

# EPIC 211964830: A transiting multi-planet system in the Praesepe open cluster

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## ABSTRACT

Planet host stars with well-constrained ages provide a rare window to the time domain of planet formation and evolution. The NASA *K2* mission has enabled the discovery of the vast majority of known planets transiting stars in clusters, providing a valuable sample of planets with known ages and radii. We present the discovery of two planets transiting EPIC 211964830, an M2 dwarf in the intermediate age (600–800 Myr) Praesepe open cluster (also known as the Beehive Cluster, M44, or NGC 2632), which was observed by *K2* during Campaign 16. The planets have orbital periods of 5.8 and 19.7 days, and radii of 2.2 and 2.7  $R_{\oplus}$ , respectively, and their equilibrium temperatures are 496 and 331 K, making this a system of two warm sub-Neptunes. When placed in the context of known planets orbiting field stars of similar mass to EPIC 211964830, these planets do not appear to have significantly inflated radii, as has previously been noted for some cluster planets. As the second known system of multiple planets transiting a star in a cluster, EPIC 211964830 should be valuable for testing theories of photoevaporation in systems of multiple planets. Follow-up observations with current near-infrared (NIR) spectrographs could yield planet mass measurements, which would provide information about the mean densities and compositions of small planets soon after photoevaporation is expected to have finished. Follow-up NIR transit observations using *Spitzer* or large ground-based telescopes could yield improved radius estimates, further enhancing the characterization of these interesting planets.

**Key words:** exoplanets – transits – observations – imaging

## 1 INTRODUCTION

The great wealth of data from large exoplanet surveys is a powerful tool for statistical studies of planet formation and evolution. For example, the large number of transiting planets, mostly discovered by the *Kepler* mission, has enabled the discovery of detailed structure in the observed planetary radius distribution (Fulton et al. 2017; Van Eylen et al. 2018; Fulton & Petigura 2018; Berger et al. 2018), which had been predicted by theories of planetary evolution via photoevaporation (e.g. Owen & Wu 2013; Lopez & Fortney 2014). The observed properties of planets are intrinsically dependent on the properties of their host stars; indeed, the phrase “know thy star, know thy planet” has become ubiquitous in the field of exoplanet science.

Besides the necessity of host star characterization for obtaining planet properties from indirect measurements, the comparison of planet properties with those of their hosts has long been a source of great interest (e.g. Fischer & Valenti 2005; Petigura et al. 2018), as the discovery of a causal relationship would provide a rare glimpse of the mechanisms underpinning planet formation and the processes sculpting them thereafter. However, the vast majority of known planet host stars are of uncertain age, so planet demographics and occurrence rates have been largely unexplored in the time domain. Planets orbiting stars in clusters thus present a rare opportunity for investigations of planet properties as a function of time.

Most of the first known planets orbiting cluster stars were discovered by the radial velocity (RV) method (e.g. Sato et al. 2007; Lovis & Mayor 2007; Quinn et al. 2012, 2014; Malavolta et al. 2016). However, an inherent limitation of the RV method is that most planets discovered in this way do not transit their host stars, so their radii are unknown and the measured masses are lower limits. By extending the *Kepler* mission to the ecliptic plane, the *K2* mission (Howell et al. 2014) has enabled the discovery of the vast majority of transiting planets in clusters (Obermeier et al. 2016; Pepper et al. 2017; David et al. 2016a; Mann et al. 2016b, 2017; Gaidos et al. 2017; Ciardi et al. 2018; Mann et al. 2018; Livingston et al. 2018a), including the youngest known transiting planet (David et al. 2016b; Mann et al. 2016a).

We present here the discovery of two planets transiting EPIC 211964830, a low mass star in the Praesepe open cluster. We identified two sets of transits in the *K2* photometric data collected during Campaign 16, then obtained high resolution adaptive optics (AO) imaging of the host star. Precise photometry and astrometry from the *Gaia* mission (Gaia Collaboration et al. 2016), along with archival data, enable the characterization of the host star and facilitate the interpretation of the transit signals. We combine the results of detailed light curve analyses and host star characterization to determine the planetary nature of the transit signals, as well as constrain fundamental properties of the two small planets. EPIC 211964830 is now the second known transiting multi-planet system in a cluster, offering a rare glimpse into the time domain of planet formation and evolution; its discovery thus significantly enhances a crucial avenue for testing theories of migration and photoevaporation. The transit detections and follow-up observations that enabled this discovery are the result of an international collaboration called

KESPRINT. While this manuscript was in preparation Rizzuto et al. (2018) announced an independent discovery of this system. Given the rarity of transiting multi-planet cluster systems, it is not surprising that multiple teams pursued follow-up observations of this valuable target.

This paper is organized as follows. In Section 2, we describe the *K2* photometry and high resolution imaging of the host star, as well as archival data used in our analysis. In Section 3 we describe our transit analyses, host star characterization, planet validation, and dynamical analyses of the system. Finally, we discuss the properties of the planet system and prospects for future studies in Section 4, and we conclude with a summary in Section 5.

## 2 OBSERVATIONS

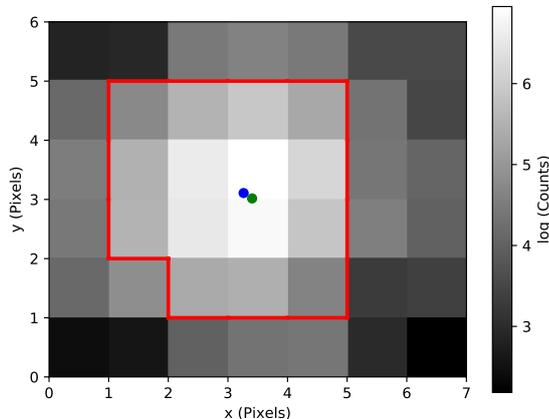
### 2.1 K2 photometry

EPIC 211964830 (also known as Cl\* NGC 2632 JS 597, 2MASS J08452605+1941544, and *Gaia* DR2 661167785238757376) was one of 35,643 long cadence (LC) targets observed during Campaign 16 of the *K2* mission, from 2017-12-07 23:01:18 to 2018-02-25 12:39:52 UT. EPIC 211964830 was proposed as a LC target by GO programs 16022 (PI Rebull), 16031 (PI Endl), 16052 (PI Stello), and 16060 (PI Agueros). The data were downlinked from the spacecraft and subsequently calibrated and made available on the Mikulski Archive for Space Telescopes<sup>1</sup> (MAST). We describe our light curve preparation and transit search procedures in detail in Livingston et al. (2018b). In brief, we extracted photometry from the *K2* pixel data with circular apertures and applied a correction for the systematic caused by the pointing drift of *K2*, similar to the approach described by Vanderburg & Johnson (2014). For EPIC 211964830, we selected an aperture with a 1.8 pixel radius (see Figure 1), as this resulted in the corrected light curve with the lowest levels of noise on six hour timescales. We then searched the light curve for transits using the Box-Least-Squares algorithm (BLS, Kovács et al. 2002), and identified two candidate planets with signal detection efficiency (Ofir 2014) values of 11.0 and 10.1. The light curve and phase-folded transits of EPIC 211964830 are shown in Figure 2. Subsequent modeling described in Section 3.1 yielded transit SNR (Livingston et al. 2018b) values of 14.0 and 15.9 for the inner and outer planet candidates, respectively. We identified an outlier most likely caused by residual systematics in the light curve and excluded it from our transit analysis (see gray data point in lower right panel of Figure 2).

