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A fresh view on carbon radio recombination lines powered by LOFAR

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TIED-ARRAY HOLOGRAPHY WITH LOFAR

5.1. INTRODUCTION

A radio telescope works by combining the signals received by the elements that constitute its aperture (a reflecting surface in the case of a dish, or an array of antennas in the case of a phased array). In order to maximize the sensitivity of the telescope towards a particular direction, the signals arriving from that direction must be combined in phase, i.e., the time difference between the signals received by different aperture elements must be zero. In the case of a dish this is accomplished by shaping the reflecting surface in such a way that all the signals arrive at the receiver at the same time. And in a phased array, by introducing instrumental time delays between its elements to compensate for the time of arrival of the signal at the antennas (e.g., Thompson et al., 2017).

Deviations from a perfect phase alignment when the signals are combined lead to a loss in the efficiency of the telescope (e.g., Ruze, 1952, 1966; D’Addario, 2008). These phase misalignments can be caused by the telescope itself, or they can be produced in the path between the source of the signals and the telescope. An example of the former are phase differences caused by misaligned panels in a reflector (e.g., Baars et al., 2007), or uncorrected cable delays in a phased array.

Different methods to reduce phase misalignments between the elements of an aperture have been developed. These include photogrammetric measurements (e.g., Wiktowy et al., 2003), direct measurement of the aperture distribution (e.g., Chen et al., 1998; Naruse et al., 2009), holographic measurements (e.g., Napier & Bates, 1973; Bennett et al., 1976; Scott & Ryle, 1977; Baars et al., 2007; Hunter et al., 2011) and calibration using astronomical sources (e.g., Fomalont & Perley, 1999; Intema et al., 2009; Thompson et al., 2017; Rioja et al., 2018). This chapter focuses on the holographic measurement of the aperture illumination of a large phased array telescope.

Since the work of Scott & Ryle (1977), holographic measurements have been used to calibrate the dishes of the Very Large Array (VLA, e.g., Kesteven, 1993; Broilo, 1993),

the Atacama Large Millimeter Array (ALMA, e.g., Baars et al., 2007) and the Green Bank Telescope (GBT, e.g., Hunter et al., 2011), to study the primary beam response of the Westerbork radio telescope (WSRT, e.g., Popping & Braun, 2008) and the Allen telescope array dishes (ATA, e.g., Harp et al., 2011), and to characterize the beam and aperture of the Low Frequency Array (LOFAR) stations (Brentjens et al. in prep.). All these measurements have been restricted to the study of apertures $\lesssim 100$ m in diameter.

In the regime of low-frequency and large apertures holographic measurements are particularly challenging. At low-frequencies the ionosphere will introduce additional time delays depending on its total electron content (TEC, e.g., Intema et al., 2009). To accurately measure the intrinsic phase errors between the elements of the phased array without ionospheric distortion, the phased array must be smaller than the diffractive scale of the ionosphere. Night time observations of the ionosphere at 150 MHz show that its diffractive scale is between 30 and 3 km (Mevius et al., 2016).

LOFAR operates at frequencies between 10 MHz and 240 MHz (van Haarlem et al., 2013). This frequency range is covered by two different types of antennas: low band antennas (LBA, 10–90 MHz) and high band antennas (HBA, 120–240 MHz). The HBA antennas are combined in a 4×4 tile with an analog beamformer. The antennas and tiles are grouped into stations, and stations are further combined to form an array. For the core stations of LOFAR, the LBA stations consist of 96 antennas, while the HBA stations have 48 tiles split into two fields. Of the 96 antennas in a core LBA station the available electronics permits only 48 to be actively beamformed. There are 24 stations in the core of LOFAR. The core stations are connected via fiber to a central clock, thus their signals can be added coherently to form a telescope with a maximum baseline of 3.5 km. The stations in the innermost 350 m are known as the *Superterp*.

Each LOFAR LBA dipole observes the entire sky, while the HBA tiles have a field of view (FoV) of 30° at 150 MHz. Since the signals from the antennas and tiles are combined digitally, the stations can point towards multiple directions within their FoV simultaneously (e.g., Barton, 1980; Steyskal, 1987). When the signals from different stations are added together coherently, a phased array (known as a tied-array) is formed. This enables LOFAR to form multiple tied-array beams (TABs) that point towards different directions.

5.2. METHOD

We want to determine the time delays for the array formed by the stations in LOFAR's core. We will refer to the tied-array formed by these stations as the array under test (AUT). In order to determine the time delays, we start from a map of its complex-valued beam B .

The basic procedure used to measure B is the same as that employed by Scott & Ryle (1977), with a difference in its implementation. In their work, a raster scan was used to map the region around the bright unresolved source. Here, we take advantage of LOFAR's multi-beaming capability to map the region around the bright unresolved source. Using multiple TABs the whole region is mapped simultaneously, and there is always a TAB pointing towards the bright unresolved source. Besides speeding up the

process by a factor equal to the number of simultaneous beams, this allows continuous calibration of the AUT as well as the reference stations by always having a TAB at the central calibrator source.

At the frequencies at which LOFAR operates, the Milky Way is bright and it will distort the observed map of B . To reduce the contribution from the Milky Way to the measurements, we use a reference station which produces a baseline that resolves out large scale Galactic structure (e.g., Colegate et al., 2015). The contribution from smaller bright sources (e.g., Cassiopeia A or Cygnus A) cannot be completely resolved out, and is reduced by limiting the field of view (FoV) through time and frequency smearing (e.g., Bridle & Schwab, 1999). Moreover, the AUT and the reference station will see different portions of the ionosphere, which will introduce an additional time delay between them. The effects of the different ionosphere seen by the AUT and the reference station are calibrated using the bright point source.

