to search again for a global minimum by using a large library of structures as starting point. But since we had only one experimental NMR constraint we cannot exclude abnormal structures that give the right CS value (see ESI†).

Results

F-NMR analysis of Fmoc-[4-F]-L-phenylalanine-OH

In order to setup optimal conditions and validate the ¹⁹F CP-MAS experiments we first analyzed a sample of powdered Fmoc-[4-F]-L-phenylalanine-OH (Fmoc-[4-F]-L-Phe). The results we obtained for

Table 1 Components of the ¹⁹F CS bond tensor $\delta^{\text{F-C}}$ and the CS polarization tensor $A^{\text{F-C}}$. The tensors are assumed to be diagonal in the bond coordinate system with the zz component in F–C direction, the yy component perpendicular to the π plane of the sp^2 carbon and xx within the π plane

Component	xx	yy	zz
¹⁹ F CS bond tensor [ppm]	−76.98	−203.50	−117.18
CS polarization tensor $A^{\text{F-C}}$ [ppm Hartree ^{−1}]	−1393.1	−2110.8	450.8

the chemical shift principal components are shown in Table 2. We found isotropic and anisotropic chemical shifts of $\delta_{\text{iso}} = -113.43 \pm 0.01$ ppm and $\delta_{\text{aniso}} = 61.0 \pm 0.5$ ppm with an asymmetry parameter of $\eta = 0.74 \pm 0.02$, corresponding to a span and skew of $\Omega = 114.0 \pm 1.0$ ppm and $\kappa = -0.21 \pm 0.01$, respectively. These values are very close to the $\delta_{\text{iso}} = -110.8$ ppm, $\delta_{\text{aniso}} = 58.9$ ppm and $\eta = 0.74$, reported for the close analog [4-F]-L-Phe by the group of Ulrich.⁹³ This validates our set-up for the subsequent ¹⁹F-NMR studies reported below.

F-NMR analysis of free 11-*cis* 12-F retinal in liposomes

The F-NMR properties of the free ligand were determined for the oxime derivative, since the type of headgroup (free carbonyl or amine derivative) does hardly influence the chemical shift of the 12-F substituent,^{49,50,52} and the oxime is much less prone to oxidation and side reactions with lipids. The 11-*cis* 12-F-retinnoxime was prepared by addition of 11-*cis* 12-F-retinal, dissolved in DMF, to a suspension of asolectin in buffer A, containing 50 mM hydroxylamine, at a molar ratio of lipid to retinoid of about 5 to 1. The liposomes were then pelleted by centrifugation and subsequently packed into the rotor as described above.

NMR analysis of the ¹⁹F CP-MAS spectra shown in Fig. 3 resulted in the data presented in Tables 2 and 3. We found an isotropic and anisotropic chemical shift of $\delta_{\text{iso}} = -104.7 \pm 0.2$ ppm and $\delta_{\text{aniso}} = 46.2 \pm 1.2$ ppm together with an asymmetry parameter of $\eta = 0.59 \pm 0.09$. These values correspond to a span of $\Omega = 83.0 \pm 2.9$ ppm and a skew of $\kappa = -0.34 \pm 0.08$. The isotropic and anisotropic shifts are slightly different but close to the ones found for Fmoc-[4-F]-L-Phe (Table 2) and for other aromatic fluorine sites in amino acids.⁹³ The isotropic shift agrees quite well with the value of -107 ppm reported for 11-*cis* 12-F-retinal in CDCl₃ or CD₂Cl₂ solution.^{49,94} The ¹⁹F resonance in 11-*cis* 12-F-retinnoxime has a linewidth of 4.3–4.6 kHz, much larger than the 0.4 kHz of Fmoc-[4-F]-L-Phe (Table 2), indicating that the 11-*cis* 12-F-retinnoxime molecules in the liposomes are not isotropic but enjoy some conformational freedom and experience a range of different environments resulting in a spread of chemical shifts.

F-NMR analysis of 12-F rhodopsin

In order to obtain ¹⁹F spectra of 12-F rhodopsin in the dark state, the sample was brought from storage in dry-ice into the magnet under dim red light (Schott RG645 longpass filter). The resulting CP-MAS fluorine spectra obtained at 10 and 12 kHz spinning frequency are shown in Fig. 4. The linewidth of about 1.8 kHz indicates a well-ordered conformation of the chromophore in 12-F rhodopsin in contrast to the case of the free 11-*cis* 12-F-retinnoxime in liposomes. From the ¹⁹F spinning-sideband pattern we were able to derive an isotropic and anisotropic chemical shift of $\delta_{\text{iso}} = -95.00 \pm 0.06$ ppm and $\delta_{\text{aniso}} = 53.3 \pm 0.9$ ppm together with an asymmetry parameter of $\eta = 0.82 \pm 0.03$ (Table 2), corresponding to a span and skew of $\Omega = 101.7 \pm 2.0$ ppm and $\kappa = -0.15 \pm 0.03$, respectively. The isotropic chemical shift of the 12-F site exhibits a significant ligation shift of +9.7 ppm in 12-F rhodopsin relative to the free 11-*cis* 12-F-retinnoxime (Table 3). This isotropic shift and ligation shift are

