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# Methanol ice formation threshold in dense clouds and dark cores revisited

CH<sub>3</sub>OH is a central molecule for interstellar complex organic molecule formation. Theoretical, experimental, and observational studies have shown that it is formed abundantly by the hydrogenation of CO-rich ice, deep inside cold, dark cores (formation threshold of  $A_V > 9$  mag). This scenario does not seem uniformly applicable, however, because observations have shown a deficiency of the CH<sub>3</sub>OH ice abundance, by at least a factor of 3, in a significant number of dark core and dense cloud sight-lines tracing depths well above 9 mag. Also, at shallower depths, CH<sub>3</sub>OH may still be formed, but at lower abundances, via, e.g., CH<sub>4</sub> abstraction reactions (Qasim et al. 2018).

In this study, we aim to help address these questions by re-analyzing previously published data (Boogert et al. 2011), as well as new L-band spectra of a sample of dense cloud and dark core sight-lines toward background stars. Focus was put into increasing the sensitivity to CH<sub>3</sub>OH abundances, enabling an improved study of abundance variations as a function of cloud/core depth. The detection and absorption profile analysis of weak ice features in dense clouds and dark cores, such as those of CH<sub>3</sub>OH, are limited by the contamination with photospheric absorption lines of the background stars. We demonstrate a method to reduce photospheric contamination by dividing over the spectra of un-reddened stars from the NASA Infrared Telescope Facility database.

In a sample of 41 stars behind quiescent dense cloud material and isolated dense cores we report 1 new CH<sub>3</sub>OH ice detection, bringing the total to 8. The CH<sub>3</sub>OH ice abundances are 3.4-21% relative to H<sub>2</sub>O, corresponding to  $1.6 \times 10^{-6} - 8.8 \times 10^{-6}$  relative to H<sub>2</sub>. A linear fit to these detections yields a CH<sub>3</sub>OH ice detection threshold of only  $13.5 \pm 7.8$  mag. This is within the uncertainty range from Boogert et al. (2011).

Significant non-detections are reported in 34 sight-lines. After correction for the photospheric lines, and at sufficiently high S/N, the CH<sub>3</sub>OH ice abundances are limited by the contamination with the 3.47  $\mu$ m ice feature (likely due to NH<sub>3</sub>.H<sub>2</sub>O hydrates). Still, the CH<sub>3</sub>OH ice abundance upper limits (rel-

ative to  $A_V$  and  $H_2O$ ) are low across the entire observed detection threshold  $A_V$  range of 5.1–46.0 mag:  $<0.34\%$  relative to  $H_2O$  at  $A_V = 5.1\text{--}16.8$  mag, and  $<0.21\%$  relative to  $H_2O$  at  $A_V = 19.6\text{--}46.0$  mag.

These results put strict observational constraints on the formation efficiency of  $CH_3OH$  ice. A full understanding requires the measurement of the CO ice abundance, as well as ice maps to trace the CO freeze out and  $CH_3OH$  formation rate in individual clouds and cores. This will be possible with the upcoming James Webb Space Telescope (JWST). Our method of the photospheric line correction of background star spectra can be applied to the JWST observations.

## 5.1 Introduction

Grain surface ice chemistry is initiated in the translucent phase of interstellar clouds (density of  $\sim 10^3$   $H_2$  molecules  $cm^{-3}$ ), and is further accelerated at densities of  $\sim 10^4 - 10^5$   $cm^{-3}$  in dense clouds and dark cores (Pontoppidan 2006; Boogert et al. 2015). At a cloud depth corresponding to a visual extinction ( $A_V$ ) of  $\sim 1.6$  mag, a  $H_2O$ -rich ice is present, in which the ice matrix is composed of simple species such as  $NH_3$ ,  $CO_2$ , and  $CH_4$  (Boogert et al. 2015). At  $\sim 3$  mag, some CO ice grows to form an apolar layer (Chiar et al. 1995). At greater depths of  $A_V \sim 9$  mag ( $10^5$   $cm^{-3}$ ), CO freezes out nearly completely (Jørgensen et al. 2005; Pontoppidan 2006; Boogert et al. 2015), and is hydrogenated to form methanol ( $CH_3OH$ ), as confirmed by a number of laboratory, modeling, and observational studies (Watanabe & Kouchi 2002; Fuchs et al. 2009; Cuppen et al. 2009; Boogert et al. 2011; Wirström et al. 2011). A list of the physical parameters of translucent, dense clouds, and dark cores is provided in Table 5.1.

