

The Mid-Infrared Instrument for the *James Webb Space Telescope*, VI: The Medium Resolution Spectrometer

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ABSTRACT. We describe the design and performance of the Medium Resolution Spectrometer (MRS) for the *JWST*-MIRI instrument. The MRS incorporates four coaxial spectral channels in a compact opto-mechanical layout that generates spectral images over fields of view up to $7.7 \times 7.7''$ in extent and at spectral resolving powers ranging from 1300 to 3700. Each channel includes an all-reflective integral field unit (IFU): an “image slicer” that reformats the input field for presentation to a grating spectrometer. Two 1024×1024 focal plane detector arrays record the output spectral images with an instantaneous spectral coverage of approximately one third of the full wavelength range of each channel. The full $5\text{--}28.5\ \mu\text{m}$ spectrum is then obtained by making three exposures using gratings and pass-band-determining filters that are selected using just two three-position mechanisms. The expected on-orbit optical performance is presented, based on testing of the MIRI Flight Model and including spectral and spatial coverage and resolution. The point spread function of the reconstructed images is shown to be diffraction limited and the optical transmission is shown to be consistent with the design expectations.

Online material: color figures

1. INTRODUCTION

The rationale for, capabilities of, and scientific context of the mid-infrared instrument (MIRI) on *JWST* are described in Rieke et al. (2015; hereafter Paper I) and the overall optical, thermal,

mechanical/structural electronic and control aspects of the design are summarized in Wright et al. (2015; hereafter Paper II). This paper describes in more detail the MIRI Medium Resolution Spectrometer (MRS), which is an Integral Field Unit (IFU)

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TABLE 1
SPATIAL AND SPECTRAL PARAMETERS FOR THE FOUR MRS CHANNELS

		Channel	1	2	3	4
Slice width		arcseconds	0.176	0.277	0.387	0.645
Number of slices		—	21	17	16	12
	λ_{short}^a	Pixels	1.405	1.452	1.629	2.253
Slice width at detector						
	λ_{long}^b	Pixels	1.791	1.821	2.043	2.824
Pixel size along slice		Arcseconds/pixel	0.196	0.196	0.245	0.273
FOV (across \times along slices)		Arcseconds	3.70×3.70	4.71×4.52	6.19×6.14	7.74×7.95
Sub-band A						
Wavelength range	$\lambda_{\text{min}} - \lambda_{\text{max}}^{c,d}$	μm	4.87–5.82	7.45–8.90	11.47–13.67	17.54–21.10
Resolution		$\lambda/\Delta\lambda$	3320–3710	2990–3110	2530–2880	1460–1930
Sub-band B						
Wavelength range	$\lambda_{\text{min}} - \lambda_{\text{max}}$	μm	5.62–6.73	8.61–10.28	13.25–15.80	20.44–24.72
Resolution		$\lambda/\Delta\lambda$	3190–3750	2750–3170	1790–2640	1680–1770
Sub-band C						
Wavelength range	$\lambda_{\text{min}} - \lambda_{\text{max}}$	μm	6.49–7.76	9.91–11.87	15.30–18.24	23.84–28.82
Resolution		$\lambda/\Delta\lambda$	3100–3610	2860–3300	1980–2790	1630–1330

^a λ_{short} is the shortest wavelength of operation for the corresponding channel.

^b λ_{long} is the longest wavelength of operation for the corresponding channel.

^c λ_{min} is the shortest wavelength of operation for the corresponding sub-band within a channel.

^d λ_{max} is the longest wavelength of operation for the corresponding sub-band within a channel.

spectrometer that covers the wavelength range from 5 to 28.5 μm at spectral resolving powers of a few thousand. The MRS consists of four channels that have coaligned fields-of-view, simultaneously observing the wavelength ranges listed in Table 1 with individually optimized IFUs, collimators and gratings.

§ 2 of this paper provides a description of the optical design, including the rationale for choosing the IFU concept and its impact on how observations are carried out. § 3 then describes the expected on-orbit optical performance of the MRS as measured during cryogenic testing, and with a description of the procedure used to construct calibrated spectral data cubes from the raw measured images. The impact of particular characteristics of the MRS, including spectral fringing and stray light are also discussed here.

2. OPTICAL DESIGN

2.1. Rationale

An IFU based design was preferred to a long-slit design for the MIRI Spectrometer for the following reasons.

First, for point source observations, the need to center the source in a narrow slit (via a peak-up procedure) is relaxed, simplifying and accelerating the target acquisition procedure. There is an additional benefit that there is no loss of light at the slices (“slit losses” in a conventional long-slit spectrometer). The MRS slice width is set to be less than or equal to the FWHM of the diffraction limited point spread function at the slicing mirror. An equivalent long-slit spectrometer would vignette the light outside this region, losing about 50% of the total. The amount of lost light could be reduced by making the slit wider,

but doing so would in turn reduce the spatial and spectral resolution and could also decrease the signal-to-noise ratio because of increased background radiation.

Second, from a scientific perspective, the wavelength range covered by the MIRI spectrometer is sufficiently broad (a factor of nearly six) that different emission mechanisms may dominate in different regions of the spectrum. In cases where these mechanisms do not share a common center (e.g., stellar output compared with infrared re-emission in a starburst galaxy), a simple slit spectrograph poses a dilemma in placing the “source” on the slit. An IFU implements three-dimensional (3D) spectroscopy, which solves this problem by giving accurately registered spatially resolved spectroscopy over the entire field.

These considerations, combined with the wavelength coverage and resolving power requirements defined by the *JWST* mission science goals, and the mass, volume, and electrical power limitations set by the *JWST* spacecraft environment, have resulted in a system with the characteristics summarized below and illustrated in block diagram form in Figure 1.

As shown in the figure, the full 5–28.5 μm wavelength range is divided within the Spectrometer Pre-Optics (SPO) into four simultaneous spectral channels (Kroes et al. 2005; Kroes et al. 2010), each with its own IFU (Lee et al. 2006) and using a simple scheme of pass-band separation by dichroic filters whose design is described in Hawkins et al. (2007) and Wells et al. (2004). We denote the short and long wavelength limits of each channel as λ_{short} and λ_{long} . Each channel serves an optimized spectrometer. This separation provides several benefits: (1) it enables the use of diffraction gratings in first order, allowing each to be used near peak efficiency around the blaze wavelength, and

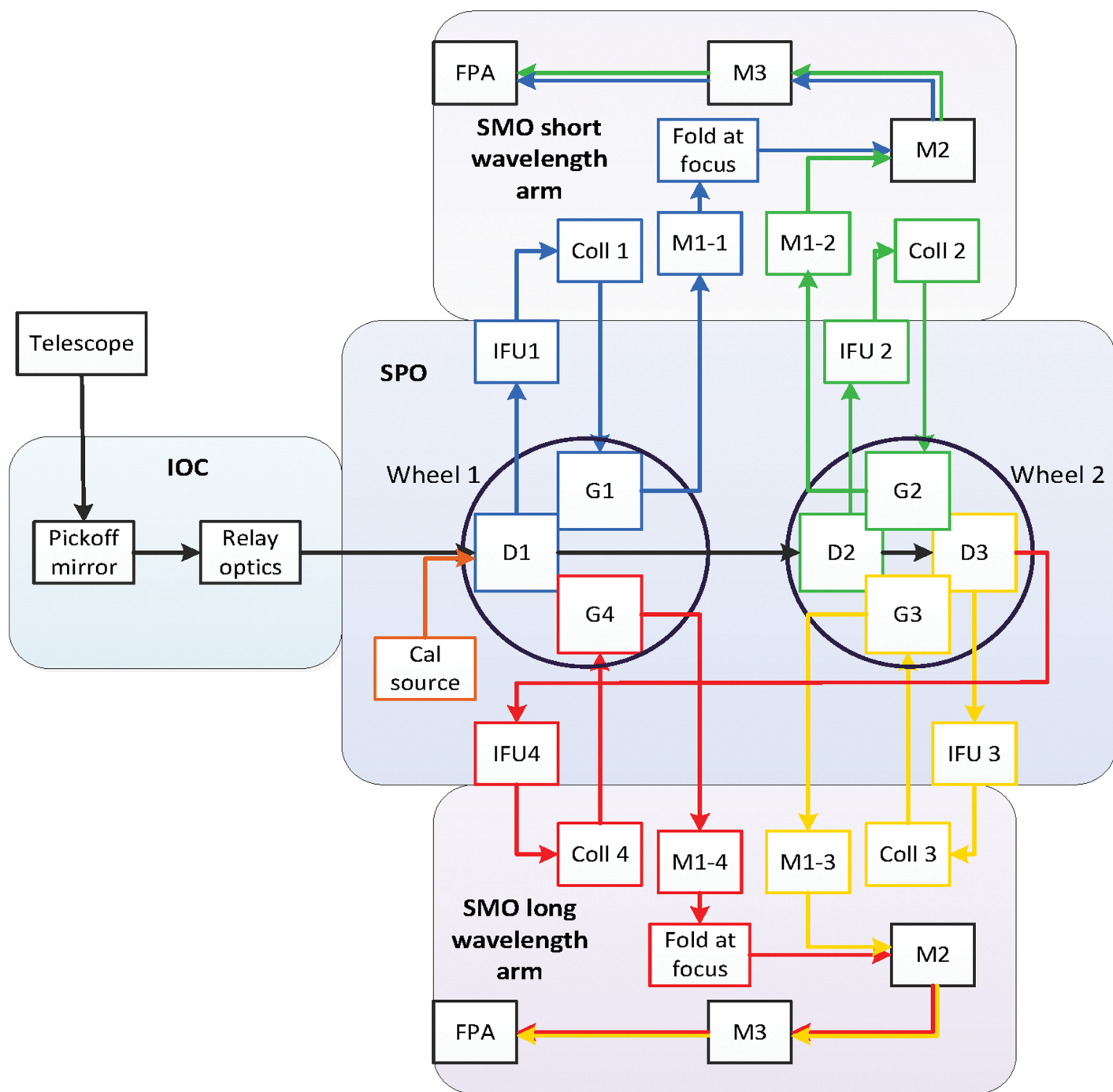


FIG. 1.—Block diagram of the MRS, showing the main optical functions. A detailed description can be found in the text. See the electronic edition of the *PASP* for a color version of this figure.

(2) it also allows the IFU slice widths to be tailored to the wavelength-scaled FWHM in each channel.

Electrical and thermal constraints limit the MRS to two 1024×1024 Si:As detector arrays (Ressler et al. 2015, hereafter Paper VIII), resulting in the spectrometer being split into two sets of optics, each with its own detector array, one for the two

short-wave channels and one for the two long-wave channels. In each case, the spectra of two wavelength channels are imaged simultaneously onto the left and right halves of a detector array.

The product of spatial and spectral coverage that can be achieved in a single instantaneous exposure is ultimately set by the number of detector pixels. For the MRS, this results

in a single exposure providing a spectral sub-band that covers only one third of each channel. Full wavelength coverage then requires three exposures, with a pair of mechanisms being used to select the gratings and dichroics in each case. We refer to the short and long wavelength limits of each sub-band as λ_{\min} and λ_{\max} . This design choice allows the grating performance to be optimized over a narrow wavelength range ($\lambda_{\max}/\lambda_{\min} \sim 1.2$). The gratings only have to work over 20% of the 1st order near the blaze wavelength and the dichroics do not have to be used near the cross over between reflection and transmission.

We will now step through the optical train as shown in Figure 1: the Input Optics and Calibration (IOC) modules pick off the MRS FOV from the *JWST* focal surface and pass it on to the SPO. The SPO spectrally splits the light into the four spectrometer channels and spatially reformats the rectangular fields of view into slits at the entrance of the Spectrometer Main Optics (SMO). The SMO comprises fixed optics that collimate the light for presentation to diffraction gratings mounted in the SPO and then reimages the dispersed spectra onto the two focal plane arrays. Key features of this optical system are described individually in the following sections. The spectral and spatial coverage of the MRS is summarized in Table 1.