### 2.2 Subaru/IRCS adaptive optics imaging

On UT 2018 June 14, we obtained high resolution adaptive optics (AO) imaging of EPIC 211964830 with the IRCS instrument mounted on the 8.2 meter Subaru telescope on Mauna Kea, HI, USA. The AO imaging utilized the target stars themselves as natural guide stars. We adopted the fine sampling mode (1 pix  $\approx$  20 mas) and five-point dithering, and a total exposure time of 300 seconds was spent

<sup>1</sup> <https://archive.stsci.edu/k2/>



**Figure 1.** *K2* “postage stamp” of EPIC 211964830 with a 1.8 pixel ( $\sim 7''$ ) photometric aperture overlaid in red. The green circle indicates the current position of the target in the EPIC, and the blue circle is the center of the flux distribution.

for EPIC 211964830. The full width at the half maximum (FWHM) of the target image was  $\sim 0''.22$  after the AO correction. Following Hirano et al. (2016), we performed dark current subtraction, flat fielding, and distortion correction before finally aligning and median combining the individual frames. In this manner we produced and visually inspected  $16'' \times 16''$  combined image, which we then used to compute a  $5\text{-}\sigma$  contrast curve following the procedure described in Hirano et al. (2018). We show the resulting contrast curve with a  $4'' \times 4''$  image of EPIC 211964830 inset in Figure 3.

### 2.3 Archival imaging

To investigate the possibility of a present-day chance alignment with a background source, we queried  $1' \times 1'$  POSS1 images centered on EPIC 211964830 from the STScI Digitized Sky Survey.<sup>2</sup> The proper motion of EPIC 211964830 is large enough that the imaging from 1950 does not show any hint of a background source at its current position (see Figure 4).

### 2.4 Literature data

To characterize the host star, we began by gathering literature data, including broadband photometry, astrometry, and physical parameters (see Table 2). We sourced the parallax, proper motion,  $G$ ,  $B_p$ ,  $R_p$  band magnitudes, effective temperature  $T_{\text{eff}}$ , and radius  $R_\star$  of EPIC 211964830 from *Gaia* DR2 (Gaia Collaboration et al. 2016, 2018a), as well as optical and infrared photometry from the SDSS (Ahn et al. 2012), Pan-STARRS (Chambers et al. 2016), UKIDSS (Lawrence et al. 2007), 2MASS (Cutri et al. 2003), and All *WISE* (Cutri & et al. 2013) catalogs.

## 3 ANALYSIS

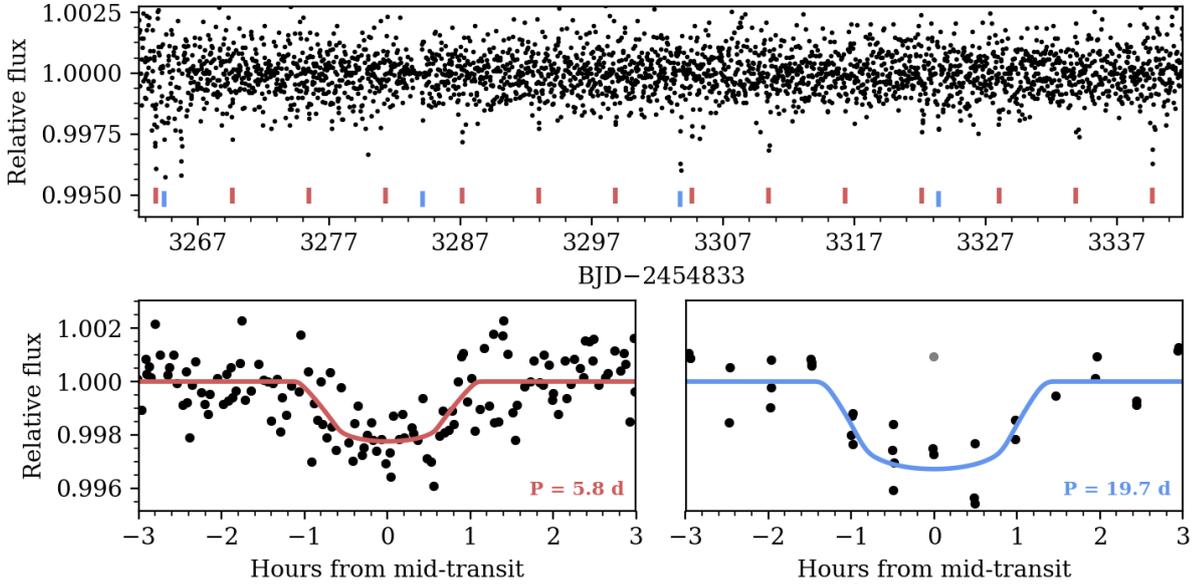
### 3.1 Transit modeling

To model the transits, we first subtracted long term trends caused by stellar variability or instrument systematics using a cubic spline with knots every 0.75 days. We adopted a Gaussian likelihood function and the analytic transit model of Mandel & Agol (2002) as implemented in the Python package *batman* (Kreidberg 2015), assuming a linear ephemeris and quadratic limb darkening. For Markov Chain Monte Carlo (MCMC) exploration of the posterior probability surface, we used the Python package *emcee* (Foreman-Mackey et al. 2013). To reduce unnecessary computational expense, we only fit the light curves in  $4 \times T_{14}$  windows centered on the individual mid-transit times. During MCMC we allowed the free parameters: orbital period  $P_{\text{orb}}$ , mid-transit time  $T_0$ , scaled planet radius  $R_p/R_\star$ , scaled semi-major axis  $a/R_\star$ , impact parameter  $b \equiv a \cos i/R_\star$ , and quadratic limb-darkening coefficients ( $q_1$  and  $q_2$ ) under the transformation of Kipping (2013). We also fit for the logarithm of the Gaussian errors ( $\log \sigma$ ) and a constant out-of-transit baseline offset, which was included to minimize any potential biases in parameter estimates arising from the normalization of the light curve. We imposed Gaussian priors on the limb darkening coefficients, with mean and standard deviation determined by Monte Carlo sampling an interpolated grid of the theoretical limb darkening coefficients tabulated by Claret et al. (2012), enabling the propagation of uncertainties in host star effective temperature  $T_{\text{eff}}$ , surface gravity  $\log g$ , and metallicity  $[\text{Fe}/\text{H}]$  (see Table 2).

