The Complete Local Volume Groups Sample - III. Characteristics of group central radio galaxies in the Local Universe

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ABSTRACT

Using new 610 MHz and 235 MHz observations from the Giant Metrewave Radio Telescope (GMRT) in combination with archival GMRT and Very Large Array (VLA) survey data we present the radio properties of the dominant early-type galaxies in the low-richness sub-sample of the Complete Local-volume Groups Sample (CLoGS; 27 galaxy groups) and provide results for the radio properties of the full CLoGS sample for the first time. We find a high radio detection rate in the dominant galaxies of the low-richness sub-sample of 82% (22/27); for the full CLoGS sample the detection rate is 87% (46/53). The group-dominant galaxies exhibit a wide range of radio power, $10^{20} - 10^{25}$ W Hz⁻¹ in the 235 and 610 MHz bands, with the majority (53%) presenting point-like radio emission, 19% hosting currently active radio jets, 6% having remnant jets, 9% being diffuse and 13% having no detected radio emission. The mean spectral index of the detected radio sources in the 235-610 MHz frequency range is found to be $\alpha_{235}^{610} \sim 0.68$, and $\alpha_{235}^{1400} \sim 0.59$ in the 235–1400 MHz one. In agreement with earlier studies, we find that the fraction of ultra-steep spectrum sources $(\alpha > 1.3)$ is ~4%, mostly dependent on the detection limit at 235 MHz. The majority of pointlike systems are found to reside in dynamically young groups, whereas jet systems show no preference between spiral-rich and spiral-poor group environments. The mechanical power of the jet sources in the low-richness sample groups is estimated to be $\sim 10^{42} - 10^{44}$ erg s⁻¹ with their black hole masses ranging between $2 \times 10^8 - 5 \times 10^9$ M_{\odot}. We confirm previous findings that, while radio jet sources tend to be associated with more massive black holes, black hole mass is not the decisive factor in determining jet activity or power.

Key words: galaxies: groups: general — galaxies: active — galaxies: jets — radio continuum: galaxies

1 INTRODUCTION

The bulk of galaxies and baryonic matter in the local Universe is found in galaxy groups (Geller & Huchra 1983; Fukugita et al. 1998; Eke et al. 2005). Typically, they extend less than a Mpc with a mass range of $10^{12.5} - 10^{14} M_{\odot}$ (Huchra & Geller 1983),

about an order of magnitude less massive than galaxy clusters. Galaxy groups exhibit shallow gravitational potential wells, with their members being at close distances and at low relative velocities. These parameters are essential in driving the transformation processes (e.g., mergers and tidal interactions) of galaxy evolution (e.g., Alonso et al. 2012). Galaxies are thought to be 'preprocessed' in the group environment before they become parts of clusters (e.g. van den Bosch et al. 2014; Haines et al. 2018) hence

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the association between rich cluster galaxies and their evolution is intimately linked to the group environment (e.g. Bekki et al. 1999; Moss & Whittle 2000). Since the properties of galaxies and their evolution depends on their local environment, galaxy groups are an ideal environment where the study of galaxy formation and evolution is of uttermost significance (e.g., Forbes et al. 2006; Sun 2012).

Many groups maintain extensive halos of hot gas with short central cooling times (e.g., O'Sullivan et al. 2017), which can fuel both star formation and active galactic nuclei (AGN). However, the cool gas can either be removed from the central region via AGN outflows (Alexander et al. 2010; Morganti et al. 2013) or heated up, eventually suppressing star formation and galaxy growth. The imprints of such a heating mechanism ('feedback') capable of balancing the radiative losses in central dominant early-type galaxies, are evident via the interaction of radio-loud AGN activity and the surrounding hot X-ray gas, leading to excavation of cavities (e.g., Mc-Namara et al. 2005; McNamara & Nulsen 2007). Their low masses mean that galaxy groups are the environment in which AGN feedback mechanism may have the greatest impact on the evolution of the hot diffuse gas.

It has been suggested that in most groups and clusters, feedback operates in a near-continuous 'bubbling' mode where thermal regulation is relatively gentle (Bîrzan et al. 2008, 2012; Panagoulia et al. 2014). However, there are cases that AGN feedback can also manifest via extreme AGN outbursts, potentially shutting down the central engine for long periods (e.g., in Hydra A, MS 0735+7421, NGC 4261, NGC 193 & IC 4296; Rafferty et al. 2006; Gitti, Brighenti & McNamara 2012; Gitti et al. 2007; O'Sullivan et al. 2011; Kolokythas et al. 2015, 2018; Grossova et al. 2019). Such powerful outbursts may have a significant effect on the development of groups, and it is therefore vital to study the energy output of radio galaxies in an unbiased set of low-mass systems in order to understand the impact of AGN in the group regime.

It is well known that galaxies at the centres of evolved groups are generally early-types, lacking strong ongoing star formation (e.g., Vaddi et al. 2016), with their properties depending strongly on the group halo mass (De Lucia & Blaizot 2007; Skibba & Sheth 2009; Gozaliasl et al. 2018), and the galaxy density (i.e., compact vs. loose groups). Brightest group early-type galaxies (BGEs) as well as brightest cluster galaxies (BCGs), exhibit different surface brightness profiles and scaling relations from field or satellite massive galaxies (e.g., Graham et al. 1996; Bernardi et al. 2007) and play an important role in the evolution of galaxy groups. They are also found to exhibit different morphologies, star formation rates, radio emission, and AGN properties compared to satellite galaxies of the same stellar mass. Their privileged position near or at the centres of the extended X-ray emitting hot intra-group medium (IGM) and dark matter halo makes them suitable laboratories for constraining cosmological models, for studying the growth history of massive galaxies (e.g., von der Linden et al. 2007; Liu et al. 2009; Stott et al. 2010; Liang et al. 2016), their super-massive black holes (e.g., Rafferty et al. 2006; Sabater et al. 2019), and the connection between galaxy properties such as stellar kinematics and the larger environment (Loubser et al. 2018).

The radio detection rate of massive galaxies is found to be $\sim 30\%$ (e.g., Best et al. 2005; Shabala et al. 2008), but this is dependent on the sensitivity of the survey used and the redshift range considered. Studies of the BGEs/BCGS of groups and clusters in the local Universe finding detection rates of $\sim 80-90\%$ (e.g., Magliocchetti & Brüggen 2007; Dunn et al. 2010; Kolokythas et al. 2018) and this can even rise to 100% for the most massive systems (Sabater et al. 2019). It has been suggested that, in addition to fu-

elling of AGN by cooling from the IGM, the connection between radio AGN and higher galaxy densities (Lilly et al. 2009; Bardelli et al. 2010; Malavasi et al. 2015) may be driven by large scale merging (e.g., infall of groups into clusters) or by 'inter-group' galaxygalaxy interactions and mergers (Miles et al. 2004; Taylor & Babul 2005). In such mergers and interactions gas can be channelled to the central AGN resulting in radio emission and the launching of jets. However, although AGNs dominate the radio emission from these massive galaxies, star formation may contribute in the less radio luminous objects, whose radio morphologies are often unresolved (e.g., Smolčić, et al. 2017).

In this paper we present results from the study of the radio properties of the dominant galaxies of the 27-group low-richness subset of the Complete Local-Volume Groups Sample (CLoGS), including new Giant Metrewave Radio Telescope (GMRT) 235 and 610 MHz observations of 25 systems. The CLoGS sample and the X-ray properties of the high-richness groups are described in more detail in O'Sullivan et al. (2017, hereafter Paper I) while the radio properties of the BGEs in the high-richness sub-sample are described in Kolokythas et al. (2018, hereafter Paper II). This paper continues the work described in paper II, presenting the properties of the central radio sources in the low-richness sub-sample and discussing the whole CLoGS sample for the first time. We examine the connection between the group environment and the dominant radio galaxies, the contribution of star formation on the radio emission of pointlike radio sources, and provide a qualitative comparison of the radio emission that BGEs exhibit in high and lowrichness group sub-samples. The paper is organized as follows: In Section 2 we briefly present the sample of galaxy groups, in Section 3 we describe the GMRT observations and the radio data analysis, in Section 4 we present the radio detection statistics of the BGEs, and in Section 5 their radio properties including information on the contribution of star formation on the radio emission. Section 6 contains the discussion of our results for the CLoGS sample as a whole, focusing on the detection statistics, the properties of the radio sources, their environment and their energetics. The summary and the conclusions are given in Section 7. Radio images and information on the central galaxies of this sample work are presented in Appendices A to C. Throughout the paper we adopt the Λ CDM cosmology with $H_o = 70 \text{ km s}^{-1} \text{ Mpc}^{-1}$, $\Omega_m = 0.27$, and $\Omega_{\Lambda} =$ 0.73. The radio spectral index α is defined as $S_{\nu} \propto \nu^{-\alpha}$, where S_{ν} is the flux density at the frequency ν . In general we quote 1σ uncertainties, and the uncertainties on our radio flux density measurements are described in Section 3.2.

2 THE COMPLETE LOCAL-VOLUME GROUPS SAMPLE

The Complete Local-Volume Groups Sample (CLoGS) is an optically selected sample of 53 groups in the local universe (D \leq 80 Mpc), that is collected from the relatively not deep, all-sky Lyon Galaxy Group catalog (LGG; Garcia 1993). Paper I provides a detailed description of the sample, its selection criteria, and the X-ray properties of the high–richness sub-sample, and we therefore only provide a brief summary of the selection here. The sample is statistically complete in the sense that, to the completeness limit of the LGG sample, it contains every group which meets our selection criteria. It is intended to be a representative survey of groups in the local universe including studies of their radio, X-ray and optical properties.

Constraints on the selection of CLoGS groups include: i) a

membership of at least 4 galaxies, ii) the presence of ≥ 1 luminous early-type galaxy ($L_B > 3 \times 10^{10} L_{\odot}$) and iii) a declination of $>-30^{\circ}$ in order to certify that the groups are observable from the GMRT and Very Large Array (VLA). The group membership was extended and refined using the HyperLEDA catalog (Paturel et al. 2003), based on which the group mean velocity dispersion and richness R parameter were estimated (R; number of member galaxies with log $L_B \ge 10.2$). Systems with R > 10 were not included in the sample as they were already known galaxy clusters, and systems with R = 1 were also excluded, as they were too poor to provide results on their physical parameters that would be reliable. From this process, we obtained a 53-group statistically complete sample, that was divided into two sub-samples: i) the 26 high-richness groups with R = 4-8 (see Paper II) and ii) the 27 low-richness groups with R = 2-3.

Distances to the CLoGS group-dominant galaxies are estimated from their recession velocities, corrected for Virgocentric flow. Inaccuracies in these distances will affect calculated values such as luminosities, but since we mainly consider relations between quantities presented on logarithmic scales, these differences will be small and will not significantly alter our findings (e.g., a distance error of 20 per cent would give a luminosity error of only 0.2 dex).

3 **OBSERVATIONS AND DATA ANALYSIS**

GMRT observations 3.1

Excluding six systems (NGC 315, NGC 524, NGC 1407, NGC 3325, NGC 3665 and NGC 4697; see Table 1) for which archival data were available, the low-richness sample of galaxy groups were observed using the GMRT in dual 235/610 MHz frequency mode during observing cycle 21, from 2011 Nov - 2012 Aug. Each target was observed at both 235 and 610 MHz with the upper side band correlator (USB) for an average of \sim 4 hours on source. The total observing bandwidth at both frequencies is 32 MHz with the effective bandwidth at 235 MHz being ~ 16 MHz.

At 610 MHz the data were obtained in 512 channels with a spectral resolution of 65.1 kHz for each channel, whereas at 235 MHz the data were obtained in 256 channels and a spectral resolution of 130.2 kHz for each channel. A detailed summary of the observations can be found in Table 1.

3.2 Data analysis

The data were processed using the SPAM pipeline¹ (Interna 2014). SPAM is a PYTHON based extension to the NRAO Astronomical Image Processing System (AIPS) package which includes directiondependent calibration, radio frequency interference (RFI) mitigation schemes, along with imaging and ionospheric modeling, adjusting for the dispersive delay in the ionosphere. We provide here only a brief description of the SPAM pipeline to account for the data analysis procedure followed. For more details regarding the SPAM pipeline and the algorithms of the SPAM package see Interna et al. (2009, 2017).

The SPAM pipeline is run in two parts. In the first, preprocessing stage, the pipeline converts the raw LTA (Long Term Accumulation) format data collected from the observations into

pre-calibrated visibility data sets for the total of the pointings observed (UVFITS format). The second stage converts these precalibrated visibility data of each pointing into a final Stokes I image (FITS format), via several repeated steps of (self)calibration, flagging, and wide-field imaging.

The final full resolution images, corrected for the GMRT primary beam pattern, provide a field of view of $\sim 1.2^{\circ} \times 1.2^{\circ}$ at 610 MHz and $\sim 3^{\circ} \times 3^{\circ}$ at 235 MHz having a mean sensitivity for the data we observed (1 σ noise level) of ~0.09 mJy/beam at 610 MHz and \sim 0.64 mJy/beam at 235 MHz (see Table 1). The measured sensitivities achieved from the analysis at both frequencies are in line with expectations from previous GMRT experience. The theoretical (thermal) noise values for our observations are 29 μ Jy/beam for 610 MHz and 80 μ Jy/beam for 235 MHz². We note that the theoretical sensitivity is dependent on the square root of the time on source, and that there is therefore a variation in the sensitivity between different targets at the same frequency, as well as a difference in quality between the older archival hardware correlator data and the newer observations using the software correlator. The main source of the residual noise in our final images comes either from calibration uncertainties in the form of phase errors from rapidly varying ionospheric delays (especially at the lowest frequencies) or from dynamic range limitation due to limited data quality, calibration and image reconstruction, that is mostly revealed by the presence of bright sources in the field. We also note that the full resolution of the GMRT is $\sim 6^{\prime\prime}$ at 610 MHz and $\sim 13^{\prime\prime}$ at 235 MHz with the *u-v* range at 610 MHz being $\sim 0.1 - 50$ k λ and at 235 MHz $\sim 0.05-21$ k λ .

The flux density scale in the images was set from the available flux calibrators in each observing session (3C 48, 3C 286 and 3C 147) using the models from Scaife & Heald (2012). We adopt a flux density uncertainty of 5% at 610 MHz and 8% at 235 MHz, representing the residual amplitude calibration errors (Chandra et al. 2004).

Out of the 27 BGEs in the CLoGS low-richness sub-sample, 25 were analyzed in this study at both 610 and 235 MHz using the SPAM pipeline. The GMRT observations and images for NGC 315 at both 235 and 610 MHz are drawn from the earlier study of Giacintucci et al. (2011) and for NGC 1407 from the detailed follow-up work of Giacintucci et al. (2012), with the analysis of the observations being described in detail in those studies. In addition, we note that the GMRT data for NGC 5903 system were also analysed by the standard procedure followed as described in Kolokythas et al. (2015, 2018) and the results are found to be consistent with results from the SPAM pipeline output.

1400 MHz data were drawn primarily from the NRAO VLA Sky Survey (NVSS, Condon et al. 1998) and the study of Brown et al. (2011), that included measurements from the NVSS, Green Bank Telescope (GBT), and Parkes Radio Telescope, and from Condon et al. (2002) for NGC 252, NGC 1106 and NGC 5127.

RADIO DETECTIONS OF BRIGHTEST GROUP 4 GALAXIES IN THE LOW-RICHNESS SAMPLE

The radio sources detected in the BGEs of each group were identified from the 235 and 610 MHz GMRT images and the available NVSS catalog data. After that, we determined their morphology,

 $^{^1\,}$ For more information on how to download and run SPAM see http: //www.intema.nl/doku.php?id=huibintemaspam

² Calculated using the rms noise sensitivity equation in $\S2.1.6$ from the GMRT Observer's Manual: http://gmrt.ncra.tifr.res.in/gmrt_hpage/Users /doc/manual/Manual_2013/manual_20Sep2013.pdf

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Table 1. Details of the new GMRT observations analyzed for this study, along with information on archival data analyzed by the authors or by Giacintucci et al. (2011, marked *a*) and Giacintucci et al. (2012, marked *b*). For each source the first line displays the details for the 610 MHz and the second line for the 235 MHz. The columns give the LGG (Lyon Groups of Galaxies) number for each group, the BGE name, observation date, frequency, time on source, beam parameters and the rms noise in the resulting images.

