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Origami metamaterials : design, symmetries, and combinatorics

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APPENDIX A

FOLD ANGLES

In this appendix we will derive closed form expressions for the relations between the fold angles of a generic 4-vertex. Expressions in the literature are either in implicit form [48, 76], are for flat foldable vertices only [16, 35, 76], or fail to clearly distinguish between the two possible discrete folding branches [16, 48]. The derivation shown here is originally by Rémi Menaut, and was shortened by Scott Waitukaitis. In addition we will show how the fold angles of a given 4-vertex and its supplement relate to each other.

A Euclidean 4-vertex consists of four rigid plates with sector angles α_i connected by four folds or hinges, where $\sum \alpha_i = 2\pi$ and we assume that all sector angles are unequal and smaller than π (Fig. A.1.A). The non-flat, folded states are characterized by the folding angles ρ_i , defined as the deviation from in-plane alignment between adjacent plates i and $i + 1$ (modulo 4). 4-vertices are equivalent to non-intersecting spherical mechanisms, allowing to represent their folded state accordingly (Fig. A.1.B). The fold angles are equal to the angle between the great circles at the point they meet, see ρ_1 in Fig. A.1.B, where we note that ρ_1 here is positive, as it is oriented counterclockwise.

Folding Branches

It was first shown by Huffman that a folded Euclidean 4-vertex will always have one fold whose sign is *unique*, i.e. the angle is opposite in sign from the other folding angles [28, 48, 58]. We call these folds *odd folds*, and these odd folds always straddle a common *odd plate*. A necessary and sufficient condition for the sector angle of the odd plate is the inequality

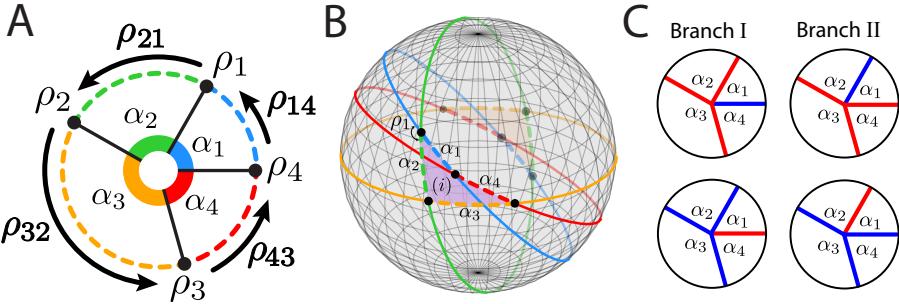


FIGURE A.1: (A) A generic 4-vertex with sector angles α_i , fold angles ρ_i , and fold operators $\rho_{i+1,i}$. (B) An origami vertex (i) can be modeled as a spherical mechanism. Dashed lines trace out the two vertices related by mirror symmetry, whereas the solid lines trace out supplemented vertices. (C) The four possible Mountain-Valley (colored red and blue respectively) arrangements of a generic Euclidean 4-vertex.

$\alpha_i + \alpha_{i+1} < \alpha_{i+2} + \alpha_{i+3}$. A generic 4-vertex always has two odd folds, which straddle a common *odd plate* [28]; we define our vertices such that ρ_4 and ρ_1 are the odd folds, and α_1 is the odd plate. Together with the $\{\rho_i\} \leftrightarrow \{-\rho_i\}$ symmetry, this yields four distinct mountain-valley patterns for a given 4-vertex, shown in Fig. A.1.C. We denote the folding branches where ρ_4 or ρ_1 has the opposite sign by I and II respectively.

Folding Operators

Along a given branch, 4-vertices have one continuous degree of freedom, and the relations between folding angles are anti-symmetric $\rho_i(-\rho_j) = -\rho_i(\rho_j)$ and bijective; we define folding operators $\rho_{i+1,i}^{I,II}$ which map the fold angles adjacent to plate i : $\rho_{i+1,i}^{I,II}(\rho_i) = \rho_{i+1}$, and suppress the index I and II when possible. Here we use the relations between the folding angles to show that the folding operators of a vertex, $\rho_{i+1,i}$ and its supplement, $\rho'_{i+1,i}$ are related as $\rho'_{i+1,i} = -\rho_{i+1,i}$.