2.2. Telescope Focal Plane to Integral Field Units

The field-of-view of the MRS is adjacent to the MIRI imager field and picked off from the *JWST* focal surface using the MIRI Pickoff Mirror (POM), which is common to both (see Paper II). The sky and the *JWST* pupil are reimaged by the IOC so that there is a pupil image at the SPO input and a sky image further on. A cold stop that is 5% oversized with respect to the *JWST* pupil is placed at the SPO input pupil for stray light control.

The light is directed toward the first dichroic filter via a fold mirror placed 10 mm beyond the pupil. The focal plane is

formed 535 mm beyond the pupil, providing the long path length and narrow beam waist needed between the pupil and the inputs of the four Integral Field Units (IFUs), for mounting a chain of dichroics to divide the light among the four spectral channels.

A hole in the fold mirror, sized to be smaller than the footprint of the telescope central obscuration, acts as an aperture for the injection of light from the on-board Spectrometer Calibration Unit (SCU), shown in Figure 2. The SCU provides spatially and spectrally uniform blackbody illumination for flux calibration and pixel flat fielding functions, using as its light source a tungsten filament heated to a temperature of ~ 1000 K by the application of an 8 mA drive current. The filament is mounted inside a nonimaging flux concentrator that generates spatially uniform focal plane illumination at the exit port of a 25 mm diameter reflective hemisphere (described in Glasse et al. [2006]). Two cadmium telluride lenses within the SCU then reimage the exit port onto the IFU input focal planes. The lenses provide a pupil image that coincides with the hole in the input fold mirror. This hole is positioned to lie within the footprint of the central obscuration of the *JWST* primary mirror and so has no impact on the science beam. In this way, the SCU can provide flood illumination of the full MRS field of view without any need for mechanisms or additional optical elements.

The overall layout of the dichroic assembly is shown in Figure 3, with the input fold mirror and SCU situated at the left-hand end (but not shown). The three dichroics needed to divide the spectral band among the four spectrometer channels for one of the three sub-bands are indicated as D1, D2, and D3 in Figures 1 and 3. Taking Sub-band A as an example, the required reflective band for dichroic D1 is the wavelength range for Sub-band A in Channel 1, while its transmission band needs to extend from the short end of Sub-band A in Channel 2 to the long end of Sub-band A in Channel 4. The bands are listed in

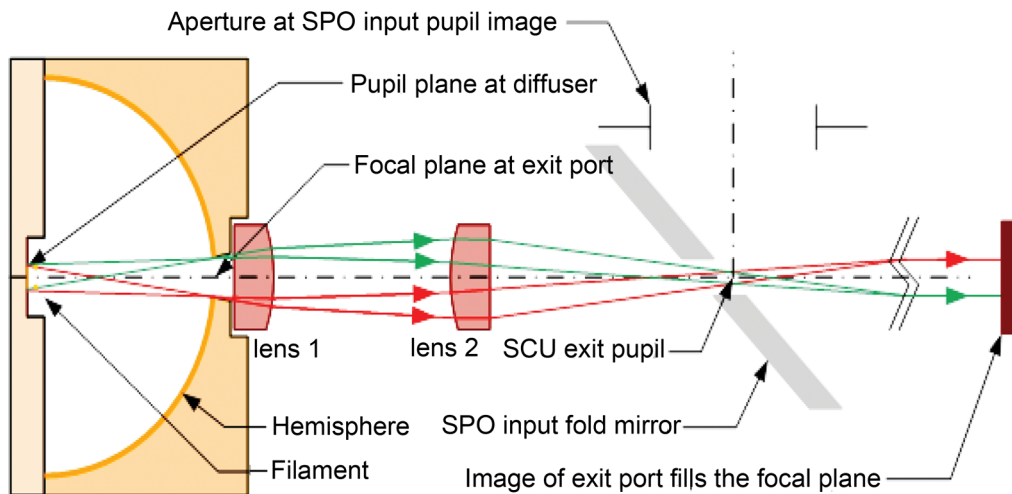


FIG. 2.—Optical Schematic (not to scale) of the Spectrometer Calibration Unit (SCU). Calibration light is injected via the obscuration at the center of the *JWST* pupil where it is imaged at the input fold mirror of the Spectrometer Pre-Optics (SPO). See the electronic edition of the *PASP* for a color version of this figure.

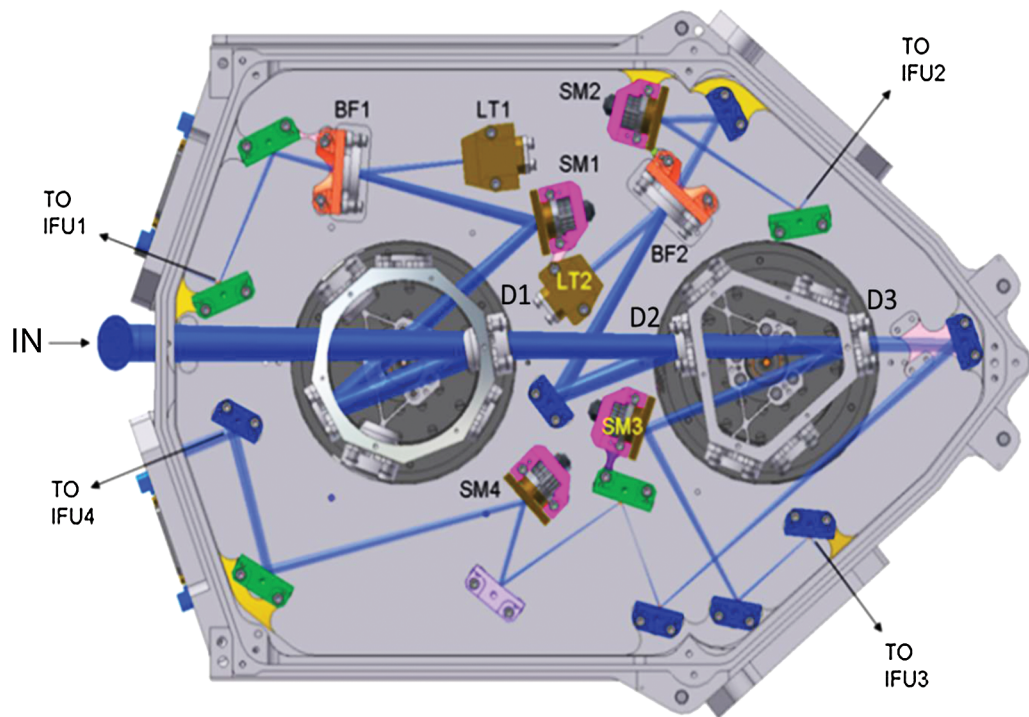


FIG. 3.—Layout of the Dichroics and Fold Mirrors for all channels in the Spectrometer Pre-Optics (SPO). The position of the input pupil and fold mirror are labeled IN. The locations of blocking filters (BF), light traps (LT), powered mirrors (SM), and dichroics (D) are shown. After the light has been divided into the appropriate spectral ranges, it is output to the integral field units (IFU) for the four spectrometer channels. See the electronic edition of the *PASP* for a color version of this figure.

Table 2 for all nine dichroics. In all cases, the mean reflectivity is above 0.95 and the mean transmission is above 0.74.

All of the gratings work in first diffractive order, so additional blocking is needed to reject second and higher orders. Because the dichroics work in series, it is possible to use the combined blocking of dichroics 1 and 2 to remove the need for blocking filters in Channels 3 and 4. For Channels 1 and 2, dedicated blocking is provided by the fixed filters shown as BF1 and BF2 in Figure 3, with light traps LT1 and LT2 absorbing unwanted reflections.

TABLE 2 REFLECTIVE AND TRANSMISSIVE BANDS OF THE MIRI DICHOIC BEAM-SPLITTERS		
Identification	Reflection band (μm) $R \geq 94\%$	Transmission band (μm) $T \geq 74\%$
Dichroic 1A	4.84–5.83	7.40–21.22
Dichroic 1B	5.59–6.73	8.55–24.73
Dichroic 1C	6.45–7.77	9.87–28.5
Dichroic 2A	7.40–8.91	11.39–21.22
Dichroic 2B	8.55–10.29	13.16–24.73
Dichroic 2C	9.87–11.88	15.20–28.5
Dichroic 3A	11.39–13.68	17.45–21.22
Dichroic 3B	13.16–15.80	20.34–24.73
Dichroic 3C	15.20–18.25	23.72–28.5

The path length required to reach the input of Channel 4 is greater than the 535 mm discussed above, so the light transmitted by the final dichroic is reimaged via an intermediate focal plane to the entrance of the Channel 4 IFU.

The 21 mm diameter dichroic filters are mounted on two wheels. First is the nine-sided wheel A containing the three Channel 1 dichroic filters and three flat mirrors to direct the light toward Channel 1. The second, six-sided wheel, contains the six dichroics needed to divert the light into Channels 2 and 3. Each dichroic filter is mounted onto a diamond machined facet on the wheel that provides the required alignment accuracy and reduces the magnitude of any print-through of surface form errors from the wheel to the filter substrate. The filters are held in place with a spring loaded bezel with a clear aperture of 17 mm compared with the coated area of 17.4 mm. The bezel prevents light from reaching the uncoated area of the filter, which might result in unfiltered light entering the optical path. The blocking filters are mounted directly to the SPO chassis with a stand and bezel, similar to the mounting arrangement for the wheels. The mechanisms that carry the dichroic wheels also carry the corresponding wheels with diffraction gratings for each sub-band (three per channel), as discussed in § 2.4. By arranging for the gratings to be mounted on the same mechanisms as their band-selecting dichroics, the number of moving mechanisms within the MRS is kept to two.

2.3. The Integral Field Units

Before the light in the four channels reaches the gratings it is sliced and reformatted by the Integral Field Units and so we describe these next. The parameters pertaining to the spatial and spectral coverage and sampling of the spectrometer, which are largely determined by the design of the four IFUs, are given in Table 1. Spatially, the image is sampled in the dispersion direction by the IFU slicing mirrors and in the slice direction by the detector pixels. Spectrally, the width of the slices defines the spectrometer entrance slit and the width of the image (in pixels) of the slice at the detector defines the width of the spectral sample. Inspection of Table 1 then shows that Channels 1–3 are all slightly undersampled spectrally, with slice widths less than two pixels at the detector.

The bandwidths of the spectral channels listed in Table 1 were chosen such that the ratio $\lambda_{\text{long}}/\lambda_{\text{short}}$ in all four channels was the same: $\lambda_{\text{long}}/\lambda_{\text{short}} = (28.3/5)^{1/4} = 1.54$. This means that, if each channel is assigned the same number of spectral pixels, the resolving power in each will be approximately the same. Each sub-band then has a width that is slightly larger than one-third of the bandwidth of its parent channel, where the excess provides sufficient overlap of adjacent spectra to allow their concatenation. In practice, the overlap is typically 10–15% of the spectral range, depending on the specific sub-bands and position in the field-of-view.

For full Nyquist sampling of the telescope PSF, the four IFU slice widths should ideally be matched to the half width at half maximum intensity of the PSF at the shortest wavelength in each channel. For the *JWST* telescope pupil, this would equate to $a_{\text{diff}} = 0.088 (\lambda_{\text{short}}/5 \mu\text{m})''$. However, the large increase in the beam size produced by diffraction at the IFU slicer mirror would then necessitate large apertures for the spectrometer optics to avoid vignetting. To control the size, mass and optical aberrations of the spectrometer, the starting point for determining the slice widths was therefore to set them equal to $2a_{\text{diff}}$, with fully sampled PSFs to be achieved using two or more pointings of the telescope.