Table 5.1: Physical conditions of the different regions of interstellar clouds, as noted in Table II of van Dishoeck et al. (1993). (a) Dust temperatures are much lower for all noted regions.

ISM region	Density ( $cm^{-3}$ )	Gas temperature <sup>a</sup> (K)	$A_V$ (mag)
Translucent	500-5000	15-50	1-5
Dense cloud	$10^2 - 10^4$	$\gtrsim 10$	$\gtrsim 2$
Cold, dark core	$10^4 - 10^5$	$\approx 10$	5-25

Hydrogen addition and abstractions reactions in the pathway to  $CH_3OH$  also lead to an assortment of complex organic molecules (COMs), such as glycolaldehyde, ethylene glycol, and methyl formate (Chuang et al. 2016), in addition to glycerol (Fedoseev et al. 2017), propanal and 1-propanol (Qasim et al. 2019b). During the dispersal of the envelopes of the Young Stellar Objects (YSOs), the temperature rises, which increases the mobility of ice species and molecular complexity (e.g., Herbst & van Dishoeck 2009). At reduced dust extinctions, enhanced UV-irradiation of  $CH_3OH$  ice leads to various COMs as well (Öberg et al. 2009). COMs formed in the ice are released into the gas-phase, however COMs may also be formed in the gas-phase starting from simple molecules that were liberated from the ice (Charnley et al. 1992). Low abundances of

gas phase COMs are also detected in cold ( $T = 10$  K) molecular clouds, where  $\text{CH}_3\text{OH}$  ice is a likely precursor (Soma et al. 2018).

$\text{CH}_3\text{OH}$  is a key interstellar molecule to complex chemistry, so understanding its prevalence and formation conditions are important. Observations have not provided a consistent formation model. At similar  $A_V$ , the  $\text{CH}_3\text{OH}$  ice abundance relative to  $\text{H}_2\text{O}$  varies between detections of 10% and  $3\sigma$  upper limits of 3% (Boogert et al. 2011, 2015). Perhaps the CO freeze out is limited in certain clouds and cores. For background star observations of dense clouds, a sight-line with a high  $A_V$  may have a long pathlength at low density. Lower densities result in higher dust temperatures, which would then decrease the freeze-out and hydrogenation rates of CO (Cuppen et al. 2009). However, this lack of  $\text{CH}_3\text{OH}$  at high  $A_V$  has been observed in numerous dark cores (Boogert et al. 2011), and dark cores have high  $A_V$  and high densities. Alternatively, CO may be consumed into other reactions, as also proposed in Shimonishi et al. (2016). The CO hydrogenation chain includes many other products besides  $\text{CH}_3\text{OH}$ , as illustrated in Figure 1 of Fedoseev et al. (2017). The low  $\text{CH}_3\text{OH}$  abundance deep in some clouds and cores could also be due to destruction either by internal UV-photons (Öberg et al. 2009) and/or H-abstraction reactions (Goumans & Kästner 2011; Chuang et al. 2016). On the other hand,  $\text{CH}_3\text{OH}$  formation should not be limited to the formation threshold of  $\sim 9$  mag (Boogert et al. 2015), as other formation routes should be relevant at shallower cloud depths. For example, in the laboratory investigation by Qasim et al. (2018), it was found that the  $\text{CH}_4 + \text{OH}$  formation route to  $\text{CH}_3\text{OH}$  should proceed on grain mantles, although at a factor of 20 lower efficiency compared to H addition in CO ices.  $\text{CH}_4$  is formed at low extinctions, and there is tentative observational evidence that it is enhanced at cloud edges where not all C has been included in CO yet (Öberg et al. 2008).