We consider a folded state of a 4-vertex, and aim to express all fold angles as function of ρ_4 . We schematically represent the arc lengths and dihedral angles of the folded states as seen on the Euclidean sphere by diagrams such as Fig. A.2.B,C. We now consider the arc length λ_{41} between folds ρ_3 and ρ_1 , depicted in Fig. A.1.B,C. Using the spherical law of cosines,

we obtain:

$$\cos \lambda_{34} = \cos \alpha_4 \cos \alpha_1 - \sin \alpha_4 \sin \alpha_1 \cos \rho_4 . \quad (\text{A.1})$$

The arc length λ_{34} is part of two spherical triangles, one with dihedral angles $\pi - \sigma_4$, $\pi - \sigma_1$ and $\pi - \sigma_3$, and the other with $\pi - \tau_1$, $\pi - \tau_2$ and $\pi - \tau_3$. Making use of the shorthand notation,

$$A(a, b, c) \equiv \arccos \left(\frac{\cos a \cos b - \cos c}{\sin a \sin b} \right) , \quad (\text{A.2})$$

and repeatedly using the spherical law of cosines, the dihedral angles σ_i and τ_i are,

$$\begin{array}{ll} \sigma_1 = A(\lambda_{34}, \alpha_1, \alpha_4) & \tau_1 = A(\alpha_2, \lambda_{34}, \alpha_3) \\ \sigma_3 = A(\alpha_4, \lambda_{34}, \alpha_1) & \tau_2 = A(\alpha_3, \alpha_2, \lambda_{34}) \\ \sigma_4 = A(\alpha_1, \alpha_4, \lambda_{34}) & \tau_3 = A(\lambda_{34}, \alpha_3, \alpha_2). \end{array}$$

These are all functions or ρ_4 through their dependence on λ_{34} .

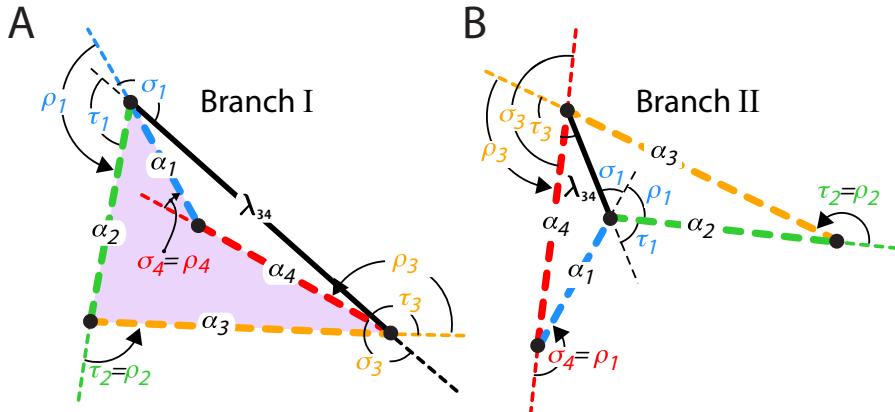


FIGURE A.2: (A) Simplified diagram of a 4-vertex (i) folded in Branch I, as in Fig. A.1.B. (B) Simplified diagram of a 4-vertex folded in Branch II.

We obtain the folding angles from σ_i and τ_i ; taking care regarding the relative signs of the fold angles on each branch, the exact equations for

$\rho_4 > 0$ are,

$$\rho_1^{I/II} = -\pi + \sigma_1 \mp \tau_1, \quad (\text{A.3})$$

$$\rho_2^{I/II} = \mp \tau_2, \quad (\text{A.4})$$

$$\rho_3^{I/II} = -\pi + \sigma_3 \mp \tau_3, \quad (\text{A.5})$$

where the minus sign in \mp corresponds to branches I , and the plus sign in \mp to branch II . Because of reflection symmetry in the flat-state plane, these equations are antisymmetric: $[\rho_i^{I/II}(\rho_4 < 0) = -\rho_i^{I/II}(\rho_4 > 0)]$. Similarly, we can obtain expressions for any fold angle as function of any other fold angle.