Optimum spatial sampling, which minimizes the number of telescope pointings needed to fully sample all four channels simultaneously in the across-slice direction, is achieved by having a single pointing offset equal to $n_i + 1/2$ slice widths for all channels where n_i is a different integer for each channel. This led us to adopt a set of slice widths that follow a scheme where the slices are factors of 1, 11/7, 11/5, and 11/3 wider than the narrowest 0.176'' slice. These values closely match the increasing FWHM of the PSF in four channels and allow full Nyquist sampling to be achieved using a single “diagonal” offset whose magnitude in the across slice direction is equal to 11/2 times the width of the narrowest slice (0.167'') and $(N + 1/2)$ detector pixels in the along slice direction (where N is an integer). With these slice widths fixed the design was integrated to, as far as possible, meet the following criteria: (1) each channel would occupy half a 1024×1024 pixel detector with borders to allow

for tolerances in alignment and gaps between slices to avoid crosstalk; (2) the FOV of each channel should be as large as possible and approximately square; and (3) the long and short wavelength arms of the spectrometer should be identical (in practice one is the mirror image of the other). This iteration resulted in the parameters in Table 1.

The IFU design was developed from the IFU deployed in *UIST*, a 1–5 μm spectrometer for *UKIRT* (Ramsay-Howat et al. 2006). Other examples of all-reflecting image slicers are the *S*pectrometer for *I*nfrared *F*aint-*F*ield *I*maging (*SPIFFI*) (Tecza et al. 2000), and the slicer in the *G*emini *N*ear-*I*nfrared *S*pectrograph (*GNIRS*) (Dubbeldam et al. 2000). The *MIRI* (and *UIST*) IFU design allows excellent control over stray light by providing through-apertures for baffling. It consists of an entrance pupil, an input fold mirror, an image slicer mirror, a mask carrying exit pupils for the individual sliced images, a mask carrying slitlets for the individual images, and an array of reimaging mirrors behind the slitlets. These components (except for the last) are shown in Figure 4. The design is all-reflecting and constructed entirely of aluminum, and hence is well-suited for operation in the infrared and at cryogenic temperatures, as discussed in Paper II. Todd et al. (2003) and Lee et al. (2006) describe the manufacture of the slicer.

The optical path through the IFUs begins with the four toroidal mirrors, which comprise the anamorphic preoptics (*APO*) module (not shown in Fig. 4). The *APO* reimages an area of up to 8 by 8'' of the input focal plane onto the image slicer mirror, with anamorphic magnification. The image slicer (at the relayed image plane) consists of a stack of thin mirrors angled to divide the image along the dispersion direction of the spectrometer into separate beams. In the across slice (dispersion) direction one slice width is matched to the FWHM of the *JWST* PSF pattern at the shortest operating wavelength for the IFU. Along the slice direction, the magnification is chosen to provide the required

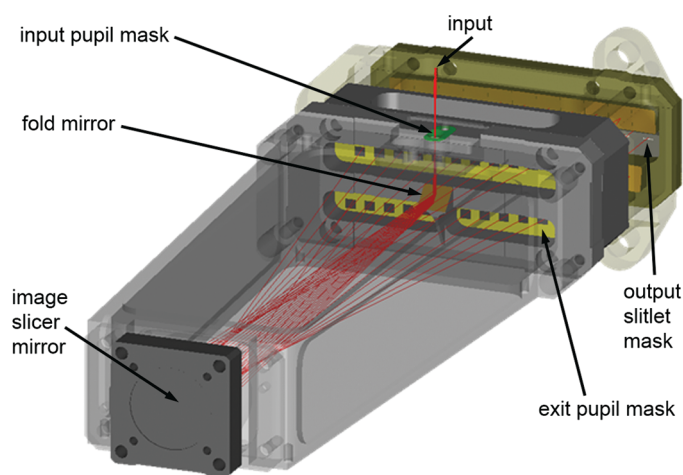


FIG. 4.—Three-dimensional view of the Channel 3 IFU. See the electronic edition of the *PASP* for a color version of this figure.

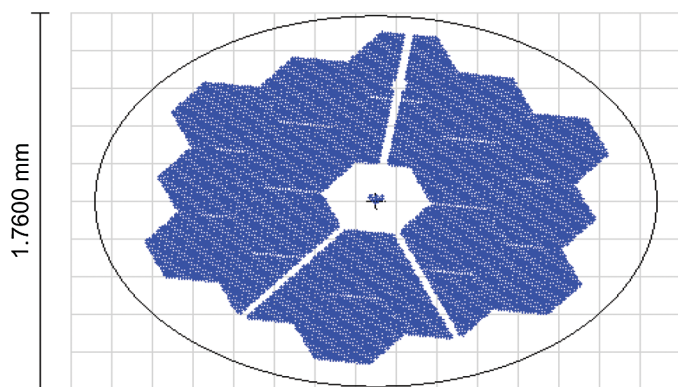


FIG. 5.—Rotated and anamorphically magnified pupil image at the entrance to the IFU of Channel 3, plotted on a square grid. See the electronic edition of the *PASP* for a color version of this figure.

plate scale at the detector. These two magnifications are not the same and so the APO exit pupil is elliptical, as illustrated in Figure 5 for Channel 3, where the footprint of the *JWST* pupil is shown in blue. The light exits the IFU through individual pupil masks for each beam, then through individual slitlets (see Fig. 4). Reimaging mirrors behind the slitlets relay the beam to the input of the appropriate spectrometer. That is, each IFU takes the rectangular FOV and transforms it into a pattern of slitlets (as shown in Fig. 4) that are reimaged onto the entrance aperture for its corresponding grating spectrometer.

The slicing mirror comprises a number of spherical optical surfaces, diamond turned onto a common substrate. This manufacturing approach was used previously for the IFU in GNIRS (Dubbeldam et al. 2000). Figure 6 is a photograph of the image slicer for Channel 1, which has 21 slicer mirrors arranged about the center of the component with 11 slices on one side and 10 on the other. Each mirror is offset to direct its output beam toward the corresponding reimaging mirror. The design for channels 2–4 is similar but with correspondingly fewer and larger slices as the wavelength increases, as listed in Table 1. The mirrors are arranged in a staircase-like manner to aid manufacturing. Each of the slices reimages the IFU entrance pupil at a scale of $\approx 1 : 1$ and because the centers of curvature of the mirrors are offset laterally with respect to each other, the pupils that the slices produce are separated, as can be seen in Figure 4.

The slicing mirrors truncate the anamorphically magnified image of the sky along the dispersion direction, as shown in Figures 7a–7c. This truncation results in the pupil image being diffraction broadened in the dispersion direction as shown in Figure 7d. To reduce vignetting, the baffle at this pupil image and subsequent optical components (including the diffraction gratings) are oversized. This is illustrated in Figure 7d where the modeled diffraction footprints for three field positions along a slice are overlaid to appear as the blurred green patch in the image. The nominal geometric pupil is shown in red and the oversized mask aperture is shown as a green rectangle.

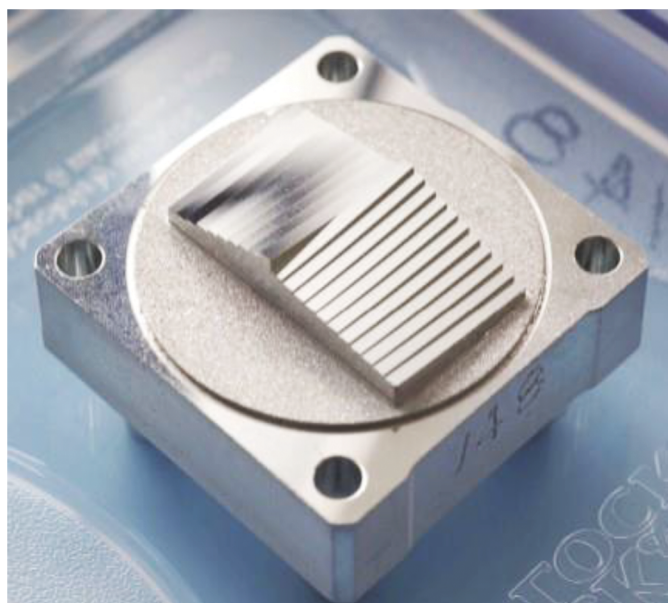


FIG. 6.—Photograph of the Channel 1 Image Slicer Mirror (Courtesy: Cranfield University). The slices are 1 mm wide and 12 mm long. See the electronic edition of the *PASP* for a color version of this figure.

The reimaging mirror array forms images of the slicing mirrors as input slits of the spectrometer and also images the pupils at infinity, making the output telecentric. Finally, a roof mirror redirects light from the reimaging mirrors so that the output

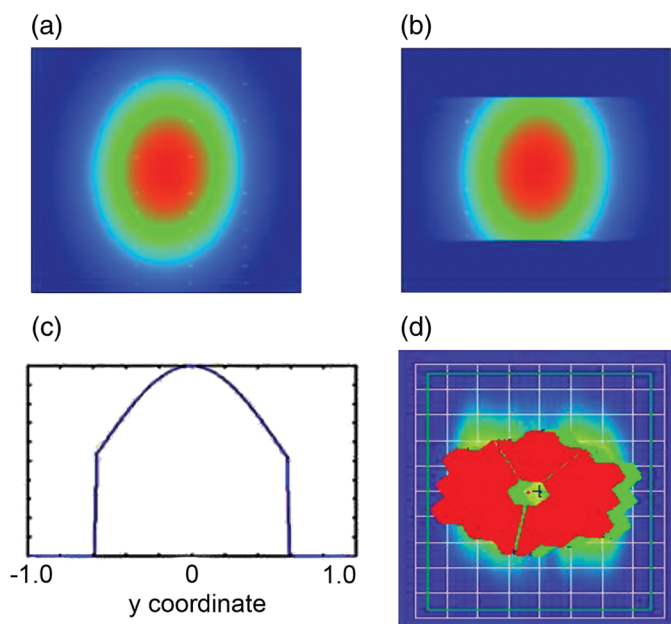


FIG. 7.—(a, b): Input and truncated images on the imager slicer mirror. (c): Cross section through the peak of the truncated image. (d): Diffraction broadened pupil images (see text for details). See the electronic edition of the *PASP* for a color version of this figure.

consists of the two lines of staggered slitlets relayed from those seen in Figure 4. A slotted mask at the position of the IFU output slits defines the field of view along the individual slices, while in the spectral direction the slots are oversized to ensure they do not vignette the image, but still act as baffles to reduce scattered light between the slices.

The numbers of detector columns covered by the four IFUs are approximately 476, 464, 485, and 434 for channels 1–4, respectively. The size of the gap between individual slits for channel 1 is set to be approximately equal to the diameter of the first dark diffraction ring in the telescope PSF (about 4 pixels). For channels 2–4 the size of the gap is slightly greater than the diameter of the first dark ring.

2.4. The Spectrometer Main Optics (SMO)

The SMO (whose layout is shown in Fig. 8) comprises four grating spectrometers in two arms. The development and test of these spectrometers are described by Kroes et al. (2005) and Kroes et al. (2010). Each of them performs three functions: collimation of the telecentric output beams of one of the four IFUs, dispersion of the collimated beam, and imaging of the resulting spectrum onto one half of one of the two focal plane arrays. One

of the two spectrometer arms includes the two short wavelength channels (1 and 2), and the other the long wavelength channels (3 and 4).