Observational constraints on  $\text{CH}_3\text{OH}$  formation models are only based on a handful of detections in dark cores (with and without embedded YSOs) (Boogert et al. 2011; Chiar et al. 2011; Pontoppidan et al. 2003). This limited sample size is in part due to the low sensitivity that is available from current observational facilities, as well as the lack of ice mapping capabilities. The James Webb Space Telescope (JWST), which is expected to be launched in 2021, will significantly improve on both aspects.

Ices in dense clouds and dark cores are typically observed in spectra of background field stars, which are typically cool stars (K- and M-giants). The detection of the ice bands, as well as their absorption profile analysis, is often limited by photospheric line contamination. Here, we present a method to increase the sensitivity of previously published and new observations of the C-H symmetric stretching mode of  $\text{CH}_3\text{OH}$  ice at  $3.537 \mu\text{m}$  by dividing out photospheric lines using spectra of un-reddened stars from the NASA Infrared Telescope Facility (IRTF) database (Rayner et al. 2009). Section 5.2 overviews the parameters of the dust/ice observations and data reduction. All observations were taken towards *background stars*. Section 5.3 details the steps involved in reducing the photospheric lines and measuring the  $\text{CH}_3\text{OH}$  column densities. Sections 5.4 and 5.5 show the application of the method to increase the sensitivity of the data and accuracy of the data analysis in order to better constrain the formation parameters of  $\text{CH}_3\text{OH}$ . Section 5.6 outlines the main findings and proposed work for future studies.

## 5.2 Observations and data reduction

41 different background stars were observed in the atmospheric L–band with the NASA IRTF/Spex (Rayner et al. 2003) and Keck 2/NIRSPEC (McLean et al. 1998) spectrometers, tracing the quiescent material in 12 different dense clouds and isolated dense cores (Table 5.2). Some cores contain embedded YSOs, and it is also not certain whether the dense cloud sight-lines trace high densities ( $\sim 10^4$ ), or low densities in long pathlengths. Note that some targets were observed with both Spex and NIRSPEC. The targets were selected to trace molecular material at a wide range of detection threshold extinctions  $A_V = 5.1\text{--}48.0$  mag, as determined from model fits to 1–5  $\mu\text{m}$  broad-band photometry and spectroscopy (Boogert et al. 2011, and in preparation). The spectra of 25 targets were previously published in Boogert et al. (2011).

Table 5.2 also lists the observational parameters, including the spectrometer used, wavelength range covered, the standard star used to reduce telluric contamination, the date of the observation, and the signal-to-noise (S/N) of the continuum emission near the wavelength of the  $\text{CH}_3\text{OH}$  ice band at 3.537  $\mu\text{m}$  (3.460 – 3.530  $\mu\text{m}$ ; a region with relatively few telluric lines).

The Spex instrument was used in the LongXD(1.9) mode, with a slit width of 0.8 arcseconds, resulting in a resolving power  $R = \lambda/\Delta\lambda = 938$ . Data reduction was performed using the software package SpexTool (Cushing et al. 2004) and Xtellcor (Vacca et al. 2003). Keck/NIRSPEC was used in the long-slit mode with a slit width of 0.57 arcseconds, resulting in a resolving power of  $\sim 1600$ . The data reduction procedure for these observations is discussed in Boogert et al. (2011). The wavelength calibration for both Spex and NIRSPEC is based on telluric emission lines. The radial velocity of the targets was not corrected for.

## 5.3 Photospheric line correction and column determination

A Python script was created to reduce contamination by photospheric absorption lines from the interstellar dust and ice spectra by fitting un-reddened template spectra from the IRTF spectral library (Rayner et al. 2009). The best-fitting templates for each source,  $\chi_\nu^2$  (reduced- $\chi^2$ ), and the  $\text{CH}_3\text{OH}$  column densities derived from the photosphere-corrected spectra are presented in Table 5.3. The following steps were used in this procedure:

- 1) The spectral resolution of the template spectra ( $R = 2500$  in the L–band; slit width of 0.3 arcseconds) was matched to those of the observed spectra by applying boxcar smoothing. Typically, a boxcar value of 3 was found to be sufficient for the lower resolution Spex spectra, and no smoothing was needed for the NIRSPEC spectra.