Group Name	BGE	Observation	Frequency	On source	Beam, P.A.	rms
LGG		Date	(MHz)	Time (minutes)	(Full array, $'' \times '', \circ$)	(mJy beam ⁻¹)
		0011.37	(10	4.55	x aa x aa x a a x a a x a a x a a x	0.14
6	NGC 128	2011 Nov	610	157	$5.32 \times 4.06, -2.96$	0.16
10		2011 Nov	235	157	$16.15 \times 12.07, 74.34$	1.90
12	NGC 252	2011 Nov	610	171	$5.46 \times 4.78, -44.72$	0.06
		2011 Nov	235	171	$16.16 \times 14.17, 77.73$	0.60
78	NGC 1106	2011 Nov	610	171	$5.44 \times 4.08, -2.47$	0.10
		2011 Nov	235	171	$14.83 \times 10.31, -9.06$	0.60
97	NGC 1395	2011 Nov	610	168	$9.59 \times 3.80, 37.03$	0.10
		2011 Nov	235	168	$28.74 \times 9.20, 39.36$	0.70
113	NGC 1550	2011 Dec	610	187	$5.66 \times 3.96, 88.52$	0.04
		2011 Dec	235	187	$14.28 \times 10.65, 88.65$	0.45
126	NGC 1779	2011 Dec	610	124	$4.68 \times 3.77, 34.90$	0.05
		2011 Dec	235	124	$12.67 \times 9.68, 25.13$	0.45
138	NGC 2292	2011 Dec	610	137	$7.58 \times 3.59, 36.19$	0.05
		2011 Dec	235	137	$20.11 \times 9.22.31.86$	0.42
167	NGC 2768	2012 Jan	610	209	$6.26 \times 3.65.49.03$	0.03
107	1100 2700	2012 Jan	235	209	16.31×9.49 36.10	0.30
177	NGC 2011	2012 Jan 2012 Jan	610	196	$5.22 \times 3.70, 60.15$	0.50 0.10^{1}
177	NOC 2911	2012 Jan	235	190	$11.58 \times 10.21, 30.02$	0.10
222	NCC 2612	2012 Jan	233	190	$11.00 \times 10.21, 39.92$	0.30
232	NGC 5015	2012 Jan	010	212	5.79 × 5.05, 55.90	0.03
255	NGG 2022	2012 Jan	235	212	16.09 × 8.77, 22.22	0.26
255	NGC 3923	2012 Mar	610	156	$7.33 \times 4.02, -16.06$	0.05
		2012 Mar	235	156	$20.77 \times 10.24, -10.22$	0.45
329	NGC 4956	2012 Apr	610	199	$5.24 \times 3.96, 24.60$	0.05
		2012 Apr	235	199	$12.54 \times 12.04, 11.54$	0.48
341	NGC 5061	2012 May	610	200	$6.54 \times 3.68, 28.89$	0.05
		2012 May	235	200	$22.42 \times 9.75, 38.61$	0.50
350	NGC 5127	2012 Apr	610	200	$6.56 \times 4.42, -30.97$	0.25
		2012 Apr	235	200	$13.51 \times 11.68, -27.35$	0.65
360	NGC 5322	2012 May	610	199	$5.07 \times 4.00, -0.06$	0.05
		2012 May	235	199	$15.30 \times 11.31, -49.08$	0.70
370	NGC 5444	2012 May	610	160	$4.42 \times 3.99, 62.22$	0.20
		2012 May	235	160	$12.90 \times 10.79, -61.85$	0.80
376	NGC 5490	2012 May	610	160	$4.14 \times 3.55, 34.12$	0.25
		2012 May	235	160	$11.88 \times 8.69.35.44$	1.85
383	NGC 5629	2012 May	610	227	$4.95 \times 3.96.80.80$	0.05
202	1100002	2012 May	235	227	12.58×10.71 83.01	0.40
308	NGC 5903	2012 May 2012 May	610	181	$5.78 \times 3.81.30.02$	0.40
570	100 5705	2012 May	235	181	$3.10 \times 0.01, 32.02$ $20.10 \times 0.54, 45.07$	0.00
457	NGC 7252	2012 May	233 610	100	$5 18 \times 4.26 20.12$	0.00
437	NOC 7252	2012 Aug	225	199	$14.95 \times 10.02 \ 10.52$	0.10
462	NCC 7277	2012 Aug	255	199	14.25 × 10.05, 19.32	0.00
403	NGC /3//	2012 Aug	610	217	5.41 × 3.92, 20.75	0.04
		2012 Aug	235	217	$15.68 \times 9.97, 17.05$	0.40
Archival data						
14	NGC 315 ^a	2008 Feb	610	380	$5.20 \times 5.00, 61.00$	0.10
		2008 Aug	235	280	$15.0 \times 15.0, 0.00$	0.70
23	NGC 524	2006 Mar	610	248	$5.52 \times 4.15, 89.20$	0.18
		2006 Mar	235	248	$11.79 \times 9.89,66.70$	2.10
100	NGC 1407 ^b	2009 Nov	610	270	$8.20 \times 4.40, 42.00$	0.05
		2009 Nov	235	270	$16.10 \times 10.90.36.00$	0.25
205	NGC 3325	2004 May	610	63	$10.52 \times 3.78,60.46$	0.31
200		2004 May	235	63	$32.57 \times 12.12.57.37$	38.0
236	NGC 3665	2009 Feb	610	248	5.08×4.10 35.38	0.17
250	1100 5005	2007 Feb	225	240	$12.00 \times 4.10, 55.50$ $12.04 \times 0.52, 22.00$	0.78
214	NGC 4607	2009 FCU	235 610	240	12.04 × 3.00, 32.99 5 11 × 1 16 - 71 50	0.78
514	1100 4097	2000 Jan	010	341	$0.44 \times 4.40, -/4.38$	0.07
		2006 Jan	235	347	$14.00 \times 11.32, -64.90$	1.14

 1 The rms around the source area here is 0.2 mJy

calculated their flux densities (see Table 2) and their largest linear size. The GMRT images of the CLoGS low–richness groups at both frequencies along with more detailed information on these sources can be found in Appendix A and B.

We find a radio detection rate of 82% (22 of 27 BGEs) for the low-richness CLoGS sample, considering both GMRT frequencies and the 1.4 GHz surveys, with five galaxies being undetected at any of our three radio frequencies (NGC 2292, NGC 3325, NGC 4956, NGC 5444 and NGC 5629).

Of these 5 undetected systems, NGC 5444 has previously been identified with a 4C source, but our higher resolution data shows that it is not related to the BGE (for more details see Section A20 and Figure B20 for the radio image). NGC 3325 has a very high upper limit due to the high noise level from the analysis of the short archival observations available (see Table 2).

If we take into consideration the detection statistics only from the GMRT data, we find that 70% of the BGEs are detected at 235 and/or 610 MHz (19 of 27 BGEs) with 3 BGEs (NGC 128, NGC 3613 and NGC 4697) being detected only at 1.4 GHz data. Figure 1, presents the flux density distribution of all 53 CLoGS BGEs along with their upper limits (both high and low–richness sub-samples) at 610 MHz with figure 2 presenting the distribution at 235 MHz. We observe that the majority of the BGEs have flux densities in the range 10 - 100 mJy, but that the 610 MHz flux density distribution of the low–richness sub-sample is shifted towards lower flux densities compared to the high–richness systems.

An upper limit of $5 \times$ the r.m.s. noise level in each image is used for the undetected radio sources, with the limits falling in the range $\sim 0.3-10$ mJy for both sub-samples and GMRT frequency ranges. However, while in figure 1 at 610 MHz we observe a similar number of radio non-detected BGEs (upper limits) between both sub-samples, from figure 2 at 235 MHz the number of radio non detected BGEs in the low–richness sub-sample is found to be almost twice than the number of non detections in the high–richness sub– sample.

As in paper II, the limiting sensitivity of our sample is estimated based on the typical noise level of our images and the maximum distance for the observed groups, which is 78 Mpc. The mean r.m.s from the analysis of the low-richness sub-sample (excluding the archival data) is ~90 μ Jy beam⁻¹ at 610 MHz and ~640 μ Jy beam⁻¹ at 235 MHz. This means that the limiting sensitivity in the low-richness sub-sample using the automated pipeline SPAM is $> 3.3 \times 10^{20}$ W Hz⁻¹ at 610 MHz or $> 2.3 \times 10^{21}$ W Hz⁻¹ at 235 MHz which is similar to the sensitivity power achieved for the high-richness sub-sample (> 2.9×10^{20} W Hz⁻¹ at 610 MHz or > 2.2×10^{21} W Hz⁻¹ at 235 MHz; Paper II). We note here that the limiting sensitivities are representing on average the total low-richness sub-sample and there may be individual radio sources detected at $>5\sigma$ level of significance at lower powers in nearby groups. The equivalent limit for NVSS 1400 MHz power sensitivity at CLoGS groups distance and level of significance for comparison is $> 1.7 \times 10^{21}$ W Hz⁻¹.

5 RADIO PROPERTIES OF THE BRIGHTEST GROUP GALAXIES

We examine here the radio properties of the detected radio sources in the central galaxies of the low–richness sample. As in the high– richness sample, the sources present a range in size, power and morphology. The GMRT images of the groups along with more details for these sources are presented in Appendices A and B.



Figure 1. Flux density distribution of the CLoGS BGEs at 610 MHz, comparing the low-richness sub-sample (grey columns) with the high-richness systems (cyan solid line) from Paper II. The upper limits of the non detected BGEs in the low-richness sub-sample are shown in magenta dot-dashed line whereas the equivalent upper limits in the high-richness sub-sample are shown in blue dashed line.



Figure 2. Same as in figure 1, but for the flux density distribution at 235 MHz.

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Table 2. Radio flux densities, radio power and spectral indices of the sources in the low–richness sub-sample. The columns list the BGE name, redshift, flux density of each source at 235 and 610 MHz, the 235–610 MHz spectral index, the flux density at 1.4 GHz (drawn from the literature), the 235–1400 MHz spectral index, and the radio power at 235, 610 and 1400 MHz. All upper limits shown here from our analysis are $5 \times r.m.s$. Five galaxies show no radio emission detected at 235, 610 or 1400 MHz. The references for the 1.4 GHz flux densities and the GMRT measurements from previous works are listed at the bottom of the table.

Source	Redshift (z)	$\begin{array}{c} \mathrm{S}_{235MHz} \\ \pm 8\% (\mathrm{mJy}) \end{array}$	$\begin{array}{c} \mathbf{S}_{610MHz} \\ \pm 5\% \text{ (mJy)} \end{array}$	$\substack{ \alpha^{610MHz}_{235MHz} \\ (\pm 0.04) }$	S _{1.4GHz} (mJy)	$\alpha^{1400MHz}_{235MHz}$	$\frac{P_{235MHz}}{(10^{23}{\rm WHz}^{-1})}$	$\frac{P_{610MHz}}{(10^{23}{\rm WHz}^{-1})}$	$\frac{P_{1.4GHz}}{(10^{23}{\rm WHz}^{-1})}$
NGC 128	0.0141	<9.5	<0.8	_	1.5 ± 0.5^{a}	_	<0.041	<0.003	0.006
NGC 252	0.0165	≤3.0	0.9	-	2.5^{b}	-	≪0.019	0.006	0.016
NGC 1106	0.0145	435.0	221.0	0.71	132 ± 4^{b}	$0.67 {\pm} 0.04$	2.122	1.078	0.645
NGC 1395	0.0057	≤3.5	3.6	-	$1.1 {\pm} 0.5^{a}$	-	≤0.002	0.002	0.0006
NGC 1550	0.0124	223.0	62.0	1.34	17 ± 2^a	$1.44 {\pm} 0.06$	0.752	0.209	0.057
NGC 1779	0.0111	4.6	2.6	0.60	$5.4 {\pm} 0.6^{c}$	-0.09 ± 0.05	0.011	0.006	0.013
NGC 2292	0.0068	≤2.1	≤0.3	-	-	-	≤0.002	≤0.0003	-
NGC 2768	0.0045	13.1	11.5	0.14	14 ± 1^a	-0.04 ± 0.07	0.008	0.007	0.009
NGC 2911	0.0106	69.0	58.0	0.18	56 ± 2^a	$0.12 {\pm} 0.04$	0.166	0.139	0.135
NGC 3613	0.0068	≤1.3	≼0.2	-	$0.3 {\pm} 0.3^{a}$	-	≤0.002	≤0.0002	0.0004
NGC 3923	0.0058	5.0	2.1	0.91	1 ± 0.5^a	$0.90 {\pm} 0.22$	0.002	0.001	0.0005
NGC 4956	0.0158	≤2.4	≤0.3	-	-	-	≤ 0.014	≤ 0.002	-
NGC 5061	0.0069	≤2.5	0.8	-	-	-	≤0.002	0.001	-
NGC 5127	0.0162	5690	1630	1.32	1980^{b}	$0.59 {\pm} 0.04$	35.464	10.159	12.278
NGC 5322	0.0059	134.6	106.0	0.25	79.3 ± 2.8^{c}	$0.30 {\pm} 0.04$	0.135	0.106	0.080
NGC 5444	0.0131	$\leqslant 12^d$	$\leqslant 1^d$	-	- *	-	≤0.051	≤ 0.004	-
NGC 5490	0.0162	1140	815	0.36	1300 ± 100^{a}	-0.07 ± 0.05	6.804	4.864	7.839
NGC 5629	0.0150	≤2.0	≤0.3	-	-	-	≤0.011	≤ 0.002	-
NGC 5903	0.0086	1830	970	0.68	321.5 ± 16.0^d	$0.99 {\pm} 0.04$	2.829	1.500	0.498
NGC 7252	0.0160	76.6	41.5	0.66	25.3 ± 1.2^{c}	$0.63 {\pm} 0.04$	0.397	0.215	0.132
NGC 7377	0.0111	2.7	2.1	0.27	3.4 ± 0.5^c	-0.13 ± 0.07	0.007	0.005	0.009
Archival data/	Previous wo	rk							
NGC 315 ^e	0.0165	15411	≥2500	≤ 1.91	1800 ± 100^{a}	$1.20 {\pm} 0.04$	99.691	16.172	11.474
NGC 524	0.0080	≤10.5	2.0	-	$3.1{\pm}0.4^{a}$	-	≼0.015	0.003	0.004
NGC 1407 ^f	0.0059	945	194	1.66	38 ± 2^{f}	$1.80 {\pm} 0.04$	0.603	0.124	0.024
NGC 3325	0.0189	≼191	≤1.55	-	-	-	≤1.449	≤0.012	-
NGC 3665	0.0069	225.2	149.0	0.43	113.2 ± 3.8^{c}	$0.39 {\pm} 0.04$	0.275	0.182	0.139
NGC 4697	0.0041	≼5.7	≤0.35	-	0.6 ± 0.5^{a}	-	≤0.002	≤0.0001	0.0002

^a Brown et al. (2011), ^b Condon et al. (2002), ^c Condon et al. (1998), ^d Upper limit was calculated from the maximum peak flux density on the position of the BGE as the phase calibration errors are too large due to a very bright 4C radio source nearby. The detection of the system in this case is limited by the dynamical range due to the very strong background source right next to the position of the BGE, ^e Giacintucci et al. (2011), ^f Giacintucci et al. (2012)
 *Condon et al. (1998) (NVSS) and Brown et al. (2011) report for NGC 5444 a flux density at 1.4 GHz of ~660±20 which is a 4C nearby source

5.1 Radio morphology

The radio structures found in the central galaxies have extents ranging from a few kpc (galactic scale) to hundreds of kpc (group scale). Table 3 lists the radio morphology for each source. We apply the radio morphology scheme described in Paper II, classifying our radio sources from our own GMRT 235 and 610 MHz analysis, the 1.4 GHz NVSS survey and from earlier studies on some of our BGEs. The purpose of this classification is to distinguish between the different classes of radio emission and provide information on their environment. Systems may belong to more than one category; when that is the case, the system is classified according to its most extended and prominent state/class.

The morphological classes are:

- (i) point-like (unresolved) sources,
- (ii) diffuse sources with no clear jet-lobe structure,
- (iii) small-scale jets (<20 kpc),
- (iv) large-scale jets (>20 kpc), and
- (v) non-detections.

A subset of the large-scale jet sources have been shown by previous studies to be *remnant* systems, where the AGN is quiescent, but emission from the aging jets or lobes is still visible.

Table 3 summarizes the radio properties of the BGEs in the low–richness sample, showing their largest radio linear size, their radio morphology, their power at 235 MHz (see § 5.3 below for more details) and an estimate of the energy injected into the IGM by the inflation of radio lobes by jet sources ($P_{\rm cav}$) calculated from the radio power at 235 MHz using the scaling relation from O'Sullivan et al. (2011, see also §6.4):

$$\log P_{cav} = 0.71 \,(\pm 0.11) \log P_{235} + 1.26 \,(\pm 0.12), \qquad (1)$$

where $P_{\rm cav}$ and P_{235} are in units of $10^{42}~erg~s^{-1}$ and $10^{24}~W~Hz^{-1}.$

We find that in the low–richness sample 14 central galaxies exhibit point-like radio emission, 2 host small-scale jets, 5 host large-scale jets (of which 1, NGC 1407, is a remnant jet; see below for details) with 5 BGEs being undetected at the frequencies we investigated, and only 1 system presenting a diffuse radio structure.