The operator ρ_{14} follows directly from Eq. A.3. Using these explicit expressions, we now show that $\rho'_{i+1,i} = -\rho_{i+1,i}$. For the supplemented vertex, we modify the sector angles from α_i to $\alpha'_i = \pi - \alpha_i$. First, note that the expression for $\cos(\lambda_{34})$, Eq. A.1, remains identical under this transformation, making use of the identity $\cos(-x) = \cos(x)$ and $\sin(-x) = -\sin(x)$. Second, note that we can write:

$$\begin{aligned} \sigma'_1 &= A(\lambda_{34}, \alpha'_1, \alpha'_4), \\ &= \arccos\left(\frac{-\cos \lambda_{34} \cos \alpha_1 + \cos \alpha_4}{\sin \lambda_{34} \sin \alpha_1}\right) \\ &= \pi - A(\lambda_{34}, \alpha_1, \alpha_4), \\ &= \pi - \sigma_1 \end{aligned} \quad (\text{A.6})$$

using the identity: $\arccos(-x) = \pi - \arccos(x)$. Likewise, we have $\tau'_1 = \pi - \tau_1$. We therefore find (dropping the branch notation),

$$\begin{aligned} \rho'_1 &= -\pi + \sigma'_1 \mp \tau'_1, \\ &= -\pi + (\pi - \sigma_1) \mp (\pi - \tau'_1), \\ &= \pi - \sigma_1 \pm \tau_1, \\ &= -\rho_1. \end{aligned} \quad (\text{A.7})$$

Hence, $\rho'_{14} = -\rho_{14}$, and we can trivially extend this argument to show that $\rho'_{i+1,i} = -\rho_{i+1,i}$.

4-VERTEX AS A SPHERICAL MECHANISM

The operator symmetry $\rho'_{ij} = -\rho_{ij}$ derived in appendix A, can also be derived graphically. As shown before, a 4-vertex can be modeled as a spherical mechanism, which can be represented on the surface of a sphere (Fig. A.1.B). In Fig. B.1.A we schematically represent a 4-vertex (i) in a folded configuration, by using a pseudo-Mercator projection. Extending the arcs of vertex (i), we obtain four directed great circles. The intersection of circle i and $i + 1$ –indicated by the black circles– correspond to the hinges of spherical linkage (i), whereas their respective angle corresponds to fold angle ρ_i . We define the fold angles ρ_i as positive when ρ_i turns counterclockwise. The grey circles indicate the antipodal points of the black points of vertex (i). Along each directed circle we name the four arc lengths: $\alpha_i, \bar{\alpha}_i, \dot{\alpha}_i, \tilde{\alpha}_i$, where $\dot{\alpha}_i = \alpha_i$, $\bar{\alpha}_i = \tilde{\alpha}_i = \pi - \alpha_i$, and $\alpha_i + \bar{\alpha}_i + \dot{\alpha}_i + \tilde{\alpha}_i = 2\pi$ (also see Fig. B.1.A). Furthermore, any pair of great circles intersects at two locations, and because they are great circles, the angles around an intersection point are identical in magnitude to the angles around its antipodal point.

We first show that we can derive the relationship between the fold angles (and fold operators) of the four related vertices (i), (ii), (iii) and (iv) of Fig. 2.2. Our goal is to relate the fold angles of vertices (ii), (iii) and (iv) to the fold angles ρ_i of vertex (i) – shaded pink in Fig. B.1.A. First, vertex (ii) –shaded orange in Fig. B.1.A– can be found by connecting the four antipodal (grey) nodes, which consists of arc lengths $\dot{\alpha}_i$. When we

consider the fold angles of this vertex (running clockwise), we see that they are all oppositely oriented with respect to those of vertex (i), and therefore each pick up a minus sign. Thus, if a vertex can be in a configuration with folding angles $\{\rho_i\}$, it can also be in a configuration with folding angles $\{-\rho_i\}$, consistent with $\rho_j(-\rho_i) = -\rho_j(\rho_i)$

Second, we consider vertex (iii), which is shaded purple in Fig. B.1.B. This vertex consists of arc lengths: $\tilde{\alpha}_1, \tilde{\alpha}_2, \tilde{\alpha}_3, \tilde{\alpha}_4$, running clockwise. We see that in this case, only the fold angle around the antipodal (gray) nodes are reversed. The same holds true for vertex (iv) ($\bar{\alpha}_1, \bar{\alpha}_2, \bar{\alpha}_3, \bar{\alpha}_4$), which is shaded orange in Fig. B.1.B. As the resultant fold angles are alternating in sign for both vertex (iii) and (iv), and we find: $\rho_{i+1}(\rho_i) = -\rho_{i+1}(\rho_i)$, or $\rho'_{ij} = -\rho_{ij}$, using the operator notation. We finally note that vertex (i) in Fig. B.1.A depicts a vertex folded on branch I, where ρ_4 is opposite to the three other folds, but the relations above also hold on branch II (where ρ_1 is opposite in sign).