Table 2 Experimental and calculated results of ¹⁹F chemical shift tensors^a

Sample	Component	Linewidth ^b [Hz]	Fraction Lorentzian ^b [%]	Haerberlen		Maryland		$\delta_{11} \geq \delta_{22} \geq \delta_{33}$		
				δ_{iso} [ppm]	η	Ω [ppm]	κ	δ_{11} [ppm]	δ_{22} [ppm]	δ_{33} [ppm]
Experiments										
Fmoc-[4F]-Phe-OH		384 ^c /407 ^c	29 ^c /27 ^c	-113.43 ± 0.01	0.74 ± 0.5	114.0 ± 1.0	-0.21 ± 0.01	-52.4 ± 0.5	-121.5 ± 0.5	-166.4 ± 0.6
11- <i>cis</i> 12-F-retinoxime		4365 ^c /4624 ^d	62 ^c /38 ^d	-104.7 ± 0.2	0.59 ± 1.2	83.0 ± 2.9	-0.34 ± 0.08	-58.4 ± 1.2	-114.2 ± 2.1	-141.4 ± 2.3
12-F rhodopsin (dark state)		1837 ^c /1785 ^d	85 ^c /73 ^d	-95.00 ± 0.06	53.3 ± 0.9	101.7 ± 2.0	-0.15 ± 0.03	-41.7 ± 0.9	-99.9 ± 0.9	-143.4 ± 1.2
12-F rhodopsin (illuminated)	Meta II	2785 ^d	100 ^d	-130.9 ± 0.1	54.0 ± 4.5	108.1 ± 9.2	0.00 ± 0.06	-76.8 ± 4.5	-130.9 ± 2.2	-184.9 ± 5.0
	Liberated retinal	3273 ^d	15 ^d	-122.5 ± 0.1	50.1 ± 4.4	97.0 ± 9.8	-0.10 ± 0.15	-72.4 ± 4.4	-125.6 ± 4.6	-169.4 ± 6.2
Calculations										
12-F rhodopsin (dark state)	Glu181 protonated			-114.0	60.0	110.1	-0.27	-54.0	-124.0	-164.1
12-F rhodopsin (dark state)	Glu181 unprotonated			-110.6	63.3	110.5	-0.44	-47.3	-126.7	-157.8
12-F rhodopsin (illuminated)	Meta II			-125.6	-62.1	117.2	0.18	-70.5	-118.5	-187.7

^a Errors are presented as standard deviation, calculated as explained in the Experimental section. ^b Linewidth of the centerband and fraction Lorentzian resulting from the lineshape fitting. ^c At 10 kHz MAS frequency. ^d At 12 kHz MAS frequency. ^e At 13 kHz MAS frequency.

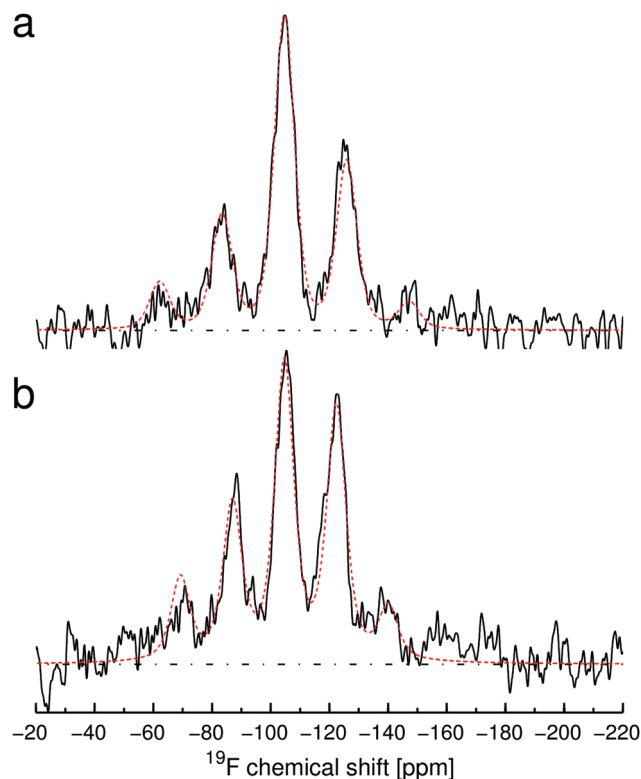


Fig. 3 Solid lines: experimental ^{19}F CP-MAS spectra of 11-*cis* 12-F-retinoyl in asolectin liposomes. Red dashed lines: best-fit results from fitting the chemical shift tensor components. The sample spinning frequency is given by (a) 12 kHz and (b) 10 kHz. The dashed dotted line indicates the baseline.