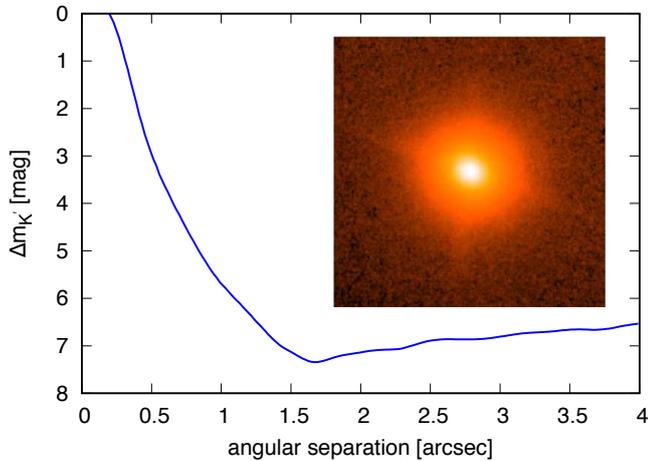
We refined initial parameter estimates from BLS by performing a preliminary nonlinear least squares fit using the Python package *lmfit* (Newville et al. 2014), and then initialized 100 “walkers” in a Gaussian ball around the least squares solution. We ran MCMC for 5000 steps and visually inspected the chains and posteriors to ensure they were smooth and unimodal, and we computed the autocorrelation time<sup>3</sup> of each parameter to ensure that we had collected 1000’s of effectively independent samples after discarding the first 2000 steps as “burn-in.” We also performed transit fits allowing for eccentricity, but found it to be poorly constrained by the light curve: the upper limits are  $e_b < 0.79$  and  $e_c < 0.87$  (95% confidence). We show the joint posterior distributions of  $\rho_\star$ ,  $b$ , and  $R_p/R_\star$  for both planets in Figure 5, derived from the MCMC samples obtained as described above. Because we imposed no prior on the mean stellar density, we can confirm that the mean stellar densities derived from the transits of each planet agree with each other and with the density we derive for the host star in Section 3.2. The mean stellar densities from the transit fits of planets b and c are  $6.49^{+3.85}_{-4.21}$  and  $9.23^{+5.75}_{-5.60}$   $\text{g cm}^{-3}$ , respectively. These values are in good agreement with each other and with our independent determination of EPIC 211964830’s mean stellar density  $\rho_\star = 6.61 \pm 0.32$   $\text{g cm}^{-3}$ , which provides additional confidence that the observed transit signals both originate from EPIC 211964830. Having confirmed this agreement, we perform a final MCMC analysis assuming a circular orbit and including a Gaussian prior on the mean stellar density.

<sup>2</sup> [http://archive.stsci.edu/cgi-bin/dss\\_form](http://archive.stsci.edu/cgi-bin/dss_form)

<sup>3</sup> <https://github.com/dfm/acor>



**Figure 2.** *K2* photometry for EPIC 211964830 with transits of the planets indicated by tick marks (top), and the same photometry phase-folded on the orbital period of each planet (bottom). The best-fitting transit models are shown in red and blue for planets b and c, respectively, which also correspond to the color of the tick marks in the top panel.



**Figure 3.**  $5\text{-}\sigma$  background sensitivity limit (blue curve) and inset  $4''\times 4''$  image of EPIC 211964830 (inset). The x-axis is angular separation from EPIC 211964830 in arcseconds, and the y-axis is differential magnitude in the  $K_s$  band.

We report the median and 68% credible interval of the resulting marginalized posterior distributions in Table 1.

### 3.2 Stellar characterization

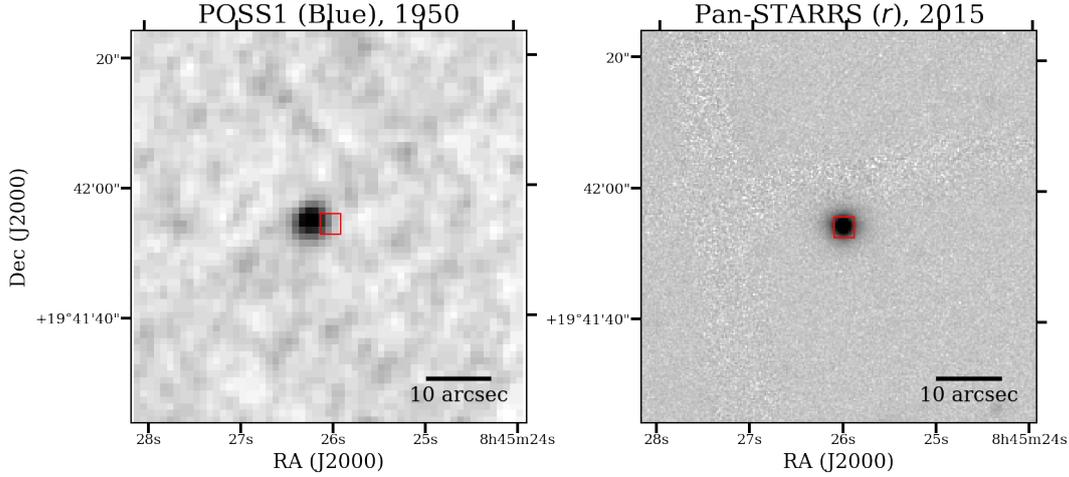
EPIC 211964830 is an M2 dwarf star in the Praesepe open cluster (Jones & Stauffer 1991; Kraus & Hillenbrand 2007; Wang et al. 2014; Gaia Collaboration et al. 2018b). Estimates for the age of Praesepe lie in the range 600–800 Myr (e.g. Kraus & Hillenbrand 2007; Fossati et al. 2008; Brandt