The details behind the calibration method will be presented in Brentjens et al. in prep. Here we summarize the rationale behind the calibration because it is an important aspect of this work. Following the Measurement Equation formalism (Hamaker et al., 1996), we have that the visibility generated by cross-correlating the signals from the AUT and the reference station is

$$V_b = J_{\text{AUT},b} E J_{\text{ref}}^\dagger, \quad (5.1)$$

where E represents the coherency matrix formed by the pure sky visibilities, and $J_{\text{AUT},b}$ and J_{ref} are the Jones matrices (Jones, 1941) of the AUT and the reference station respectively. The subscript b represents the TABs formed with the AUT. The $\dagger$ symbol denotes taking the conjugate transpose of the corresponding matrix. The calibration consists of finding the inverse of the visibility of the central beamlet, V_c^{-1} , and right multiplying all the visibilities with it. After doing this, for the central beam $\tilde{V}_{b=c} = V_c V_c^{-1} = 1$, where 1 represents the identity matrix. For the remaining directions $\tilde{V}_i = V_i V_c^{-1} = J_{\text{AUT},i} J_{\text{AUT},c}^{-1}$. That is, the calibrated visibility for the i -eth beam only depends on the Jones matrix of the AUT, and not on the sky brightness distribution. These calibrated visibilities map B .

From the observed map of B we determine the amplitude and phase over the aperture of the AUT, A . In the far-field approximation, and for a coplanar array, they are related by (e.g., D’Addario, 1982; Baars et al., 2007; Thompson et al., 2017),

$$B(l, m) \propto \iint A(p, q) e^{2\pi i (pl + qm) \frac{v}{c}} dp dq, \quad (5.2)$$

where i denotes the imaginary unit, c is the speed of light, v is frequency, p and q are orthogonal coordinates in the aperture plane, and l and m are the direction cosines measured with respect to p and q . For LOFAR, the (p, q) coordinate system has its origin at the center of the aperture and it lies in the plane of the station, or in this case, the plane of the superterp stations. The phase of $A(p, q)$ is set by e.g. uncalibrated errors in the clock distribution, cable length, antenna position, and ionospheric phase variations across the aperture.

5.3. OBSERVATIONS

5.3.1. LOFAR HOLOGRAPHY OBSERVATIONS

Table 5.1 summarizes the observations. Each TAB recorded complex voltages in two orthogonal polarizations (X and Y) at $5.12 \mu\text{s}$ time resolution in ten 195.3125 kHz wide spectral windows. The data were subsequently ingested into the LOFAR long term archive. The calibrator sources are selected to be small compared to the size of the TAB, as well as the fringe spacing of the baselines between the AUT and the reference stations. The former prevents systematic distortions in the measured beam, while the latter guarantees high signal-to-noise ratios (SNR) on the baselines towards the reference stations.

The complex valued beam maps were measured on a regular hexagonal grid, 1.32 and 0.37 deg across for LBA and HBA, respectively. The map size is limited by the number of TABs, spectral windows and stations that the beamformer, COBALT (Broekema et al., 2018), can process simultaneously. Per Fourier relation Equation 5.2 this implies a spatial resolution in the aperture plane of 270 m (HBA at 174 MHz) and 170 m (LBA at 68 MHz), comparable to the diameter of the *Superterp* (350 m). To avoid aliasing in the aperture domain, the separation between TABs was set at λ/D at the highest frequency, and kept constant for lower frequencies.

The required integration time is set by the error on the phase in the aperture plane, $\Delta\phi$, (D’Addario, 1982)

$$\Delta\phi \approx \frac{\pi D}{4\sqrt{2}d\text{SNR}_{\text{bm}}}, \quad (5.3)$$

where D is the telescope diameter, d the spatial resolution on the aperture plane and SNR_{bm} the signal-to-noise ratio in the complex-valued beam map, i.e., the ratio between the peak response of the array and the root-mean-squared (rms) over the complex-valued beam map, $\text{SNR}_{\text{bm}} = I/\sigma$. σ can be estimated from (e.g., Napier & Crane, 1982)

$$\sigma = \sqrt{(\text{SEFD}_{\text{CS}}/N_{\text{CS}})(\text{SEFD}_{\text{RS}})/(\sqrt{\Delta\nu\Delta t})}, \quad (5.4)$$

where N_{CS} is the number of stations in the AUT, SEFD_{CS} and SEFD_{RS} are the system equivalent flux density (SEFD) of a core station and a reference station, respectively, $\Delta\nu$ is the bandwidth and Δt the integration time. For the LBA the SEFD of each antenna field is $\approx 30 \text{ kJy}$ at 60 MHz and for the HBA $\approx 3 \text{ kJy}$ at 150 MHz (van Haarlem et al., 2013). Thus, for an observation of 3C196 with integration times of 600 s at 60 MHz and 60 s at 150 MHz we can determine the time delays with errors of 1.8 ns and 0.4 ns, respectively.

For a baseline of 52 km, an integration time of 600 s and an observing frequency of 37 MHz, time and bandwidth smearing (e.g., Bridle & Schwab, 1999) will reduce the response to sources at $45'$ from the map center by 5% and 70%, respectively. For the same baseline and an integration time of 60 s these sources will have their apparent fluxes reduced by 23% at 115 MHz.

Table 5.1.: Observations used in this work.[†]

Observation ID [‡]	Antenna	# antenna fields	Start time	Duration (s)	# TABs	TAB spacing (arcmin)	FoV (degrees)	Source
L658168	HBA ^b	46 ^c	June 14 13:40:00 UT	60	169	1.6	0.37	3C147
L658158	HBA ^b	46 ^c	June 14 13:30:00 UT	60	169	1.6	0.37	3C196
L650445	LBA ^d	24 ^e	April 19 09:20:00 UT	600	271	5	1.32	3C48
L645357 ^a	LBA ^d	24 ^e	March 20 19:45:00 UT	600	271	5	1.32	3C196

[†] The stations used as reference and their distance from the center of the array are; RS210 65 km; RS509 59 km; RS310 52 km; and DE605 226 km.

[‡] For each observation ID the three/four previous odd values contain the observations with the reference stations.

^a For this observation DE605 was not used as a reference.

^b Data was recorded for 10 spectral windows centered at; 108.0078125, 114.84375, 125.1953125, 136.9140625, 147.65625, 150.1953125, 154.1015625, 166.6015625, 178.90625 and 189.0625 MHz.

^c All core stations, except CS024.

^d Data was recorded for 10 spectral windows centered at; 36.71875, 45.1171875, 50.0, 51.5625, 53.125, 56.0546875, 61.5234375, 65.234375, 67.7734375 MHz.

^e All core stations.