- 2) A cross-correlation was applied to detect differences of the central wavelengths of the photospheric lines between the template and background star spectra. A wavelength shift was subsequently applied to the science spectrum, correcting for small wavelength calibration errors as well as differences in the radial velocity. In the wavelength range relevant for this work ( $\sim 3.34\text{--}3.71$   $\mu\text{m}$ ), shifts of  $< 70$  km/s were typically applied, but sometimes as large as 170 km/s.

Table 5.2: Observational parameters. Extinction  $A_V$  calculated assuming  $A_V/A_K = 8$  (Cardelli et al. 1989). B2011: published in Boogert et al. (2011).

Source 2MASS J	Cloud/Core	Instrument	$\lambda$ $\mu\text{m}$	Standard star (SPT)	Date	Continuum S/N yyyy/mm/dd	Notes 3.537 $\mu\text{m}$
04215402+1530299	IRAM 04191	NIRSPEC	2.82-4.14	HR 1251 (A0.5V)	2006/10/09	45	B2011
17111501-2726180	B 59	NIRSPEC	2.41-4.14	HR 6141 (B3V)	2005/06/28	54	B2011
17111538-2727144	B 59	NIRSPEC	2.41-4.14	HR 6141 (B3V)	2005/06/28	55	B2011
17112005-2727131	B 59	NIRSPEC	2.38-4.14	HR 5993 (B1V)	2006/07/07	52	B2011
17155573-2055312	L 100	NIRSPEC	2.38-4.14	HR 5993 (B1V)	2006/07/06	54	B2011
17160467-2057072	L 100	NIRSPEC	2.39-4.14	HR 6946 (B2II)	2005/06/29	41	B2011
17160860-2058142	L 100	NIRSPEC	2.38-4.14	HR 5993 (B1V)	2006/07/06	76	B2011
18140712-0708413	L 438	NIRSPEC	2.82-4.14	HR 7141 (A5V)	2006/10/09	79	B2011
18160600-0225539	CB 130-3	NIRSPEC	2.82-4.14	HR 7141 (A5V)	2006/10/09	65	B2011
18165296-1801287	L 328	NIRSPEC	2.41-4.14	HR 6946 (B2II)	2005/06/28	64	B2011
18165917-1801158	L 328	NIRSPEC	2.41-4.14	HR 6946 (B2II)	2005/06/28	50	B2011
18170376-0815070	L 429-C	NIRSPEC	3.00-3.87	HR 6963 (B9V)	2018/06/01	43	
18170426-1802408	L 328	NIRSPEC	2.06-4.14	HR 6946 (B2II)	2005/06/29	52	B2011
18170429-1802540	L 328	NIRSPEC	2.39-4.09	HR 6946 (B2II)	2005/06/29	56	B2011
18170470-0814495	L 429-C	NIRSPEC	2.38-4.14	HR 7141 (A5V)	2006/07/07	62	B2011
18170957-0814136	L 429-C	NIRSPEC	2.38-4.14	HR 7141 (A5V)	2006/07/07	63	B2011
18171181-0814012	L 429-C	NIRSPEC	2.06-4.14	HR 7141 (A5V)	2006/10/09	43	B2011
18171366-0813188	L 429-C	NIRSPEC	2.06-4.14	HR 7141 (A5V)	2006/10/09	39	B2011
18171700-0813504	L 429-C	NIRSPEC	3.00-3.87	HR 6963 (B9V)	2018/06/01	54	
18172690-0438406	L 483	NIRSPEC	2.38-4.14	HR 7141 (A5V)	2006/10/09	59	B2011
18275901+0002337	Serpens MC	IRTF	1.92-4.19	HR 7141 (A5V)	2010/05/22	50	
		NIRSPEC	2.77-4.23	HR 7141 (A5V)	2009/10/12	62	
18282010+0029141	Serpens MC	IRTF	1.92-4.19	HR 7141 (A5V)	2010/05/24	28	