**Table 1.** Planet parameters

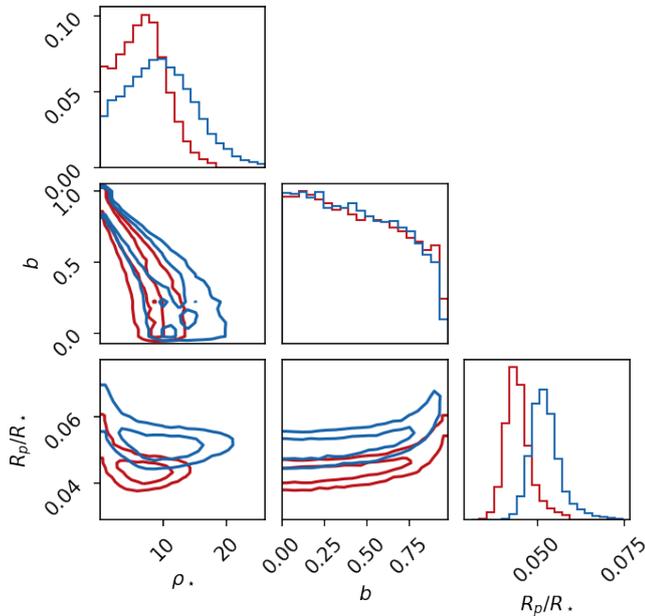
Parameter	Unit	Planet b	Planet c
<i>Free parameters</i>			
$P$	days	$5.840002^{+0.000676}_{-0.000662}$	$19.660302^{+0.003496}_{-0.003337}$
$T_0$	BJD	$2458102.59177^{+0.00428}_{-0.00523}$	$2458117.09169^{+0.00485}_{-0.00447}$
$R_p$		$0.04318^{+0.00275}_{-0.00259}$	$0.05164^{+0.00368}_{-0.00354}$
$a$	$R_\star$	$22.84^{+0.36}_{-0.38}$	$51.30^{+0.82}_{-0.84}$
$b$	—	$0.40^{+0.16}_{-0.23}$	$0.55^{+0.12}_{-0.20}$
$\log(\sigma)$	—	$-6.89^{+0.06}_{-0.06}$	$-7.05^{+0.11}_{-0.10}$
$q_1$	—	$0.51^{+0.11}_{-0.10}$	$0.51^{+0.12}_{-0.10}$
$q_2$	—	$0.25^{+0.03}_{-0.03}$	$0.25^{+0.03}_{-0.03}$
<i>Derived parameters</i>			
$R_p$	$R_\oplus$	$2.231^{+0.151}_{-0.145}$	$2.668^{+0.201}_{-0.194}$
$T_{\text{eq}}$	K	$496 \pm 10$	$331 \pm 7$
$a$	AU	$0.05023^{+0.00042}_{-0.00043}$	$0.11283^{+0.00095}_{-0.00097}$
$i$	deg	$89.01^{+0.58}_{-0.40}$	$89.38^{+0.22}_{-0.13}$
$T_{14}$	hours	$1.884^{+0.118}_{-0.149}$	$2.618^{+0.271}_{-0.233}$
$T_{23}$	hours	$1.701^{+0.137}_{-0.176}$	$2.256^{+0.325}_{-0.302}$
$R_{p,\text{max}}$	$R_\star$	$0.05114^{+0.01380}_{-0.00825}$	$0.07426^{+0.02527}_{-0.01822}$

& Huang 2015), which is consistent with a recent estimate using data from *Gaia* DR2 of  $\log(\text{age}) = 8.85^{+0.08}_{-0.06}$  by Gaia Collaboration et al. (2018b). Because the analysis of Brandt & Huang (2015) accounts for rotation, their older age estimate of  $790 \pm 60$  Myr is likely to be more accurate than earlier determinations, but we adopt the full range to be conservative.

As a preliminary assessment of the stellar parameters of EPIC 211964830, we built the spectral energy distribution (SED; Fig.6) of EPIC 211964830 using the optical and infrared magnitudes listed in Table 2. We did not include the All *WISE* W3 and W4 magnitudes because the former has a signal-to-noise ratio of  $\text{SNR}=3.7$ , while the latter is an up-



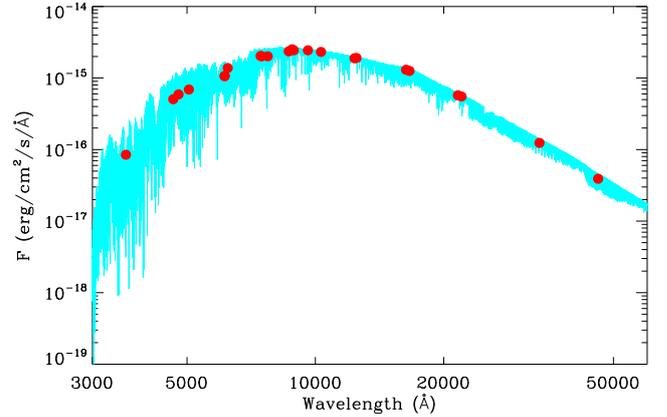
**Figure 4.** Archival imaging from POSS1 (left) and Pan-STARRS (right), with the position of EPIC 211964830 indicated by a red square.



**Figure 5.** Joint posterior distributions of  $\rho_*$ ,  $b$ , and  $R_p/R_*$  without using a prior on stellar density, with 1- and 2- $\sigma$  contours. As in Figure 2, planet b is in red and planet c is in blue.

per limit. We used the web-tool *VOSA*<sup>4</sup> (Version 6; Bayo et al. 2008) to compare the SED to the grid of *BT-Settl* synthetic model spectra of very-low-mass stars (Allard et al. 2012). *VOSA* is a virtual observatory tool specifically designed to derive stellar fundamental parameters (e.g., effective temperature, metallicity, gravity, luminosity, interstellar extinction) by comparing the observed SED to theoretical models. We found that EPIC 211964830 has an effective temperature of  $T_{\text{eff}} = 3500 \pm 50$  K, a surface gravity of  $\log g = 5.00 \pm 0.25$  (cgs), and a metallicity of  $[M/H] = 0.30 \pm 0.15$  dex. Assuming a normal value for the total-to-selective extinction

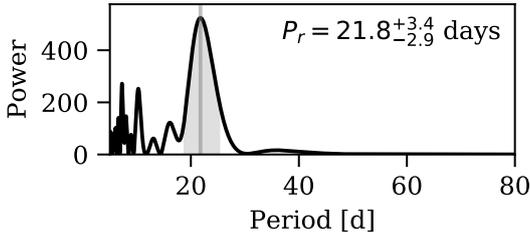
<sup>4</sup> <http://svo2.cab.inta-csic.es/theory/vosa>.



**Figure 6.** Spectral energy distribution of EPIC 211964830. The red circles mark the observed fluxes as derived from the optical and infrared magnitudes listed in Table 2. The best fitting *BT-Settl* is overlotted with a light blue thick line.

( $R = A_V/E(B - V) = 3.1$ ), we derived an interstellar extinction of  $A_V = 0.03 \pm 0.03$  mag. We note that both metal content and extinction are consistent with the average values measured for other member stars of the Praesepe open cluster (see, e.g., Boesgaard et al. 2013; Yang et al. 2015). We used *Gaia* DR2 parallax to determine the luminosity and radius of EPIC 211964830. Following Luri et al. (2018), we accounted for systematic errors in *Gaia* astrometry by adding 0.1 mas in quadrature to the parallax uncertainty of EPIC 211964830 from *Gaia* DR2. Assuming a black body emission at the star's effective temperature, we found a luminosity of  $L_\star = 0.0329 \pm 0.0014 L_\odot$  and a radius of  $R_\star = 0.493 \pm 0.018 R_\odot$ .