NGC 3665 and NGC 5322 are the two systems that exhibit a small-scale jet morphology and both present symmetrical structures and thin, straight jets (see Figures B13 and B19) with similar sizes (see Table 3). The four currently-active large-scale jet systems are all Fanaroff-Riley type I (FR I; Fanaroff & Riley 1974) radio galaxies with their jet/lobe components extending from several tens of kpc (e.g., 33 kpc; NGC 1550) to several hundreds of kpc (1200 kpc; NGC 315) away from their host galaxy.

Looking at the large-scale jet systems in more detail, NGC 1550 presents an unusually asymmetric FR I radio morphology with the eastern lobe roughly half as far from the optical centre of the galaxy as its western counterpart (see Figure B6). Our GMRT data provides a considerably clearer view of the radio morphology than previous studies (e.g., Dunn et al. 2010), revealing the previously unresolved lobes and a sharp z-bend in the western jet which produces an offset between the east and west jet axes. We will discuss this system in greater detail in a forthcoming paper (Kolokythas et al., in prep.).

In NGC 5490 we also observe an asymmetric morphology with the eastern jet detected to \sim 50 kpc while the opposite side of the source breaks up into detached clumps of radio emission (see Figure B21). NGC 5127 on the other hand, presents a normal symmetric double jet morphology (see Figure B18) with NGC 315 hosting the giant FR I radio galaxy B2 0055+30 which has been extensively investigated at multiple frequencies and angular resolutions (e.g., Bridle et al. 1976, 1979; Venturi et al. 1993; Mack et al. 1997; Worrall et al. 2007; Giacintucci et al. 2011). NGC 315 presents two asymmetric jets, with one (northwest) being brighter and appearing bent backwards and the opposite (southeast) jet appearing much fainter and intermittent with the radio source having a total linear size of ~1200 kpc (see Giacintucci et al. 2011).

NGC 1407 is the only radio galaxy in the low–richness sample classed as remnant radio source, as the radio spectral age analysis in the study of Giacintucci et al. (2012) revealed that a faint \sim 300 Myr old, ultra-steep spectrum ($\alpha = 1.8$) radio plasma of \sim 80 kpc surrounds the central jet source, which is a typical product of former AGN activity and characteristic of a dying radio galaxy.

The only diffuse radio source in the low–richness sample, appears in the galaxy group NGC 5903 (see Figure B23). It has also been examined by several earlier studies at many different frequencies (e.g., Gopal-Krishna 1978; Gopal-Krishna et al. 2012; O'Sullivan et al. 2018). It hosts a ~75 kpc wide diffuse, steep-spectrum ($\alpha_{150}^{612} \sim 1.03$; O'Sullivan et al. 2018) radio source whose origin may be through a combination of AGN activity and violent galaxy interactions (see A23). We note that NGC 5903 could also be categorised as a remnant system, but due to its complex origin and current state we consider it more conservative to class this system as diffuse.

5.2 Radio Spectral index

Where possible, we estimate the spectral indices of each radio source in the frequency ranges 235–610 MHz and 235–1400 MHz. The radio spectral index of a source will, over time, steepen owing to synchrotron and inverse Compton losses, provided that there is no new source of electrons and no additional energy injection. Spectral index is therefore, in the absence of complicating factors, an indicator of source age.

Examining the radio sources in the low–richness sample, we find that only just over half the BGEs (15/27 galaxies; 56%) are detected at both GMRT frequencies at 235 and 610 MHz. At this frequency range we find that four sources present steep radio spec-

tra of $\alpha_{235}^{610} > 1$ (NGC 1550, NGC 5127, NGC 315 and NGC 1407) with NGC 315 and NGC 1407 presenting ultra-steep spectra of $\alpha_{235}^{610} \ge 1.91$ and $\alpha_{235}^{610} = 1.66$ respectively (see Table 2). However, in the study of Giacintucci et al. (2011) the flux density of NGC 315 at 610 MHz is quoted as underestimated, therefore the spectral index calculated must be considered an upper limit. For the same galaxy, Mack et al. (1997), using the WSRT at the same frequency (609 MHz) reported a flux density of 5.3 Jy, thus reducing the spectral index value of the radio source in NGC 315 to $\alpha_{235}^{609} = 1.11$. For NGC 1407, Giacintucci et al. (2012) reported that the radio spectral index value at the 235–610 MHz range was calculated after subtracting contributions in the flux densities at both frequencies from an inner double (young jets associated with the NGC 1407 AGN) and two point radio sources whose positions fall within the large-scale diffuse emission of the old radio lobes.

The rest of the 11 radio sources present α_{235}^{610} that ranges from very flat values of ~0.1 to typical radio synchrotron spectra of ~0.9. We find that the total mean spectral index value at the GMRT frequency range, leaving the two overestimated steep spectrum outliers out (NGC 315 and NGC 1407), is $\alpha_{235}^{610} = 0.60 \pm 0.16$ (for 13/15 radio sources).

Examining the radio spectral index distribution between 235 and 1400 MHz we find that four radio sources (NGC 1779, NGC 2768, NGC 5490 and NGC 7377), three of which exhibit a relatively flat spectral index in the 235–610 MHz range (all except NGC 1779 which has $\alpha_{235}^{610} = 0.60$), have flux densities greater at 1400 MHz than at 235 MHz, giving an inverted spectral index in the 235–1400 MHz range (see Table 2). Three of these four systems present weak point-like radio sources at both frequencies (NGC 1779, NGC 2768 and NGC 7377), hence the deviation from a powerlaw of the observed flux density can be attributed to selfabsorption of the lower frequency emission at 235 MHz as a result of the 'cosmic conspiracy' (see Cotton et al. 1980).

For the FR I radio jet system in NGC 5490, the observed inverted spectral index in the 235–1400 MHz range ($\alpha_{235}^{1400} =$ -0.07) arises primarily from the difference in detected morphology between the two frequencies. The bright source close to NGC 5490 limits the sensitivity we were able to achieve at 235 MHz and probably prevented us from detecting all the corresponding radio emission from the source that has been picked up at 1400 MHz. The mean value of α_{235}^{1400} that we calculate for the 15/27 BGEs in the low-richness sample is 0.57 ± 0.16.

In Table 4 we list the mean values of the radio spectral indices for the point-like, small-scale and large-scale jet radio morphologies in the low-richness sample. We neglected to include values for remnant jets and diffuse sources since these classes consist of only 1 system each. We find that large-scale jet systems present the steepest mean indices ($\alpha_{235}^{610} = 1.23 \pm 0.08$ and $\alpha_{235}^{1400} = 0.79 \pm 0.08$) in the low-richness sample, with smallscale jet systems exhibiting the flattest spectra, having similar values at both frequency ranges we examined. The point-like radio sources are found to exhibit typical radio spectrum indices with the mean values being comparable within the uncertainties at both frequency ranges. In large-scale jet systems, the steeper mean spectral index seen in the 235-610 MHz frequency range can be attributed to low number statistics and detection sensitivity between the two frequency ranges, as NGC 315 presents an upper limit for $lpha_{235}^{610}$ due to the flux density underestimation at 610 MHz (Giacintucci et al. 2011) and NGC 5490 presents an inverted spectrum at 235-1400 MHz reducing the mean at this frequency range (see below for more on NGC 5490).

For the systems analysed in this paper where it was possi-

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ble to separate their extended emission from the main core at both 235 and 610 MHz, we calculated the spectral indices for their core and the extended components separately. We created images with matched resolution, cellsize, uv range and image pixel size at these two frequencies. Table 5 lists the spectral index values for the 4 systems that meet these criteria. In NGC 3665 we find that the spectral index does not differ between the two components. For two systems (NGC 1550 and NGC 5127) the spectral index in their cores is flatter, indicating either aging along their jets, self-absorption, or the presence of free-free emission. For NGC 5490 we find that the core (which is clearly detected) presents a fairly typical spectral index value of 0.55 ± 0.04 but that the extended emission has a very flat value of 0.02 ± 0.04 . The extended low-surface brightness emission presents different morphology between 235 MHz and 610 MHz, so the flat spectral index we measure is likely biased, and probably does not reflect the true index of the extended component. A more accurate measurement would require higher signal-to-noise ratio data and low resolution images to trace the extended emission at both frequencies.

5.3 Radio power in CLoGS BGEs

As in paper II, we calculate the radio power P_{ν} , for the low-richness sample BGEs as:

$$P_{\nu} = 4\pi D^2 (1+z)^{(\alpha-1)} S_{\nu}, \qquad (2)$$

where D is the distance to the source, α is the spectral index in the 235–610 MHz regime, z the redshift and S_{ν} is the flux density of the source at frequency ν . For the systems that were detected at only one frequency the typical spectral index value of $\alpha = 0.8$ was used for the calculation (Condon 1992).

We find that in the low-richness sample the radio power of the detected radio sources in the BGEs is in the range $10^{19} - 10^{25}$ W Hz^{-1} . Figure 3, presents the radio power distribution of the BGEs from the low-richness sub-sample at both 235 and 610 MHz. The majority of the galaxies in figure 3 have low radio powers in the range between $10^{20} - 10^{21}$ W Hz⁻¹ and $10^{22} - 10^{23}$ W Hz⁻¹ with only one BGE exhibiting power in the range $10^{21}\,-\,10^{22}$ W Hz⁻¹ (NGC 1779). The CLoGS low-richness sample contains two high-power sources with $P_{235MHz} > 10^{24} \text{ W Hz}^{-1}$, both of which host large-scale bright radio jets (NGC 315 and NGC 5127). The radio power upper limits of the non detected radio sources at 610 MHz extend to lower values $(10^{19} - 10^{20} \text{ W Hz}^{-1})$ than at 235 MHz, with the majority of the radio power upper limits at both frequencies seen between $\sim 5 \times 10^{19} - 5 \times 10^{21} \text{ W Hz}^{-1}$. Only NGC 3325 presents an upper limit of an order of magnitude higher $(\sim 10^{23} \text{ W Hz}^{-1})$ compared to rest of the BGEs, due to the very high noise level in the image produced from the available archival data. We note that previously reported typical values of radio power at 235 and 610 MHz in BCGs range between $\sim 10^{23} - 5 \times 10^{26}$ W Hz^{-1} (see e.g. Table 2, Yuan et al. 2016). The equivalent radio power statistics for the full sample will be discussed in \S 6.

5.3.1 Radio-loudness in low-richness CLoGS sample BGEs

Following Best et al. (2005), we consider as radio-loud systems with $P_{1.4GHz} > 10^{23}$ W Hz⁻¹. In the low–richness sample we find that three BGEs have radio power higher than this value: NGC 315, NGC 5127 and NGC 5490 (see Table 6). While NGC 5903 and NGC 1106 have radio powers greater than 10^{23} W Hz⁻¹ at



Figure 3. Radio power distribution of CLoGS BGEs in low-richness sample at 610 MHz (cyan columns) and 235 MHz (solid magenta line), with the upper limits at 610 MHz (blue dashed line) and at 235 MHz (black dotted line) also shown here.

610 MHz their corresponding flux density at 1.4 GHz is lower than the radio-loud limit, and we do not include them in the radio-loud category. In addition, the sensitivity limit of the NVSS survey and Brown et al. (2011) study guarantees that all the undetected systems at 1.4 GHz have radio powers below this limit. In § 6 we discuss the equivalent statistics for the radio-loud systems in the full CLoGS sample.

5.3.2 235 MHz radio power vs largest linear size

It is well known that the linear size of FR I radio galaxies is proportional to their radio power (Ledlow et al. 2002). As in Paper II, we examine the relation between the 235 MHz power of our sources against their largest linear size (LLS). The relation was examined for the radio sources that were morphologically resolved, with their largest linear size being calculated across the longest extent of the detected radio emission, at the frequency at which the emission is most extended.

The resolved radio sources of the CLoGS low–richness central galaxies span a broad spatial scale from ~15 kpc (small scale jets; NGC 5322) to > 1200 kpc (large scale jets; NGC 315) with the corresponding 235 MHz radio powers being in the range of ~ 10^{21} W Hz⁻¹ to ~ 10^{25} W Hz⁻¹.

Figure 4 shows the relation for the full CLoGS sample, along with the group systems from the study of Giacintucci et al. (2011). We find that our CLoGS radio sources are in agreement with the linear correlation between size and power found by Ledlow et al. (2002) and Giacintucci et al. (2011). In addition, our CLoGS BGEs range of luminosities and size distribution is in line with the measured values of radio sources from the recent LOFAR AGN study of Hardcastle et al. (2016) at 150 MHz over a greater redshift range (0-0.8) (see their fig. 18). We find that there is no significant dif-

Table 3. Morphological properties of CLoGS low-richness groups and their central radio sources. For each group we note the LGG number, the BGE name, the angular scale, the largest linear size (LLS) of the radio source, measured from the 235 MHz radio images unless stated otherwise, the radio morphology class, and lastly the energy output of any radio jets, estimated from the 235 MHz power using equation 1. The errors on the energy output were calculated based on the errors from scaling relation 1.

Group	BGE	Scale (kpc/")	LLS (kpc)	Radio morphology	Energy output (radio) $(10^{42} \text{ erg s}^{-1})$
LGG 6	NGC 128	0.291	$\leqslant 3^a$	point	-
LGG 12	NGC 252	0.349	≼3	point	-
LGG 14	NGC 315	0.354	1200^{b}	large-scale jet	$93.13^{+64.97}_{-38.27}$
LGG 23	NGC 524	0.165	$\leqslant 2^{c}$	point	-
LGG 78	NGC 1106	0.310	≤14	point	-
LGG 97	NGC 1395	0.102	≪3	point	-
LGG 100	NGC 1407	0.112	80^b	remnant	-
LGG 113	NGC 1550	0.257	33	large-scale jet	$2.89^{+0.02}_{-0.03}$
LGG 126	NGC 1779	0.218	≪4	point	-
LGG 138	NGC 2292	0.145	-	-	-
LGG 167	NGC 2768	0.112	≪3	point	-
LGG 177	NGC 2911	0.218	≤7	point	-
LGG 205	NGC 3325	0.388	-	-	-
LGG 232	NGC 3613	0.156	$\leqslant 3^a$	point	-
LGG 236	NGC 3665	0.156	16	small-scale jet	$1.42^{+0.16}_{-0.18}$
LGG 255	NGC 3923	0.097	≼3	point	-
LGG 314	NGC 4697	0.087	$\leqslant 3^a$	point	-
LGG 329	NGC 4956	0.344	-	-	-
LGG 341	NGC 5061	0.136	$\leqslant 2$	point	-
LGG 350	NGC 5127	0.349	244	large-scale jet	$44.74^{+23.06}_{-15.22}$
LGG 360	NGC 5322	0.141	15^c	small-scale jet	$0.86^{+0.15}_{-0.19}$
LGG 370	NGC 5444	0.291	-	-	-
LGG 376	NGC 5490	0.344	120 (30 at 610 MHz)	large-scale jet	$13.84^{+3.65}_{-2.89}$
LGG 383	NGC 5629	0.325	-	-	-
LGG 398	NGC 5903	0.175	65^d	diffuse	-
LGG 457	NGC 7252	0.320	≤14	point	-
LGG 463	NGC 7377	0.223	≼5	point	-

^a Measured from the 1.4 GHz image, ^b Giacintucci et al. (2011), ^c Measured from the 610 MHz image, ^d O'Sullivan et al. (2018)

Table 4. Mean spectral indices α_{235}^{610} and α_{235}^{1400} for the different radio morphology classes in the low–richness sample. The last column shows the number of sources used in calculating the means.

Radio Morphology	$\mathrm{Mean}\ \alpha^{610}_{235}$	$\mathrm{Mean} \ \alpha^{1400}_{235}$	No of sources
Point-like	$0.50 {\pm} 0.11$	$0.29 {\pm} 0.11$	7
Small-scale jet	$0.34{\pm}0.06$	$0.35 {\pm} 0.06$	2
Large-scale jet	$1.23{\pm}0.08$	$0.79{\pm}0.08$	4

ference in the spread of our group radio sources between the low and high–richness sample and confirm across group and cluster environments the linear correlation for the radio sources is valid and stands for almost 4 orders of magnitude at 235 MHz.