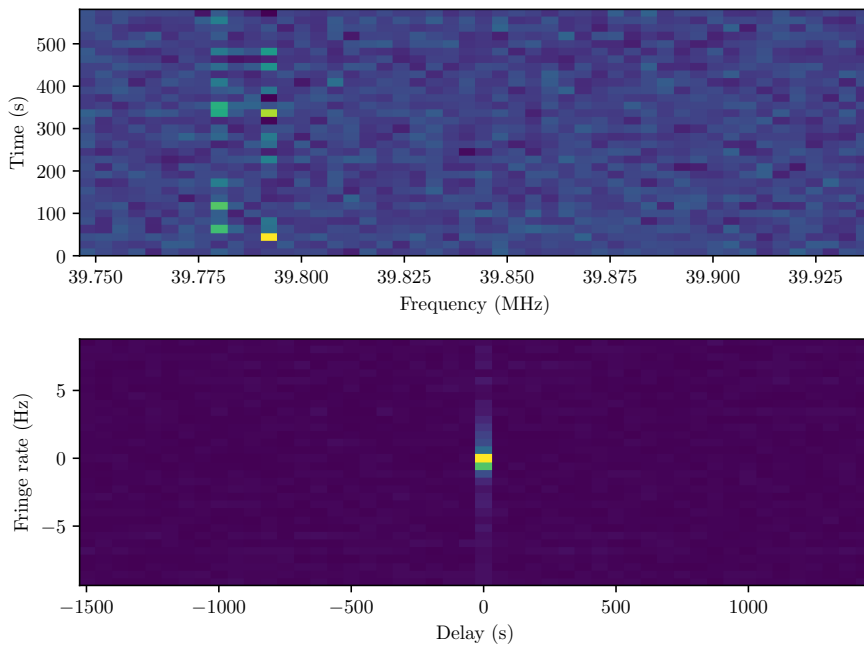


Figure 5.1.: Waterfall plots of the complex visibilities and their Fourier transform. The *top panel* shows the amplitude of the visibilities. RFI can be seen at 39.78 MHz. The *bottom panel* shows the Fourier transform of the visibilities. If the visibilities are dominated by a point source in the array tracking center, then this should be a delta function centered at 0 delay and 0 fringe rate. In this case there is some time variability in the data which produces a spread along the fringe rate.

5.3.2. FROM RAW VOLTAGES TO BEAM MAPS

To obtain a complex-valued map of the array beam we cross-correlate the voltage from the AUT with that of the reference station. This is done using an FX correlator (e.g., Thompson et al., 2017) implemented in *python*. To channelize the time series data from each spectral window we use a polyphase filter bank (PFB, e.g., Price, 2016) with a Hann window to alleviate spectral leakage and scalloping losses. We produce spectra of 64 3 kHz channels, and a time resolution of 327.68 μ s. This enables us to, at a later stage, flag narrowband radio frequency interference (RFI) without losing the time sample. In each spectral window we discard 25% of the channels at the edges, leaving a bandwidth of 146.48 kHz per spectral window. The two orthogonal polarizations are combined to produce four cross-correlation products, i.e. XX, XY, YX and YY products.

Before proceeding, we check that the bright unresolved source in the map center dominates the signal. In this case, a time delay versus fringe rate plot will show a peaked response in the center of the diagram. An example of such diagram is presented in Figure 5.1.

After cross-correlation, the visibilities are time averaged to ensure that their SNR is large enough (> 3) for calibration. For the HBA observations we average to a time resolution of 0.4 s, which results in a SNR of 6. For the LBA, which has a lower sensitivity and is more severely affected by the ionosphere (Figure 5.2), the averaging times are longer. For L645357 we average to 20 s and for L650445 — 5 s.

After time averaging, we remove visibilities affected by radio frequency interference (RFI) in the frequency-time domain. We use a SumThreshold method (AOFlogger, Offringa, van de Gronde, & Roerdink, 2012) on each TAB, polarization and spectral window independently. For the LBA and HBA observations the fraction of flagged data is $\approx 5\%$. After RFI flagging we average each spectral window in frequency to a single 146.48 kHz channel.

The flagged visibilities are then calibrated by multiplying by the inverse of the Jones matrix of the central beam (Brentjens et al. in prep.). This has the effect of removing most of the undesired systematic effects present in the data, such as the dependence of the observed visibilities on the sky brightness distribution, beam pattern of the reference station, or ionospheric delays between reference station and AUT.

After calibration, we further average the visibilities in time to one time sample with a duration of one minute for the HBA and ten minutes for the LBA (Table 3.1). After averaging in time, we are left with one calibrated complex visibility for each polarization (XX, XY, YX, and YY), spectral window and TAB. These calibrated complex visibilities map the complex-valued beam. Finally, we compute the inverse-variance weighted mean beam maps, averaged over all reference stations.

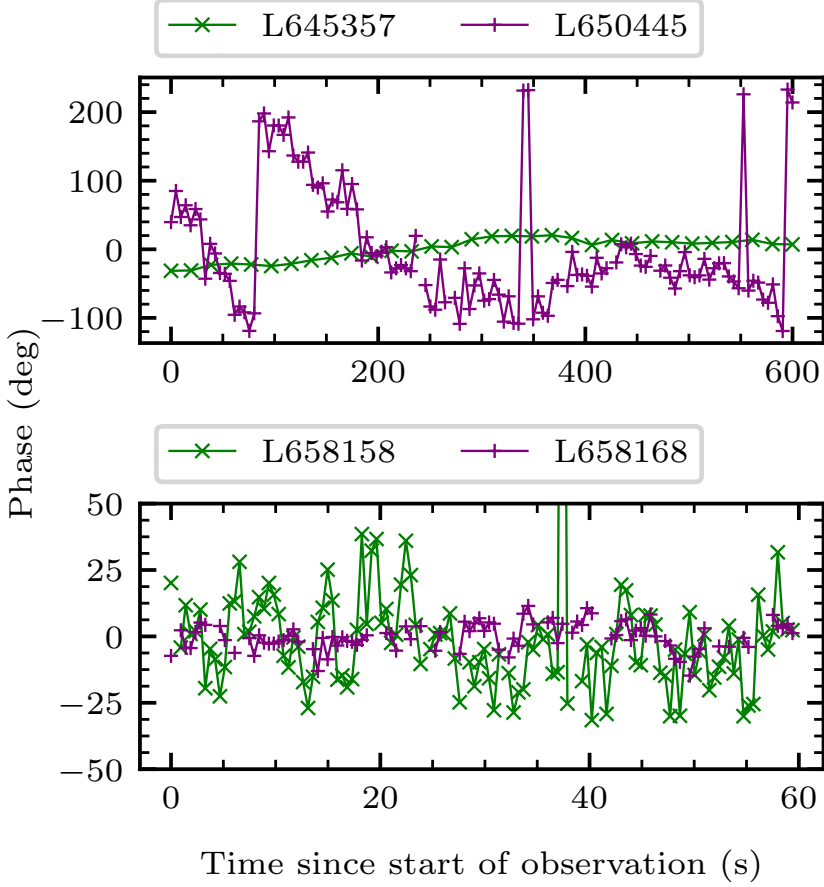


Figure 5.2.: Phase of the visibilities at the map center. The *top* panel shows the observed phase for the LBA observations. The *bottom* panel shows the observed phase for the HBA observations.