To obtain the final set of stellar parameters we use in this work, we utilized the *isochrones* (Morton 2015a) Python interface to the *Dartmouth* stellar evolution models (Dotter et al. 2008) to infer stellar parameters using the 2MASS *JHKs* photometry and *Gaia* DR2 parallax (with augmented uncertainty to account for systematics as above). *isochrones* uses the *MultiNest* (Feroz et al. 2013) algo-



**Figure 7.** Lomb-Scargle periodogram of the K2 light curve of EPIC 211964830.

rithm to sample the posteriors of fundamental stellar properties of interest, and resulted in the following constraints: effective temperature  $T_{\text{eff}} = 3660^{+80}_{-45}$  K, surface gravity  $\log g = 4.783 \pm 0.012$  (cgs), metallicity  $[\text{Fe}/\text{H}] = -0.013 \pm 0.180$  dex, radius  $R_{\star} = 0.473 \pm 0.011 R_{\odot}$ , mass  $M_{\star} = 0.496 \pm 0.013 M_{\odot}$ , extinction ( $A_V$ ) =  $0.301 \pm 0.162$  mag, and distance =  $187.0 \pm 4.0$  pm. We opted not to include a prior on the metallicity of EPIC 211964830 based on its cluster membership, as the resulting stellar parameter uncertainties may not accurately reflect intrinsic variability of metallicity within the Praesepe birth nebula. The posteriors were consistent with the results of our SED analysis and with Praesepe membership; metallicity is poorly constrained, but is consistent with that of Praesepe Boesgaard et al. ( $[\text{Fe}/\text{H}] = 0.12 \pm 0.04$ ; 2013). Most posteriors appeared roughly symmetric and Gaussian, so we list the median and standard deviation in Table 2; the  $T_{\text{eff}}$  posterior was asymmetric, so we list the median and 68% credible region instead. We note that these values are in moderate disagreement with the stellar parameters computed by Huber et al. (2016), which may be due to the lack of a parallax constraint in their analysis, but may also reflect a systematic bias for low mass stars, which has been attributed to their choice of stellar models (Dressing et al. 2017). We note that these estimates are consistent with the *Gaia* DR2 values for EPIC 211964830 ( $T_{\text{eff}} = 3422^{+478}_{-22}$  K,  $R_{\star} = 0.54^{+0.01}_{-0.12} R_{\odot}$ , distance =  $186.573 \pm 2.105$  pc).

The light curve of EPIC 211964830 exhibits clear quasi-periodic rotational modulation, which is characteristic of surface magnetic activity regions moving in and out of view as the star rotates around its axis. We measured the rotation period using two different methods. After masking the transits from the *K2* light curve and subtracting a linear trend, we computed the Lomb-Scargle periodogram, from which we derived a stellar rotation period of  $21.8^{+3.4}_{-2.9}$  days by fitting a Gaussian to the peak (see Figure 7). In Figure 8 we show a Gaussian Process (GP) fit to the light curve using a quasi-periodic kernel (e.g. Haywood et al. 2014; Grunblatt et al. 2015; Dai et al. 2017), from which we measured a rotation period of  $22.2 \pm 0.6$  days. Although the GP-based approach is in agreement and yields higher precision, we conservatively adopt the Lomb-Scargle estimate due to its larger error bars. The rotational modulation of EPIC 211964830 has a similar period and amplitude to K2-95, an M3 dwarf in Praesepe hosting a transiting sub-Neptune (Obermeier et al. 2016). Using the gyrochronology relation of Angus et al. (2015) we find that the measured stellar rotation period is consistent with the age of Praesepe.