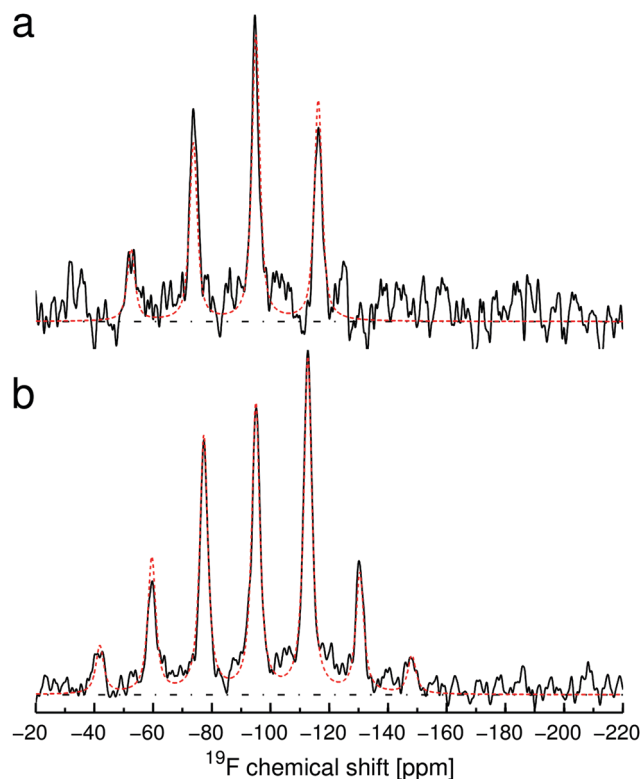


Fig. 4 Solid lines: experimental ^{19}F CP-MAS spectra of 12-F rhodopsin (dark state). Red dashed lines: best-fit results from fitting the chemical shift tensor components. The sample spinning frequency is given by (a) 12 kHz and (b) 10 kHz. The dashed dotted line indicates the baseline.

in good agreement with data reported before for 12-F rhodopsin solubilized in detergent micelles: -94.2 ppm and about $+13$ ppm, respectively.^{53,94} We now show that the large ligation shift rests on very significant shifts in the δ_{22} and δ_{11} chemical shift tensor elements (Tables 2 and 3).

F-NMR analysis of the late photoproducts of 12-F rhodopsin

Analysis of the ^{19}F CP-MAS spectra of the late photoproducts of 12-F rhodopsin shown in Fig. 5 resulted in two sets of resonances with about similar total intensity but differing in centerband linewidth (Table 2). This is expected since in the time required to pack the sample into the rotor slow decay of 12-F Meta II will be proceeding and in the range of 30–50% will have decayed, while the signals of the remaining 12-F rhodopsin and 12-F Meta I are too low to be detectable above the noise level. During the decay of 12-F Meta II, the all-*trans* 12-F-retinal is liberated from the

binding site and then partly diffuses freely in the lipid matrix and partly becomes randomly bound to protein and lipid amino groups.³² Once the rotors are packed, they are stored and analyzed in a frozen state, which will halt further decay of Meta II. As further outlined in the discussion, we expect that, similar to the difference observed between the free 11-*cis* 12-F-retinoyl and 12-F rhodopsin, the liberated retinal will have a larger linewidth and a smaller anisotropy than the retinal still restricted in the binding site of Meta II. This difference is again observed between the two sets of resonances, and we have assigned them accordingly (Table 2). Since the signal-to-noise ratio in this experiment was not as good as in the rhodopsin analysis, a shorter fully independent second experiment was performed. This corroborated the data of the first experiment ± 0.2 ppm for Meta II and ± 3 ppm for the liberated retinal. However, it had a much poorer signal-to-noise ratio and therefore only the data of the first experiment have been presented in Table 2.

Table 3 Ligation shifts for the ^{19}F chemical shift tensors upon binding of 11-*cis* 12-F-retinal in rhodopsin and all-*trans* 12-F-retinal in Meta II. Shifts and errors are propagated from Table 2. Data in bold type represent significant differences

Sample	Ligation shift in chemical shift tensors upon Table 2 (free subtracted from bound)							
	$\Delta\delta_{\text{iso}}$ [ppm]	$\Delta\delta_{\text{aniso}}$ [ppm]	$\Delta\eta$	ΔQ [ppm]	$\Delta\kappa$	$\Delta\delta_{11}$ [ppm]	$\Delta\delta_{22}$ [ppm]	$\Delta\delta_{33}$ [ppm]
Bound versus free 11- <i>cis</i> 12-F-retinal	$+9.7 \pm 0.2$	$+7.1 \pm 1.5$	$+0.23 \pm 0.10$	$+18.8 \pm 3.5$	$+0.20 \pm 0.09$	$+16.8 \pm 1.5$	$+14.3 \pm 2.3$	-2.0 ± 2.6
Bound versus free all- <i>trans</i> 12-F-retinal	-8.4 ± 0.1	$+4.0 \pm 6.3$	$+0.13 \pm 0.20$	$+11.1 \pm 13.4$	$+0.10 \pm 0.16$	-4.5 ± 6.3	-5.3 ± 5.1	-15.5 ± 7.9