5.4 Star formation contribution in CLoGS radio sources

Although the source engine of radio jet systems is most certainly an AGN, it is not clear whether that is the case for the majority of the radio sources in the low–richness sample, which present unresolved (point-like) radio morphology. Unless a luminous radio source is present ($P_{1.4GHz} > 10^{23}$ W Hz⁻¹), it is possible in these systems to confuse a compact central star forming region with AGN emission, as the beam sizes correspond to areas a few kiloparsecs across. Therefore, we examine the possibility of stellar contribution **Table 5.** Spectral indices in the 235–610 MHz range for the cores and extended emission of those radio sources in our sample whose morphology is resolved at both frequencies. Columns show the group LGG number, the BGE name, and the spectral index α_{235}^{610} of the core and the surrounding emission in each source.

Group BGE		Core α_{235}^{610}	Surrounding α_{235}^{610}
LGG 113 LGG 236 LGG 350 LGG 376	NGC 1550 NGC 3665 NGC 5127 NGC 5490	0.94 ± 0.04 0.39 ± 0.04 0.77 ± 0.04 0.55 ± 0.04	$\begin{array}{c} 1.36{\pm}0.04\\ 0.44{\pm}0.04\\ 1.33{\pm}0.04\\ 0.02{\pm}0.04\end{array}$

to the radio emission of the low-richness sample BGEs and in particular for those systems which exhibit low radio luminosities. We note here that in NGC 5903, the only system in the low-richness sample that presents a diffuse radio structure, the radio emission likely arises from a combination of AGN jet activity and galaxy interactions driven, and star formation cannot be a significant contributor to the radio luminosity of this system (see O'Sullivan et al. 2018).

We estimate the expected contribution to the radio flux density from star formation using two methods. Firstly, we use Far-Ultraviolet (FUV) fluxes from the Galaxy Evolution Survey GR6



Figure 4. Radio power at 235 MHz plotted against the largest linear size for the full sample of CLoGS radio sources. Different symbols indicate the radio morphology of the radio sources. The points circumscribed by square box are representing BGEs from the high–richness sample, while those without the box are from the low–richness one. Open diamonds indicate systems in the Giacintucci et al. (2011) sample, for comparison.



Figure 5. Radio power at 1.4 GHz in relation to molecular gas mass for the 14 point-like BGEs in the low–richness sample. Black triangles indicate systems with detected molecular gas and rhombus with arrows indicates systems with 3σ upper limits on molecular gas mass. The dashed red line represents the expected radio emission from star formation, assuming the molecular gas mass to star formation rate scaling relation of Gao & Solomon (2004).

Table 6. Radio-loud AGN ($P_{1.4GHz} > 10^{23}$ W Hz⁻¹) in the BGEs of the CLoGS low-richness sample. Columns show the LGG number, the BGE name and the radio power at 1.4 GHz.

Group	BGE	${P_{1.4GHz}\over (10^{23}{ m WHz^{-1}})}$
LGG 14	NGC 315	11.5 ± 0.6
LGG 350	NGC 5127	12.3 ± 0.6
LGG 376	NGC 5490	$7.8 {\pm} 0.6$

catalog³ (Martin 2005) to estimate the star formation rate (SFR) from the calibration of Salim et al. (2007). We then calculate the expected 610 MHz radio emission from star formation at that rate using the relation of Garn et al. (2009) (more details on the star formation in CLoGS BGEs will be presented in a future paper; Kolokythas et al., in prep.).

Table 7 lists the calculated SFR from FUV and the estimated contribution to 610 MHz radio power, for the galaxies where star formation could potentially affect the radio flux density measurements in the low-richness sample (point-like systems). The only exception is NGC 1779, where no FUV flux was available. We consider that star formation dominates the radio emission if the expected radio power at 610 MHz of a BGE is >50% of its actual P_{610MHz} value. We find that in 6/13 systems, the luminosity from the detected radio emission exceeds by over an order of magnitude the expected radio luminosity from the FUV SFR, making these galaxies AGN dominated. In three galaxies, (NGC 1395, NGC 3613 and NGC 7377) we find that star formation may contribute 20-40% in the radio emission, in two galaxies may contribute significantly (50-60%; NGC 3923 and NGC 5061) and in two systems (NGC 252 and NGC 4697) it appears to be the dominant source of radio emission with the expected radio luminosity from the FUV SFR exceeding or being similar to the observed radio luminosity. A plausible explanation to why the expected radio luminosity from the FUV SFR exceeds the observed radio luminosity, is that different phases of star formation can be traced by radio and FUV wavelengths - FUV from the largest young stars, radio from supernovae - therefore, a change in SFR over time could lead to a disagreement.

However, using as star formation estimate the FUV wavelength we probe star formation from only young massive stars and cannot account for internal extinction due to dust. In addition, the calibration from Salim et al. (2007) is only valid when the SFR remains constant over the lifetime of the UV-emitting stars $(<10^{8} \text{ yr})$. Thus the expected radio from SFR_{FUV} can hold for systems that do not present extreme recent starbursts or dust obscured star forming regions. NGC 7252 is likely an example of dust-obscured star formation in a post-starburst merger (see A24). We find SFR_{*FUV*} \sim 0.04 M_{\odot} yr⁻¹, considerably less than the SFR estimated from the integrated far infrared (FIR) flux (8.1 M_{\odot} yr⁻¹, O'Sullivan et al. 2015). However, as noted by George et al. (2018), radio and H α SFR estimates for NGC 7252 (Schweizer & Seitzer 2013) agree within a factor of ~ 2 with the FIR rate, and the implied degree of extinction is not unreasonable for such a gas-rich system. Modelling of the star formation history of the galaxy suggests peak SFRs of tens of M_{\odot} yr⁻¹ (Chien & Barnes 2010) within the last 200-600 Myr, and the discrepancy may also indicate a more recent decline in SFR, with the FUV tracing the youngest stars. While the

³ http://galex.stsci.edu/GR6/?page = mastform

Table 7. Star formation rates calculated from FUV, and expected radio power at 610 MHz owing to star formation for the CLoGS low–richness sample pointlike radio sources. For each system we list the SFR_{FUV}, the expected radio power at 610 MHz $P_{610 expected}$ from the calculated SFR_{FUV}, the radio power at 610 MHz and 1.4 GHz ($P_{610 MHz}$, $P_{1.4 GHz}$), the molecular mass values (M_{H2}) drawn from O'Sullivan et al. (2018b) and the radio morphology of the source.

Group LGG	BGE	$\frac{\text{SFR}_{FUV}}{(10^{-2} \text{ M}_{\odot} \text{ yr}^{-1})}$	$\frac{P_{610\ expected}}{(10^{21}\ W\ Hz^{-1})}$	$P_{610 MHz} (10^{21} \mathrm{W Hz^{-1}})$	$P_{1.4 GHz} (10^{21} \text{ W Hz}^{-1})$	${M_{H2} \over (10^8 \ M_{\odot})}$	Morphology
6	NGC 128	4.9	0.13	1.25 ^a	0.64	1.76±1.5	point
12	NGC 252	30.7	0.69	0.60	1.55	$6.29 {\pm} 0.79$	point
23	NGC 524	2.7	0.05	0.30	0.43	$1.90{\pm}0.23$	point
78	NGC 1106	26.7	0.94	107.80	64.50	6.81 ± 0.33	point
97	NGC 1395	4.2	0.07	0.20	0.06	< 0.27	point
126	NGC 1779	-	-	-	1.31	$4.57 {\pm} 0.60$	point
167	NGC 2768	2.6	0.06	0.70	0.89	$0.18 {\pm} 0.01$	point
177	NGC 2911	3.2	0.11	13.90	13.53	$2.66 {\pm} 0.31$	point
232	NGC 3613	1.4	0.02	0.07^{a}	0.04	< 0.46	point
255	NGC 3923	4.2	0.06	0.10	0.05	< 0.29	point
314	NGC 4697	3.1	0.04	0.05^{a}	0.02	< 0.07	point
341	NGC 5061	2.7	0.04	0.08	0.04^{b}	< 0.43	point
457	NGC 7252	36.8	1.29	21.50	13.16	58.00 ± 8.70	point
463	NGC 7377	6.0	0.13	0.50	0.86	$4.74 {\pm} 0.44$	point

^{*a*} The P_{610 MHz} was calculated by extrapolating the 610 MHz flux density from the available 1.4 GHz emission, using a spectral index of 0.8 ^{*b*} The P_{1.4 GHz} was calculated by extrapolating the 1400 MHz flux density from the available 610 MHz emission, using a spectral index of 0.8

high radio/FUV ratio might have been an indicator of a radio AGN, the low limit on the AGN X-ray luminosity (Nolan et al. 2004) and its radio morphology strongly suggest the radio emission from this system is dominated by star formation rather than an active nucleus (though there is evidence that the opposite was true in the past, Schweizer & Seitzer 2013).

In order to test whether the star formation contribution in the radio emission of point-like BGEs is consistent with typical SFRs in star forming galaxies, we also examine the expected radio power from star formation at 1.4 GHz, by converting molecular gas mass to SFR using the scaling relation of Gao & Solomon (2004). Following O'Sullivan et al. (2018b) in figure 5 we show the relation between the radio power at 1.4 GHz and the molecular gas mass (M_{H2}) for the 14 point-like radio BGEs (see also Table 7). The dashed red line indicates the expected radio emission from star formation for the measured M_{H2} . We find that along with NGC 7252, five more BGEs (NGC 128, NGC 252, NGC 524, NGC 1779 and NGC7377) fall close to the line for the expected radio emission from star formation along with the systems that present upper limits in molecular gas mass. The systems with upper limits (NGC 1395, NGC 3613, NGC 3923, NGC 4697 and NGC 5061) are seen about the line, indicating that they may be star formation dominated or have a significant contribution from star formation to their radio emission. In NGC 4697 the expected radio emission from SFR_{FUV} is >50% hence we consider that this galaxy's radio emission is star formation dominated. The radio emission from three BGEs (NGC 1106, NGC 2768 and NGC 2911) is clearly far above the line, suggesting that their radio emission is AGN dominated.

While estimating the expected SFR from the molecular gas mass is not a direct measurement like that provided by the FUV flux, we find that both methods suggest that three systems are AGN dominated: NGC 1106, NGC 2768 and NGC 2911. We find inconsistent results between the two methods for NGC 524, NGC 7252 and NGC 128, which were classed as AGN dominated from their SFR_{FUV}, but both methods produce similar results for the systems where star formation may make a significant contribute to the radio emission. We know that NGC 7252 is a starburst system and we

 Table 8. Radio morphology occurrence for the BGEs in total CLOGS sample. The numbers for the high–richness sample are drawn from Paper II.

Radio Morphology	Low-richness	High-richness	CLoGS total
Point-like	52% (14/27)	54% (14/26)	53% (28/53)
Non-detection	19% (5/27)	8% (2/26)	13% (7/53)
Large-scale jet	15% (4/27)	8% (2/26)	11% (6/53)
Diffuse emission	4% (1/27)	14% (4/26)	9% (5/53)
Small-scale jet	7% (2/27)	8% (2/26)	8% (4/53)
Remnant jet	4% (1/27)	8% (2/26)	6% (3/53)
Overall	82% (22/27)	92% (24/26)	87% (46/53)

hence we consider its radio emission star formation dominated as well as in NGC 524 as is shown by figure 5. For NGC 128, where P_{610MHz} was estimated by the 1.4 GHz emission using a spectral index of 0.8, we consider that star formation may contribute significantly to the radio emission as is indicated by figure 5.

We therefore in summary, consider that in the low–richness sample, star formation most likely dominates the radio emission in 5/14 of the BGEs with point-like radio emission. It is possible that star formation may have a significant impact on our measurements of radio properties in six more (NGC 128 and the five BGEs with upper limits on molecular gas mass). We find that with certainty, the radio emission is AGN dominated in 3/14 point-like BGEs.

6 DISCUSSION

6.1 Detection statistics in CLoGS

Brightest group and cluster galaxies (BGGs and BCGs) in the local Universe are more likely to be detected in radio than the wider galaxy population. However, matching group and cluster samples to perform a comparison one has to be prudent, as very often different selection criteria have been applied (redshift, luminosity etc). The most recent statistical study by Hogan et al. (2015) using different X-ray selected cluster samples including systems out to redshift z~0.4, reports a radio detection percentage for BCGs of $60.3 \pm 7.7\%$ (from the ROSAT ESO Flux Limited X-ray/Sydney University Molonglo Sky Survey (REFLEX-SUMSS) sample, Böhringer et al. 2004), $62.6\pm5.5\%$ (REFLEX-NVSS sample), and $61.1 \pm 5.5\%$ from the extended Brightest Cluster Sample (eBCS; Ebeling et al. 2000). BCGs in the X-ray selected galaxy clusters of Ma et al. (2013) and Kale et al. (2015) (from the Extended Giant Metrewave Radio Telescope -GMRT- Radio Halo Survey; EGRHS; Venturi et al. 2007, 2008) present a radio detection rate of ~52% and ~48% respectively, from the NVSS (> 3 mJy) and NVSS/FIRST catalogs.

In what may be the most relevant studies for CLoGS, Dunn et al. (2010) used a sample of nearby early-type galaxies, most of which are in the centres of groups or clusters, and found a detection rate of $\sim 81\%$ (34/42) using NVSS and SUMSS data. Bharadwaj et al. (2014) report a $\sim 77\%$ (20/26) detection rate for the X-ray selected groups sample of Eckmiller et al. (2011) using the NVSS, SUMSS and VLA Low frequency Sky Survey (VLSS, 74 MHz) radio catalogs.

In the CLoGS low-richness sample, considering only the 1.4 GHz data, we find a detection rate of \sim 78% (21/27 BGEs) with a similar detection rate of $\sim 81\%$ (21/26 BGEs) being found for the high-richness sample in Paper II. In total, the detection percentage at 1.4 GHz for the complete CLoGS sample is found to be \sim 79% (42/53 BGEs) which is in good agreement with the general trend of the high detection rate of brightest galaxies in the local Universe found by the studies mentioned above. We also note that CLoGS brightest group galaxies present a higher detection rate compared to those from clusters, as the average detection rate value for the BCGs from Hogan et al. (2015) is $\sim 61.3 \pm 11\%$. However, we note here that the latter study included cluster systems at a much greater volume ($z\sim0.4$) than our CLoGS groups ($\sim20x$ farther) along with systems from the southern hemisphere, where the detection limit from SUMSS is ~ 6 mJy ($\sim 2.5x$ higher than NVSS). Hence, we treat this only as indicative of the relative detection statistics between groups and clusters.

The total radio detection rate for the CLoGS sample at any of the three radio frequencies (235, 610 and 1400 MHz) is \sim 87% (46/53 BGEs) which drops to \sim 79% (42/53 BGEs) if we consider only the GMRT data. In a comparison between the CLoGS two group sub-samples, the radio detection rate was found to be slightly higher for the high–richness sample (24/26 BGEs, \sim 92%; Paper II) than for the low–richness groups (\sim 82% 22/27 BGEs), but the difference is not statistically significant.

Examining the radio morphology appearance in CLoGS sample as a total, we find that more than half of the radio sources (\sim 53%; 28/53) present point-like radio emission, followed by the BGEs with non-detections (\sim 13%; 7/53). Large-scale jets are present in \sim 11% (6/53) of the BGEs, \sim 9% (5/53) host a diffuse radio source, \sim 8% (4/53) a small-scale jet, and \sim 6% (3/53) a remnant of an old radio outburst. In Table 8 we show the occurrence of radio morphology in total and in both sub-samples. While both sub-samples present a similar rate for small-scale and remnant jet morphologies, large-scale jet emission appears at higher rate in high-richness sample, with the most prominent difference between the two sub-samples being the higher rate of undetected radio sources presented in the low–richness sample (\sim 19%) compared to the high–richness one (\sim 8%).

We find that, in total, 6/53 of our CLoGS BGEs can be considered as radio-loud following the definition by Best et al. (2005) $(P_{1.4GHz} > 10^{23} \text{ W Hz}^{-1})$. This gives a percentage of ~11% radio loud galaxies for CLoGS groups which is similar to the ~13%



Figure 6. Spectral index α of CLoGS different radio morphologies in the frequency range 235–610 MHz, in relation to radio power at 235 MHz (P_{235MHz}). The radio morphology of each source is indicated by the symbols, with the groups from the high–richness sample being marked by open squares.

fraction of radio loud galaxies found in low-mass clusters and groups ($10^{13} < M_{200} < 10^{14.2}$) by Lin & Mohr (2007).