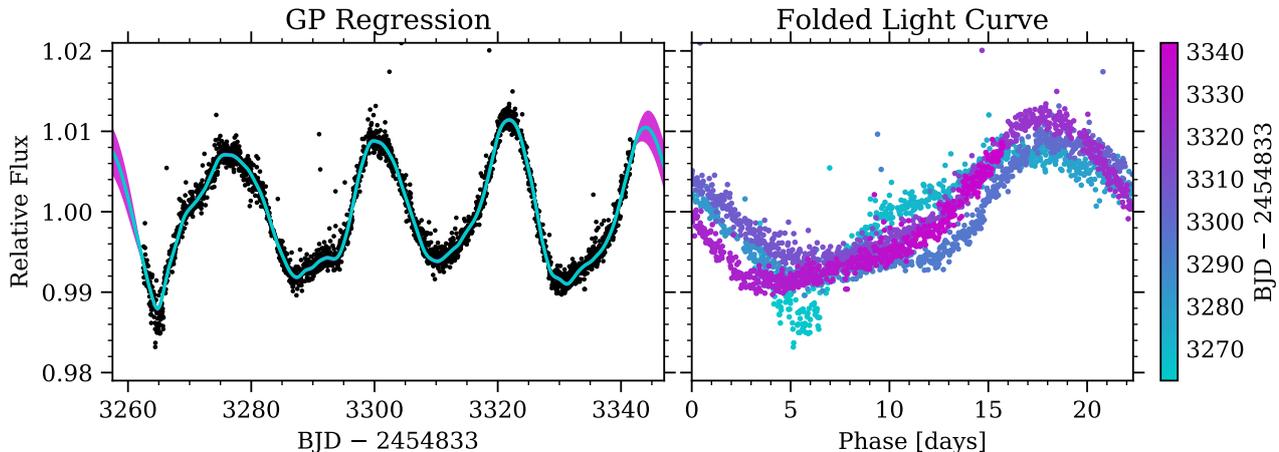
**Table 2.** Stellar parameters

Parameter	Unit	Value	Source
<i>Astrometry</i>			
$\alpha$ R.A.	deg	131.358352378	<i>Gaia</i> DR2
$\delta$ Dec.	deg	19.698400987	<i>Gaia</i> DR2
$\pi$	mas	$5.3598 \pm 0.0605$	<i>Gaia</i> DR2
$\mu\alpha$	$\text{mas yr}^{-1}$	$-37.900 \pm 0.095$	<i>Gaia</i> DR2
$\mu\delta$	$\text{mas yr}^{-1}$	$-13.079 \pm 0.061$	<i>Gaia</i> DR2
<i>Photometry</i>			
$K_p$	mag	15.318	EPIC
$B_p$	mag	$16.946 \pm 0.006$	<i>Gaia</i> DR2
$R_p$	mag	$14.538 \pm 0.002$	<i>Gaia</i> DR2
$G$	mag	$15.663 \pm 0.001$	<i>Gaia</i> DR2
$u$	mag	$19.994 \pm 0.036$	Sloan/SDSS
$g$	mag	$17.499 \pm 0.005$	Sloan/SDSS
$r$	mag	$16.089 \pm 0.004$	Sloan/SDSS
$i$	mag	$14.963 \pm 0.004$	Sloan/SDSS
$z$	mag	$14.374 \pm 0.004$	Sloan/SDSS
$g$	mag	$17.260 \pm 0.006$	Pan-STARSS
$r$	mag	$16.075 \pm 0.002$	Pan-STARSS
$i$	mag	$14.965 \pm 0.003$	Pan-STARSS
$z$	mag	$14.471 \pm 0.002$	Pan-STARSS
$y$	mag	$14.221 \pm 0.004$	Pan-STARSS
$Z$	mag	$13.848 \pm 0.002$	UKIDSS
$J$	mag	$12.997 \pm 0.002$	UKIDSS
$H$	mag	$12.393 \pm 0.001$	UKIDSS
$K$	mag	$12.157 \pm 0.001$	UKIDSS
$J$	mag	$13.047 \pm 0.025$	2MASS
$H$	mag	$12.386 \pm 0.022$	2MASS
$K_s$	mag	$12.183 \pm 0.020$	2MASS
$W1$	mag	$12.048 \pm 0.023$	All <i>WISE</i>
$W2$	mag	$11.978 \pm 0.023$	All <i>WISE</i>
$W3$	mag	$11.317 \pm 0.294$	All <i>WISE</i>
$W4$	mag	8.173	All <i>WISE</i>
<i>Physical parameters</i>			
$T_{\text{eff}}$	K	$3660^{+80}_{-45}$	This work
$\log g$	cgs	$4.783 \pm 0.012$	This work
$[\text{Fe}/\text{H}]$	dex	$-0.013 \pm 0.180$	This work
$M_{\star}$	$M_{\odot}$	$0.496 \pm 0.013$	This work
$R_{\star}$	$R_{\odot}$	$0.473 \pm 0.011$	This work
$\rho_{\star}$	$\text{g cm}^{-3}$	$6.610 \pm 0.322$	This work
$A_V$	mag	$0.301 \pm 0.162$	This work
distance	pc	$187.0 \pm 4.0$	This work
$P_r$	days	$21.8^{+3.4}_{-2.9}$	This work

### 3.3 Validation

Transiting planet false positive scenarios typically involve an eclipsing binary (EB) blended with a brighter star within the photometric aperture. If an EB's mass ratio is close to unity, then the primary and secondary eclipses will have the same depth, and in such a case the dilution from the brighter star will make these eclipses shallower and thus more similar to planetary transits. In such a scenario, the EB's orbit must also be circular such that the eclipses mimic the regular periodicity of planetary transits. Another possibility is an extreme EB mass ratio, in which case the (diluted) secondary eclipses would be small enough that they could be below the detection limit of the photometry. Because of the large ( $4''$ ) pixel scale of the *Kepler* photometer, blended EB scenarios are not rare, and must therefore be properly accounted for. Such a false positive scenario could be caused by the chance alignment of a background source (BEB), or by a hierarchical triple system (HEB), the relative frequencies of which depend on the density of sources in the vicinity of the candidate host star.

To investigate the possibility of a BEB false positive scenario, we utilize the observed transit geometry in conjunc-



**Figure 8.** Gaussian Process fit to the light curve of EPIC 211964830 with transits removed (left), and the same light curve folded on the maximum a posteriori rotation period of 22.3 days (right; color of datapoints correspond to time).

tion with a simulated stellar population appropriate for the line of sight to EPIC 211964830. The eclipse depth of an EB can in principle reach a maximum of 100%, which sets a limit on the faintness of any putative background sources that could be responsible for the observed signals. Using Equation 1 of Livingston et al. (2018b) and the observed transit depths, this corresponds to  $Kp \approx 22$  mag. Using a simulated stellar population in the direction of EPIC 211964830 from TRILEGAL Galaxy model (Girardi et al. 2005), the expected frequency of sources brighter than this limit is very low, at  $\sim 0.07$  for a  $7''$  photometric aperture (see Figure 1). Indeed, the non-detection of any background sources in our AO image (see Figure 3) and the POSS1 image from 1950 (see Figure 4) is consistent with the expectation of zero such sources from the Galaxy model.

If, on the other hand, the observed signals are actually the result of a HEB scenario, we must instead consider the possibility that EPIC 211964830 is actually a bound triple star system. In order for the eclipsing component to have a negligible impact on the observed SED (see Figure 6), it would need to be composed of stars with much lower masses than EPIC 211964830. However, from the observed transit geometry we have  $3\text{-}\sigma$  upper limits on the radius ratio of 9% and 15% for the inner and outer planets, respectively, using Equation 21 of Seager & Mallén-Ornelas (2003) (see  $R_{p,\max}$  in Table 1). Radius ratios below this limit would involve either an eclipsing component in the planetary mass regime or an occulted component that would contribute non-negligible flux to the combined SED and thereby have observable signatures. Perhaps most importantly, the existence of two periodic transit-like signals from the same star is *a priori* more difficult to explain with non-planetary scenarios, because the BEB and HEB scenarios consistent with the observed signals would require vanishingly infrequent chance alignment or higher stellar multiplicity. Indeed, candidates in systems of multiple transiting planets have been shown to have a very low false positive rate (Lissauer et al. 2012), and are thus essentially self-validating.

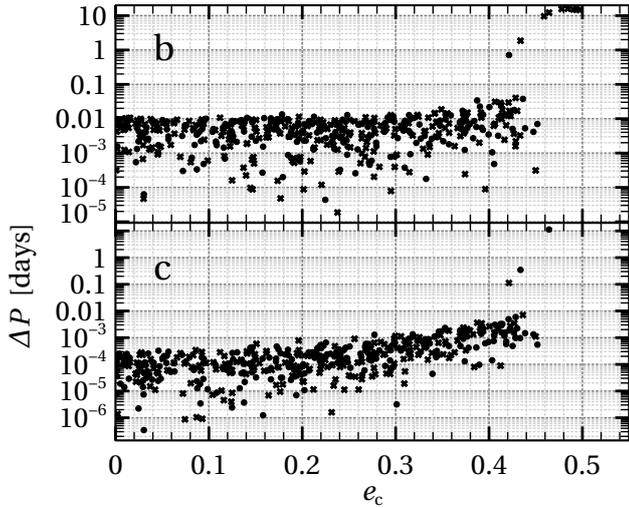
Besides these qualitative considerations, we also computed the false positive probabilities (FPPs) of the planet

candidates of EPIC 211964830 using the Python package *vespa* (Morton 2015b). *vespa* employs a robust statistical framework to compare the likelihood of the planetary scenario to likelihoods of several astrophysical false positive scenarios involving eclipsing binaries, relying on simulated eclipsing populations based on TRILEGAL. The FPPs from *vespa* for planets b and c are 0.007% and 0.012%, respectively, well below the standard validation threshold of 1%. Moreover, these FPPs are overestimated due to the fact that *vespa* does not account for multiplicity: Lissauer et al. (2012) demonstrated that a candidate in a system with one or more additional transiting planet candidates is 25 times more likely to be a planet based on multiplicity alone. Therefore, in addition to the qualitative arguments above, the planet candidates also quantitatively warrant validation; we conclude that EPIC 211964830 is thus the host of two *bona fide* transiting planets.