		NIRSPEC	2.77-4.23	HR 6744 (A5E)	2009/10/12	37	
18282631+0052133	Serpens MC	IRTF	1.92-4.19	HD 163336 (A0V)	2010/05/20	89	
		NIRSPEC	2.77-4.23	HR 7141 (A5V)	2009/10/12	116	
18284038+0044503	Serpens MC	IRTF	1.92-4.19	HD 161868 (A1V)	2010/05/21	34	
		NIRSPEC	2.77-4.23	HR 7141 (A5V)	2009/10/12	49	
18284797+0037431	Serpens MC	IRTF	1.92-4.19	HR 7141 (A5V)	2010/05/24	32	
		NIRSPEC	2.77-4.23	HR 7141 (A5V)	2009/10/12	50	
18285266+0028242	Serpens MC	NIRSPEC	2.75-4.21	HR 7141 (A5V)	2007/07/05	31	
18290316+0023090	Serpens MC	IRTF	1.92-4.19	HD 161868 (A1V)	2010/05/2	34	
		NIRSPEC	2.77-4.23	HR 7141 (A5V)	2009/10/12	63	
18290436+0116207	Serpens MC	IRTF	1.92-4.19	HR 7141 (A5V)	2010/05/20	41	
		NIRSPEC	2.77-4.23	HR 7141 (A5V)	2009/10/12	71	
18291619+0045143	Serpens MC	NIRSPEC	2.77-4.23	HR 7141 (A5V)	2009/10/12	55	
18291699+0037191	Serpens MC	IRTF	1.92-4.19	HR 6629 (A1V)	2010/05/23	19	
		NIRSPEC	2.77-4.23	HR 7141 (A5V)	2009/10/12	41	
18292528+0003141	Serpens MC	IRTF	1.92-4.19	HR 7141 (A5V)	2010/05/20	53	
		NIRSPEC	2.77-4.23	HR 7141 (A5V)	2009/10/12	62	
18294108+0127449	Serpens MC	IRTF	1.92-4.19	HR 7141 (A5V)	2010/05/23	26	
		NIRSPEC	2.77-4.22	HR 7141 (A5V)	2009/10/12	72	
18300061+0115201	Serpens MC	IRTF	1.92-4.19	HR 6744 (A0V)	2010/05/22	41	
		NIRSPEC	2.77-4.23	HR 7141 (A5V)	2011/04/18	47	
		NIRSPEC	2.82-4.14	HR 7141 (A5V)	2009/10/12	47	B2011
18300085+0017069	Serpens MC	IRTF	1.92-4.19	HR 7141 (A5V)	2010/05/23	110	
		NIRSPEC	2.77-4.23	HR 7141 (A5V)	2009/10/12	51	
18300896+0114441	Serpens MC	IRTF	1.92-4.19	HD 163336 (A0V)	2010/05/22	38	
		NIRSPEC	2.77-4.22	HR 7141 (A5V)	2009/10/12	64	
19201597+1135146	CB 188	NIRSPEC	2.41-4.14	HR 7724 (A1V)	2005/09/20	54	B2011
19201622+1136292	CB 188	NIRSPEC	2.06-4.14	HR 7724 (A1V)	2005/09/20	56	B2011
19214480+1121203	L 673-7	NIRSPEC	2.06-4.14	HR 7141 (A5V)	2006/07/07	62	B2011
21240517+4959100	L 1014	NIRSPEC	1.49-4.14	HR 7929 (B2V)	2005/06/28	65	B2011
21240614+4958310	L 1014	NIRSPEC	1.49-4.14	HR 7929 (B2V)	2005/06/28	72	B2011
22063773+5904520	L 1165	NIRSPEC	1.49-4.14	HR 8585 (A1V)	2006/10/09	117	B2011

3) After interpolating the modified template spectrum to the source spectrum, the source spectrum was divided by the template to produce a "residual" spectrum.

4) The "residual" spectrum was divided by a smooth baseline, constructed by applying a Savitzky-Golay filter. Subsequently, the standard deviation was measured in the range of 3.6 - 4.0  $\mu\text{m}$  (i.e., the range at which photospheric lines are strongly present, and telluric and ice features are relatively weak).