We find that ~9 \pm 4% (5/53) of the CLoGS BGEs are currently hosts of double-lobed radio galaxies (note that this excludes the lobeless double-jet system NGC 5490). This is slightly higher than the ~8 \pm 5% (2/26; Paper I) found from the high–richness sub-sample but comparable within uncertainties. Huang & Chen (2010) find that the equivalent percentage of BCGs having double-lobed radio morphology, using a catalog of ~13000 clusters from the maxBCG sample, is ~4%. It should also be noted that the remnant radio sources in our sample (NGC 1167, NGC 1407 and NGC 5044) all appear to be the remains of large-scale jet-lobe systems. Although this could suggest that double-lobed radio sources are more common in galaxy groups than clusters, this result needs to be treated with caution, since it is unclear how sensitive the cluster samples are to such steep spectrum remnant sources.

6.2 Spectral index of CLoGS BGEs

The spectral index of a radio source provides insight on the age and the nature of the mechanisms that give rise to radio emission (Lisenfeld & Völk 2000). In AGN, the most important mechanisms at high frequencies are inverse Compton losses and synchrotron emission which is produced by relativistic electrons spiralling down magnetic field lines. The cores of AGN usually present a flat radio spectrum as a result of self-absorption at low frequencies rather than a flat electron energy distribution, with a steep spectrum being observed from the extended emission (lobes) of an AGN.

For extended radio galaxies that lie in the centre of coolcore clusters and groups, the spectral index distribution is generally found to be 1 - 2 in the 235–610 MHz frequency range (Giacintucci et al. 2011). The spectral indices of these sources are thought to be steep owing to the confinement of the relativistic radio plasma by the surrounding intra-cluster medium (e.g., Fanti et al. 1995; Murgia et al. 1999). Bornancini et al. (2010), using a sample of maxBCG clusters and radio data from the Faint Images of the Radio Sky at Twenty-cm survey (FIRST; 1400 MHz), the NVSS (1400 MHz), the Westerbork Northern Sky Survey (WENSS; 325 MHz) and the Green Bank 6 cm Survey (GB6; 4850 MHz), calculated the mean value of the spectral index for BCGs between 325 MHz and 1.4 GHz to be $\alpha_{325}^{1400} = 0.65$. However, the authors did not select by radio morphology, and the inclusion in the spectral index calculation of radio point sources as well as jet-lobe systems probably explains the flatter mean index, showing that different morphological types of radio sources have an impact on the mean value of the spectral index.

In the CLoGS sample we are able to measure the 235-610 MHz spectral index in 34/53 BGEs. The mean spectral index for these radio sources is $\alpha_{235}^{610} = 0.68 \pm 0.23$, with a standard deviation of 0.55, whereas in the 235–1400 MHz frequency range, the mean radio spectral index for 33/53 BGEs is $\alpha_{235}^{1400} =$ 0.59 ± 0.26 with a standard deviation of 0.43. The uncertainties on the 1400 MHz flux for two galaxies, NGC 3923 and NGC 5982, were excluded from the error calculation of the mean α_{235}^{1400} as their high uncertainties provided for their low flux densities by Brown et al. (2011) are equivalent to their value (≤ 1 mJy see Table 2 and Paper II) which was increasing the average error on the mean by over a factor of 2.

The mean 235–610 MHz spectral index for CLoGS sample is significantly flatter than the value found by Giacintucci et al. (2011), probably owing to the inclusion of sources with a wide range of morphologies in our study. If we take into consideration the spectral indices only from the CLoGS large–scale and remnant jet systems in order to match with the extended group radio sources examined by the sample of Giacintucci et al. (2011), we find that $\alpha_{235}^{610} = 1.13 \pm 0.12$, in much better agreement. The mean α_{235}^{1400} of 0.59 ± 0.26 found for all our CLoGS systems is very close to the mean spectral index found by Bornancini et al. (2010) ($\alpha_{325}^{1400} = 0.65$) using a similar frequency range.

In the recent study of de Gasperin et al. (2018), using data extracted from the NVSS (1400 MHz) and a re-imaged version of the TGSS (TIFR GMRT Sky Survey, 147 MHz; Intema et al. 2017), the weighted average spectral index distribution α_{147}^{1400} for over a million extra-galactic radio sources was found to be 0.79 with a standard deviation of 0.24 (for $S_{\nu} \propto \nu^{-\alpha}$). This is in agreement with previous results on extra-galactic radio sources which found mean spectral indices α_{147}^{1400} in the range 0.75 – 0.8 (Ishwara-Chandra et al. 2010; Intema et al. 2011). Although these studies include radio sources from various environments and redshifts, the CLoGS mean spectral index is consistent (within uncertainties) with their mean spectral indices.

Table 9 shows the mean spectral index values in the 235–610 MHz and 235–1400 MHz bands for the different radio morphology classes in the full CLoGS sample. We find that point-like systems present a mean spectral index of $\alpha_{235}^{610} = 0.47 \pm 0.16$, diffuse radio sources a value of $\alpha_{235}^{610} = 0.65 \pm 0.09$ with jet systems (both small and large scale jets) presenting a steeper index of $\alpha_{235}^{610} = 0.81 \pm 0.13$. The mean 235-1400 MHz spectral index in jet systems is $\alpha_{235}^{1400} = 0.62 \pm 0.14$ with the mean value for point-like systems being $\alpha_{235}^{1400} = 0.43 \pm 0.26$.

As expected, jet systems show a somewhat steeper mean spectral index than point-like sources. In figure 6 we present the spectral index α_{235}^{610} distribution of CLoGS BGEs, in relation to their radio power at 235 MHz (P_{235MHz}). We observe that point-like systems fall mainly about the calculated mean spectral index distribution. We also find that 50% (3/6) of the large–scale jet systems exhibit steep α_{235}^{610} values with the other half having values close to the mean. While small–scale jet systems present small deviations in their spectral index values, we see that large–scale systems are divided between flat and very steep spectral α_{235}^{610} indices.

Table 9. Mean spectral indices α_{235}^{610} and α_{235}^{1400} for the different radio morphology classes in the 53-group CLoGS sample.

Radio Morphology	Mean α^{610}_{235}	Mean α^{1400}_{235}
Point-like	$0.47 {\pm} 0.16$	$0.43 {\pm} 0.26$
Small-scale jet	$0.52 {\pm} 0.08$	$0.47 {\pm} 0.08$
Remnant jet	$1.38 {\pm} 0.07$	1.11 ± 0.07
Large-scale jet	$1.00 {\pm} 0.10$	$0.72 {\pm} 0.11$
Diffuse emission	$0.65 {\pm} 0.09$	$0.72 {\pm} 0.10$

This divide in the large-scale jet sources might be explained by the different evolutionary stages of the jet systems (younger sources with flatter spectral indices, vs. older or remnant sources with aged electron populations) or by the properties of the environment that the jets propagate into (e.g., bending lobes or confinement of electrons due to higher density environment). From figure 6 we also find a set of six radio sources that present a very steep spectrum, (> 1), in the 235–610 MHz frequency range, with two of them being naturally remnant radio sources, which as expected present the largest deviation from the mean in our sample, another three sources being large–scale jet systems and the remaining one being of diffuse radio morphology.

Interna et al. (2011) categorized as ultra-steep spectrum (USS) those sources with spectral indices steeper than 1.3, and found that the fraction of USS found in the 153–1400 MHz frequency range is 3.8% (16/417). This is in good agreement with earlier studies (e.g., 3.7%; Sirothia et al. 2009). In the CLoGS sample, we find that only 2 radio sources present a spectral index steeper than 1.3 in the 235–1400 MHz frequency range (NGC 1407 and NGC 1550). This gives a USS fraction in good agreement with larger samples (3.8%), but the result is strongly dependent on the detection limit at 235 MHz.

6.3 Environmental properties

We investigate the relation between the radio morphology of a BGE and its near environment by examining the fraction of spiral galaxies in each group. The number of spiral galaxies that a group possesses is known to be connected to its dynamical age, as older systems will have had more opportunities for galaxy interaction and/or merging to drive morphological transformation.

We classify as spiral-rich (i.e, dynamically young) galaxy groups with a spiral fraction $F_{sp} > 0.75$ (Bitsakis et al. 2010). In Table 10 we present the fraction of spiral galaxies in each group for the low–richness sample. The morphologies of the galaxies were drawn from the HyperLEDA⁴ catalogue, with morphological T-type <0 galaxies being classified as early-type, and T-type ≥ 0 (or unknown) galaxies classed as late-type. As in Paper II, we define spiral fraction F_{sp} as the number of late-type galaxies over the total number of group members.

We find that $19/27 (70\pm9\%)$ of the groups in the low-richness sample are spiral-rich and $8/27 (30\pm9\%)$ are spiral-poor, whereas for the high-richness sample in Paper II we found that 11/26 $(42\pm10\%)$ of the groups are spiral-rich and $15/26 (58\pm10\%)$ spiral-poor. Declining spiral fraction with richness may be an indication that galaxies undergo gradual morphological transformation if rich groups emerge over an extended period of time, or it may be a manifestation of galaxy downsizing, where galaxy-sized

⁴ http://leda.univ-lyon1.fr/



Figure 7. Spiral fraction F_{sp} plotted against 235 MHz radio power, with symbols indicating radio morphology of the BGE. Groups from the high-richness sample are marked with open squares.

density peaks riding on larger-scale density fluctuations destined to form rich groups collapse early and rapidly to form preferentially spheroidal galaxies.

Examining the CLoGS sample in total, 30/53 (57%) of the groups are spiral-rich and 23/53 (43%) are spiral-poor. Bitsakis et al. (2014), using a sample of 28 Hickson compact groups (HCGs, Hickson 1992), found that 46% (13/28) were spiral-rich and 54% (15/28) spiral-poor. Although Bitsakis et al. (2014) used a sample of compact groups which are likely to have higher galaxy interaction rates and thus evolve more rapidly than our systems, the spiral-rich fraction found in Bitsakis et al. (2014) is similar to our high–richness sample. This again demonstrates that richer groups are a conducive environment for galaxy evolution.

In Fig 7 we show the relation between F_{sp} and the radio power at 235 MHz (P_{235MHz}) for all CLoGS systems whose BGEs are detected at 235 MHz. For the low–richness sample we find that all but one of the point-like systems appear in galaxy groups with higher spiral fractions (6/7; 86%), with a mean spiral fraction of 0.84±0.14. Jet systems in the low-richness sub-sample are also found to reside more often in a spiral-rich environment (4/6; 67%) having a mean spiral fraction of 0.77±0.17.

Looking at CLoGS groups as a whole, we find that the majority (75%) of radio point-like systems detected at 235 MHz (12/16) reside in groups with high spiral fractions (dynamically young) with a mean spiral fraction of 0.78 ± 0.10 . Jet systems on the other hand, are found to be equally distributed between spiral-rich (5/10; 50%) and spiral-poor environments (5/10; 50%). However, we find the mean F_{SP} for jet sources (0.73 \pm 0.14) is similar to that for point radio sources. Given the large uncertainties, this result cannot be regarded as statistically significant and larger sample of groups are required in order to confirm it. However, the even appearance of jet systems in both spiral rich and poor groups suggests that the dynamical evolution of the galaxy population of a group has little influence on the radio jet activity of the BGE. This is of interest given the connection between jets and intra-group medium properties found in paper I, where we found that more massive groups were more likely to have declining central temperature profiles and that the presence of jets was correlated with these cool cores and other evidence of rapid gas cooling. We might have expected that more dynamically evolved (i.e., spiral-poor) groups, which are likely to

Table 10. Spiral fraction for the low–richness CLoGS groups. Spiral fraction (F_{sp}) is defined as the ratio of the number of late-type and unknown morphology galaxies to the total number of group members.

LGG	BGE	Spiral fraction
	F_{sp}	
LGG 6	NGC 128	0.71
LGG 12	NGC 252	0.67
LGG 14	NGC 315	0.70
LGG 23	NGC 524	0.64
LGG 78	NGC 1106	0.87
LGG 97	NGC 1395	0.79
LGG 100	NGC 1407	0.50
LGG 113	NGC 1550	0.78
LGG 126	NGC 1779	0.92
LGG 138	NGC 2292	0.64
LGG 167	NGC 2768	0.89
LGG 177	NGC 2911	0.96
LGG 205	NGC 3325	0.96
LGG 232	NGC 3613	0.90
LGG 236	NGC 3665	0.86
LGG 255	NGC 3923	0.78
LGG 314	NGC 4697	0.87
LGG 329	NGC 4956	0.90
LGG 341	NGC 5061	0.86
LGG 350	NGC 5127	0.95
LGG 360	NGC 5322	0.84
LGG 370	NGC 5444	0.79
LGG 376	NGC 5490	0.50
LGG 383	NGC 5629	0.95
LGG 398	NGC 5903	0.78
LGG 457	NGC 7252	0.89
LGG 463	NGC 7377	0.60

be more massive, would more commonly host central jet sources. Of the remnant jet sources, 2/3 (NGC 5044 and NGC 1407) reside in dynamically old groups ($F_{SP} \sim 0.5$) while the third, NGC 1167, is the BGE of an X-ray faint group with a high spiral fraction (0.75). Diffuse radio sources, similar to jet systems, occur in groups with a wide range of spiral fractions (0.57-0.91) with 3/5 residing in high spiral fraction groups. Our X-ray analysis of the low-richness sample is not yet complete, and will be reported in a later paper, but our preliminary analysis suggests that 10/13 jet systems in the full CLoGS sample are hosted by X-ray bright groups with group-scale X-ray halos. As reported in paper I, NGC 1167 is the BGE of an X-ray faint group with little or no hot gas, and in the low-richness sample we find two jet sources in systems with only galaxy-scale gas halos. This suggests that even in low mass, less evolved groups, a hot gas rich environment is preferable for jets (e.g., McNamara & Nulsen 2007). We examine the properties of jet-hosting BGEs in more detail in the next section (\S 6.4).

6.4 Power output and properties of jet activity in CLoGS

As in Paper II, the energy output of jet sources in the low-richness sample is estimated using the $P_{cav} - P_{235}$ scaling relation by O'Sullivan et al. (2011). This relation estimates the cavity power of a jet system based on the BGE's radio power at 235 MHz allowing the use of radio data as a proxy to estimate the AGN mechanical power output of radio jet systems. It has the benefit of providing with an estimate of the mechanical power output in the absence of direct calculations of cavity power (P_{cav}) from the X-rays. Table 3 shows the results from the calculations of the radio energy output. We find that the energy output values for the jet systems in the low–richness sample range from $\sim 10^{42} - 10^{44}$ erg s⁻¹, which is common for galaxy groups, but roughly an order of magnitude greater than the range found for the high–richness sample. The giant radio galaxy NGC 315 produces the difference in the upper bound of the range, and given the limited scale of our sample, the low and high–richness sub-samples are likely comparable.

For the BGEs in the CLoGS sample that present either an ongoing jet radio emission (large and small-scale jets; 10/53) or a past one (remnant jets; 3/53) we examine whether a correlation exists between the mass of the central black hole (BH) and their radio activity at 235 MHz. We calculate the BH masses for all CLoGS systems using the correlation of Ferrarese & Merritt (2000), estimating the mass of the BH from the central velocity dispersion of the host galaxy. This correlation is expressed as:

$$log_{10}\left(\frac{M_{BH}}{M_{\odot}}\right) = \alpha + blog_{10}\left(\frac{\sigma_o}{200 \ km \ s^{-1}}\right) \tag{3}$$

where M_{BH} is the mass of the black hole in solar masses (M_{\odot}) and σ_o is the central velocity dispersion. The central velocities dispersions for 23/53 of our BGEs were drawn from the recent kinematic study of Loubser et al. (2018) with the values for the remaining 30 obtained from the HyperLeda⁵ catalogue. We estimated the BH masses for all CLoGS BGEs (see Table C1) using the values for the parameters a = 8.39 ± 0.06 and b = 5.20 ± 0.36 given by the fit of the velocity dispersion–black hole mass relation in McConnell & Ma (2013).

We find that the BGE systems with 235 MHz detections (34/53) have BH masses in the range of ${\sim}10^8-5{\times}10^9~M_{\odot}$, with the majority of the systems having masses ${\leqslant}10^9~M_{\odot}$. Figure 8 shows the relation between the radio power at 235 MHz (P_{235MHz}) and the BH mass for different radio morphologies. Jet systems present in total a BH mass range of $0.2{\times}10^9-5{\times}10^9~M_{\odot}$, with an average of ${\sim}2.3{\pm}2.7{\times}10^9~M_{\odot}$ for the large–scale systems and ${\sim}1.07{\pm}1.16{\times}10^9~M_{\odot}$ for the small–scale.