### 3.4 Dynamical stability

Given the large separation between the two planets, the system is manifestly Hill stable. Assuming that the orbits are circular, their separation is about 25 times their mutual Hill radius, much larger than the threshold value of  $3.46R_H$  (Gladman 1993; Chambers et al. 1996; Deck et al. 2013). Using the angular momentum deficit criterion of Petit et al. (2018), we find that the eccentricity of the outer planet must be less than  $e_c = 0.4$  to ensure the stability of the system.

We use the probabilistic mass-radius relation of Wolfgang et al. (2016) to estimate the masses of the planets given their measured radii, yielding  $7.7 \pm 2.3$  and  $9.5 \pm 2.7 M_\oplus$  for planets b and c, respectively. We use the MERCURIUS integrator of the REBOUND package (Rein & Liu 2012; Rein & Tamayo 2015) to simulate the orbital evolution of 500 realizations. Consistent with the planets' orbital inclinations from the measured transit geometry, we set their mutual inclination to zero and sampled the eccentricity of the outer planet between 0 and 0.5. Figure 9 shows the difference between the initial orbital periods and the final ones, after 2 Myr of integration. For all systems with  $e_c < 0.45$  the



**Figure 9.** Difference between the initial and final orbital periods of the planets in the  $N$ -body simulations as a function of the initial eccentricity of the outer planet  $c$ . Top panel: period difference of the inner planet  $b$ . Bottom panel: period difference of the outer planet  $c$ . Circles and crosses represent a period increase and decrease with respect to the initial one, respectively.

difference in orbital periods remain well below 0.01 days. For  $e_c > 0.45$ , the perturbations between the two planets lead to instability and chaotic evolution. Therefore,  $e_c \approx 0.45$  is a robust upper limit for the eccentricity of the outer planet.

We also calculate the eccentricity damping timescale for planet  $b$  using the tidal model of Hut (1981). Since EPIC 211964830 is a low mass star with a convective envelope, we can compute the tidal timescale  $k/T$  due to the tides raised on the star from the stellar structure parameters (Zahn 1977). We use the stellar models of the PARSEC stellar evolution code (Bressan et al. 2012; Chen et al. 2015; Fu et al. 2018), and derive  $k/T = 0.34 \text{ yr}^{-1}$ . Considering only the tides raised on the star, we find that the circularization timescale is much longer than the age of the system for all  $e_b < 0.9$ . We also take into account the tides raised on the planet  $c$  using the tidal quality factor  $Q = nT/3k$ , where  $n$  is the mean motion of planet  $c$ . Note that the stellar  $k/T$  corresponds to a quality factor of  $\sim 350$ , which tells us that any  $Q > 350$  makes the planetary tide less efficient than the stellar tide. Since the tidal quality factor for gas giants is expected to be  $Q \gg 10^2$ , we conclude that the inner planet could be still undergoing tidal circularization.

#### 4 DISCUSSION

Assuming a bond albedo of 0.3, the equilibrium temperatures of planets  $b$  and  $c$  are  $496 \pm 13$  and  $331 \pm 8$  K, respectively, making EPIC 211964830 a system of two warm sub-Neptunes. Although such planets have been found in large numbers by previous surveys (e.g. *Kepler*), the number orbiting cluster stars is extremely small. EPIC 211964830 is thus an important system because it significantly improves the statistics for demographic studies of cluster planets. Furthermore, prior to this discovery only one member of a cluster was known to host multiple transiting planets (K2-136; Mann et al. 2018; Livingston et al. 2018a; Ciardi et al.

2018). EPIC 211964830 is thus a unique laboratory for studies of system architectures as a function of time. We place EPIC 211964830 in the context of the general exoplanet population, as well as other cluster planets, by plotting planet radius as a function of host star mass in Figure 10, using data from a query of the NASA Exoplanet Archive<sup>5</sup> (Akeson et al. 2013). From this perspective, the planets of EPIC 211964830 do not appear to have significantly inflated radii, as has previously been a matter of speculation for cluster systems (e.g. K2-25, Mann et al. 2016b; K2-95, Obermeier et al. 2016). It is worth noting, however, that K2-25 and K2-95 have lower masses than EPIC 211964830, and the radii of planets orbiting higher mass host stars in both Hyades and Praesepe appear less inflated. Two cluster planets buck this apparent trend: K2-33 and K2-100. However, K2-33 may still be undergoing radial Kelvin–Helmholtz contraction due to its young age (5–10 Myr; David et al. 2016b), and K2-100 is much more massive ( $M_\star = 1.18 \pm 0.09 M_\odot$ ; Mann et al. 2017). The radii of the planets orbiting EPIC 211964830 lend support to this trend, and thus to the hypothesis that radius inflation results from higher levels of X-ray and ultraviolet (UV) flux incident upon planets orbiting lower mass stars; the absence of such a trend for field stars may tell us something about the timescales of radial relaxation after early-stage X-ray/UV flux from low mass stars diminishes.

Planets orbiting cluster stars are expected to have large eccentricities and large mutual inclinations, if perturbations from cluster members are efficient. While we cannot yet constrain the eccentricity of the outer planet, it is safe to assume that the system is coplanar. Even in the hypothesis of the presence of an outer, inclined, non-transiting planet, produced during a stellar encounter in the early life of the cluster, perturbations from the outer planet would have propagated inward, altering the inclinations of the inner planets. Cai et al. (2018) show that planetary systems in the outskirts of the cluster (i.e. outside its half-mass radius) are unlikely to have been perturbed by passing stars. We compute the distance of EPIC 211964830 from the center of Praesepe, using the cluster center coordinates derived by Khalaj & Baumgardt (2013) and the coordinates of EPIC 211964830 from *Gaia* DR2. We find that EPIC 211964830 lies at  $4.365 \pm 0.206$  pc projected distance ( $8.8 \pm 4.2$  unprojected) from the cluster center, well outside the half-mass radius of the Praesepe cluster (3.9 pc, Khalaj & Baumgardt 2013). This suggests that perturbations from other stars have likely played a minor role in shaping the planetary system of EPIC 211964830.