We will now step through the optical path from the IFUs to the detectors. The IFU output beam for each channel is collimated and its light directed toward the corresponding diffraction gratings as shown in Figure 9. Each spectrometer arm uses six gratings (two wavelength channels, three sub-band exposures). Figure 8 then shows the optical paths from the gratings to the detectors. The dispersed beams are imaged by three-mirror-anastigmat (TMA) camera systems (M1–M2–M3). In each spectrometer arm the channels have separate, but identical, M1 camera mirrors that provide intermediate images of the spectra between M1 and M2. Folding flats at these intermediate focus positions reflect the Channel 1 and Channel 4 beams such that the combined (Ch. 1 + 2) and (Ch. 3 + 4) beam pairs are imaged onto opposite halves of the detectors by common M2 and M3 mirrors. The symmetry through the center plane of the camera optics and the positioning of the optics allow for an opto-mechanical design of two mirror-imaged boxes, with identical optics and almost identical structures. The combination of symmetry and an all aluminum design (optics as well as structures) allowed the very strict alignment and stability requirements to

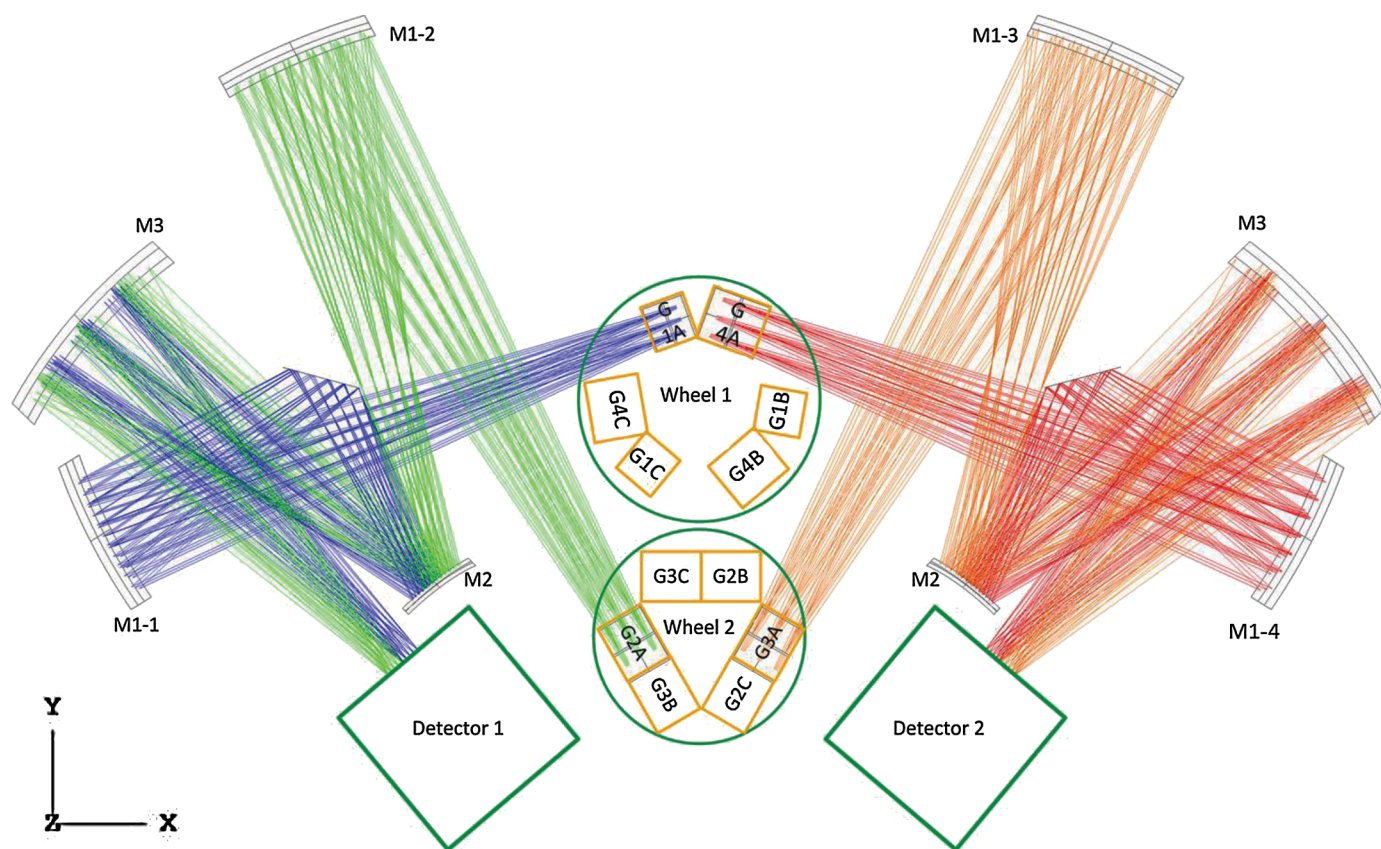


FIG. 8.—SMO cross section of the camera symmetry plane in the spatial direction. Light arrives at the gratings ($G1A$, $G2A$, $G3A$, $G4A$) from the IFU output collimating mirrors. See the electronic edition of the *PASP* for a color version of this figure.

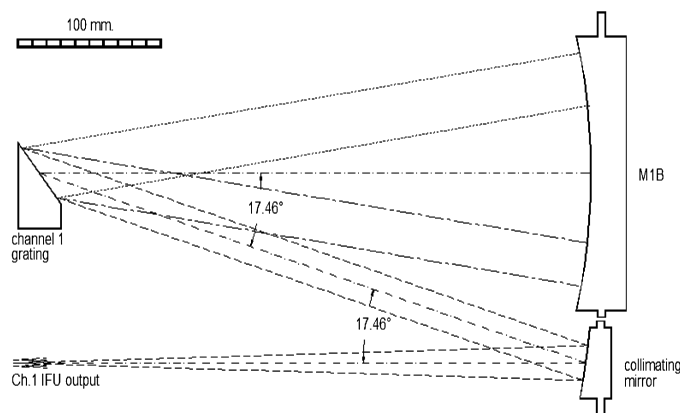


FIG. 9.—Cross section along the dispersion direction of the collimator, grating, and camera mirror M1 of the SMO.

be met fully with an adjustment-free mounting (i.e., positional accuracy by manufacturing accuracy only except for the focus shim at the detector interface), a major advantage regarding the total design, manufacture and test effort.

The mirror substrates are all light-weighted aluminum with diamond turned optical surfaces that are gold coated for maximum reflectivity at the MIRI operating wavelengths. They are mounted through holes in the housing walls using stress-relieving lugs under light-tight covers to prevent any high temperature radiation from the instrument enclosure (at ~ 40 K) reaching the detectors. The two sets of gratings for the A, B, and C sub-bands of the two channels are mounted on a single wheel, as shown in Figure 10, that is mounted on a mechanism that also carries the dichroics wheel on a lower level. The mechanisms are located in the SPO and each mechanism is rotated by

a single actuator to give the correct combination of dichroics and gratings across all channels for each exposure. The gratings are master rulings on aluminum substrates and are gold coated. Table 3 lists the design parameters of the gratings. All gratings operate in first order. The angular values quoted in Figure 1 are calculated for the (virtual) nominal input beam, originating from the center of the IFU output area.

2.5. Stray Light Control

The SPO is designed to be a well-baffled optical system. To intercept stray light from the telescope the first optical element in the MRS is a cold stop placed at the entrance pupil. A number of features within the MRS help to control stray light: (1) stops at each pupil and sky image in the optical train; (2) light traps where significant stray light occurs (0th order from the gratings); (3) labyrinths at the edges of apertures between modules and tight control of all external apertures, e.g., electrical feed-throughs; (4) black coating of all surfaces not in the main optical path, which is an inorganic black anodizing applied to a roughened surface of aluminum (the process was carried out by Protection des Metaux, Paris); and (5) avoidance of surfaces at grazing incidence near the main optical path. In addition, for each slice of each IFU there are output pupil and slit masks.

All the diamond turned aluminum mirrors in MIRI were specified to have a surface roughness of <10 nm RMS. This results in a total integrated scatter per surface of $<0.06\%$; non-sequential optical modeling indicated that this level of scatter should not significantly degrade the PSF of the MRS. All mirrors have an overcoated gold surface with reflectivity $>98.5\%$ in the MIRI wavelength range.

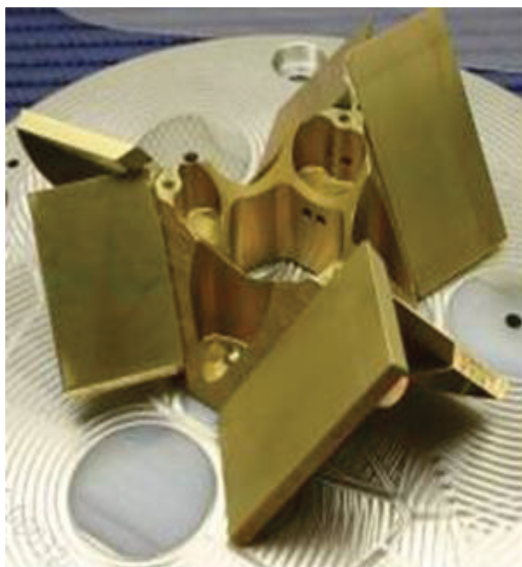


FIG. 10.—Grating Wheel Assembly (Wheel 1). See the electronic edition of the *PASP* for a color version of this figure.

TABLE 3

MRS GRATING DIMENSIONS: GROOVE LENGTH; RULED WIDTH;
ANGLE OF INCIDENCE; AND ANGLE OF DIFFRACTION

Band	Grating constant (mm ⁻¹)	Length (mm)	Width (mm)	θ_{inc} (deg)	θ_{diff} (deg)
Channel 1					
A	266.67	28	44	55.46	29.2
B	230.77				
C	200.00				
Channel 2					
A	171.43	30	63	54.46	28.5
B	148.45				
C	128.57				
Channel 3					
A	112.06	30	63	54.46	28.5
B	96.97				
C	83.96				
Channel 4					
A	71.43	34	64	53.46	27.2
B	61.28				
C	52.55				

All gratings have a blaze angle of 45°.

Stray light analysis shows that the extensive baffling combined with the low scattering optical surfaces and blackened structure should reduce unwanted light (cross talk between channels, degraded PSF and out-of-beam light at the output image) to levels that are estimated to be less than 0.1% of the background radiation included in the science beam.

3. MEASURED PERFORMANCE

3.1. The Spectral Image

Due to the nature of an IFU spectrograph with its slicing and dispersing optics, the resulting detector images are not straightforward to analyze: Spatial, spectral, and with it photometric information, are spread over the entire detector array. Further, the gratings are used in an optically fast, non-Littrow configuration. The resulting anamorphic magnification varies from the short to long wavelength limit for each sub-band which, when combined with the curved spectral images of the dispersed slices on the detector as shown in Figure 11, makes the optical distortion significant and complex. The anamorphic and slicing optics add other components of distortion that vary the plate scale in the along-slice direction for each slice individually. This leads to a very complicated variation of the spatial plate scales and finally to a dependency of the spatial and spectral axes on each other and on the location on the detector array.

To enable scientific studies based on MRS data, the flux measured in each detector pixel needs to be associated with a wavelength and location on the sky. Since the IFU provides two spatial dimensions, each detector pixel corresponds to a position within a three-dimensional cube with two spatial and one spectral dimension. Due to the optical distortion, the edges of these cubes are neither orthogonal nor constant in length. Consequently, we have developed a process (called image or cube reconstruction), which allows a transformation of the detector pixels onto orthogonal cubes.