We observe that 4/6 large–scale jet systems present a BH mass $>2\times10^9$ M_☉ with the hosts of 2/6 large–scale systems exhibiting a BH mass an order of magnitude less (NGC 193 and NGC 5127; $\sim 2\times10^8$ M_☉). The most extended large–scale radio jet system in CLoGS (~ 1200 kpc; NGC 315) presents a BH mass estimate of $\sim 3.8\times10^9$ M_☉. The recent study of Horellou et al. (2018) reports a BH mass of $\sim 1.6\times10^9$ M_☉ for the BGE that hosts the giant radio source 'Double Irony' (~ 1100 kpc), which is equivalent to the BH mass that we find for NGC 315. We observe that 5/8 systems with BH mass $\geq 10^9$ M_☉ are large–scale jet systems. For the systems that have no radio emission detected at any frequency (7/53) we find that their BH mass ranges $\sim 0.4\times10^8 - 9\times10^8$ M_☉ with an average of $2.8\pm2.5\times10^8$ M_☉ which is on order of magnitude lower than the host BGEs of large–scale / jet systems.

On the other hand, the majority of small–scale jet systems (3/4) presents BH masses ${\sim}2{-}5{\times}10^8~M_{\odot}$ with only one system, NGC 1060, which is known as a dynamically interacting group (O'Sullivan et al. 2017) having a BH mass of ${\sim}3.1{\times}10^9~M_{\odot}$. We also note that only 10 BGEs present BH masses ${>}10^9~M_{\odot}$ with all of them being known to be embedded in an extended X-ray environment (Paper II and Sun et al. (2003) for NGC 1550).

From figure 8 we observe a mild trend that suggests higher BH masses are linked with highest radio powers, as found by pre-



Figure 8. Black hole mass (M_{BH}) of different radio morphologies, in relation to radio power at 235 MHz (P_{235MHz}). The radio morphology of each group can be seen in different symbols, with the host BGEs other than jet systems being marked as 'rest of BGEs' in triangles.

vious studies (e.g., Franceschini et al. 1998; Liu et al. 2006). However, we do not see a strong correlation between radio power at 235 MHz and the BH mass of the host BGEs in jet systems. There is a broad dispersion in radio power for a given BH mass for our CLoGS BGEs suggesting that a certain threshold value of BH mass may be necessary (but not on its own sufficient) to produce a powerful radio source, as suggested by a study of nearby field and group elliptical galaxies by Vaddi et al. (2016). This may indicate that jet activity depends more closely on environmental processes and the fuel that is available, rather than the mass of the host BGE's BH. In the diverse environment of galaxy groups we find that BH mass is not the main drive for the occurrence of jet emission, in agreement with the recent study of Sabater et al. (2019) in a sample of nearby AGN in the LOFAR Two-meter Sky Survey (LoTSS). Environmental mechanisms such as interactions, radiative cooling of the IGM, and precipitation of cold molecular gas are most likely prominent drivers for jet emission in galaxy groups. In addition to the effects of the large scale environment, the efficiency of gas accretion and the BH spin likely also play a role in the large radio power dispersion seen for our group BGES at a given BH mass (e.g., Baum et al. 1995).

7 CONCLUSIONS

In this paper we have presented the GMRT radio observations of the brightest group early-type galaxies from the low–richness CLoGS sample at 235 and 610 MHz. Combining with the study of the high–richness sample from Paper II we present results and statistics from the complete CLoGS sample. The CLoGS low–richness sample is found to have a radio detection rate of 82% (22 of 27 BGEs) at 235, 610 or 1400 MHz, with three radio sources being characterized as radio loud ($P_{1.4 GHz} > 10^{23}$ W Hz⁻¹ - NGC 315, NGC 5127 and NGC 5490). The overall detection rate of the CLoGS sample as a whole is 87% (46 of 53 BGEs), a result which confirms the high radio detection rate of dominant galaxies in groups or clusters in the nearby Universe. We find that CLoGS group-dominant radio galaxies exhibit radio powers in the range of $10^{20} - 10^{25}$ W Hz⁻¹.

Examining the radio morphology, we find that the majority of the groups in the CLoGS sample, ~53% (28/53) have a point-like radio morphology while ~13% (7/53) of groups are not detected in radio at any frequency. ~25% (13/53) of the groups are found to be hosting currently or recently active small/large scale jets and the remaining ~9% (5/53) host diffuse radio sources. The size of the extended CLoGS radio sources spans over 3 orders of magnitude from ~ 10 - 1200 kpc.

For the very low-powered unresolved point-like radio sources in the CLoGS low-richness sample using the expected radio emission from SFR_{*FUV*} and the expected radio emission from the conversion of molecular gas to star formation, we find that 3/14 are AGN dominated (NGC 1106, NGC 2768 and NGC 2911), and 5/14 sources (NGC 128, NGC 1395, NGC 3613, NGC 3923 and NGC 5061) probably have a significant contribution (20-50%) to their radio emission from star formation. In 6/14 point-like radio sources (NGC 252, NGC 524, NGC 1779, NGC 4697, NGC 7252 and NGC7377) we find that the stellar population likely dominates the radio emission.

In the CLoGS low–richness sample only 15/27 galaxies are detected at both 610 and 235 MHz with 2/15 sources exhibiting ultra-steep radio spectra with $\alpha_{235}^{610} > 1.3$ (NGC 315 and NGC 1407). The mean spectral index value for 13/15 sources is $\alpha_{235}^{610} = 0.60 \pm 0.16$ with the equivalent mean value in the 235–1400 MHz frequency range measured to be $\alpha_{235}^{1400} = 0.57 \pm 0.16$. In the full CLoGS sample we are able to measure a spectral index for 34/53 BGEs, and the mean spectral index is $\alpha_{235}^{610} = 0.59 \pm 0.26$. Jet systems are found to exhibit steeper mean $\alpha_{235}^{610} = 0.59 \pm 0.26$. Jet systems are found to exhibit steeper mean α_{235}^{610} than point-like systems. Roughly 4% of our sources have ultra-steep spectral indices ($\alpha_{235}^{610} > 1.3$), in good agreement with the earlier studies of Intema et al. (2011) and Sirothia et al. (2009).

Considering the group environments in which the BGEs are located, we find that 57% (30/53) of the CLoGS groups are spiralrich with 43% (23/53) being spiral-poor. Comparing the mean spiral fraction (F_{sp}) to BGE radio morphology, we find that the majority (75%) of radio point-like systems are found in dynamically young group systems (mean spiral fraction of 0.78 ± 0.10), with jet systems appearing equally in both dynamically young and old group systems, meaning that groups may form/host jets at any stage of their evolution. Our three remnant jet systems are also split between spiral-rich and spiral-poor groups, and BGEs which host diffuse, non-jet radio sources appear in groups with a wide range of spiral fractions.

We estimate the mechanical power output of the jet sources in the CLoGS low–richness sample from their 235 MHz luminosity, and find it to be in the range $\sim 10^{42} - 10^{44}$ erg s⁻¹. This is typical for galaxy groups, but about an order of magnitude greater than the range of powers measured in our high-richness sub-sample. A combined X-ray and radio analysis will be needed to explore the reliability and origin of this difference.

Lastly, for all radio detected CLoGS groups, we find only a mild positive trend between the BH mass of the BGE and its radio power at 235 MHz. We do not find a correlation between the central BH mass and jet activity. Although a certain minimum black hole mass is necessary to form jets, we also suggest that it is not the most essential prerequisite for the formation of a jet, in agreement with other recent studies (e.g., Vaddi et al. 2016; Sabater et al. 2019).

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APPENDIX A: INFORMATION ON INDIVIDUAL GROUPS

In this appendix we provide notes on the 25/27 individual groups of the low–richness sample analyzed in this study, the BGEs and their radio sources, where relevant. The GMRT data for NGC 315 and NGC 1407 at both 235 and 610 MHz are drawn from the earlier studies of Giacintucci et al. (2011) and Giacintucci et al. (2012).

A1 LGG 6 / NGC 128

NGC 128 has a peculiar ('peanut-shape') bulge morphology that kinematic studies attribute to the current strong dynamical interactions with its close neighbours NGC 127 and NGC 130 (Jarvis B. 1990). A tidal bridge of cold gas links NGC 128 and NGC 127, and the three galaxies have narrow spread in systemic velocity ($\leq 200 \text{ km s}^{-1}$; D'Onofrio et al. 1999).

A faint central point source was detected at both 1.4 GHz (Brown et al. 2011) and single dish 2.38 GHz (Dressler & Condon 1978), with flux densities of 1.5 mJy and 5 ± 3 mJy respectively. However, we note that the Arecibo 2.38 GHz value is not trustworthy as its poor spatial resolution (>5 arcmin) means that it cannot separate NGC 128 from its neighbours. Despite the previous low power detection at 1.4 GHz, NGC 128 is undetected in our GMRT observations, but NGC 127 hosts a radio point source at both 610 and 235 MHz. Figure B1 shows the radio contours that reveal the detection associated with NGC 127.

A2 LGG 12 / NGC 252

Previous 1.4 GHz (Condon et al. 2002) and single dish 2.38 GHz observations of NGC 252 found only a tentative detection with flux density of 2.5 mJy and 8 ± 4 mJy (Dressler & Condon 1978) respectively. We detect a weak radio point source associated with NGC 252 at 610 MHz (one of the weakest in the sample; see fig B2), and place an upper limit of <3.0 mJy at 5σ level of significance for the 235 MHz flux (Table 2).

A3 LGG 23 / NGC 524

X-ray observations by the *ROSAT* PSPC (Mulchaey et al. 2003) and more recently with *Chandra*-ACIS (Kharb et al. 2012) show that NGC 524 group contains an extended X-ray hot gas component associated with NGC 524. Previous 1.4 GHz (NVSS Condon et al. 1998) and single dish 2.38 GHz observations of NGC 524 detected a radio source of flux density 3.1 mJy and 10 ± 5 mJy (Dressler & Condon 1978) respectively. We detect a radio point source associated with NGC 524 only at 610 MHz (see fig B3), with an upper limit of <10.5 mJy at 5 σ significance for 235 MHz (Table 2).

A4 LGG 78 / NGC 1106

NGC 1106 is the BGE of the LGG 78 group and is identified as a Compton-thick dust obscured AGN Tanimoto et al. (2016). The NVSS (Condon et al. 2002) detected a strong point-like 1.4 GHz radio source with a flux density of 132 mJy. Our images at 610 MHz and 235 MHz (see figure B4) show a compact radio point-like source at both GMRT frequencies that coincides with the central region of the optical galaxy. The spectral index of the associated AGN radio emission is $\alpha_{235}^{610} \sim 0.7$.

A5 LGG 97 / NGC 1395

LGG 97 is part of the Eridanus supergroup and its BGE, NGC 1395, is X-ray luminous (O'Sullivan et al. 2003), with the extended diffuse X-ray emission being detected in both *ROSAT* PSPC (Omar & Dwarakanath 2005, out to \sim 135 kpc) and *XMM-Newton* (Escudero et al. 2018, out to \sim 66 kpc) observations.

Previous 1.4 GHz observations detected a weak point radio source with a flux density of 1.1 mJy (Brown et al. 2011). We detect a radio point-like source associated with NGC 1395 only at 610 MHz (see fig B5), with an upper limit of <3.5 mJy (5σ significance) at 235 MHz (Table 2).

A6 LGG 113 / NGC 1550

NGC 1550 is the BGE of the X-ray bright LGG 113 group, is classified as an optically overluminous early-type galaxy (OLEG; Vikhlinin et al. 1999) or potential fossil group, and is considered to be a product of multiple earlier mergers (Kawaharada et al. 2009). Previous X-ray observations by *ROSAT All-Sky Survey* (RASS; Beuing et al. 1999) detected a very bright X-ray source (RX J0419+0225) associated with the system, with more recent X-ray observations by *Chandra* (Sun et al. 2003) and *XMM-Newton* (Kawaharada et al. 2009) showing that the hot IGM extends out to at least ~200 kpc from the centre of the galaxy.

Previous 1.4 GHz observations from both NVSS (Condon et al. 1998) and Brown et al. (2011) detected a relatively weak radio source with a flux density of 16.6 mJy and 17 mJy respectively. The source presents an asymmetric double peaked radio morphology with the first peak coincident with the optical centre of NGC 1550 and the second one being observed as a lobe-like extension to the west (Dunn et al. 2010).

Our GMRT radio images of the source are shown in figure B6. We detect NGC 1550 at both frequencies with much higher angular resolution and sensitivity than prior studies. We find an asymmetric jet-lobe structure along an east-west axis, with the centre of the eastern lobe close to the stellar body of the galaxy and the second lobe expanding ~ 20 kpc to the west. At 610 MHz the western jet has a bend or kink to the south before continuing relatively straight to the lobe. This system will be examined in more detail in a future paper (Kolokythas et al., in prep.).

A7 LGG 126 / NGC 1779

The NVSS shows a weak 1.4 GHz point source in NGC 1779 with a flux density of 5.4 ± 0.6 mJy (Condon et al. 1998). We also detect a weak radio point source associated with NGC 1779 at both GMRT frequencies (see fig B7), with the spectral index in the 235–610 MHz range being ~0.60 and in the 235–1400 MHz range appearing to be inverted (α_{235}^{1400} ~-0.09; see Table 2).

A8 LGG 138 / NGC 2292

The S0 NGC 2292 is part of the interacting galaxy pair NGC 2292/2293. Earlier studies of this system (e.g., Barnes et al. 1999; Stickel et al. 2005) have found an extended, dusty tidal HI ring surrounding the pair. The galaxy was not previously detected in radio continuum emission and our GMRT observations find no radio detection at either frequency (fig. B8). The flux density upper limit set for this galaxy at 5σ level of significance is 0.3 mJy at 610 MHz and 1.9 mJy at 235 MHz. However, we do detect a point

source in the companion NGC 2293 at both 235 MHz (\sim 2.3 mJy) and 610 MHz (\sim 1.9 mJy).

A9 LGG 167 / NGC 2768

NGC 2768 hosts a polar ionized O[II] gas disk, rotating perpendicular to the stellar component and thus likely the product of a past merger (Fried & Illingworth 1994). Previous 1.4 GHz observations from the NVSS (Condon et al. 1998) and Brown et al. (2011) detected a radio source with a flux density of 15.9 ± 1 mJy and 14 ± 1 mJy respectively. We detect a compact radio point source at both GMRT frequencies (see fig B9). The spectral index of the associated compact radio emission appears to be flat in the 235–610 MHz regime, $\alpha_{235}^{610} \sim -0.04$ (Table 2).

A10 LGG 177 / NGC 2911

NGC 2911 is a low-ionization nuclear emission region (LINER) galaxy (Prieto et al. 2010). VLA observations at 4.9 and 14.9 GHz with sub-arcsecond resolution revealed a ~500 pc jet with a steep spectral index ($\alpha_{4.9}^{14.9}$ ~1.1 Mezcua et al. 2014). Lower resolution 1.4 GHz observations (Brown et al. 2011) reveal a radio source with a flux density of 56±2 mJy. We detect a radio point-like source associated with NGC 2911 at both GMRT frequencies (see fig B10) presenting hints of extension. An elongated asymmetric structure of ~7 kpc is observed only at 235 MHz, north-west from the centre whereas at 610 MHz the source appears to be slightly elongated but unresolved. We note here that the repetitive patterns observed around the source at 610 MHz, are artifacts owing to residual calibration errors. The spectral index of the associated compact radio emission appears to be flat in both 235–610 MHz ($\alpha_{235}^{610} \sim 0.18$), and 235–1400 MHz regimes ($\alpha_{235}^{1400} \sim 0.12$; Table 2).

A11 LGG 205 / NGC 3325

Earlier studies of NGC 3325 found no detection in the radio band, with only a tentative single dish 2.38 GHz detection of 1 ± 3 mJy (Dressler & Condon 1978) being reported. The galaxy is undetected in our GMRT observations at both frequencies (fig. B11). The flux density upper limit set for this galaxy at 5σ significance is <1.55 mJy at 610 MHz and <191 mJy at 235 MHz. The extremely high upper limit set at 235 MHz is a result of the short observing time (~60 mins) of the available archival data.

A12 LGG 232 / NGC 3613

A very faint central radio source was detected at 1.4 GHz in NGC 3613 using the VLA (Brown et al. 2011), with a flux density of 0.3 ± 0.3 mJy, the faintest in our sample. Our GMRT observations find no radio detection at either frequency (fig. B12). The flux density upper limit set for this galaxy at 5σ significance is <0.2 mJy at 610 MHz and <1.1 mJy at 235 MHz.

A13 LGG 236 / NGC 3665

NGC 3665 is the BGE of the X-ray luminous LGG 236 loose group of galaxies, which exhibits an unusually low velocity dispersion (\sim 217 km s⁻¹ Heldson et al. 2005). The galaxy contains a significant cold gas reservoir (Young et al. 2011; Alatalo et al. 2013) but has only low-level star formation (Xiao et al. 2018). NGC 3665 hostss a relatively weak double-jet radio source (Parma et al. 1986; Nyland et al. 2016) with its jet axis nearly perpendicular to the major axis of the central molecular gas disk (Onishi et al. 2017).