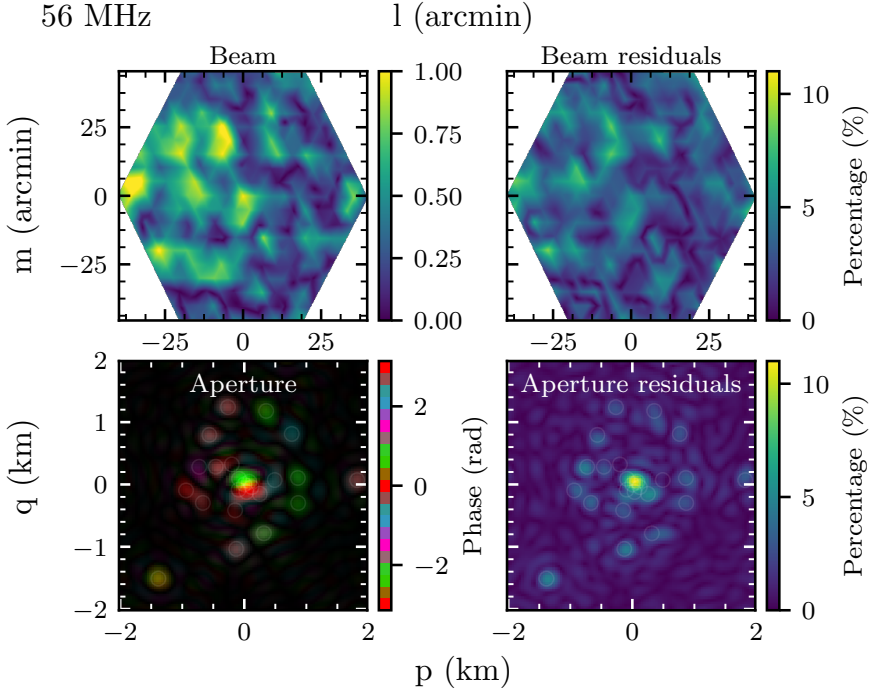


Figure 5.3.: Voltage beam and aperture maps for the LBA at 52 MHz derived from the observation L645357. *Top left:* Amplitude of the voltage beam. *Top right:* Amplitude of the residuals after subtracting the beam model from the observed voltage beam (*top left*). The beam model is generated using the derived phases for the stations in the AUT. *Bottom left:* Aperture map obtained by taking the Fourier transform of the voltage beam (inverse Fourier transform of Equation 5.2). The colorscale represents the phase and the brightness the amplitude at each location in the map. The *white circles* show the location of the stations in the AUT. *Bottom right:* Amplitude of the Fourier transform of the beam residuals (*top right*).

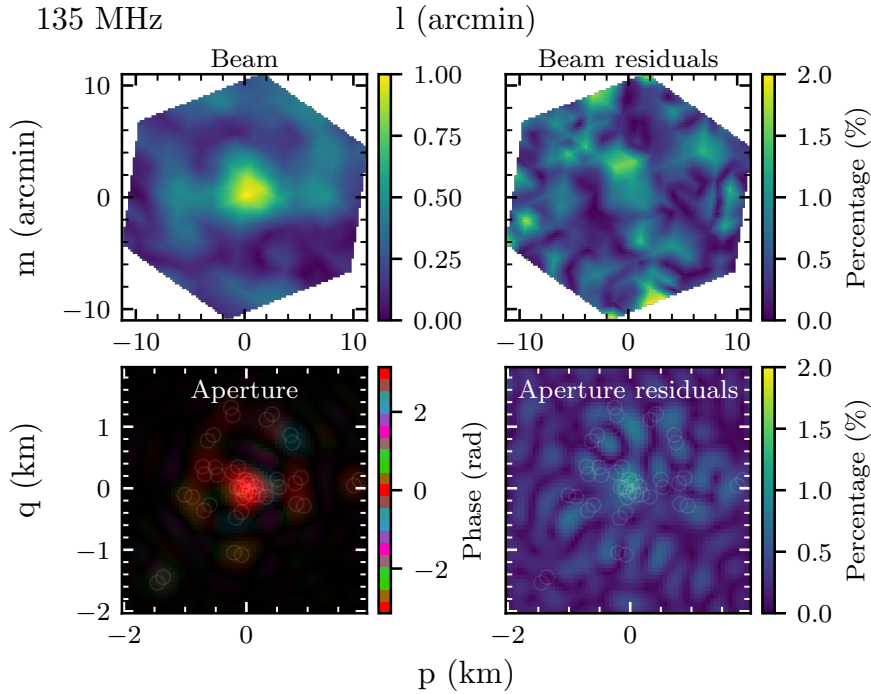


Figure 5.4.: Voltage beam and aperture maps for the HBA at 135 MHz derived from the observation L658168. *Top left:* Amplitude of the voltage beam. *Top right:* Amplitude of the residuals after subtracting the beam model from the observed voltage beam (*top left*). The beam model is generated using the derived phases for the stations in the AUT. *Bottom left:* Aperture map obtained by taking the Fourier transform of the voltage beam (inverse Fourier transform of Equation 5.2). The colorscale represents the phase and the brightness the amplitude at each location in the map. The *white circles* show the location of the stations in the AUT. *Bottom right:* Amplitude of the Fourier transform of the beam residuals (*top right*).