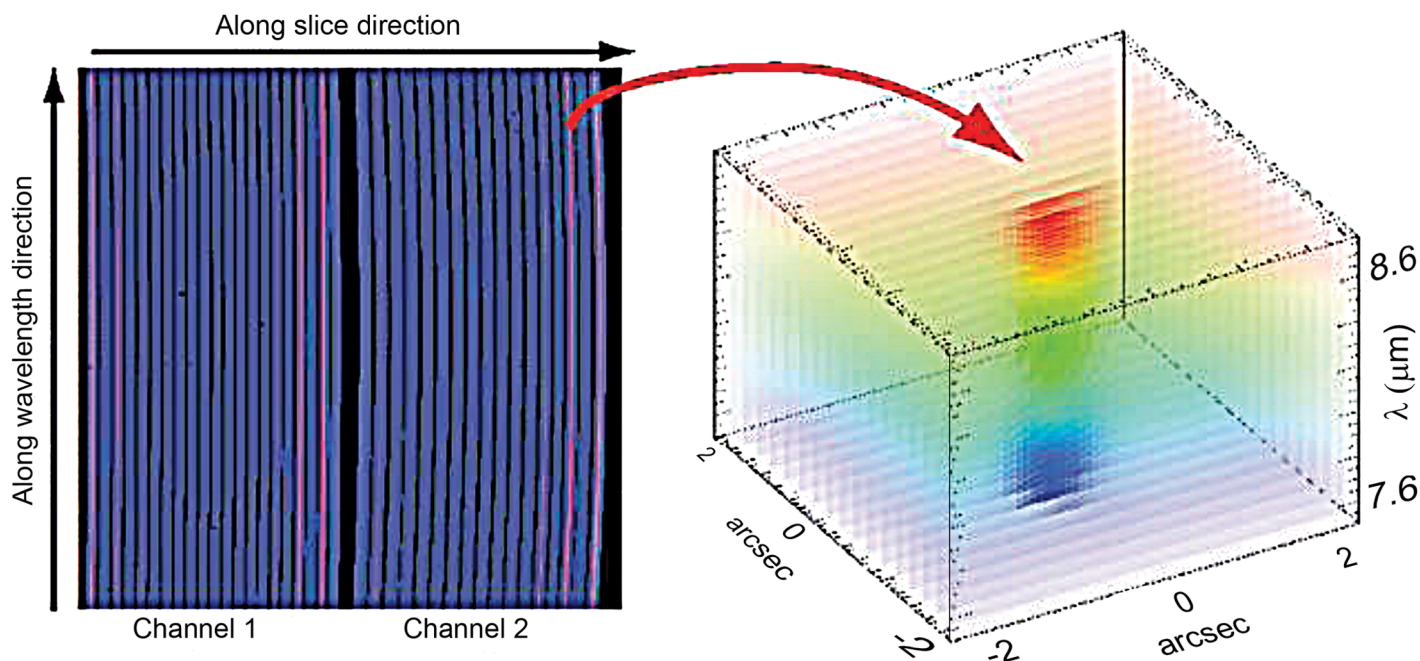


FIG. 11.—Illustration of the reconstruction process: The detector image for two channels is reconstructed into a 3D cube with two spatial dimensions (α along the slice and β across the slice) and one spectral dimension (λ). See the electronic edition of the *PASP* for a color version of this figure.

Cube reconstruction requires a thorough understanding of the optical distortion. It is not possible to just characterize the curved images of the dispersed slices (as shown in the left plot of Fig. 11) by approximating the curvature with polynomial functions. Due to the distortion caused by the slicing optics, the spatial plate scale varies within a slice nonuniformly. The characterization of the plate scale and the association of each detector pixel with its projected location on the sky at any given wavelength is a task for the astrometric and wavelength calibration, described in the next sections. Once the sky-coordinates and wavelength are known for each detector pixel, the cube can be reconstructed as discussed by Glauser et al. (2010).

3.2. Astrometric Calibration

To describe the optical distortion on the image plane of the detectors, we began from the optical model, using ray-tracing techniques to project each detector pixel through the slicing optics backwards onto the sky. When the as-built geometry was taken into account, we achieved a high spatial accuracy, as demonstrated for the MIRI Verification Model (Glauser et al. 2010). To verify this approach, a dedicated astrometric calibration campaign was conducted during the instrument test campaign of the MIRI Flight Model.

As outlined in Glauser et al. (2010), the intraslice spatial distortion can be approximated with a second-order polynomial to accuracies of a few milliarcseconds—much better than required for the astrometric accuracy given the minimum plate scale of $0.196''$ per pixel in the along-slice direction for channel 1. We conducted the astrometric calibration by placing a broad-band point source at three field positions for each slice and each channel and recording the dispersed signature on the detector. The central pixel of the spatial profile was determined and correlated with the known position of the point source (our reference on the sky). With this method we were able to determine the spatial plate scale at any location on the detector. Due to limitations of the test setup, i.e., strong field distortion of the steerable point-source, the achievable relative astrometric accuracy was very limited. However, within an estimated upper limit of ~ 0.5 pixels ($0.098''$ for channels 1 and 2, $0.123''$ for channel 3, and $0.137''$ for channel 4), the approach of using the optical model for the reconstruction was validated.

A further result from this calibration campaign was the measurement of the FOV for each sub-band. Figure 12 shows the MRS FOV in the *JWST* coordinate frame and its position relative to the imager field. The FOV common to all MRS channels is $3.64''$ in the along-slice (α) and $3.44''$ in the across-slice direction (β).

3.3. Wavelength Calibration

To determine the absolute wavelengths for the MRS channels we conducted a series of calibration measurements during the instrument test campaign. Fabry-Pérot etalon filters were used

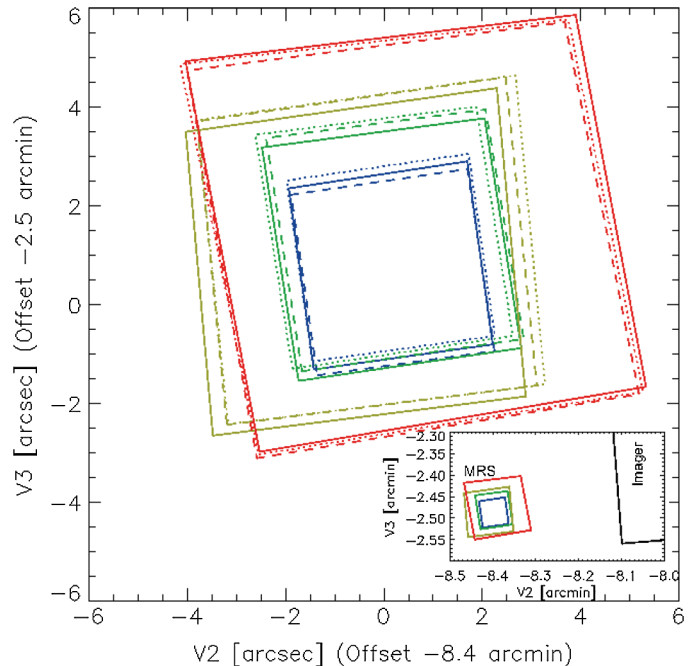


FIG. 12.—Transformed MRS fields in the *JWST* coordinate frame ($V2$, $V3$). The field borders are drawn as solid lines (Sub-band A), dashes (Sub-band B), and dots (Sub-band C) for channel 1 (blue), 2 (green), 3 (yellow), and 4 (red). The inset shows the position of the MRS FOV relative to the imager field. See the electronic edition of the *PASP* for a color version of this figure.

to create a dense pattern of unresolved spectral lines on the detector. Many tens of lines were typically visible at a signal to noise ratio of more than 50 in a single sub-band exposure. An example etalon measurement is shown in Figure 13. To provide a reference wavelength to distinguish between adjacent etalon lines, the telescope simulator used in the test was also equipped with a pair of “edge” filters, one of which cut-on at around $\lambda = 6.6 \mu\text{m}$ (in Sub-band 1C) and the other which cut off at $\lambda = 21.5 \mu\text{m}$ (in Sub-band 4B).

The absolute wavelengths of the etalon lines and edge filters had previously been measured at ambient and 77 K temperatures using a laboratory standard Fourier Transform spectrometer with a spectral resolving power of $R = 100,000$. The extrapolation of the wavelength scale to the 34 K operating temperature for the test campaign was based on published measurements (Browder & Ballard 1969; Browder & Ballard 1972; Smith & White 1975). This extrapolation resulted in a typical wavelength correction of less than 1% of the width of the MRS spectral resolution element. The repeatability of the scale after multiple mechanism reconfigurations has been measured to be 0.02 resolution elements.

The wavelength calibration process then involved the assignment of an absolute wavelength to an etalon line by fitting the calibrated spectra of the edge filters to their MRS measured spectra in Sub-band 1C and 4B. The known separation of the etalon lines was used to extend the wavelength scale across the

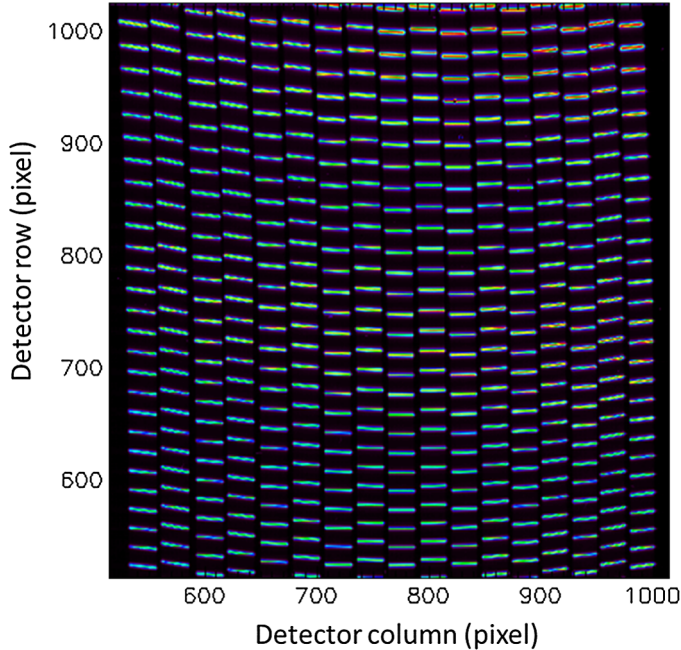


FIG. 13.—A typical wavelength calibration observation. The long wavelength half of Sub-band 2A is shown when illuminated by an extended continuum source modulated by a vacuum spaced etalon with a free spectral range (fringe spacing) of 3.1 cm^{-1} at ambient. See the electronic edition of the *PASP* for a color version of this figure.

full spectral image in each of these sub-bands. To extend the scale to other sub-bands, pairs of measured etalon spectra were coadded and the positions of unique identifying features (due to spectral beating between the two patterns) were used. A set (2 per sub-band) of second-order polynomial fits to the positions of the etalon lines was used to generate a wavelength value for the corners of all illuminated pixels.

The wavelength calibration derived in this way was encoded for use in the MRS calibration pipeline by forming six images (one for each MRS detector and all three grating wheel settings) where the image signal values were set equal to the wavelengths at the corner of each detector pixel. We note that as a result, these wavelength reference images have one more row and column than the detector images. The relative accuracy of this wavelength scale (within and between sub-bands) is estimated to be better than 0.02 spectral resolution elements which, for example, corresponds to 0.03 nm at $\lambda = 5 \mu\text{m}$. The absolute accuracy of the wavelength scale is estimated to be comparable to this figure but this will need to be confirmed during on-orbit commissioning using spectral standards.

The wavelength calibrated etalon spectra were also used to measure the spectral resolving power of the MRS. The results are shown in Figure 14, where the sub-band averaged values are shown in red and the spread of values seen across the field of view is indicated by the black band. We note that Figure 14 does not take account of the intrinsic spectral width of the etalon

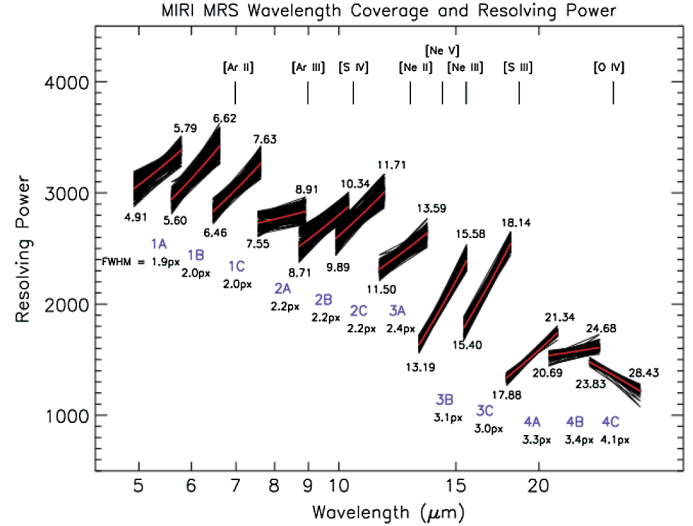


FIG. 14.—The spectral resolving power of the MRS. The black regions represent variations in the resolving power across the FOV for each sub-band, and the red lines are spatially averaged values. The wavelength ranges of each sub-band are indicated, as well as some scientifically important spectral lines. The mean width (FWHM) of the spectral resolution element in pixels is also provided for each sub-band. See the electronic edition of the *PASP* for a color version of this figure.

features, which was determined using the $R = 100,000$ calibrated spectra, described above. Initial efforts to deconvolve the intrinsic line profiles from the MRS measured spectra suggest that the resolving powers quoted in Figure 14 may be underestimated by around 10%. We have therefore used the band averaged measurements scaled by a factor 1.1 to generate the summary values quoted in Table 1.