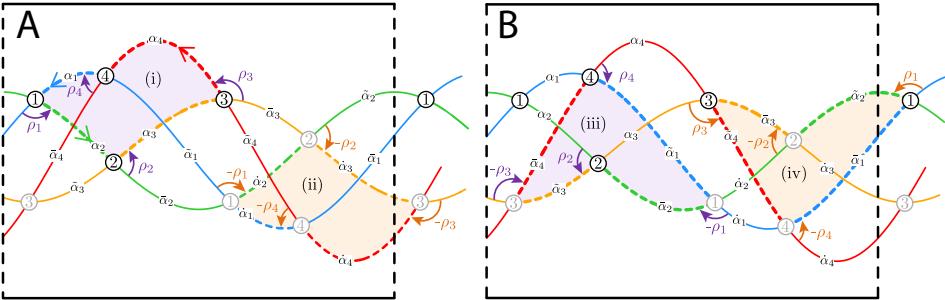


FIGURE B.1: (A) Simplified Mercator projection of a vertex (i) as shown in Fig. A.1, and its mirror image on the other side of the sphere (ii). (B) Simplified Mercator projection of two the two supplemented vertices, (iii) and (iv). Dashed line indicates periodic boundary. For details see text.

Besides the related Euclidean vertices (i)-(iv) and their representative spherical mechanisms (also termed 'folding linkages' [77]), there are an additional 12 related spherical mechanisms [78, 79]. All of these represent non-Euclidean vertices. Some of these mechanisms however, are self-intersecting, meaning they can not be converted into a 4-vertex as the plates would intersect. To study these additional spherical mechanisms in detail, we use the same pseudo-Mercator map as in Fig. B.1, in Fig. B.2. On this map we express all 16 arc lengths $(\alpha_i, \bar{\alpha}_i, \dot{\alpha}_i, \tilde{\alpha}_i)$, as well as the

angles around each node (ρ_i and ρ'_i) in terms of the arc lengths and angles of the original counterclockwise oriented vertex (i), which is depicted by the dashed line.

An example of one the 12 non-Euclidean spherical mechanisms is shown in light-blue in Fig. B.2, which we denote as $\tilde{\alpha}_1\dot{\alpha}_2\bar{\alpha}_3\alpha_4$. When we consider the magnitude of the angles between consecutive links of this mechanism, we find that they are: $\rho'_1, \rho'_2, \rho'_3, \rho'_4$. The signs of these angles can be found by comparing the orientation of these angles to those of the original vertex (vertex (i) in Fig. B.1.A), where the orientation of the mechanism itself is set by the colors of the segments (blue \rightarrow green \rightarrow yellow \rightarrow red). In this case, the orientation of the angles in the mechanism $\tilde{\alpha}_1\dot{\alpha}_2\bar{\alpha}_3\alpha_4$ are respectively: counterclockwise, counterclockwise, clockwise, and counterclockwise. When comparing this to vertex (i), we see that both ρ'_3 and ρ'_4 are oppositely oriented, which is why they obtain a minus sign. The folding angles of this mechanism are therefore $\rho'_1, \rho'_2, -\rho'_3, -\rho'_4$. An example of a self-intersecting non-Euclidean mechanism is colored green in Fig. B.2. This mechanism is denoted as $\bar{\alpha}_1\tilde{\alpha}_2\dot{\alpha}_3\dot{\alpha}_4$. Comparison of the angles of this mechanism to those of vertex (i) yields: $\rho_1, \rho'_2, -\rho_3, \rho'_4$.