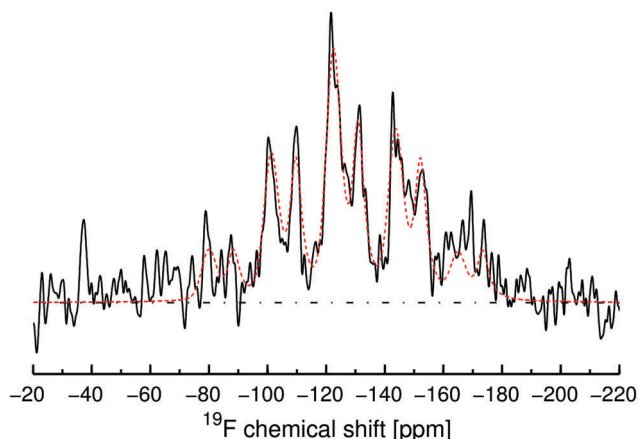


Fig. 5 Solid line: experimental ^{19}F CP-MAS spectrum of illuminated 12-F rhodopsin obtained at 12 kHz sample-spinning frequency. Red dashed lines: best-fit results from fitting the chemical shift tensor components. The dashed dotted line indicates the baseline.

The 12-F Meta II ^{19}F resonance has a linewidth of 2.8 kHz and the analysis of the spinning sideband pattern resulted in isotropic and anisotropic chemical shifts of $\delta_{\text{iso}} = -130.9 \pm 0.1$ ppm and $\delta_{\text{aniso}} = 54.0 \pm 4.5$ ppm, where the asymmetry parameter is given by $\eta = 1.00 (+0, -0.16)$ (Table 2). This corresponds to a span of $\Omega = 108.1 \pm 9.2$ ppm and a skew of $\kappa = 0.00 \pm 0.06$. The large change in the isotropic ^{19}F chemical shift of -35.9 ppm from the dark-state to the Meta-II state of 12-F rhodopsin is quite noteworthy.

The linewidth of the liberated all-*trans* 12-F-retinal is significantly broader (3.3 kHz). In this case we found an isotropic and anisotropic chemical shift of $\delta_{\text{iso}} = -122.5 \pm 0.1$ ppm and $\delta_{\text{aniso}} = 50.1 \pm 4.4$ ppm with an asymmetry parameter of $\eta = 0.87 (+0.13, -0.24)$ (Table 2), corresponding to span and skew of $\Omega = 97.0 \pm 9.8$ ppm and $\kappa = -0.10 \pm 0.15$, respectively. These values differ significantly from the ones found for the free 11-*cis* 12-F-retinonoxime in liposomes, but the isotropic shift is close to that of all-*trans* 12-F-retinal in CDCl_3 solution (-120 ppm) reported before.⁴⁹

We note that the errors in the anisotropic shifts and the asymmetry parameters are significantly larger for 12-F Meta II and liberated all-*trans* 12F-retinal compared to dark-state 12-F rhodopsin and 12-F-retinonoxime, as the experiments were solely performed at a single spinning frequency, resulting in less spinning sidebands being used in the analysis. Nevertheless, we find a significant ligation shift (-8.4 ppm) for all-*trans* 12-F-retinal upon binding to the protein, now in the Meta II state, but, remarkably, of opposite sign compared to the ligation shift observed for 11-*cis* 12-F-retinal (Table 3). Also in contrast to 11-*cis* 12-F-retinal, the major contribution to the ligation shift is primarily reflected in the δ_{33} tensor element.

Discussion

We have used a fluorine label at the C12 position of the retinylidene chromophore to probe its environment in the ground state of the rod visual pigment, rhodopsin, and in the activated state,

Meta(rhodopsin) II, by means of solid-state F-NMR. Fluorine labels are quite sensitive NMR-probes, with a broad chemical shift range, and good monitors of their local environment.^{44,45} The C12 position was selected for several reasons. First of all, it is part of the photo-active central segment of the chromophore (C-10–C-13) and its vibrational modes strongly support the 11-*cis* to 11-*trans* isomerization and subsequent conformational changes.^{10,54–58,60,61,63,64,95} In addition, our previous FTIR study shows significant changes in the vibrational properties of the C12-F derivative in the activated state, but not of the C10-F derivative.⁶⁴ Indeed, a preliminary F-NMR study on 10-F labeled octopus rhodopsin indicates small and very similar positive ligation shifts for the rhodopsin and the metarhodopsin state of $+6.4$ and $+4.5$ ppm, respectively.⁹⁶ Finally, in contrast to the C12 position, a fluorine label at C11 only shows a very small ligation shift in rhodopsin ($+2$ ppm) and would be very sensitive to bonding partners at C15 that affect the charge distribution over the polyene chain of the chromophore.^{49,50,52} This would compromise the interpretation of any changes observed in the activated state, since in rhodopsin the Schiff base linking the chromophore to Lys₂₉₈ in the protein is protonated, but it is unprotonated in Meta II.^{1,5,6,19} Furthermore, we have shown before that the 12-F analog of rhodopsin has retained the ability to activate its cognate G-protein transducin following illumination.⁶⁴ In conclusion, a fluorine label at the C12 position is most suitable to probe the differences in its microenvironment in rhodopsin and the activated state, Meta II, and correlate this with the available crystal structures.