3.4. Spatial Sampling

As outlined in § 2, the PSF of the MRS is under-sampled by design, with full sampling in both spatial and spectral dimensions requiring that the object be observed in at least two dither positions that include an offset in the across-slice direction of $1\frac{1}{2}$ times the channel 1 slice width (which corresponds to $7\frac{1}{2}$ slices in channel 2, $5\frac{1}{2}$ slices in channel 3, and $3\frac{1}{2}$ slices in channel 4). Due to the curved shape of the distorted spectrum on the detector and the variable plate scale along the individual slices, the exact dither offset in the along-slice direction is less well determined (but also less critical). Figure 15 shows the nominal MRS dither pattern to be used in a single observation to sample point sources fully, as derived during test campaigns, and as proposed for in-flight operations.

To achieve a fully sampled PSF, these dithered observations must be combined. We anticipate that this could most readily be achieved using the reconstructed cubes. To avoid loss of spatial resolving power caused by any shift- and coadding algorithm due to the rebinning of the data (for example, Fruchter & Hook [2002]), we minimize the necessary rebinning steps by

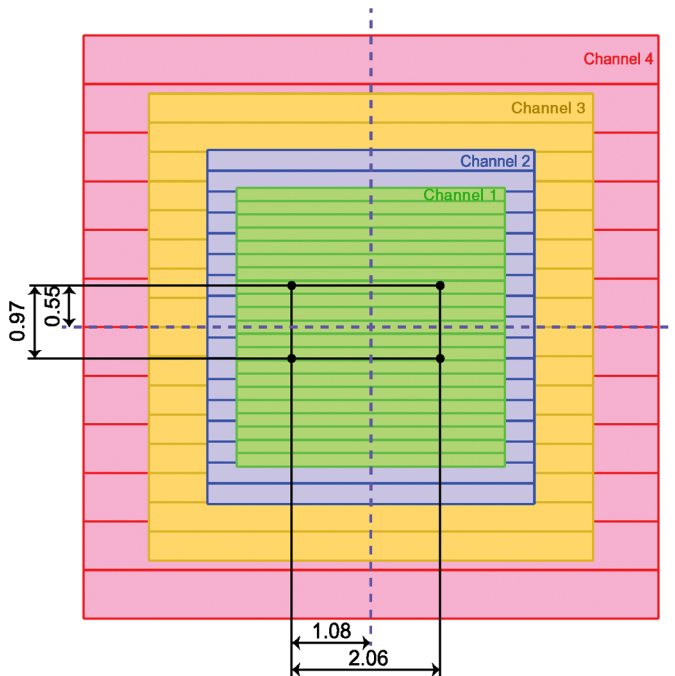


FIG. 15.—Coordinates in arcseconds of the four nominal dither positions relative to the center of the MRS field (*dashed axes*). See the electronic edition of the *PASP* for a color version of this figure.

incorporating dither offsets parallel to the slice into the reconstruction algorithm itself. This has the advantage that only one rebinning step is required from detector data to reconstructed cubes, while dither offsets in the across slice direction can be corrected and combined afterward using an interlacing method. We expect more sophisticated techniques than we have developed so far to be incorporated into the data reduction pipeline before launch.

3.5. PSF and Optical Quality

We attempted to measure the MRS PSF during flight model testing at RAL using the MIRI Telescope Simulator (MTS), described in Paper II. However, optical aberrations and vignetting in the MTS led to the generated point source being elongated and extended at short wavelengths, such that it did not provide a sufficiently point-like image. Even at longer wavelengths, where aberrations became less apparent as diffraction started to dominate, the measured PSF was broader than nominal due to vignetting in the telescope simulator.

We therefore repeated the PSF measurement on the flight model during the first cryo-vacuum test campaign at NASA-Goddard (CV1RR), where a compact and well defined point source was available. Deep exposures at 17 different locations in the MRS field were combined and reconstructed to form the PSF image for Channel 1A, shown in Figure 16. For these data, the spectral coverage was limited to a narrow ($0.125 \mu\text{m}$ wide) wavelength range around $5.6 \mu\text{m}$. A comparison with the model

PSF (pure diffraction limited Fourier-transformed *JWST* pupil) shows a very good match across the slices. In the along-slice direction, a broadening of approximately 50% is observed. Currently, possible causes considered for the broadening include scattering in the detector substrate (the detector halo effect, which is also observed in the imager and is discussed in Rieke et al. [2015], Paper VII) or a side-effect of the stray light discussed in § 3.9. In the case of the halo effect, the larger degree of broadening seen in the MRS may be due to the larger spatial sample per pixel. More detailed modeling is required to confirm the root cause.

3.6. Spectral Fringing

Spectral fringes are a common characteristic of infrared spectrometers. They originate from interference at plane-parallel surfaces in the light path of the instrument. These surfaces act as Fabry-Pérot etalons, each of which can absorb light from the source signal with a unique fringe pattern. In the infrared wavelength range, surfaces separated by a fraction of a mm up to a few cm may form very efficient etalons. The most obvious source of fringes in the MRS is the detector itself with a physical thickness of $500 \mu\text{m}$ (Paper VII). Similar fringes have also been observed in spectra measured with the MIRI Verification Model during testing and with the *Spitzer*-IRS instrument, which employs comparable (though smaller) Si:As BIB detectors (Lahuis & Boogert 2003).

For the initial analysis, we follow the formalism as defined in Kester et al. (2003) and Lahuis & Boogert (2003). The sine approximation, $I_t = A_f \cos(2\pi\omega D) + B_f \sin(2\pi\omega D)$, is used with ω being the wavenumber and D the optical thickness of the instrument component. Of primary interest to help in the identification and first characterization of the fringes is the optical thickness D . In the MRS test data three distinct fringe components are seen, with key parameters listed in Table 4.

Of the three fringe components two are directly matched to optical components in the instrument. The main fringe component (#1 in Table 4) has a derived optical thickness of approximately 3.5 mm for all sub-bands. This corresponds to the optical thickness of the detector substrate which has $d \sim 500 \mu\text{m}$ and $n \sim 3.42$, giving $D \sim 0.34 \text{ cm}$. Figure 17 gives an illustration of this fringe pattern, showing the main detector and dichroic fringes in more detail for Channel 4. The optical thickness derived for the second, high frequency, component ($D \sim 2.7 \text{ cm}$) is matched to that of the CdTe Dichroic filters (§ 2.2). For the third set of low frequency components ($D = 0.01\text{--}0.1 \text{ cm}$), no unique surfaces in the instrument are identified; instead these are likely to originate from beating between primary fringe components. The fringe variations come from the layered structure of the detector substrate and optical thickness differences between individual dichroic filters.

Figure 18 shows the peak normalized fringes over the entire wavelength range of the MRS. The two curves in Figure 18 show predictions for the expected fringe amplitudes based on

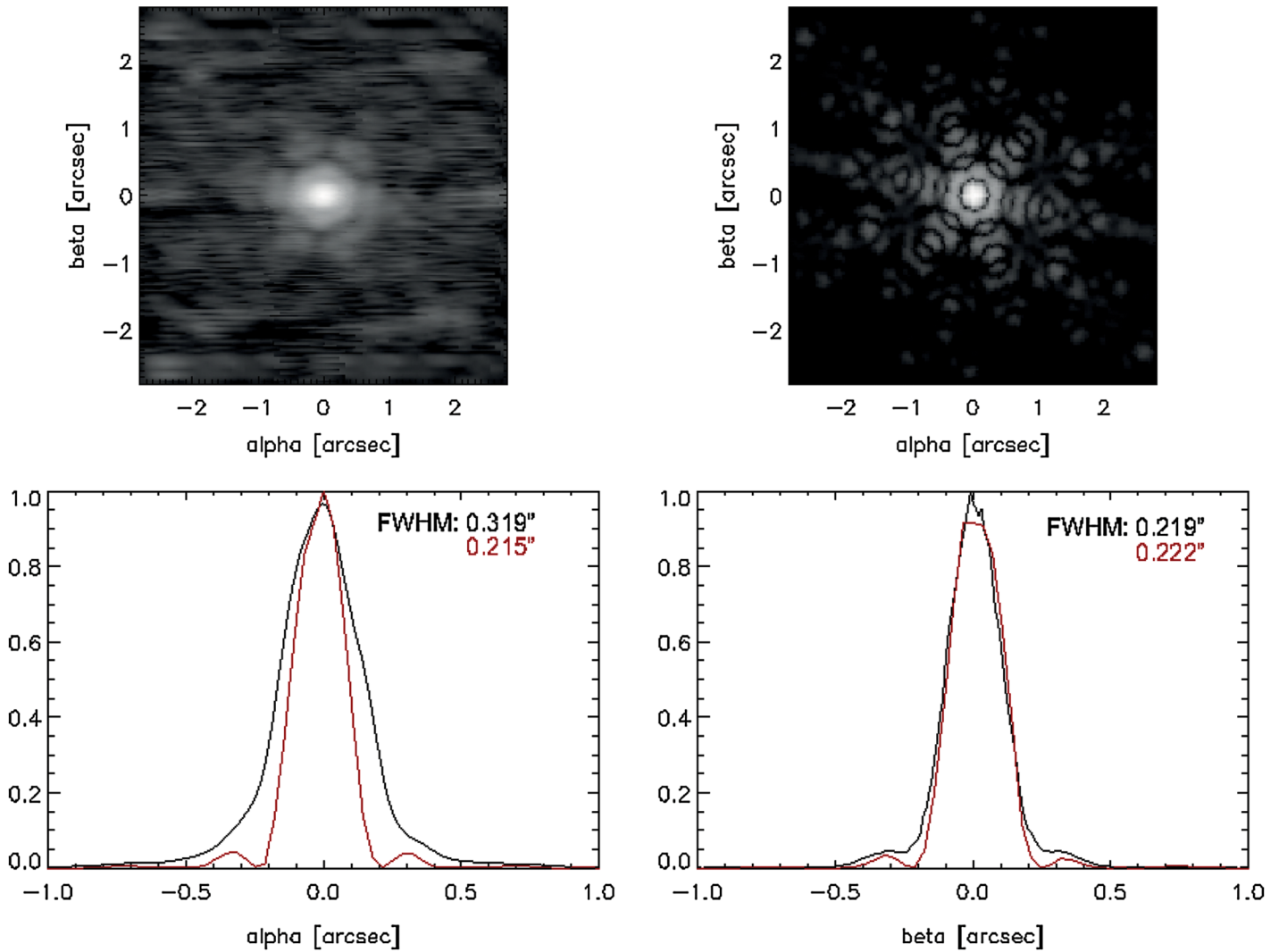


FIG. 16.—Measured PSF at $5.6\ \mu\text{m}$ (*top left*) compared with the FFT of the *JWST* pupil (*top right*) and normalized peak profiles of the measured (*black*) and modeled (*red*) for the along-slice (*bottom left*) and across-slice direction (*bottom right*). The x-axis is multiplied by two for the lower figures. See the electronic edition of the *PASP* for a color version of this figure.

representative anti-reflection (AR) coating profiles as applied to the MIRI flight detectors. The solid curve assumes a pure two-sided etalon while the dashed-dotted curve simulates a back-illuminated surface with a fully reflective front surface and photon absorption in the active layer (adopted from Woods et al. [2011]). Though not a perfect match, this approximate

detector model does reproduce the general trend and magnitude of the fringes.