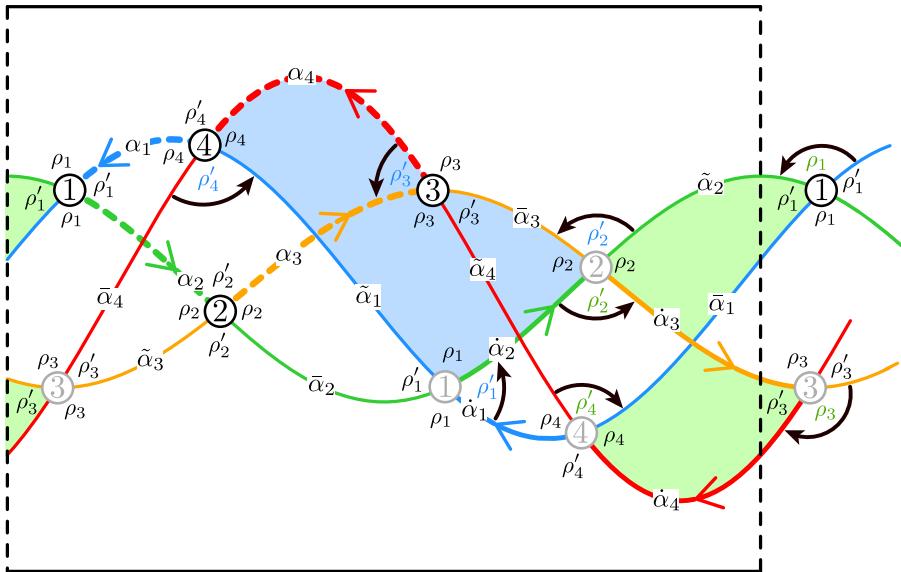


FIGURE B.2: Simplified mercator projection of the flat vertex shown in Fig. A.1 (dashed lines), folded on branch I. Dashed lines indicate periodic boundary.

In Table B.1 we list all 16 spherical mechanisms and their respective arc lengths (β_i), and fold angles (θ_i), expressed in terms of the arc lengths and fold angles of the vertex $\alpha_1\alpha_2\alpha_3\alpha_4$. Although the fold angles derived here are derived from Fig. B.2, which depicts a spherical mechanism on branch I (where the sign of ρ_4 on vertex $\alpha_1\alpha_2\alpha_3\alpha_4$ is opposite to the other three), we note that the expressions are valid for branch II as well (where the sign of ρ_1 on vertex $\alpha_1\alpha_2\alpha_3\alpha_4$ is opposite to the other three). In addition to the fold angles θ_i , we also display the sign of θ_i on both branch I and branch II. From the $\bar{\alpha}_1\tilde{\alpha}_2\dot{\alpha}_3\dot{\alpha}_4$ mechanism we know that self intersecting mechanisms have two consecutive positive fold angles, and two consecutive negative fold angles. We therefore see that mechanisms $\alpha_1\alpha_2\tilde{\alpha}_3\bar{\alpha}_4$, $\tilde{\alpha}_1\bar{\alpha}_2\alpha_3\alpha_4$, $\dot{\alpha}_1\dot{\alpha}_2\bar{\alpha}_3\dot{\alpha}_4$, and $\bar{\alpha}_1\tilde{\alpha}_2\dot{\alpha}_3\dot{\alpha}_4$ are self-intersecting on branch I. On branch II we find that $\alpha_1\tilde{\alpha}_2\bar{\alpha}_3\alpha_4$, $\tilde{\alpha}_1\dot{\alpha}_2\alpha_3\alpha_4$, $\dot{\alpha}_1\bar{\alpha}_2\tilde{\alpha}_3\dot{\alpha}_4$, and $\bar{\alpha}_1\alpha_2\alpha_3\tilde{\alpha}_4$ are self intersecting. Other than the four Euclidean vertices $\alpha_1\alpha_2\alpha_3\alpha_4$, $\tilde{\alpha}_1\bar{\alpha}_2\tilde{\alpha}_3\bar{\alpha}_4$, $\dot{\alpha}_1\dot{\alpha}_2\dot{\alpha}_3\dot{\alpha}_4$, $\bar{\alpha}_1\tilde{\alpha}_2\bar{\alpha}_3\tilde{\alpha}_4$, this leaves four spherical mechanisms that represent non-Euclidean vertices which can fold from branch I to branch II, these are: $\alpha_1\tilde{\alpha}_2\dot{\alpha}_3\bar{\alpha}_4$, $\tilde{\alpha}_1\dot{\alpha}_2\bar{\alpha}_3\alpha_4$, $\dot{\alpha}_1\bar{\alpha}_2\alpha_3\tilde{\alpha}_4$, $\bar{\alpha}_1\alpha_2\bar{\alpha}_3\dot{\alpha}_4$.