The isotropic chemical shift we measure for 12-F rhodopsin in a natural membrane environment (-95.0 ppm) agrees very well with the one reported for 12-F rhodopsin in a micellar solution (-94.2 ppm).^{53,94} This is not really surprising, since the chromophore is fixed quite tightly in the binding site of the protein and well shielded from solvating molecules.^{1,3,8,12,51,97,98} Consequently, the protein micro-environment around the chromophore will be largely unperturbed by the environment the protein is embedded in, as is also evident from the nearly identical spectral properties of rhodopsin in different solvation shells.^{1,99}

With respect to the identification of the photoproduct resonances, we have assigned the family of resonances at higher chemical shift to liberated all-*trans* retinal and the set at lower chemical shifts to Meta II. This is based upon the isotropic chemical shift of all-*trans* retinal in organic solvent (-120 ppm), as well as the lower anisotropy and asymmetry parameter together with a larger linewidth for the set at higher chemical shifts, which is also observed for the free 11-*cis* retinonoxime relative to the dark-state rhodopsin (Table 2) and likewise is expected for the liberated all-*trans* retinal relative to Meta II. In confirmation, the isotropic chemical shift of the free 11-*cis* and all-*trans* 12-F retinals in our lipid bilayer environment (-104.7 and -122.5 ppm, respectively) is close to the one observed in organic solvents (-107 ppm and -120 ppm, respectively).^{49,94}

Clearly, there is a large difference between the isotropic chemical shifts of the 11-*cis* and all-*trans* 12-F retinals in the free state, which can be traced back in all tensor elements (Table 2). Part of this difference is intrinsic to the isomeric state. For instance,

there is a 5–10 ppm shift towards higher isotropic chemical shifts of 10-F and 11-F retinal, when going from the all-*trans* to the 9-*cis* and 11-*cis* configuration, respectively.^{52,96} The additional positive shift observed for free 12-F retinal in the 11-*cis* configuration probably originates in intramolecular crowding.¹⁰⁰ The 11-*cis* isomer experiences steric interaction between the C10–H and C13–CH₃ elements, which in the free retinal is partially relieved by twisting the C12–C13 bond,^{101–103} bringing C12–F and the C13–methyl group in closer contact, which would increase the chemical shift.⁹⁴ We find a ± 2 ppm difference in the isotropic chemical of the 12-F retinals in a bilayer environment as compared to organic solvent. We would expect a shielding effect of the more hydrophobic lipid acyl chain environment in the bilayer, resulting in a small shift that indeed is observed for the all-*trans* configuration (about –2 ppm). However, the 11-*cis* isomer displays a small shift of about +2 ppm, instead. Currently, we attribute this to stabilization of a special conformation of the curved and twisted 11-*cis* retinal in the acyl chain solvation shell, possibly in combination with intermolecular crowding around the “protruding” C10–C13 structural element. In view of this, the 11-*cis* isomer in organic solvent probably represents a better model of the “free ligand” for estimating protein-induced ligation shifts.

The isotropic chemical shift of the 12-F label, we attribute to the activated state, Meta II, lies at –130.9 ppm, displaying an astonishingly large shift difference of –35.9 ppm from the ground state, rhodopsin. This is almost equally distributed over all principal components of the chemical shift (Table 2), where the differences in each component between the dark and Meta II state are significantly larger than the error of the individual tensor elements. The shift partly originates in the concomitant 11-*cis* to all-*trans* isomerization of the C11=C12 bond, as indicated above, but also stems from the large difference in ligation shift (Table 3).

Such a large shift reflects a significant change in the local environment of the 12-F label upon progressing from the ground state to Meta II. It would be very satisfying if we could link this large shift to the crystal structures of rhodopsin and Meta II.^{12,36} This quest is described in the Experimental section and the ESI.† Some relevant data are reported here. First, let us write down eqn (1) in terms of ¹⁹F isotropic fluorine chemical shifts:

$$\delta_{\text{iso}} = n\delta_{\text{iso}}^{\text{F-C}} + \delta^{\text{ref}} + n^2 A_{\text{iso}}^{\text{pol}} V^{\text{pol}} = -132.58 - 1017.74 V^{\text{pol}} \quad (2)$$

For an ideal F–C single bond we would obtain $n = 1.0$ and obtain the simple form in the second part of eqn (2). The charge distribution in the binding pocket around the 12-F–C moiety can therefore have two effects: (i) influencing the chemical shift by changing F–C bond length and (ii) causing chemical shift changes by bond polarization.