Fringe removal will be achieved using the techniques developed for and applied to ISO and *Spitzer* data (see Lahuis & van Dishoeck [2000]; Kester et al. [2003]; Lahuis & Boogert [2003]). This involves dividing the observed spectra by a fringe flat-field

TABLE 4
OBSERVED FRINGE COMPONENTS WITH BASIC PARAMETERS AND IDENTIFICATION

Component	Optical thickness (cm)	$\delta\omega$ (cm^{-1})	Fringe origin	Peak-to-peak amplitude (%)
1	0.35 ± 0.02	$2.8 (\pm 0.2)$	Detector Substrate	≤ 40
2	2.7 ± 0.2	$0.37 (\pm 0.03)$	Dichroic	≤ 10
3	0.01–0.1	10–100	Fringe beating	≤ 10

The dichroic is modeled as chemical vapor deposited CdTe.

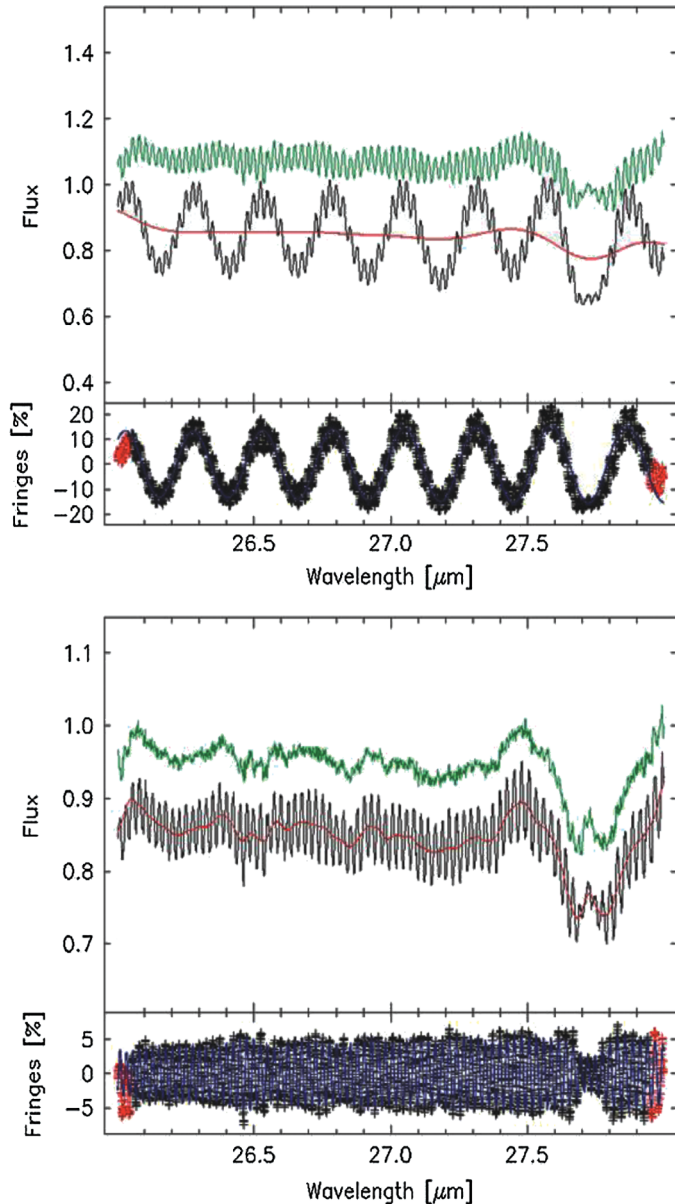


FIG. 17.—Examples of two main observed MRS fringe components. In the top panels of each plot the input data is shown in black with a fitted continuum in red. The spectrum (offset for clarity) after fringe removal is shown in green. In the lower panels the normalized spectrum and fitted fringes are shown. *Top*: the main component #1 in Channel 4 (after defringing, fringe component #2 remains); *bottom*: the high-frequency fringe component #2. See the electronic edition of the *PASP* for a color version of this figure.

followed by the removal of fringe residuals using the sine fitting method. This approach has proven to be reliable and robust for most spectra and the experience with ISO and Spitzer has shown that it allows the removal of fringe residuals down to the noise level. The main limitations with this technique are the definition of the spectral continuum in the presence of spectral features and isolating the fringe spectrum from broad molecular (vibration) rotation bands (e.g., those of C_2H_2 , HCN, CO_2 , and H_2O).

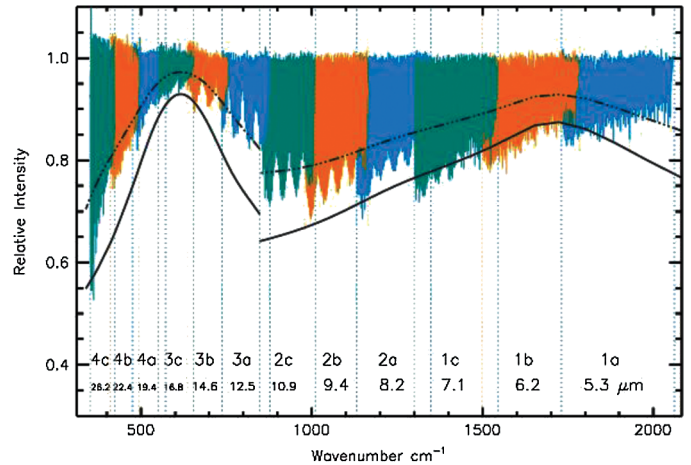


FIG. 18.—The combined fringes for each channel and sub-band. The data for each neighboring band is plotted in a different color for clarity. The main and high-frequency fringe components are not resolved on this scale. In Channels 2 and 3, an example of a low-frequency beating fringe can clearly be seen. Note: the spectrum is normalized to the peaks of the main fringe amplitude. See text for a description of the *solid* and *dashed-dotted* lines. See the electronic edition of the *PASP* for a color version of this figure.

For the MIRI IFU other effects play a role and may limit the fringe removal for point and compact source measurements. The two major effects are: (1) the illumination (and its effective angle) on the detector depends on the source morphology (position and extent); and (2) the wavelength depends on the spatial location of the point source in the IFU field. This results in (1) changes in the effective optical thickness from source to source and (2) a wavelength shift with spatial offset. Both modify the detailed fringe pattern for individual cases. Figure 19 illustrates this with

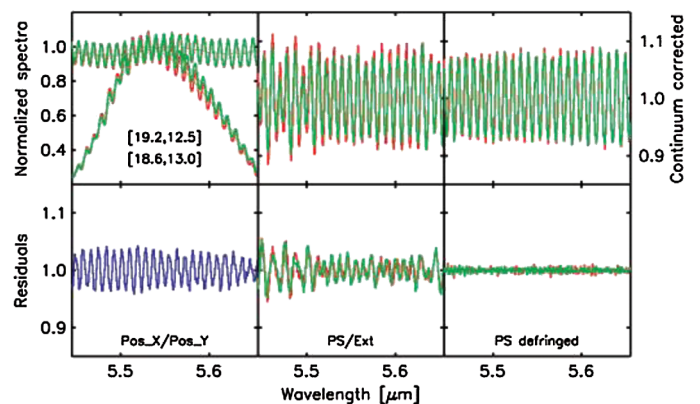


FIG. 19.—Illustration of the fringe source dependence for point and extended blackbody source spectra at two positions (*red and green spectra*) as printed in the top-left panel. *Top from left to right*: peak normalized point and extended spectra; normalized point source spectra; normalized extended blackbody spectra. *Bottom from left to right*: both point source spectra divided; point source spectra divided by extended source spectra (akin to direct flat-field correction); both normalized point source spectra defringed using a sine fitting technique. See the electronic edition of the *PASP* for a color version of this figure.

point source measurements from the CVIRR test campaign and using extended source blackbody spectra from flight model testing at RAL. Small subpixel pointing offsets are seen to have a discernable impact on the fringe pattern.

These effects can be mitigated by traditional fringe removal techniques using optimized and iterative reduction algorithms (e.g., by modifying, shifting, and stretching the reference fringe spectrum before applying it). This has been used in individual ISO and Spitzer cases and the MIRI team will work on developing optimized methods for the MRS. For MIRI an alternative model-based approach is also under study which uses the observed source morphology to define the fringe spectrum. This method will be applicable to both the MRS and the LRS, and will complement the traditional fringe removal techniques. This requires both a well-defined and well-calibrated fringe model (an ongoing MIRI team activity) and a flexible and iterative reduction pipeline (in development at STScI based on input from the MIRI team).

3.7. Pixel Flat-Field

We determined the pixel-to-pixel variation of the response for both short and long wavelength detectors, using measurements taken during testing at RAL. The illumination was provided by an external, extended source in the MTS. A number of exposures measured over the period of the whole test run and covering the full MRS wavelength range were included in this calculation.

We first defringed the data using the prescription described in § 3.6, with the results illustrated in Figure 17. We then calculated the mean value of each pixel, $f_{\text{ave}}(i, j)$, by averaging over 5×5 pixel boxes centered on pixel (i, j) and accounting for pixels close to the edges of the slices or near bad pixel clusters. The original data were then divided by these averaged values to create a map of the normalized pixel gain. We checked the distribution of the pixel-to-pixel variation for each exposure in all the available data for both detectors and all the sub-bands and found that the distribution is Gaussian with a full-width-half-maximum of $\sim 2\%$ (Fig. 20). The variation among different datasets has a standard deviation of $\leq 0.2\%$ for 16 different observations. The pixel-to-pixel variation does not appear to be wavelength dependent as both short and long wavelength detectors show the same overall flatness. This uniformity suggests that the correction of pixel response variations on orbit can be achieved using infrequent calibration measurements.

3.8. Spectrophotometric Performance

The absolute responsivity of the MRS is expressed in terms of the quantity referred to as the Photon Conversion Efficiency (PCE), which is equal to the number of electrons detected by the focal plane array for each photon incident at the MIRI entrance focal plane. The wavelength-dependent PCEs for the MRS were derived during testing at RAL. This was done by configuring

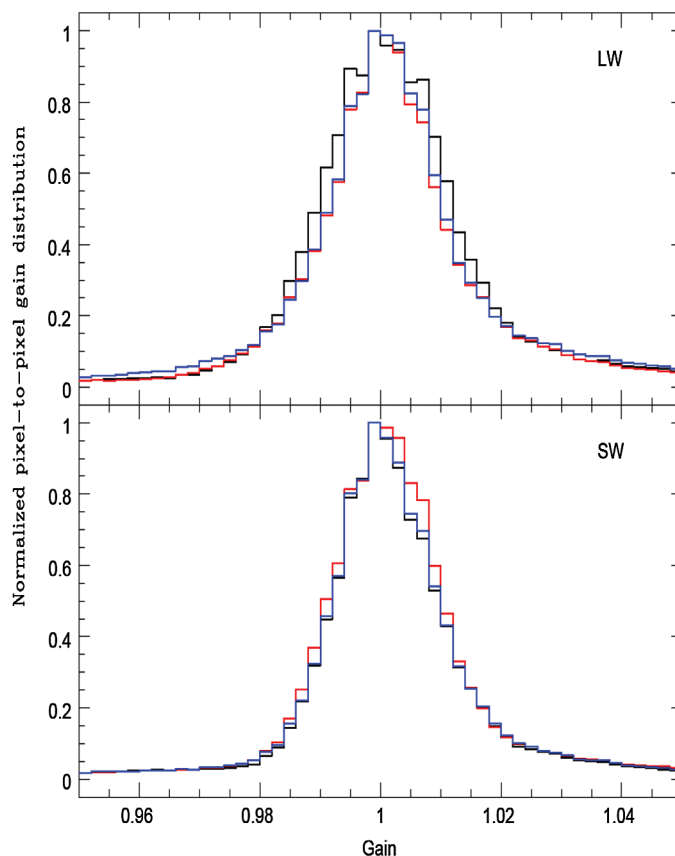


FIG. 20.—Normalized pixel-to-pixel gain distributions derived from 3 observations for SW and LW detectors (black, red, and blue). The FWHMs are 1.91% and 1.99% for the SW and LW channels, respectively. See the electronic edition of the *PASP* for a color version of this figure.

the MTS to provide extended illumination of the MIRI entrance pupil with the spectral energy distributions of blackbodies of 400, 600, and 800 K. For each MTS blackbody configuration, MRS spectra were obtained in all 12 spectral bands, together with background measurements using the blank position in the MTS filter wheel.