The best-fitting template was selected to be the one resulting in the lowest standard deviation in the "residual" spectrum divided by a smooth baseline. Also, a range of best-fitting spectral types was determined by selecting those with  $\chi^2_{\nu}$  values within a factor of 2 of the minimum, although a precise spectral classification of the background stars is not the primary goal of this work.

After the division of the observed background star spectrum over the best fitting template spectrum, the ice features were analyzed. The well known 3.47  $\mu\text{m}$  ice absorption feature, which overlaps with the 3.53  $\mu\text{m}$   $\text{CH}_3\text{OH}$  ice feature, is readily visible in many spectra. The 3.47  $\mu\text{m}$  absorption feature is usually attributed to the  $\nu(\text{O}-\text{H}\cdots\text{N})$  mode of ammonia hydrates,  $\text{NH}_3\cdot\text{H}_2\text{O}$  (Dartois & d'Hendecourt 2001), although this is still open to interpretation (Shimonishi et al. 2016). A second or third order polynomial was fitted to the continuum surrounding this feature, in the wavelength ranges 3.34-3.36 and 3.60-3.70  $\mu\text{m}$  (see Appendix 5.7). The photospheric line-corrected background star spectrum  $F$  was converted into optical depth ( $\tau$ ) scale using this baseline  $F_0$  as follows:

$$\tau = -\ln(F/F_0) \quad (13)$$

Subsequently, two gaussians were fitted simultaneously to the overlapping 3.537  $\mu\text{m}$  and 3.47  $\mu\text{m}$  features. This decomposition method is similar to what was used for YSO spectra in Brooke et al. (1999). Instead of using a laboratory spectrum of  $\text{CH}_3\text{OH}$ , we applied a gaussian with a similar peak position and width (FWHM) as the main peak of the  $\text{CH}_3\text{OH}$  laboratory spectrum: 3.537 and 0.04  $\mu\text{m}$ , respectively. For the second gaussian, we used values of 3.47 and 0.1  $\mu\text{m}$ . The fitting procedure was performed in Python, using the GaussianModel class within the package, LmFit. An example of the decomposition of the two peaks is shown in Figure 5.1.

The  $\text{CH}_3\text{OH}$  column,  $N(\text{CH}_3\text{OH})$ , was then derived following

$$N(\text{CH}_3\text{OH}) = \frac{\tau_0 \times \text{FWHM}}{A} \quad (14)$$

where the full width at half maximum (FWHM) was fixed at 32  $\text{cm}^{-1}$  (0.04  $\mu\text{m}$ ) and the peak optical depth ( $\tau_0$ ) was derived with the Gaussian fits. The integrated band strength ( $A$ ) of  $5.6 \times 10^{-18}$   $\text{cm molecule}^{-1}$  for the C-H symmetric stretching mode of solid  $\text{CH}_3\text{OH}$  was obtained from Kerkhof et al. (1999).

## 5.4 Results

The  $\text{CH}_3\text{OH}$  ice column densities and  $3\sigma$  upper limits are presented in Table 5.3, along with the best-fitting templates for each source spectrum and



Table 5.3: CH<sub>3</sub>OH column densities in cm<sup>-2</sup> and relative to H<sub>2</sub>O ice for all observed dense cloud and dark core lines of sight after correction for the photospheric absorption lines. The upper limits have 3 $\sigma$  significance and are indicated by (<). The uncertainties of the detections are 1 $\sigma$  and are indicated by parentheses. The uncertainty range for the spectral types is based on the template stars that have  $\chi_v^2$  values within a factor 2 of the minimum.  $A_V$  and the H<sub>2</sub>O column densities are obtained from Boogert et al. (2011) and Boogert et al. (in preparation). (\*) indicates CH<sub>3</sub>OH detected in this sight-line in Boogert et al. (2011). B2011: published in Boogert et al. (2011).