In our case both rhodopsin models have similar bond C–F bond length and therefore similar values of n : 0.955 for the Meta II state and 0.966 for the dark state. Obviously the value of around –130 ppm, representing the Meta II state, would be reached in situations with low polarization energy V^{pol} where the negative and positive terms in V^{pol} nearly cancel out (see eqn (2)). To reach –95.0 ppm, representing the ground state rhodopsin, the valence must be lower than 1.0—that means a

longer F–C bond—and/or the polarization energy must display large negative values. This will be further elaborated below in the discussion of the CS tensor.

Since most ¹³C chemical shifts of retinal in rhodopsin and Meta II have been measured and assigned, first a family of structures has been generated linking these shifts with crystal structures (see ESI†). The 20 best structures from this ¹³C CS refinement are not sufficient for a global minima search for the chromophore with the isotropic ¹⁹F chemical shifts of the 12-F as constraints but nevertheless some valuable conclusions could be drawn from these ¹⁹F CS calculations (Table 2). At first it can be observed that the trend of the decrease in isotropic chemical shift of the Meta II state is correctly reproduced by the calculations but the total shift was not correctly predicted. It is also noteworthy, that the calculations with a rhodopsin where the carbonyl group of Glu 181 is not protonated perform better than when it is protonated (Table 2).

In going from the dark state to the Meta II state the 33 component is most shifted to negative values (by –41.5 ppm) and the 22 component least shifted (by –31 ppm) and the span increases by 6.4 ppm. This trend is also present in the calculations but the total effect is smaller than in the experiment and the 22 component is shifted into the opposite direction. This is not surprising since the total isotropic shift of the dark state (–95.0 ppm) was not reached in our calculations. Since also the total span is correctly reproduced in our calculations we can conclude that the calculated CS tensor correctly reflects the situation in the two pigments except for an offset. What happens with the CS tensor when the conformation of the retinal is changing is best illustrated in the synthetic powder patterns in Fig. 6. The powder patterns based on the experimentally determined principal

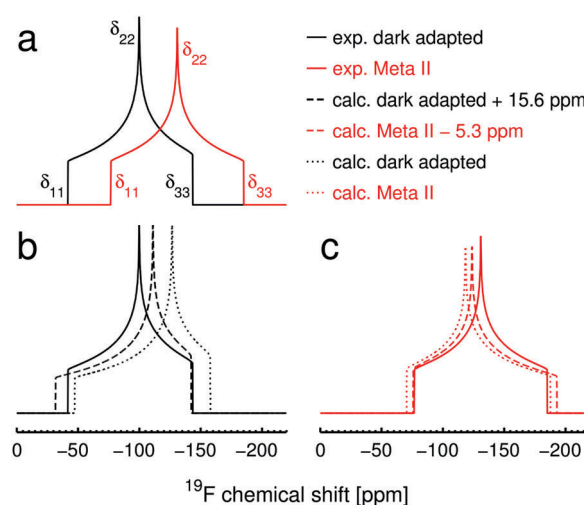


Fig. 6 Synthetic ¹⁹F powder patterns for the 12-F retinal pigments: black: dark adapted (rhodopsin) and red: light adapted (Meta II). The principal tensor components are indicated according to Table 1. (a) Comparison of experimental powder patterns for the dark adapted and Meta II state. (b) Comparison of the powder patterns for the dark adapted state based on the experimental and calculated principal components of the chemical shift tensor. (c) Comparison of the powder patterns for Meta II based on the experimental and calculated principal components.

components of the CS tensor for the dark adapted and Meta II state are directly compared in Fig. 6a. Furthermore, Fig. 6b and c compare the powder patterns based on experimentally determined and calculated principal components. We note that if we adjust all calculated principal components by +15.6 ppm for the dark adapted and −5.3 ppm for the Meta II state, we obtain excellent agreement between calculated and experimental powder patterns. See Tables S3 and S4 in the ESI† for details. This indicates that the error in the calculation of the absolute sizes of the principal components is larger than the resulting error in the chemical shift anisotropies that are based on relative sizes. The structural change in the retinal-protein complex is accompanied by a typical change from a negative skew in rhodopsin to a very symmetric powder pattern with skew near zero for the Meta II state. To explain the changes in tensor components it is useful to write down eqn (2) in terms of tensor components (also see Table 1):

$$\begin{pmatrix} \delta_{11}^{\text{F}} \\ \delta_{22}^{\text{F}} \\ \delta_{33}^{\text{F}} \end{pmatrix} = n \begin{pmatrix} -76.98 \\ -117.18 \\ -203.5 \end{pmatrix} + n^2 \begin{pmatrix} -1393.1 \\ 450.8 \\ -2110.8 \end{pmatrix} V^{\text{pol}} \quad (3)$$

The principal tensor components of $\delta^{\text{F-C}}$ and A^{pol} in eqn (3) can be written as column vectors since the tensor are diagonal within the framework of BPT. Both components δ_{11} and δ_{33} (perpendicular to the bond direction, see Fig. 7) change in such a way that the span of the tensor stays nearly constant (101.7 ppm for rhodopsin and 108.1 ppm for Meta II; Table 2). The higher span of the F bond tensor (126.5 ppm, see Table 1) indicates that the F–C bond is slightly longer ($n < 1.0$) for the retinal in the opsin binding pocket. This can readily derived

from eqn (3) since the configuration weight n acts on all tensor components in the same way.