5.4. RESULTS

5.4.1. BEAM AND APERTURE MAPS

An example of the observed beam of the LBA is presented in the top left panel of Figure 5.3. There, the main lobe of the beam is at the map center, and we can also see that there is a side lobe with a similar amplitude at $(l, m) = (-6, 25)$. This is produced by improperly calibrated time delays between stations. For the HBA (top left panel of Figure 5.4) the side lobes have amplitudes $\approx 30\%$ of the main lobe. The time delays between HBA stations are regularly calibrated using synthesis imaging observations.

The voltage beam is the Fourier transform of the aperture illumination (Eq. 5.2), shown in the bottom left panel of Figures 5.3 and 5.4. There, we can see that the amplitudes are non-zero at the location of the stations in the AUT. The amplitudes are larger in the *Superterp* because there the stations are unresolved and their amplitudes, and phases, overlap.

5.4.2. TIME DELAYS AND 0 HZ PHASE OFFSETS

To measure the phase of the stations in the AUT we pose Equation 5.2 as a linear problem. Then, we use least squares parameter estimation to determine the amplitude and phase at the locations of the stations. The phases derived are not meaningful on their own, as an interferometer only measures relative phases (e.g., Jennison, 1958). To remove the arbitrary offset from the phases we reference them with respect to one of the stations in the AUT.

From the referenced phases we can recover the time delay, τ , and the 0-Hz phase offset, ϕ_0 , of each station. These are related to the phase, ϕ , by the linear relation $\phi = 2\pi\nu\tau + \phi_0$. An example of the observed phases and their best fit linear relation are presented in Figure 5.5. There we can see that the phases show a linear relation with frequency and that the errorbars on the phases become larger for stations closer to the array center.

Examples of the measured τ and ϕ_0 for the HBA stations derived from the L658168 observations are shown in Figure 5.6. For the *Superterp* stations, CS002 to CS007, the errorbars are a factor of three larger than for the rest of the stations. This is a consequence of the larger phase errors obtained for the *Superterp* stations (see Figure 5.5). This is also reflected in the larger aperture residuals at the *Superterp* (bottom right panel of Figure 5.4). For the stations outside the *Superterp*, the errors on τ have a mean value of 1.4 ± 1.2 ns and 1.2 ± 0.9 ns for the XX and YY polarizations, respectively. For the observation L658158 the same stations have errors on τ with a mean of 3.9 ± 1.7 ns and 3.7 ± 1.7 ns for the XX and YY polarizations, respectively. Since the flux density of 3C196 is a factor of 1.2 larger than that of 3C147, the larger errors on τ for L658158 are produced by the larger phase fluctuations in this observation (Figure 5.2).

The measured time delays and 0-Hz phase offsets for the LBA stations derived from L645357 and L650445 are presented in Figure 5.7. In both observations the derived values of τ and ϕ_0 agree to within 3σ , albeit in L645357 they have a factor of 12 higher SNR. This is because for the L645357 observations the ionosphere over the array

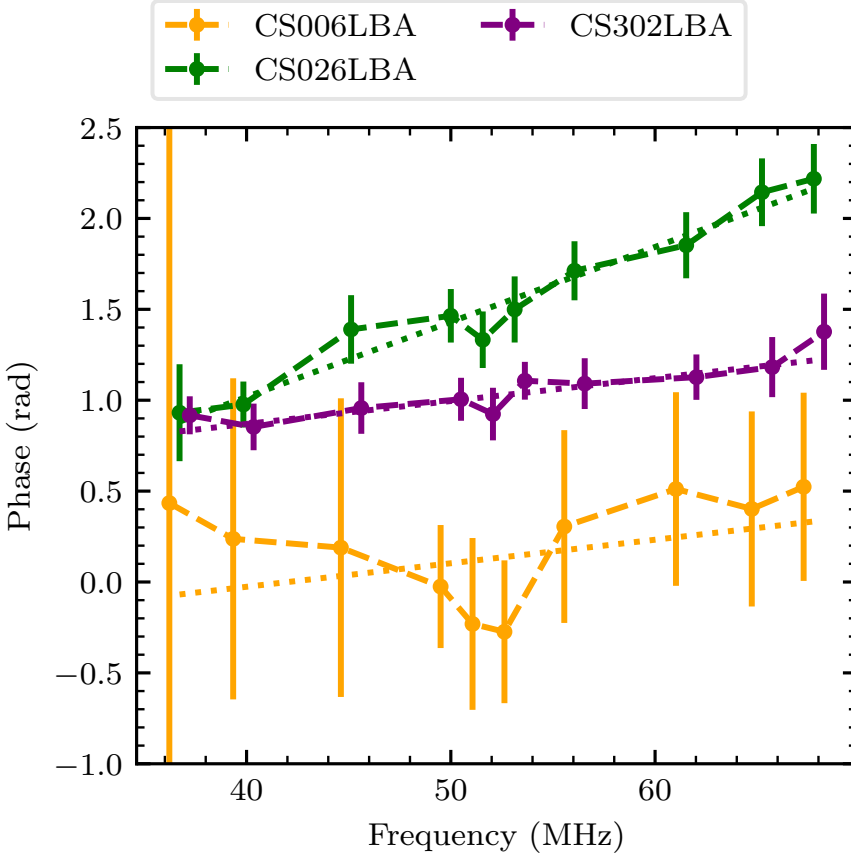


Figure 5.5.: Example of the measured phases as a function of frequency for three LBA stations. The measured phases are shown with *dashed lines* and *error bars*, while the best fit lines are shown with *dotted lines*. CS302LBA is at a distance of two km from the array center, CS026LBA at 870 m and CS006LBA 126 m, part of the *superterp*. These correspond to values derived from observation L645357.