Previous 1.4 GHz observations from NVSS (Condon et al. 1998) detected a radio jet source with a flux density of 113.2 \pm 3.8 mJy. Our GMRT observations confirm the symmetric small-scale jet radio morphology at both frequencies, with the largest linear size of the source being ~16 kpc at 235 MHz (see fig. B13). At 235 MHz the source appears to be wider and more extended, whereas at 610 MHz it presents a symmetric thin-jet structure.

A14 LGG 255 / NGC 3923

NGC 3923 is a shell galaxy, with multiple shells of stars in its halo and is considered to be a product of a minor merger, (e.g., Dupraz & Combes 2010; Hernquist & Quinn 1988). Previous 1.4 GHz VLA (Brown et al. 2011) and single dish 5 GHz observations from Parkes radio telescope found only a tentative detection with a flux density of 1±0.5 mJy and 1.5±1.5 mJy (Disney & Wall 1977) respectively. We detect a weak radio point source at both GMRT frequencies (fig B14), with a spectral index of $\alpha_{235}^{610} \sim 0.91$.

A15 LGG 314 / NGC 4697

NGC 4697 is the BGE of the X-ray faint LGG 314 group of galaxies (ROSAT; Sansom et al. 2000). Previous 1.4 GHz VLA (Brown et al. 2011) observations detect only a very weak radio source with a flux density of 0.6 ± 0.5 mJy. Our GMRT observations find no radio detection at either frequency (fig. B15) for this galaxy, with 5σ flux density upper limits of <0.35 mJy at 610 MHz and <5.7 mJy at 235 MHz.

A16 LGG 329 / NGC 4956

The S0 NGC 4956 was not previously detected in the radio band and our GMRT observations find no radio detection at either frequency (fig. B16). The 5σ flux density upper limits are <0.3 mJy at 610 MHz and <2.4 mJy at 235 MHz.

A17 LGG 341 / NGC 5061

NGC 5061 is the dominant elliptical galaxy of the X-ray faint LGG 341 group of galaxies (Beuing et al. 1999). It is thought to have formed through several dry minor and major merging events (Dullo & Graham 2014). The galaxy was not previously detected at 1.4 GHz, and only a single dish tentative detection from Parkes telescope at 5 GHz is reported, with a flux density of 1 ± 2.8 mJy (Disney & Wall 1977). We detect a weak radio point source associated with NGC 5061 only at 610 MHz (fig B17) with a 5σ flux density upper limit of <2.5 mJy at 235 MHz.

A18 LGG 350 / NGC 5127

NGC 5127 hosts the FR I radio source B2 1321+31. Westerbork Synthesis Radio Telescope (WSRT) observations Klein et al. (1995) revealed that the there is little variation in the 610–10700 MHz spectral along the jet major axis but the spectral index steepens from the jet centre to its edge. The mean spectral index of the jet is $\alpha_{610}^{10700} \sim 0.56$ with the steepest indices found in the west lobe ($\alpha_{610}^{10700} \sim 1.00$). We detect the B2 1321+31 at both GMRT frequencies (fig B18). It has a symmetric double jet radio morphology with the total size of the source being \sim 244 kpc. At 235 MHz we see extensive lobes and diffuse radio emission that surrounds the main jet, whereas at 610 MHz the source appears to be thin, with only the linear jet emission being detected along with some clumpy emission in the west lobe. The eastern jet appears to bend to the north as it expands into the lobe.

A19 LGG 350 / NGC 5322

NGC 5322 is a LINER (Baldi & Capetti 2009) with a central stellar mass deficit (Carollo et al. 1997), a counter-rotating nuclear stellar and gaseous disk (e.g., Bender 1988; Bender & Surma 1992; Dullo et al. 2018), and hosts a central radio jet source (Feretti et al. 1984; Hummel, van der Hulst & Dickey 1984). In a recent study, Dullo et al. (2018) suggested that the most likely formation scenario for NGC 5322 is through a sequence of several initial 'dry' major mergers followed up by the subsequent accretion of a metal-rich gaseous satellite, with the latter event considered to have created the counter-rotating cold gas reservoir in the centre that fueled the AGN with gas and powered the radio source.

Previous 1.4 GHz VLA observations show an FR I radio morphology extending over ~20" (Hummel, van der Hulst & Dickey 1984). Dullo et al. (2018) using sub-arcsecond resolution 1.5 GHz e-MERLIN observations report that the inner jet structure has the same position angle as the extended one observed by the VLA. In our GMRT observations the source has a symmetric small-scale jet morphology at both frequencies, with the largest linear size of the source being ~15 kpc at 610 MHz (see fig. B19). At 235 MHz the source is poorly unresolved and only slightly elongated, whereas at 610 MHz the thin-jet structure is much clearer, aligned roughly north-south. The small-scale jet presents a very flat spectral index of $\alpha_{235}^{610} \sim 0.25$ indicating a young radio AGN.

A20 LGG 370 / NGC 5444

Previous 1.4 GHz VLA observations of NGC 5444 (Brown et al. 2011) report a detection with a flux density of 660 ± 20 mJy. However, this flux actually arises from the nearby radio source 4C +35.32, and not from NGC 5444 (see Brandie et al. 1976). Figure B20 shows that our GMRT observations find no radio detection at either frequency associated with NGC 5444. 4C +35.32 is detected as a strong source with a projected separation of only \sim 14 kpc (\sim 1') from NGC 5444. Owing to the presence of this source, the detection of the BGE here is dynamical range limited. Therefore, the upper limit for this galaxy is set from the maximum peak flux density at the position of NGC 5444, and is <1 mJy at 610 MHz and <12 mJy at 235 MHz.

A21 LGG 376 / NGC 5490

NGC 5490 hosts an FR I radio source (Condon & Broderick 1988) in which high resolution (milliarcsecond) VLBA observations at 5 GHz detect a polarized core-jet structure towards the stronger eastern large-scale jet (Xu et al. 2000; Kharb et al. 2012). VLA 1.4 GHz observations of the source (Brown et al. 2011) show a total flux density of 1300 ± 100 mJy. Our GMRT observations detect the source at both frequencies but with different morphology (fig B21). At 235 MHz we see an asymmetric radio core-jet morphology with only the eastern jet being visible as a continuous line of emission. The western jet is much shorter, but detached clumps of radio emission at larger radii presumably mark its path. At this frequency, the total size of the source is ~120 kpc. At 610 MHz, we observe only the central ~30 kpc jet which is linear and thin. The radio source presents an inverted spectral index of $\alpha_{235}^{1400} \sim -0.07$ in the 235–1400 MHz range and a relatively flat spectral index of $\alpha_{235}^{610} \sim 0.36$ in the 235–610 MHz one. The inverted spectral index can be attributed to the fact that we do not detect the same morphology at different frequencies.

A22 LGG 383 / NGC 5629

The S0 NGC 5629 was previously undetected in the radio band and our GMRT observations also find no radio detection at either frequency (fig. B22). The 5σ flux density upper limits for this galaxy are <0.3 mJy at 610 MHz and <2.0 mJy at 235 MHz.

A23 LGG 398 / NGC 5903

The radio continuum, HI and X-ray properties of the NGC 5903 group are described in detail in O'Sullivan et al. (2018). The group has a ~0.65 keV IGM extending out to at least 145 kpc (O'Sullivan et al. 2018), and hosts an extended (~40 by 75 kpc) steep-spectrum diffuse radio source associated with NGC 5903 (Gopal-Krishna 1978; Gopal-Krishna et al. 2012). The spectral index of the radio emission is $\alpha_{150}^{612} \sim 1.03 \pm 0.08$, suggesting a radiative age greater than 360 Myr. The most likely origin scenario for the radio structure is a combination of an interaction-triggered AGN outburst and high-velocity collisions between galaxies. Our images at 610 MHz and 235 MHz (see figure B23) show the diffuse radio structure, with its brightest part coinciding with the central region of the galaxy.

A24 LGG 457 / NGC 7252

NGC 7252 (also known as the 'Atoms-for-Peace' galaxy) is the BGE in the X-ray faint LGG 457 group of galaxies (Nolan et al. 2004) and is identified as a Compton-thick dust obscured AGN (Tanimoto et al. 2016). It is considered to be the remnant of an advanced merger between two gas-rich spiral galaxies and has two clearly visible and HI-rich tidal tails (e.g., Schweizer 1982; Dupraz et al. 1990; Wang et al. 1992; Hibbard et al. 1994). The merger is estimated to have started ~600–700 Myr ago, triggering a starburst (Hibbard & Mihos 1995; Chien & Barnes 2010), and NGC 7252 is expected to evolve into an elliptical galaxy after exhausting its cold gas reservoir (Toomre 1977; Schweizer 1982).

NVSS detected a point-like radio source with a flux density of 25.3 ± 1.2 mJy (Condon et al. 1998). Our images at 610 MHz and 235 MHz (see figure B24) show a compact radio point-like source at both GMRT frequencies that coincides with the optical nucleus.

A25 LGG 463 / NGC 7377

The NVSS shows NGC 7377 to host a weak radio point source with a flux density of 3.4 ± 0.5 mJy (Condon et al. 1998). In our GMRT observations we also detect a weak radio point source at both frequencies (see fig B25). The spectral index in the 235–610 MHz range is ~0.27 and in the 235–1400 MHz range appears to be inverted (α_{235}^{1400} ~-0.13; see Table 2).

APPENDIX B: GMRT RADIO CONTINUUM IMAGES

In this appendix we provide GMRT radio and Digitized Sky Survey optical images for 25/27 of our BGEs analysed in this study. In all figures, the radio images are shown in the left-hand panels, and the optical images on the right. 235 MHz contours (green) are overlaid on the upper-right panels, 610 MHz (cyan) on the bottom-right. The white circle, when shown, indicates the optical position of the BGE.



Figure B1. LGG 6 / NGC 128. *Top left:* GMRT 235 MHz image. *Top right:* GMRT 235 MHz contours in green ($1\sigma = 1.90$ mJy beam⁻¹), overlaid on the *Digitized Sky Survey (DSS)* optical image. *Bottom left:* GMRT 610 MHz image. *Bottom right:* GMRT 610 MHz contours in cyan ($1\sigma = 160 \mu$ Jy beam⁻¹), overlaid on the *Digitized Sky Survey (DSS)* optical image. In both panels the radio contours are spaced by a factor of two, starting from 3σ level of significance. For this source the scale is 0.291 kpc arcsec⁻¹.



Figure B2. LGG 12 / NGC 252. *Top left:* GMRT 235 MHz image. *Top right:* GMRT 235 MHz contours in green ($1\sigma = 0.60$ mJy beam⁻¹), overlaid on the *Digitized Sky Survey (DSS)* optical image. *Bottom left:* GMRT 610 MHz image. *Bottom right:* GMRT 610 MHz contours in cyan ($1\sigma = 60 \mu$ Jy beam⁻¹), overlaid on the *Digitized Sky Survey (DSS)* optical image. In both panels the radio contours are spaced by a factor of two, starting from 3σ level of significance. For this source the scale is 0.349 kpc arcsec⁻¹.



Figure B3. LGG 14 / NGC 524. *Top left:* GMRT 235 MHz image. *Top right:* GMRT 235 MHz contours in green ($1\sigma = 2.10$ mJy beam⁻¹), overlaid on the *Digitized Sky Survey (DSS)* optical image. *Bottom left:* GMRT 610 MHz image. *Bottom right:* GMRT 610 MHz contours in cyan ($1\sigma = 180 \ \mu$ Jy beam⁻¹), overlaid on the *Digitized Sky Survey (DSS)* optical image. In both panels the radio contours are spaced by a factor of two, starting from 3σ level of significance. For this source the scale is 0.165 kpc arcsec⁻¹.



Figure B4. LGG 78 / NGC 1106. *Top left:* GMRT 235 MHz image. *Top right:* GMRT 235 MHz contours in green ($1\sigma = 0.60$ mJy beam⁻¹), overlaid on the *Digitized Sky Survey (DSS)* optical image. *Bottom left:* GMRT 610 MHz image. *Bottom right:* GMRT 610 MHz contours in cyan ($1\sigma = 100 \mu$ Jy beam⁻¹), overlaid on the *Digitized Sky Survey (DSS)* optical image. In both panels the radio contours are spaced by a factor of two, starting from 3σ level of significance. For this source the scale is 0.310 kpc arcsec⁻¹.



Figure B5. LGG 97 / NGC 1395. *Top left:* GMRT 235 MHz image. *Top right:* GMRT 235 MHz contours in green ($1\sigma = 0.70$ mJy beam⁻¹), overlaid on the *Digitized Sky Survey (DSS)* optical image. *Bottom left:* GMRT 610 MHz image. *Bottom right:* GMRT 610 MHz contours in cyan ($1\sigma = 100 \ \mu$ Jy beam⁻¹), overlaid on the *Digitized Sky Survey (DSS)* optical image. In both panels the radio contours are spaced by a factor of two, starting from 3σ level of significance. For this source the scale is 0.102 kpc arcsec⁻¹.



Figure B6. LGG 113 / NGC 1550. *Top left:* GMRT 235 MHz image. *Top right:* GMRT 235 MHz contours in green ($1\sigma = 0.45$ mJy beam⁻¹), overlaid on the *Digitized Sky Survey (DSS)* optical image. *Bottom left:* GMRT 610 MHz image. *Bottom right:* GMRT 610 MHz contours in cyan ($1\sigma = 40 \mu$ Jy beam⁻¹), overlaid on the *Digitized Sky Survey (DSS)* optical image. In both panels the radio contours are spaced by a factor of two, starting from 3σ level of significance. For this source the scale is 0.257 kpc arcsec⁻¹.



Figure B7. LGG 126 / NGC 1779. *Top left:* GMRT 235 MHz image. *Top right:* GMRT 235 MHz contours in green ($1\sigma = 0.45$ mJy beam⁻¹), overlaid on the *Digitized Sky Survey (DSS)* optical image. *Bottom left:* GMRT 610 MHz image. *Bottom right:* GMRT 610 MHz contours in cyan ($1\sigma = 50 \mu$ Jy beam⁻¹), overlaid on the *Digitized Sky Survey (DSS)* optical image. In both panels the radio contours are spaced by a factor of two, starting from 3σ level of significance. For this source the scale is 0.218 kpc \arccos^{-1} .



Figure B8. LGG 138 / NGC 2292. *Top left:* GMRT 235 MHz image. *Top right:* GMRT 235 MHz contours in green ($1\sigma = 0.42$ mJy beam⁻¹), overlaid on the *Digitized Sky Survey (DSS)* optical image. *Bottom left:* GMRT 610 MHz image. *Bottom right:* GMRT 610 MHz contours in cyan ($1\sigma = 50 \mu$ Jy beam⁻¹), overlaid on the *Digitized Sky Survey (DSS)* optical image. In both panels the radio contours are spaced by a factor of two, starting from 3σ level of significance. For this source the scale is 0.145 kpc \arccos^{-1} .



Figure B9. LGG 167 / NGC 2768. *Top left:* GMRT 235 MHz image. *Top right:* GMRT 235 MHz contours in green ($1\sigma = 0.30$ mJy beam⁻¹), overlaid on the *Digitized Sky Survey (DSS)* optical image. *Bottom left:* GMRT 610 MHz image. *Bottom right:* GMRT 610 MHz contours in cyan ($1\sigma = 30 \mu$ Jy beam⁻¹), overlaid on the *Digitized Sky Survey (DSS)* optical image. In both panels the radio contours are spaced by a factor of two, starting from 3σ level of significance. For this source the scale is 0.112 kpc arcsec⁻¹.



Figure B10. LGG 177 / NGC 2911. *Top left:* GMRT 235 MHz image. *Top right:* GMRT 235 MHz contours in green ($1\sigma = 0.30$ mJy beam⁻¹), overlaid on the *Digitized Sky Survey (DSS)* optical image. *Bottom left:* GMRT 610 MHz image. *Bottom right:* GMRT 610 MHz contours in cyan ($1\sigma = 200 \mu$ Jy beam⁻¹), overlaid on the *Digitized Sky Survey (DSS)* optical image. In both panels the radio contours are spaced by a factor of two, starting from 3σ level of significance at 235 MHz and at 4σ level of significance at 610 MHz due to the higher rms noise around the source at this frequency. For this source the scale is 0.218 kpc arcsec⁻¹.