The change from a very symmetric powder pattern in the case of the Meta II form (skew $\kappa = 0.00$) to a more axial pattern in the case of the dark state (skew $\kappa = -0.15$) can be explained from the different sign of the 22 component of the polarization tensor A^{pol} . That means that the 22 component (in bond direction) is shifted into the opposite field direction with the respect to 11 and 33 by bond polarization (V^{pol} see eqn (3)). We can conclude that the more axial ^{19}F powder pattern of the dark adapted state is caused by large negative polarization energy. In accordance with eqn (3) the largest shift with polarization is calculated for the 33 component because of the large 33 component of A^{pol} .

The orientation of the ^{19}F tensor with respect to the molecular framework could not be extracted from our experiments but for oriented measurements⁴⁶ this information is of central interest. As indicated in Fig. 7 the BPT calculations allow assigning the tensor directly to the F–C bond coordinate system. The 22 component of the tensor is in F–C bond direction, the 11 component in the π plane of the sp^2 carbon 12 and the 33 component perpendicular to it. With this information we can assign the principal tensor components to the bond coordinate system that is used for the definition of C–F bond CS tensor and the CS polarization tensor A^{pol} : the 22 principal component is in z direction, the 11 component in x direction and finally the 33 component in y direction.

The ligation shift we calculate for rhodopsin, *i.e.* the protein induced shift in the chemical shift of its ligand, is a positive shift of +10 to +12 ppm, depending on which reference system is used for the free 11-*cis* ligand. This shift compares well with the *ca.* +13 ppm reported before for rhodopsin solubilized in a micellar environment.^{53,94} We here show that this ligation shift mainly originates in large positive shifts in the δ_{11} and δ_{22} tensor elements (Table 2). The ligation shift in rhodopsin has been attributed to steric crowding in the ligand binding site around C12 of the retinylidene moiety.⁵³ As discussed above, our data indicate that major effects of the binding pocket are elongation and polarization of the 12-F–C bond. Fig. 8 shows the functional groups of the amino acids surrounding the retinal in the binding pocket that are in closest contact to the 12-F site. For the chemical shift optimized structure of the dark-adapted state shown in Fig. 8a we find close contacts of 2.31 and 2.4 Å to the carbonyl groups of Gly-114 and Cys-187, respectively. The oxygen of the closest water molecule is with a distance of 4.21 Å from the 12-F relatively far away.^{11,12,92,97} The bond polarization would agree with the strong C–F stretch vibration we observed in 12-F rhodopsin.⁶⁴

The huge difference in isotropic chemical shift of −35.9 ppm, we observe for the activated state, Meta II, relative to the ground state, rhodopsin, has two causes: (i) the 11-*cis* to all-*trans* isomerization shift and (ii) the fact that the ligation shift of the all-*trans* retinal is −8.4 ppm and hence in the opposite direction compared to the 11-*cis* retinal for which it is +9.7 ppm. The ligation shift in Meta II entails a chemical shift decrease, and of a magnitude that is quite unusual upon binding in the

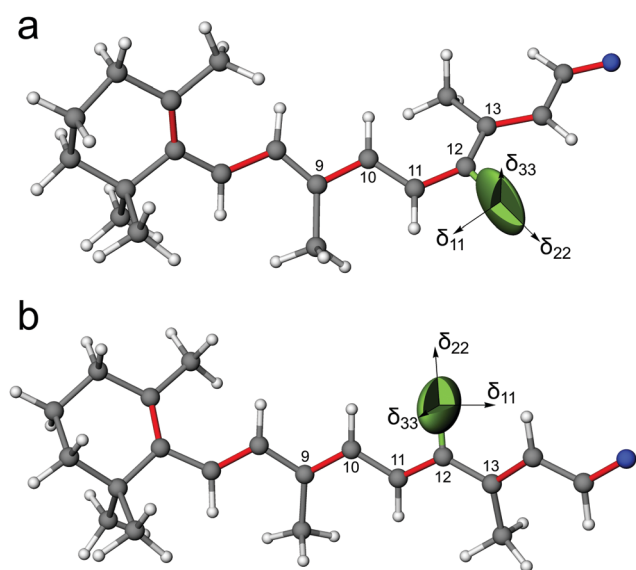


Fig. 7 Orientation of calculated ^{19}F chemical shift tensor in the 12-F pigments: (a) the dark adapted state with the 11-*cis* chromophore (rhodopsin). (b) The photo-activated state with all-*trans* chromophore (Meta II). The blue sphere represents the ϵ -amino group of the protein lysine residue with which the retinal is connected. See also Fig. 8.