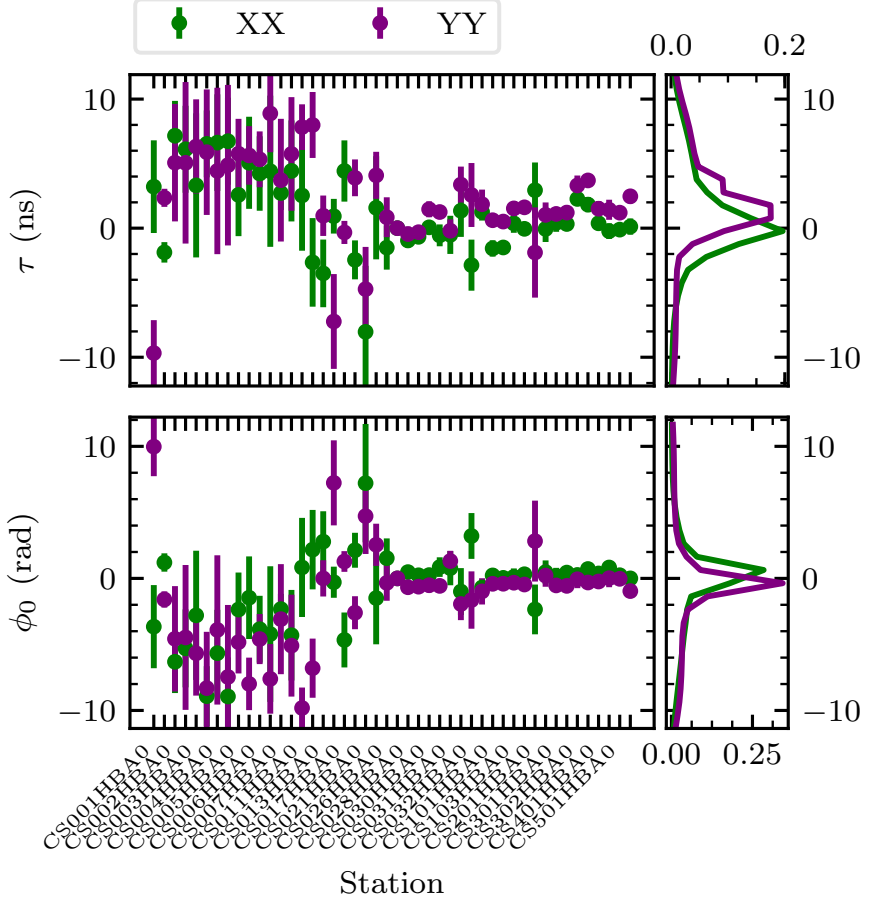


Figure 5.6.: Time delay and 0-Hz phase offsets for the HBA stations derived from observation L658168. The *top panel* shows the derived time delay τ for each station in the AUT. The time delay for CS026HBA1 is 0 because this station was used to reference the phases. The *bottom panel* shows the 0-Hz phase offset. It can be seen that the time delays and 0-Hz phase offsets are consistent between the two polarizations. For the innermost stations, CS002 to CS007, the scatter is larger because the stations are unresolved.

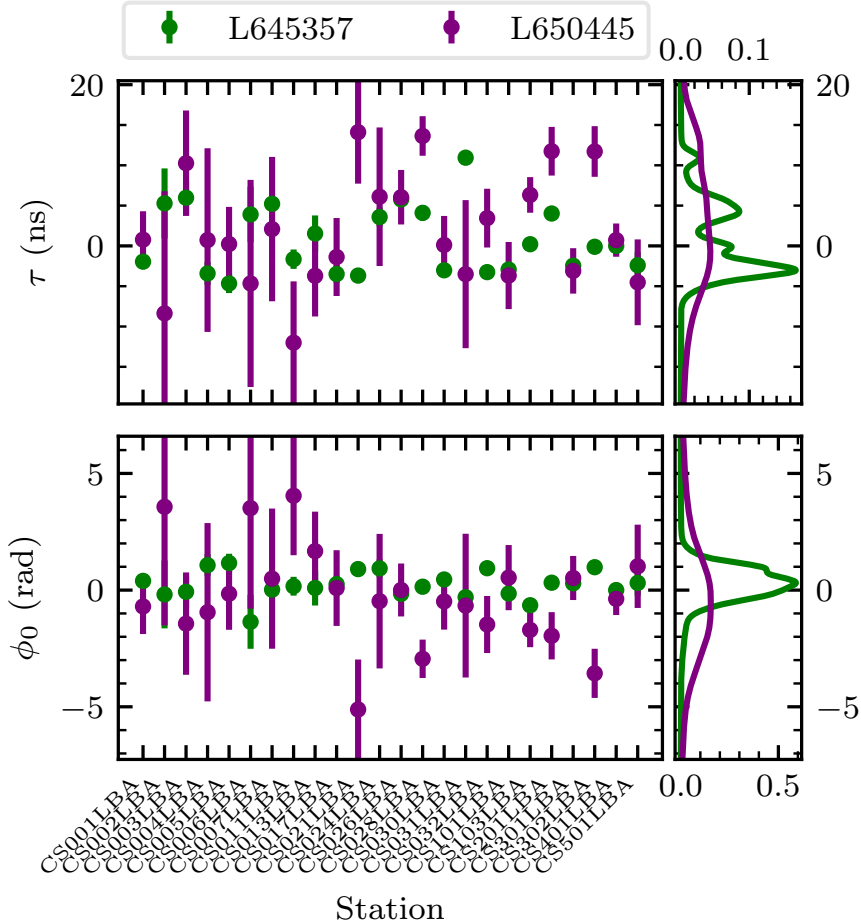


Figure 5.7.: Time delays and 0-Hz phase offsets for the LBA stations derived from holography observations. The time delays and 0-Hz phase offsets correspond to those of the XX polarization for two observations (L645357 and L650445). The *top panel* shows the derived time delay τ for each station in the AUT. The time delay for CS028 is 0 because this station was used to reference the phases. The *bottom panel* shows the 0-Hz phase offset. It can be seen that the time delays and 0-Hz phase offsets are consistent between the two observations. For the innermost stations, CS002 to CS007, the scatter is larger because the stations are unresolved.

produces a smooth slow time-varying phase rotation, while for L650445 the changes are faster and more pronounced (Figure 5.2). The time delays for the *Superterp* stations have errors a factor of four larger than for the rest of the stations.

For both HBA and LBA (Figures 5.6 and 5.7), the 0-Hz phase offsets are consistent with being zero at the 5σ level. Motivated by this, we fit a linear relation to the phases with $\phi_0 = 0$. The values of τ for the LBA stations under this assumption are presented in Figure 5.8. We can see that the derived time delays are consistent with those presented in Figure 5.7, but in this case the errorbars are smaller because there is one free parameter less and setting $\phi_0 = 0$ is a strong constraint. Using $\phi_0 = 0$ the mean value of the error of the derived time delays is 0.26 ± 0.16 ns and 0.17 ± 0.10 ns for the HBA and LBA, respectively.