Figure B11. LGG 205 / NGC 3325. *Top left:* GMRT 235 MHz image. *Top right:* GMRT 235 MHz contours in green $(1\sigma = 38 \text{ mJy beam}^{-1})$, overlaid on the *Digitized Sky Survey (DSS)* optical image. *Bottom left:* GMRT 610 MHz image. *Bottom right:* GMRT 610 MHz contours in cyan $(1\sigma = 310 \ \mu\text{Jy beam}^{-1})$, overlaid on the *Digitized Sky Survey (DSS)* optical image. In both panels the radio contours are spaced by a factor of two, starting from 3σ level of significance. The higher rms noise for this field at both frequencies is attributed to the short duration of the archival observations at hand. For this source the scale is 0.388 kpc arcsec⁻¹.



Figure B12. LGG 232 / NGC 3613. *Top left:* GMRT 235 MHz image. *Top right:* GMRT 235 MHz contours in green ($1\sigma = 0.26$ mJy beam⁻¹), overlaid on the *Digitized Sky Survey (DSS)* optical image. *Bottom left:* GMRT 610 MHz image. *Bottom right:* GMRT 610 MHz contours in cyan ($1\sigma = 30 \mu$ Jy beam⁻¹), overlaid on the *Digitized Sky Survey (DSS)* optical image. In both panels the radio contours are spaced by a factor of two, starting from 3σ level of significance. For this source the scale is 0.156 kpc \arccos^{-1} .



Figure B13. LGG 236 / NGC 3665. *Top left:* GMRT 235 MHz image. *Top right:* GMRT 235 MHz contours in green $(1\sigma = 0.78 \text{ mJy beam}^{-1})$, overlaid on the *Digitized Sky Survey (DSS)* optical image. *Bottom left:* GMRT 610 MHz image. *Bottom right:* GMRT 610 MHz contours in cyan $(1\sigma = 170 \mu\text{Jy beam}^{-1})$, overlaid on the *Digitized Sky Survey (DSS)* optical image. In both panels the radio contours are spaced by a factor of two, starting from 3σ level of significance. For this source the scale is 0.156 kpc arcsec⁻¹.



Figure B14. LGG 255 / NGC 3923. *Top left:* GMRT 235 MHz image. *Top right:* GMRT 235 MHz contours in green ($1\sigma = 0.45$ mJy beam⁻¹), overlaid on the *Digitized Sky Survey (DSS)* optical image. *Bottom left:* GMRT 610 MHz image. *Bottom right:* GMRT 610 MHz contours in cyan ($1\sigma = 50 \mu$ Jy beam⁻¹), overlaid on the *Digitized Sky Survey (DSS)* optical image. In both panels the radio contours are spaced by a factor of two, starting from 3σ level of significance. For this source the scale is 0.097 kpc arcsec⁻¹.



Figure B15. LGG 314 / NGC 4697. *Top left:* GMRT 235 MHz image. *Top right:* GMRT 235 MHz contours in green ($1\sigma = 1.14 \text{ mJy beam}^{-1}$), overlaid on the *Digitized Sky Survey (DSS)* optical image. *Bottom left:* GMRT 610 MHz image. *Bottom right:* GMRT 610 MHz contours in cyan ($1\sigma = 70 \mu$ Jy beam⁻¹), overlaid on the *Digitized Sky Survey (DSS)* optical image. In both panels the radio contours are spaced by a factor of two, starting from 3σ level of significance. For this source the scale is 0.087 kpc arcsec⁻¹.



Figure B16. LGG 329 / NGC 4956. *Top left:* GMRT 235 MHz image. *Top right:* GMRT 235 MHz contours in green $(1\sigma = 0.48 \text{ mJy beam}^{-1})$, overlaid on the *Digitized Sky Survey (DSS)* optical image. *Bottom left:* GMRT 610 MHz image. *Bottom right:* GMRT 610 MHz contours in cyan $(1\sigma = 50 \ \mu\text{Jy beam}^{-1})$, overlaid on the *Digitized Sky Survey (DSS)* optical image. In both panels the radio contours are spaced by a factor of two, starting from 3σ level of significance. For this source the scale is $0.344 \text{ kpc } \arccos^{-1}$.



Figure B17. LGG 341 / NGC 5061. *Top left:* GMRT 235 MHz image. *Top right:* GMRT 235 MHz contours in green $(1\sigma = 0.50 \text{ mJy beam}^{-1})$, overlaid on the *Digitized Sky Survey (DSS)* optical image. *Bottom left:* GMRT 610 MHz image. *Bottom right:* GMRT 610 MHz contours in cyan $(1\sigma = 50 \mu\text{Jy beam}^{-1})$, overlaid on the *Digitized Sky Survey (DSS)* optical image. In both panels the radio contours are spaced by a factor of two, starting from 3σ level of significance. For this source the scale is 0.136 kpc arcsec⁻¹.



Figure B18. LGG 350 / NGC 5127. *Top left:* GMRT 235 MHz image. *Top right:* GMRT 235 MHz contours in green ($1\sigma = 0.65$ mJy beam⁻¹), overlaid on the *Digitized Sky Survey (DSS)* optical image. *Bottom left:* GMRT 610 MHz image. *Bottom right:* GMRT 610 MHz contours in cyan ($1\sigma = 250 \mu$ Jy beam⁻¹), overlaid on the *Digitized Sky Survey (DSS)* optical image. In both panels the radio contours are spaced by a factor of two, starting from 3σ level of significance. For this source the scale is 0.349 kpc arcsec⁻¹.



Figure B19. LGG 360 / NGC 5322. *Top left:* GMRT 235 MHz image. *Top right:* GMRT 235 MHz contours in green ($1\sigma = 0.70$ mJy beam⁻¹), overlaid on the *Digitized Sky Survey (DSS)* optical image. *Bottom left:* GMRT 610 MHz image. *Bottom right:* GMRT 610 MHz contours in cyan ($1\sigma = 50 \mu$ Jy beam⁻¹), overlaid on the *Digitized Sky Survey (DSS)* optical image. In both panels the radio contours are spaced by a factor of two, starting from 3σ level of significance. For this source the scale is 0.141 kpc arcsec⁻¹.



Figure B20. LGG 370 / NGC 5444. *Top left:* GMRT 235 MHz image. *Top right:* GMRT 235 MHz contours in green $(1\sigma = 0.80 \text{ mJy beam}^{-1})$, overlaid on the *Digitized Sky Survey (DSS)* optical image. *Bottom left:* GMRT 610 MHz image. *Bottom right:* GMRT 610 MHz contours in cyan $(1\sigma = 200 \mu\text{Jy beam}^{-1})$, overlaid on the *Digitized Sky Survey (DSS)* optical image. In both panels the radio contours are spaced by a factor of two, starting from 3σ level of significance. For this source the scale is 0.291 kpc arcsec⁻¹.



Figure B21. LGG 376 / NGC 5490. *Top left:* GMRT 235 MHz image. *Top right:* GMRT 235 MHz contours in green ($1\sigma = 1.85$ mJy beam⁻¹), overlaid on the *Digitized Sky Survey (DSS)* optical image. *Bottom left:* GMRT 610 MHz image. *Bottom right:* GMRT 610 MHz contours in cyan ($1\sigma = 250 \mu$ Jy beam⁻¹), overlaid on the *Digitized Sky Survey (DSS)* optical image. In both panels the radio contours are spaced by a factor of two, starting from 3σ level of significance at 235 MHz and at 4σ level of significance at 610 MHz due to the higher rms noise around the source at this frequency. For this source the scale is 0.344 kpc arcsec⁻¹.



Figure B22. LGG 383 / NGC 5629. *Top left:* GMRT 235 MHz image. *Top right:* GMRT 235 MHz contours in green $(1\sigma = 0.40 \text{ mJy beam}^{-1})$, overlaid on the *Digitized Sky Survey (DSS)* optical image. *Bottom left:* GMRT 610 MHz image. *Bottom right:* GMRT 610 MHz contours in cyan $(1\sigma = 50 \ \mu\text{Jy beam}^{-1})$, overlaid on the *Digitized Sky Survey (DSS)* optical image. In both panels the radio contours are spaced by a factor of two, starting from 3σ level of significance. For this source the scale is 0.325 kpc arcsec⁻¹.



Figure B23. LGG 398 / NGC 5903. *Top left:* GMRT 235 MHz image. *Top right:* GMRT 235 MHz contours in green ($1\sigma = 0.60$ mJy beam⁻¹), overlaid on the *Digitized Sky Survey (DSS)* optical image. *Bottom left:* GMRT 610 MHz image. *Bottom right:* GMRT 610 MHz contours in cyan ($1\sigma = 80 \mu$ Jy beam⁻¹), overlaid on the *Digitized Sky Survey (DSS)* optical image. In both panels the radio contours are spaced by a factor of two, starting from 3σ level of significance. For this source the scale is 0.175 kpc arcsec⁻¹.



Figure B24. LGG 457 / NGC 7252. *Top left:* GMRT 235 MHz image. *Top right:* GMRT 235 MHz contours in green $(1\sigma = 0.60 \text{ mJy beam}^{-1})$, overlaid on the *Digitized Sky Survey (DSS)* optical image. *Bottom left:* GMRT 610 MHz image. *Bottom right:* GMRT 610 MHz contours in cyan $(1\sigma = 100 \ \mu\text{Jy beam}^{-1})$, overlaid on the *Digitized Sky Survey (DSS)* optical image. In both panels the radio contours are spaced by a factor of two, starting from 3σ level of significance. For this source the scale is 0.320 kpc arcsec⁻¹.



Figure B25. LGG 463 / NGC 7377. *Top left:* GMRT 235 MHz image. *Top right:* GMRT 235 MHz contours in green ($1\sigma = 0.40$ mJy beam⁻¹), overlaid on the *Digitized Sky Survey (DSS)* optical image. *Bottom left:* GMRT 610 MHz image. *Bottom right:* GMRT 610 MHz contours in cyan ($1\sigma = 40 \mu$ Jy beam⁻¹), overlaid on the *Digitized Sky Survey (DSS)* optical image. In both panels the radio contours are spaced by a factor of two, starting from 3σ level of significance. For this source the scale is 0.223 kpc arcsec⁻¹.

APPENDIX C: BGE CALCULATED PROPERTIES FOR CLOGS SAMPLE

This paper has been typeset from a $T_{\ensuremath{E}} X/I \!\! \ensuremath{\Delta} T_{\ensuremath{E}} X$ file prepared by the author.

Table C1. Ce	ntral velocity	/ dispersions and	d estimated BH	I masses for	r CLoGS BG	Es. The c	olumns list t	the BGE name	, the central	velocity	dispersion a	σ_o , the
references fro	m which σ_o	was drawn, and	the BH mass d	lerived using	g the central v	elocity d	lispersion σ_o	- BH mass re	lation from	McConn	ell & Ma (2	2013).

	e	J 1	-
BGE	$\sigma_o ({\rm km} {\rm s}^{-1})$	References (for σ_o)	$M_{BH\sigma}~(10^9~M_{\odot})$
HIGH-RICH	NESS SUB-SAMPI	LE	
NGC 193	197.6±4.8	HyperLeda	$0.23^{+0.03}_{-0.03}$
NGC 410	$311.0 {\pm} 6.0$	Loubser et al. (2018)	$2.44^{+0.84}_{-0.63}$
NGC 584	220.0 ± 6.0	Loubser et al. (2018)	$0.40^{+0.08}_{-0.06}$
NGC 677	$228.4{\pm}20.2$	HyperLeda	$0.49^{+0.10}_{-0.08}$
NGC 777	321.0±9.0	Loubser et al. (2018)	$2.87^{+1.04}_{-0.76}$
NGC 940	180.5 ± 8.7	HyperLeda	$0.14_{-0.01}^{+0.02}$
NGC 924	201.0 ± 8.0	Loubser et al. (2018)	$0.25_{-0.03}^{+0.04}$
NGC 978	231.8 ± 5.3	HyperLeda	$0.53_{-0.09}^{+0.11}$
NGC 1060	326.0 ± 6.0	Loubser et al. (2018)	$3.11^{+1.15}_{-0.84}$
NGC 1167	202.9 ± 9.9	HyperLeda	$0.27_{-0.04}^{+0.04}$
NGC 1453	312.0±6.0	Loubser et al. (2018)	$2.48^{+0.86}_{-0.64}$
NGC 1587	239.0±6.0	Loubser et al. (2018)	$0.62^{+0.14}_{-0.11}$
NGC 2563	287.0 ± 6.0	Loubser et al. (2018)	$1.61^{+0.49}_{-0.28}$
NGC 3078	243.3 ± 6.4	HyperLeda	$0.68^{+0.16}_{-0.12}$
NGC 4008	215.8±14.3	HyperLeda	$0.36^{+0.07}_{-0.06}$
NGC 4169	183.5±4.7	HyperLeda	$0.16^{+0.02}_{-0.02}$
NGC 4261	314.4 ± 4.0	Loubser et al. (2018)	$2.56^{+0.02}$
ESO 507-25	243.6 ± 6.8	HyperLeda	$0.68^{+0.16}_{-0.10}$
NGC 5044	224.9+9.1	HyperLeda	$0.45^{+0.09}_{-0.13}$
NGC 5084	199.8 ± 5.8	HyperLeda	$0.24^{+0.04}$
NGC 5153	182.6 ± 5.1	HyperLeda	$0.15^{+0.02}$
NGC 5353	300.0+4.0	Loubser et al. (2018)	$2.02^{+0.66}$
NGC 5846	195.0 ± 4.0	Loubser et al. (2018)	$0.22^{+0.03}_{-0.03}$
NGC 5982	260.0+7.0	Loubser et al. (2018)	$0.96^{+0.25}$
NGC 6658	210.0+9.0	Loubser et al. (2018)	$0.32^{+0.05}$
NGC 7619	331.0±5.0	Loubser et al. (2018)	$3.37^{+1.27}_{-0.92}$
LOW-RICHN	ESS SUB-SAMPL	E	
NGC 128	216.4±14.0	HyperLeda	$0.37^{+0.07}_{-0.06}$
NGC 252	223.6±12.4	HyperLeda	$0.44^{+0.09}_{-0.07}$
NGC 315	339.0±8.0	Loubser et al. (2018)	$3.82^{+1.48}_{-1.07}$
NGC 524	$245.0 {\pm} 5.0$	Loubser et al. (2018)	$0.71_{-0.13}^{+0.17}$
NGC 1106	143.7±7.9	HyperLeda	$0.040^{+0.001}_{-0.001}$
NGC 1395	240.1±4.2	HyperLeda	$0.63^{+0.14}_{-0.12}$
NGC 1407	265.6 ± 5.1	HyperLeda	$1.07^{+0.29}_{-0.23}$
NGC 1550	300.4 ± 5.3	HyperLeda	$2.04^{+0.67}_{-0.50}$
NGC 1779	164.0 ± 10.0	Loubser et al. (2018)	$0.09^{+0.01}_{-0.01}$
NGC 2292	139.8±21.0	HyperLeda	0.0400 + 0.0004
NGC 2768	172.0 ± 3.0	Loubser et al. (2018)	$0.11^{+0.01}_{-0.01}$
NGC 2911	239.1±7.5	HyperLeda	$0.62^{+0.01}_{-0.11}$
NGC 3325	168.9 ± 12.1	HyperLeda	$0.10^{+0.01}_{-0.01}$
NGC 3613	214.0±4.0	Loubser et al. (2018)	$0.35^{+0.01}_{-0.05}$
NGC 3665	224.0±5.0	Loubser et al. (2018)	$0.44^{+0.09}_{-0.07}$
NGC 3923	245.6 ± 4.9	HyperLeda	$0.71^{+0.17}_{-0.14}$
NGC 4697	165.2±1.6	HyperLeda	$0.09^{+0.01}_{-0.01}$
NGC 4956	144.7 ± 24.7	HyperLeda	$0.050^{+0.001}_{-0.001}$
NGC 5061	188.0 ± 4.8	HyperLeda	$0.18^{+0.02}_{-0.001}$
NGC 5127	194.0 ± 5.0	Loubser et al. (2018)	$0.21^{+0.02}_{-0.02}$
NGC 5322	230.0±3.8	HyperLeda	$0.51^{+0.03}_{-0.00}$
NGC 5444	220.7±33.1	HyperLeda	$0.41^{+0.09}_{-0.07}$
NGC 5490	353.0±6.0	Loubser et al. (2018)	$4.71^{+1.93}$
NGC 5629	257.0 ± 8.0	Loubser et al. (2018)	$0.93^{+0.23}_{-0.18}$
NGC 5903	203.6 ± 2.8	HyperLeda	$0.27^{+0.04}_{-0.04}$
NGC 7252	156.8 ± 22.0	HyperLeda	$0.070^{+0.001}_{-0.001}$
NGC 7377	174.5 ± 12.1	HyperLeda	$0.12^{+0.01}_{-0.01}$
		-	0.01