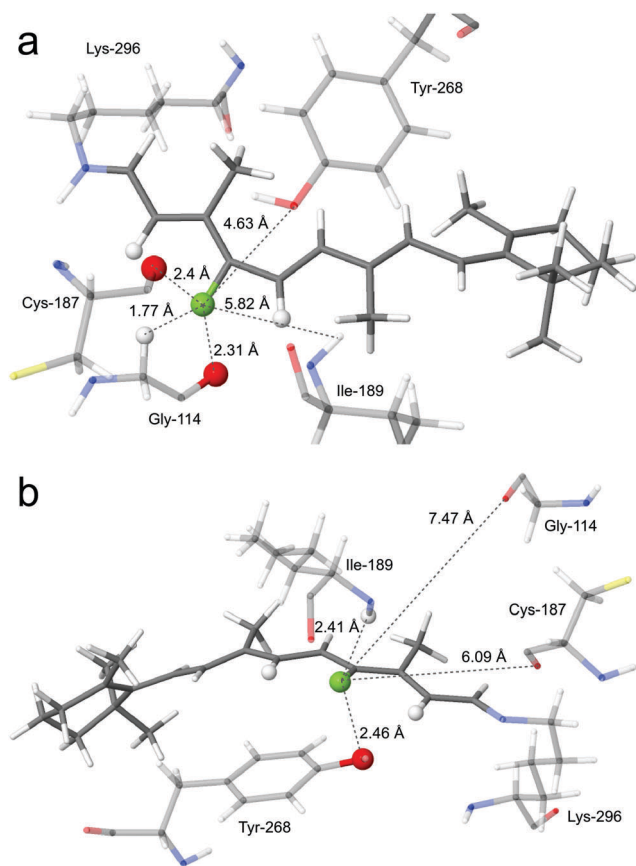


Fig. 8 Amino acids in the surrounding of the retinal in the binding pocket in the dark adapted (a) and Meta II form (b) of rhodopsin in chemical shift optimized structures. Contacts of the 12-F (shown in green) mainly responsible for the ^{19}F chemical shift are indicated. In the dark state these are contacts to the oxygen atoms of carbonyl groups of Cys-187 and Gly-144. In the Meta II form these are contacts to the hydrogen atom of the NH group of Ile-189 and the oxygen atom of the hydroxyl group of Tyr-268.

protein interior.⁴⁴ In addition, the ligation shift in Meta II is most prominently visible in a very large negative shift in the δ_{33} tensor element (Table 3 and eqn (3)). The BPT calculations provide an explanation together with the chemical shift optimized structure of Meta II shown in Fig. 8b: after the 11-*cis* to all-*trans* transformation of the retinal the 12-F-C bond points in the opposite direction in the binding pocket, hence replacing the close contacts to two carbonyl groups in the dark-state with one close contact to the hydroxyl group of Tyr-268 and one to the amide group of Ile 189 in the Meta II state with distances given by 2.48 Å and 2.41 Å, respectively. The effect of those closest functional groups on the bond polarization cancel each other to a large degree, resulting in a significant decrease of the chemical shift. The slight decrease in bond C-F length and the strong effect on bond polarization, would also explain the slight change in frequency and strong decrease in intensity observed for the C-F stretch vibration.⁴⁶ We conclude that the interactions of the retinal with the surrounding binding pocket is strongly reduced in the Meta II state. This is functional, of course, since Meta II has the purpose to mobilize protein structural elements to allow binding of the G protein, and the retinal is about to be released.

Conclusions

Using a ^{19}F label, strategically positioned at C-12 of retinal, the corresponding chemical shift and tensor elements were determined in ground state bovine rhodopsin, its active state photo-intermediate Meta II, and free retinal in a lipid environment. Opposite chemical shifts (+9.7 and −8.4 ppm) were obtained upon binding of the chromophore to obtain rhodopsin and Meta II, respectively. These are due to large changes in the δ_{11} and δ_{22} and in the δ_{33} tensor elements, respectively. An unusually large shift (35.9 ppm) was observed between rhodopsin and Meta II, to which all tensor elements contribute. This shift partly originates in the light-triggered 11-*cis* to all-*trans* isomerization of the chromophore. The other part reflects the local conformational rearrangements in the chromophore and the binding pocket. We set out to correlate these data with the available X-ray structures of rhodopsin and Meta II. However, for this purpose hydrogen atoms have to be inserted and hereto, based upon the X-ray structures, a family of structures was derived that best correlated with the ^{13}C chemical shifts of the chromophores reported in the literature. Following up on these structures, a 12-F derivative was obtained that best corresponds with the ^{19}F chemical shifts and tensor elements that were determined experimentally. The combined data indicate strong changes in the local environment of the C-12 position and a profound change in the interaction pattern of the chromophore with the protein during the transition from rhodopsin to Meta II. We propose, that using alternative ^{19}F -labels on the chromophore as well as selective fluorination of protein residues, in combination with progress in ^{13}C -labeling, will eventually lead to comprehensive structural models and detailed insight in the activation mechanism of visual rhodopsins.

Conflicts of interest

There are no conflicts to declare.

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