In Figure 5.8 we also show the time delays for the LBA stations derived from imaging observations. In these, phase solutions for each station are derived from observations of a bright calibrator source and a model of the sky brightness distribution (e.g., Fomalont & Perley, 1999). Then, the contribution to the phase from the station delays and the ionosphere are separated (e.g., van Weeren et al., 2016; de Gasperin et al., 2018). We see that the time delays derived using holography and imaging observations agree to within 3σ for 20 out of the 23 stations present in both observations. This shows that the time delays derived here, where no model of the sky brightness distribution is used, are indistinguishable from those derived in imaging observations. The interferometric time delays have smaller errorbars because they are the average of 4 hours of data.

We check that the derived time delays and 0-Hz phase offsets capture the status of the AUT by using them to simulate the array beam and comparing it with the observed one. To simulate the array beam we use the time delays and 0-Hz phase offsets shown in Figures 5.6 and 5.7 for the HBA and LBA stations, respectively. These are used to evaluate the phase of each station at a particular frequency, producing a complex valued map of the AUT. The Fourier transform of the AUT simulates the array beam. The residuals between the observed and simulated array beams are presented in the top right panel of Figures 5.3 and 5.4. The residuals in the image plane have no obvious structure and show amplitudes of $\lesssim 10\%$. This shows that we can reproduce the array beam using the derived time delays and 0-Hz phase offsets. However, the Fourier transform of the beam residuals reveals that there is significant structure in the aperture plane (bottom right panel of Figures 5.3 and 5.4). This can be noticed as a larger amplitude (11% for the LBA) at the location of the *Superterp*, for which we are not capturing the phase behavior as accurately as for the stations away from it (see also Figure 5.5).

5.5. DISCUSSION

5.5.1. CORRECTED TIME DELAYS

We use the derived time delays (Figure 5.8) to update the instrumental time delays in LOFAR's beamformer. To test the effect of updating the instrumental time delays in the beamformer we observed Cygnus A with the core stations of LOFAR. The observation was one minute long using the imaging mode, where the signals of different stations

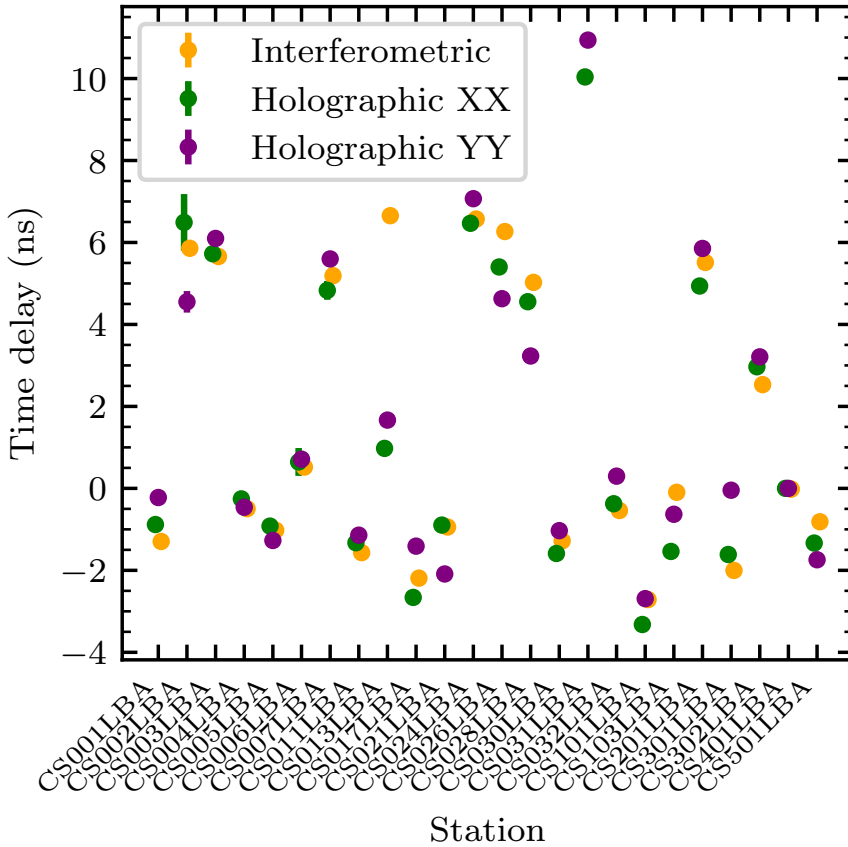


Figure 5.8.: Time delays for the LBA stations assuming that $\phi_0 = 0$. The time delays obtained from holography observations are shown for two orthogonal polarizations (XX and YY). We also show the time delays derived from imaging observations for Stokes I. The time delay for CS401LBA is 0 because this station was used to reference the phases. Station CS031LBA is missing in the imaging data because it was completely flagged.