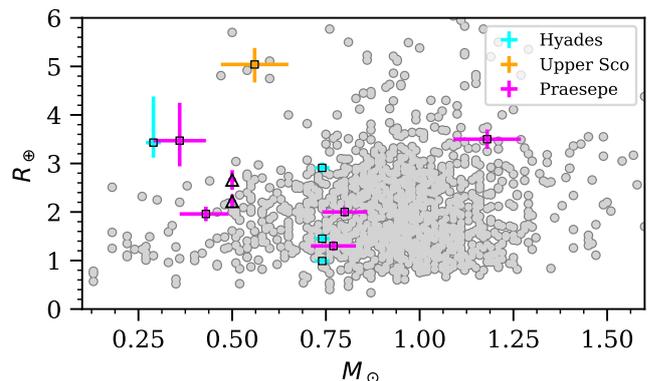
Because the planets orbiting EPIC 211964830 have a common host star history, X-ray and UV stellar flux at young ages can be controlled for, better enabling their observed radii to yield insights into atmospheric evolution due to irradiation from the host star. Additionally, the 600–800 Myr age of the system is particularly good for testing photo-evaporation theory, as this is the timescale over which photo-evaporation should have finished (Owen & Wu 2013); by this age, the radius distribution of small planets should approach that of field stars. The planet radii place them both securely above the observed gap in the radius distribution (Fulton

<sup>5</sup> <https://exoplanetarchive.ipac.caltech.edu/>

et al. 2017; Van Eylen et al. 2018; Berger et al. 2018), which suggests either that they have large enough core masses to have retained substantial atmospheres, or that photoevaporation may have played a less significant role in their evolution, or both. However, the host star’s spectral type indicates substantial X-ray/UV irradiation during the first few hundred million years, which makes it less likely that the planets could have largely escaped the effects of photoevaporation unless they had larger core masses. Indeed, the location of the bimodality has been shown to shift to smaller radii for lower mass host stars (Fulton & Petigura 2018), consistent with the expectation that smaller stars produce smaller planet cores. This implies that the planets orbiting EPIC 211964830 are more likely to have relatively massive cores and always occupied the larger radius mode. Given the age of the system, it is likely that photoevaporation is effectively over, and the planet radii will no longer undergo substantial evolution.

Systems of multiple transiting planets sometimes allow for the masses and eccentricities in the system to be measured via dynamical modeling of the observed transit timing variations (TTVs; Holman & Murray 2005; Agol et al. 2005), in which the mutual gravitational interaction between planets produces regular, measurable deviations from a linear ephemeris. To test if either planet exhibits TTVs, we used the best-fitting transit model as a template for the determination of individual transit times. Keeping all parameters fixed except the mid-transit time, we fitted this template to each transit in the data, but we did not detect any TTVs over the  $\sim 80$  days of *K2* observations. The absence of TTVs is perhaps not surprising given that the orbital periods are not especially close to a low-order mean motion resonance, with  $P_c/P_b \approx 3.367$ , about 12% outside of a 3:1 period commensurability. Although the planets do not exhibit measurable mutual gravitational interactions, the pull they exert on their host star presents an opportunity for characterization via Doppler spectroscopy.

By obtaining precise RV measurements of EPIC 211964830, it may be possible to measure the reflex motion of the host star induced by the gravity of its planets (e.g. Struve 1952; Mayor & Queloz 1995). Such measurements would yield the planet masses and mean densities, which would constrain the planets’ interior structures. The predicted masses of planets b and c, along with their orbital periods and the mass of the host star, yield expected RV semi-amplitudes values of  $4.4 \pm 1.3$  and  $3.6 \pm 1.0$   $\text{m s}^{-1}$ , respectively. However, the youth and photometric variability of EPIC 211964830 imply RV stellar activity signals larger in amplitude than the expected planet signals from optical spectroscopy. This suggests that the planets of EPIC 211964830 may be amenable to mass measurement using a high precision NIR spectrograph, such as IRD (Tamura et al. 2012) or HPF (Mahadevan et al. 2012), as the RV amplitude of stellar activity signals should be significantly lower in the infrared. Assuming no orbital obliquity, from the radius and rotation period of EPIC 211964830 we estimate low levels of rotational line broadening, with  $v \sin i$  of  $\sim 1$   $\text{km s}^{-1}$ . Prior knowledge about the star’s rotation period from *K2* should prove useful for modeling the stellar activity signal simultaneously with the Keplerian signals of the planets using a Gaussian Process



**Figure 10.** Planet radius versus host star mass of EPIC 211964830 (triangles) and a selection of other transiting planet systems in clusters (squares), as compared to the field star planet population (gray points). Besides EPIC 211964830, the data shown are from a query of the NASA Exoplanet Archive (Akeson et al. 2013). Besides EPIC 211964830, the cluster systems shown are: K2-25 and K2-136 (Hyades; Mann et al. 2016b, 2018; Livingston et al. 2018a; Ciardi et al. 2018); K2-33 (Upper Sco; David et al. 2016b; Mann et al. 2016b); K2-95, K2-100, K2-101, K2-102, K2-103, and K2-104 (Praesepe; Obermeier et al. 2016; Mann et al. 2017)

model (e.g. Haywood et al. 2014; Grunblatt et al. 2015; Dai et al. 2017).

Besides spectroscopy, follow-up NIR transit photometry of EPIC 211964830 could enable a better characterization of the system by more precisely measuring the transit geometry. Besides yielding a better constraint on the planet radius, transit follow-up would also significantly refine estimates of the planets’ orbital ephemerides, enabling efficient scheduling for any subsequent transit observations, e.g. with *JWST*. Using the *WISE* W2 magnitude in Table 2 as a proxy for *Spitzer* IRAC2, the expected transit SNR is in the range 4–8; given the systematic noise in *Spitzer* light curves, such transit measurements would be challenging, but may be feasible by simultaneously modeling the transit and systematics signals using methods such as pixel-level decorrelation (PLD; Deming et al. 2015). Furthermore, by simultaneously modeling the *K2* and *Spitzer* data, *Spitzer*’s high photometric observing cadence and the diminished effects of limb-darkening in the NIR could be leveraged to more precisely determine the transit geometry (Livingston et al., in review). NIR transit observations from the ground could also be useful, but would likely require a large aperture (e.g. 4–8 meter) telescope to yield better performance than *Spitzer*.

## 5 SUMMARY

Using data from the *K2* mission and ground-based follow-up observations, we have detected and statistically validated two warm sub-Neptunes transiting the star EPIC 211964830, which is a member of the 600–800 Myr Praesepe open cluster. Unlike several previously discovered planets orbiting lower mass stars in clusters, their radii are fairly consistent with the those of planets orbiting field stars of comparable mass to their host, suggesting that radius inflation is a function

of host star mass. The system presents opportunities for RV follow-up using high precision NIR spectrographs, which would yield the planets' densities and thereby test theories of planet formation and evolution. NIR transit photometry could more precisely measure the planet's ephemerides and transit geometry, and thus also their radii. By leveraging the known age of the system, such characterization would yield a direct view of the planets' atmospheric evolution. EPIC 211964830 joins a small but growing list of cluster planets, and is particularly valuable as it is only the second known system of multiple transiting planets in a cluster.

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