Source 2MASS J	Instrument	$A_V$	Template star <sup>1</sup> mag	$\chi_v^2$	N(H <sub>2</sub> O)		N(CH <sub>3</sub> OH)		Notes
					10 <sup>18</sup> cm <sup>-2</sup>	10 <sup>17</sup> cm <sup>-2</sup>	10 <sup>17</sup> cm <sup>-2</sup>	%H <sub>2</sub> O	
04215402+1530299	NIRSPEC	24.3	HD6903/F9 (F5-G2)	0.0021	1.84 (0.20)	<0.097	<0.53	B2011	
17111501-2726180	NIRSPEC	31.3	HD94705/M5.5 (M0-M5)	0.00051	2.35 (0.26)	<0.084	<0.36	B2011	
17111538-2727144	NIRSPEC	28.2	HD120052/M2 (M0-M3)	0.0018	2.01 (0.22)	<0.081	<0.40	B2011	
17112005-2727131	NIRSPEC	48.0	HD124897/K1.5 (G8-K5)	0.0027	3.79 (0.42)	<0.081	<0.21	B2011	
17155573-2055312	NIRSPEC	21.2	HD19058/M4 (K2-M4)	0.0013	1.63 (0.18)	<0.91	<5.58	B2011	
17160467-2057072	NIRSPEC	25.9	Gl268AB/M4.5 (M2-M4.5)	0.0027	2.13 (0.23)	<1.08	<5.07	B2011	
17160860-2058142	NIRSPEC	19.6	Gl381/M2.5 (K1-M2.5)	0.0029	1.53 (0.17)	<0.69	<4.51	B2011	
18140712-0708413*	NIRSPEC	15.1	HD9852/K0.5 (K0-K1)	0.0026	1.40(0.15)	<0.59	<4.21	B2011	
18160600-0225539	NIRSPEC	10.0	HD216946/K5 (K1-K5)	0.0028	1.06 (0.19)	<0.71	<6.70	B2011	
18165296-1801287	NIRSPEC	24.0	HD204724/M1 (K5-M3)	0.00094	1.62 (0.22)	<0.72	<4.44	B2011	
18165917-1801158	NIRSPEC	26.7	HD94705/M5.5 (M2-M5.5)	0.00044	1.85 (0.20)	<0.96	<5.19	B2011	
18170376-0815070	NIRSPEC	16.0	HD108849/M7 (M1-M7)	0.00052	0.99 (0.12)	<0.48	<4.85	B2011	
18170426-1802408	NIRSPEC	20.0	HD6903/F9 (F6-F9)	0.0014	1.09 (0.12)	<0.069	<0.63	B2011	
18170429-1802540	NIRSPEC	23.7	HD28487/M3.5 (M1-M5.5)	0.00058	1.43 (0.20)	<0.069	<0.48	B2011	
18170470-0814495	NIRSPEC	34.2	HD108849/M7 (M2-M7)	0.0014	3.93 (0.44)	2.77 (0.016)	7.05	B2011	
18170957-0814136	NIRSPEC	26.2	HD39045/M3 (M3-M7)	0.00071	2.85 (0.31)	2.91 (0.015)	10.21	B2011	
18171181-0814012	NIRSPEC	34.5	HD124897/K1.5 (K1.5-K2)	0.0012	3.81 (0.42)	5.37 (0.014)	14.10	B2011	
18171366-0813188	NIRSPEC	28.6	HD76151/G2 (G2-M0)	0.0044	3.40 (0.38)	3.85 (0.49)	11.32	B2011	
18171700-0813504	NIRSPEC	14.0	HD207076/M7 (M2-M7)	0.0023	0.68 (0.12)	<0.26	<3.82	B2011	
18172690-0438406	NIRSPEC	36.8	HD108849/M7 (M3-M7)	0.00084	4.31 (0.48)	3.48 (0.020)	8.07	B2011	
18275901+0002337	IRTF	7.2	HD214665/M4 (M3-M4)	0.00043	<0.08	<0.23	...		
	NIRSPEC		HD194193/K7 (K4-M0)	0.00063		<0.083	...		
18282010+0029141	IRTF	10.6	HD94705/M5.5 (M2-M6)	0.00023	<0.17	<0.42	...		