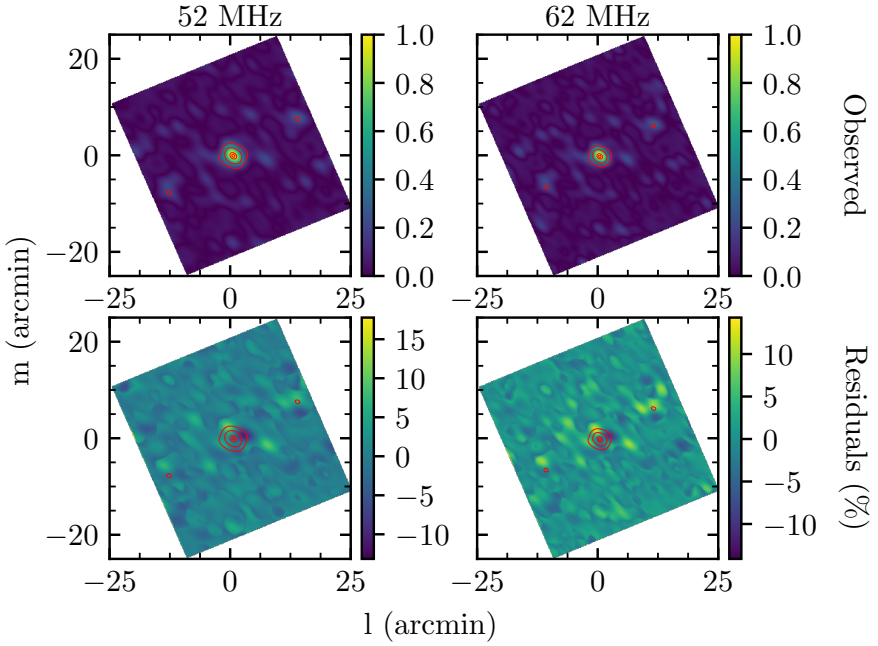


Figure 5.9.: Observed beam for the LBA. The *top row* shows the observed beam after updating the instrumental time delays in the beamformer (*colorscale*). The *red contours* show the amplitude of a beam model generated from the position of the stations with no additional time delays. They are at the 20%, 50%, 90% and 98% level with respect to the peak response. The *bottom row* shows the residuals after subtracting the beam model from the observed beam (*colorscale*). The *red contours* are the same as in the top row.

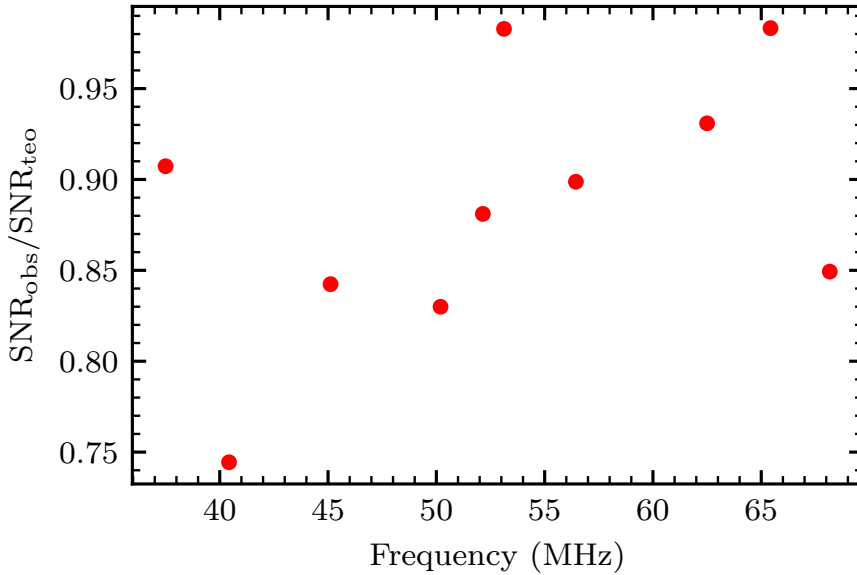


Figure 5.10.: Ratio between the observed SNR and the theoretical value.

are cross-correlated instead of added. Since the signal path between stations and the beamformer is the same in tied-array and imaging modes, these observations have a beam equivalent to the one observed using holography. Cygnus A has a size of $\approx 1'$ (e.g., McKean et al., 2016), so it will be unresolved by the LOFAR core at LBA frequencies (at 90 MHz the spatial resolution of the core is $3'.3$). Hence, a dirty image obtained from this observation will show the array beam.

A comparison between the LBA beam after the update and a model of the beam is presented in Figure 5.9. In the beam model the phase of each station is given by its location with no additional time delays. We see that the sidelobes in the observed beam are similar to those of the beam model. After subtracting the beam model the residuals are $\lesssim 20\%$.

We compare the SNR of the observations of Cygnus A with the theoretical SNR in Figure 5.10. The ratio between the observed SNR and its theoretical value has a mean of 0.88 ± 0.06 . These results show that after updating the time delays in the beamformer the array beam is close to the ideal case. The remaining differences are produced by propagation delays introduced by the ionosphere and any remaining errors in the instrumental time delays.

When using an unpolarized bright unresolved source to determine the time delays, there will be an arbitrary offset between the X and Y polarizations. When these time delays are applied to the beamformer, the offset will produce a rotation of Stokes U into Stokes V . In order to find this offset we need to observe a linearly polarized source for which the sign of the rotation measure is known. From this observation the offset

is determined from the angle between the apparent Stokes U and V (e.g., Brentjens, 2008). This step has not been performed yet.

5.5.2. IMPROVEMENTS TO HOLOGRAPHIC MEASUREMENTS WITH LOFAR

One of the main limitations of the holographic measurements presented is the spatial resolution over the telescope aperture. For a constant number of TABs, this can be improved by observing a larger portion of the beam using a mosaic while keeping a TAB at the calibrator source. Additionally, the separation between TABs could be made smaller, reducing aliasing artifacts in the aperture plane.

For the experiments presented in this work we have used only four reference stations. Outside its core, LOFAR has 14 stations within the Netherlands and 13 stations distributed all over Europe. Any of these stations can be used as reference station, as long as the baseline formed with the AUT does not resolve the source used to map the beam. This means that the improvement in the SNR of the complex-valued beam map can be up to a factor of four using the same sources. With an improvement a factor of four in the maps of the complex-valued beam the integration times could be made shorter, or more precise time delays could be derived.

The time delays derived from holographic measurements can be used to update the instrumental time delays prior to an observation with the tied-array (e.g., of a pulsar). Moreover, if the integration time required to reach ns precision can be made shorter, and the post-processing of the holographic measurements made in real time, then it would be possible to interleave holographic observations during the tied-array observations. This could be a way towards semi-real-time beam calibration for the Square Kilometer Array (SKA). Implementing a dedicated holography mode in the supercomputer that processes the raw LOFAR data is one of the next steps towards this goal.

5.6. SUMMARY

In this work we have used radio holography along with a new calibration method to characterize the time delays between LOFAR's core stations. This new calibration method consists on calibrating the measured complex-valued beam map by right multiplying by the matrix inverse of the map center. This calibration makes the observed complex-valued reception pattern independent of the sky brightness distribution.

Four (HBA) and three (LBA) reference stations were used simultaneously to produce maps of the tied array voltage beam. Using 60 s (HBA) and 600 s (LBA) long observations of 3C196, 3C147 and 3C48 we derived time delays with an error < 1 ns. We find that the main limitations in reaching nanosecond precision in the measured time delays are the condition of the ionosphere over the array and the ability to spatially resolve the array elements. LOFAR now uses the derived time delays operationally.

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