The data were processed using the standard DHAS tool (Paper II), which converts the raw readouts of the integration ramps to slopes in physical units (electrons/second). The DHAS miri cube module, which performs the reconstruction technique described in § 3.1, was then used to construct spectral cubes from the slope images, regridding the focal plane array pixel signals onto an equidistant spectral cube. The spectral cubes of the background measurements (MTS filter blank position) were subtracted from the flood illumination measurements to correct for the test facility background.

The spectral cubes produced with the DHAS miri cube routine have a fully calibrated World Coordinate System (WCS) with the plate scale and wavelength coverage of every cube pixel. When combined with an estimate for the absolute flux delivered to the entrance focal plane by the MTS (Paper II), this

TABLE 5
MEAN PHOTON CONVERSION EFFICIENCY (ELECTRON/
PHOTON) FOR ALL MRS BANDS

	Band		
	A	B	C
Channel 1	0.117	0.117	0.138
Channel 2	0.114	0.130	0.134
Channel 3	0.115	0.108	0.114
Channel 4	0.028	0.020	0.011

The PCE was averaged over all reconstructed spectral cube sample elements and all wavelengths.

allows us to calculate the photon conversion efficiency in every pixel of the spectral cube for all MRS spectral bands. Table 5 lists the mean PCEs for all MRS bands. We estimate the fractional error to be 20% in these figures, due to systematic effects, primarily in estimating the absolute flux from the MTS.

One obvious feature of Table 5 is the sharp drop in PCE from Channel 3 to Channel 4. The PCE in Channel 4 is roughly a factor of 2.5 lower than was expected from subsystem measurements of the nominal MRS optical train. The extra loss was identified (M. te Plate, ESA, private communication) as being caused by a fault in the groove profiles of the Channel 4 gratings, which cannot be corrected before launch.

Following the procedure to determine the photon conversion efficiency, we can establish a first spectrophotometric calibration. By comparing the flux conversion factors derived from blackbody measurements of the MTS at 400, 600, and 800 K, we can assess the achievable spectrophotometric calibration accuracy. Figure 21 shows the ratios of the different obtained flux

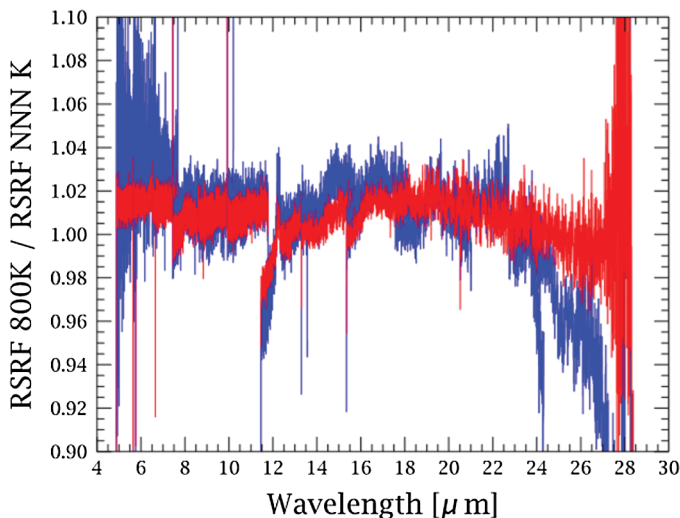


FIG. 21.—Comparison of the Relative Spectral Response function calculated using observations of the MTS black body at 400 K (blue) and at 600 K (red) with the RSRF at 800 K. See the electronic edition of the *PASP* for a color version of this figure.

conversion factors as a function of wavelength. We are encouraged by the good agreement (less than 2% variation) over a large swathe of the MRS waveband.

3.9. MRS Stray Light

As described in §§ 2.3 and 2.5, great care was taken to minimize sources of stray light and optical cross-talk within the MRS IFUs. However, a source of stray light was detected during RAL testing, which was identified as being caused by scattering in optical components within the SMO. The stray light is manifested as a signal that extends in the detector row direction. Its magnitude is proportional to that of bright illuminated regions of the spectral image, at a ratio that falls with increasing wavelength, from about 2% in Channel 1A to undetectably low levels longward of Channel 2B. Figure 21 emphasizes the impact of the stray light in Channel 1B, using the wavelength-averaged, reconstructed image of a bright source (seen at the top-left). The stray light signature is seen as the two horizontal bands, where the variation in brightness with the “alpha” coordinate is well explained by the mapping between alpha and detector row coordinate.

The development of algorithms for correction of this stray light is underway. They take advantage of the stray light being the dominant signal in the inter-slice regions of the detector, thereby making it amenable to accurate characterization. Initial indications (Fig. 22) suggest that an effective correction algorithm will be available before launch.

3.10. Spectral Leakage

Testing at RAL revealed a gap in the performance of the train of dichroics described in § 2.2. The set of filters that define the pass-band for Channel 3, Sub-band A, have an unwanted (but small) transmission peak at a wavelength of $6.1 \mu\text{m}$, which allows light in the second diffraction order to reach the detector at the position where $12.2 \mu\text{m}$ light is detected in the first diffraction order. This leak was confirmed using Fabry-Pérot etalon data and characterized to produce the leak profile plotted in Figure 23. This curve can be interpreted as the transmission profile by which the $6.1 \mu\text{m}$ spectrum of a target object should be multiplied to determine the leakage signal at $12.2 \mu\text{m}$.

There are two options for mitigating the effects of the spectral leak. First, the Channel 3A spectrum can be corrected by multiplying an observation of the $6.1 \mu\text{m}$ (Channel 1B) spectrum of the target object by the leak profile, resampling the wavelength scale of the resulting spectrum to the wavelength grid of Channel 3A and then subtracting it from the contaminated Channel 3A spectrum. This requires that a separate Channel 1B observation has been taken.

The second method would be to make the Channel 3A observation with Dichroic Wheel 1 set to use the Sub-band C dichroics (see § 2.2 to see what this means in terms of the

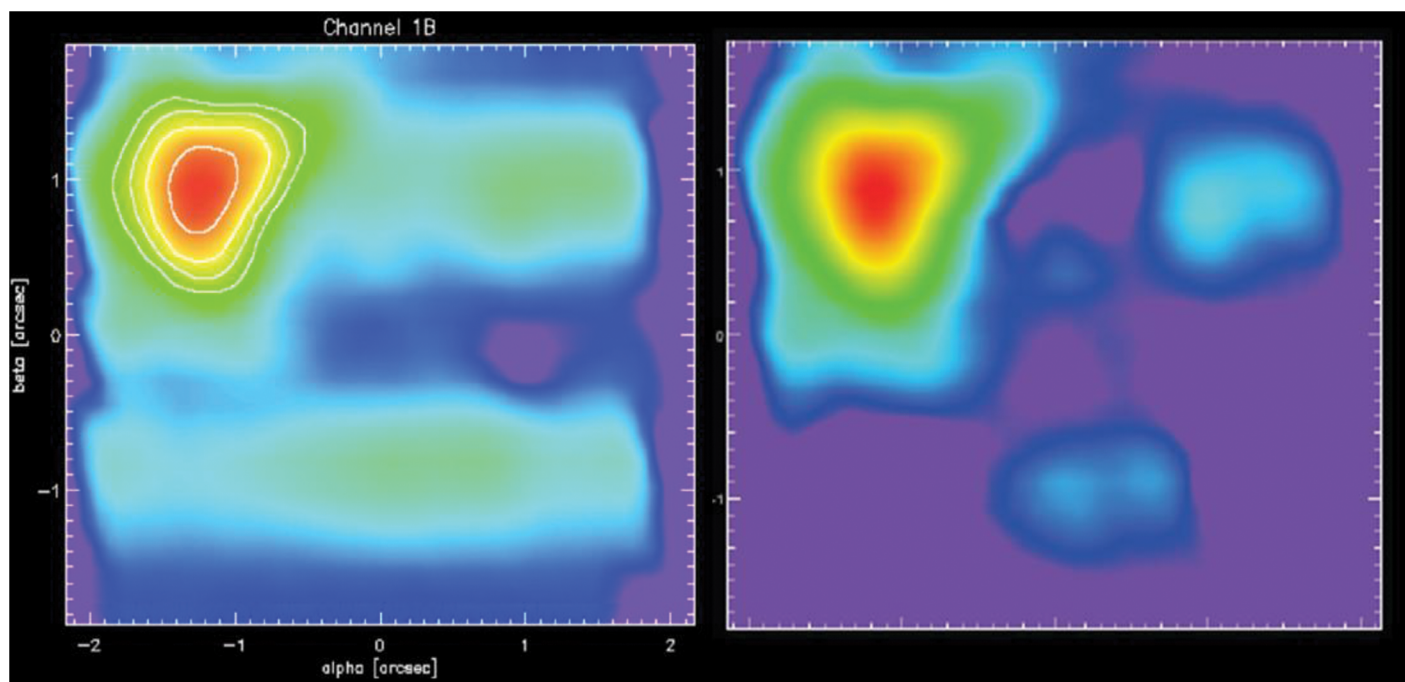


FIG. 22.—Stray light in Channel 1B. *Left*: The reconstructed image of a bright, moderately compact source is plotted on a logarithmic scale with contours at 50, 20, 10, and 5 % of the peak signal level. *Right*: The image after stray-light correction using a prototype algorithm. See the electronic edition of the PASP for a color version of this figure.

optical train). This combination of dichroics reduces the spectral leak by a factor of more than 1000 at $\lambda = 6.1 \mu\text{m}$. The unwanted side-effect of this solution is up to a factor of three loss of PCE at the short wavelength ends of Channel 2A and Channel 4A.

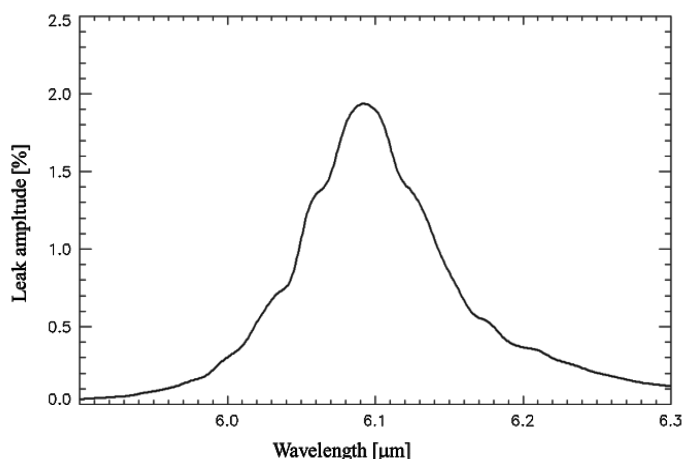


FIG. 23.—MRS Channel 3A spectral leak profile. This is the fraction of the source spectrum at $\lambda = 6.1 \mu\text{m}$ which is added to the nominal $12.2 \mu\text{m}$ spectrum.

4. CONCLUSION

We have presented the key parameters that describe the performance of the MIRI MRS spectrometer as designed or measured, in a form that both provides our best estimate of the behavior of the instrument on-orbit and also that is accessible to the prospective user. The optical design behind the parameters is presented at a level of detail that is intended to provide the astronomer with an understanding of what to expect in terms of operating restrictions and data format when planning observations.

The impact of straylight and spectral leaks in contaminating the spectral images has been discussed, along with proposals of operational and analytical techniques that should mitigate their effects. When combined with the latest sensitivity estimates (Glasse et al. 2015, Paper IX), we are confident that the MIRI MRS will meet all of its scientific objectives as part of the *JWST* Observatory.

The work presented is the effort of the entire MIRI team and the enthusiasm within the MIRI partnership is a significant factor in its success. MIRI draws on the scientific and technical expertise many organizations, as summarized in Papers I and II. A portion of this work was carried out at the Jet Propulsion Laboratory, California Institute of Technology, under a contract with the National Aeronautics and Space Administration.

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