Arc Lengths				Folding Angle				Sign (Branch I)				Sign (Branch II)			
β_1	β_2	β_3	β_4	θ_1	θ_2	θ_3	θ_4	θ_1	θ_2	θ_3	θ_4	θ_1	θ_2	θ_3	θ_4
α_1	α_2	α_3	α_4	ρ_1	ρ_2	ρ_3	ρ_4	+	+	+	-	-	+	+	+
α_1	$\tilde{\alpha}_2$	$\bar{\alpha}_3$	α_4	$-\rho'_1$	$-\rho_2$	$-\rho'_3$	ρ_4	-	-	-	-	+	-	-	+
α_1	α_2	$\tilde{\alpha}_3$	$\bar{\alpha}_4$	ρ_1	$-\rho'_2$	$-\rho_3$	$-\rho'_4$	+	-	-	+	-	-	-	-
α_1	$\tilde{\alpha}_2$	$\dot{\alpha}_3$	$\bar{\alpha}_4$	$-\rho'_1$	ρ'_2	ρ'_3	$-\rho'_4$	-	+	+	+	+	+	+	-
$\tilde{\alpha}_1$	$\bar{\alpha}_2$	$\tilde{\alpha}_3$	$\bar{\alpha}_4$	$-\rho_1$	ρ_2	$-\rho_3$	ρ_4	-	+	-	-	+	+	-	+
$\tilde{\alpha}_1$	$\dot{\alpha}_2$	$\dot{\alpha}_3$	$\bar{\alpha}_4$	ρ'_1	$-\rho_2$	ρ'_3	ρ_4	+	-	+	-	-	-	+	+
$\tilde{\alpha}_1$	$\bar{\alpha}_2$	α_3	α_4	$-\rho_1$	$-\rho'_2$	ρ_3	$-\rho'_4$	-	-	+	+	+	-	+	-
$\tilde{\alpha}_1$	$\dot{\alpha}_2$	$\tilde{\alpha}_3$	α_4	ρ'_1	ρ'_2	$-\rho'_3$	$-\rho'_4$	+	+	-	+	-	+	-	-
$\dot{\alpha}_1$	$\dot{\alpha}_2$	α_3	$\dot{\alpha}_4$	$-\rho_1$	$-\rho_2$	$-\rho_3$	$-\rho_4$	-	-	-	+	+	-	-	-
$\dot{\alpha}_1$	$\bar{\alpha}_2$	$\tilde{\alpha}_3$	$\dot{\alpha}_4$	ρ'_1	ρ_2	ρ'_3	$-\rho_4$	+	+	+	+	-	+	+	-
$\dot{\alpha}_1$	$\dot{\alpha}_2$	$\bar{\alpha}_3$	$\tilde{\alpha}_4$	$-\rho_1$	ρ'_2	ρ_3	ρ'_4	-	+	+	-	+	+	+	+
$\dot{\alpha}_1$	$\bar{\alpha}_2$	α_3	$\tilde{\alpha}_4$	ρ'_1	$-\rho'_2$	$-\rho'_3$	ρ'_4	+	-	-	-	-	-	-	+
$\bar{\alpha}_1$	$\tilde{\alpha}_2$	$\bar{\alpha}_3$	$\tilde{\alpha}_4$	ρ_1	$-\rho_2$	ρ_3	$-\rho_4$	+	-	+	+	-	-	+	-
$\bar{\alpha}_1$	α_2	α_3	$\tilde{\alpha}_4$	$-\rho'_1$	ρ_2	$-\rho'_3$	$-\rho_4$	-	+	-	+	+	+	-	-
$\bar{\alpha}_1$	$\tilde{\alpha}_2$	$\dot{\alpha}_3$	$\dot{\alpha}_4$	ρ_1	ρ'_2	$-\rho_3$	ρ'_4	+	+	-	-	-	+	-	+
$\bar{\alpha}_1$	α_2	$\tilde{\alpha}_3$	$\dot{\alpha}_4$	$-\rho'_1$	$-\rho'_2$	ρ'_3	ρ'_4	-	-	+	-	+	-	+	+

TABLE B.1: Table listing all 16 spherical mechanisms that can be linked to a generic Euclidean 4-vertex, and their fold angles.

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