	NIRSPEC		HD194193/K7 (K5-M3)	0.00059		<0.082	...		
18282631+0052133	IRTF	5.1	HD132935/K2 (G9-K2)	0.0012	<0.08	<0.16	...		
	NIRSPEC		HD2901/K2 (K1-K7)	0.00092		<0.044	...		
18284038+0044503	IRTF	12.5	HD214665/M4 (K3-M4)	0.0022	0.38 (0.13)	<0.47	<12.37		
	NIRSPEC		HD94705/M5.5 (K3-M5)	0.00060		<0.079	<2.08		
18284797+0037431	IRTF	10.6	HD214665/M4 (K5-M4)	0.00032	0.40 (0.08)	<0.32	<8.00		
	NIRSPEC		HD39801/M1-M2 (M0-M3)	0.00049		<0.073	<1.83		
18285266+0028242	NIRSPEC	36.0	HD28487/M3.5 (M3-M6)	0.0017	2.9 (0.2)	6.14 (0.48)	21.17		
18290316+0023090	IRTF	8.1	HD213893/M0 (M0-M5)	0.00030	0.23 (0.07)	<0.29	<12.61		
	NIRSPEC		HD194193/K7 (K4-M1)	0.00075		<0.060	<2.61		
18290436+0116207	IRTF	10.3	HD207991/K4 (K2-K4)	0.0035	0.18 (0.05)	<0.58	<32.22		
	NIRSPEC		HD187238/K3 (K2-M1)	0.0010		<0.040	<2.22		
18291619+0045143	NIRSPEC	6.1	HD124897/K1.5 (K1.5-K4)	0.00051	<0.10	<0.052	...		
18291699+0037191	IRTF	19.8	HD213893/M0 (M0-M4)	0.00084	1.06 (0.12)	<0.66	<6.23		
	NIRSPEC		HD95735/M2 (M0-M3)	0.0013		<0.071	<0.67		
18292528+0003141	IRTF	7.0	HD28487/M3.5 (M1-M3.5)	0.00022	<0.12	<0.15	...		
	NIRSPEC		HD39801/M1-M2 (K7-M3)	0.00061		<0.065	...		
18294108+0127449	IRTF	8.6	HD27598/M4 (K1.5-M4)	0.00023	<0.08	<0.066	...		
	NIRSPEC		HD91810/K1 (K1-K2)	0.00095		<0.066	...		
18300061+0115201	IRTF	46.0	HD236697/M0.5 (M0.5-M5.5)	0.010	3.04 (0.34)	<0.18	<0.59		
	NIRSPEC		HD108849/M7 (M3-M7)	0.00059		<0.090	<0.30		
	NIRSPEC		HD108849/M7 (M3-M7)	0.00059		<0.095	<0.31	B2011	
18300085+0017069	IRTF	5.9	HD2901/K2 (G9-K2)	0.00040	<0.10	<0.17	...		
	NIRSPEC		HD132935/K2 (K0-K2)	0.0014		<0.035	...		
18300896+0114441	IRTF	13.1	HD207991/K4 (K4-M2)	0.00033	0.50 (0.17)	<0.32	<6.40		
	NIRSPEC		HD14469/M3-M4 (K2-M3)	0.00068		<0.063	<1.26		
19201597+1135146	NIRSPEC	25.1	HD39045/M3 (K5-M3)	0.00057	1.34 (0.14)	<0.87	<6.49	B2011	
19201622+1136292	NIRSPEC	20.0	HD120477/K5.5 (K1.5-K5.5)	0.0013	1.01(0.11)	<0.062	<0.61	B2011	
19214480+1121203	NIRSPEC	16.8	HD39045/M3 (M2-M3)	0.00063	1.43 (0.16)	<0.048	<0.34	B2011	
21240517+4959100	NIRSPEC	24.8	HD28487/M3.5 (M0-M3.5)	0.00052	2.19 (0.24)	0.74 (0.21)	3.38	B2011	
21240614+4958310	NIRSPEC	12.8	HD124897/K1.5 (K1.5-K3)	0.0011	0.98 (0.11)	<0.058	<0.59	B2011	
22063773+5904520	NIRSPEC	13.9	HD10307/G1 (F5-G1)	0.0030	1.13 (0.12)	<0.044	<0.39	B2011	

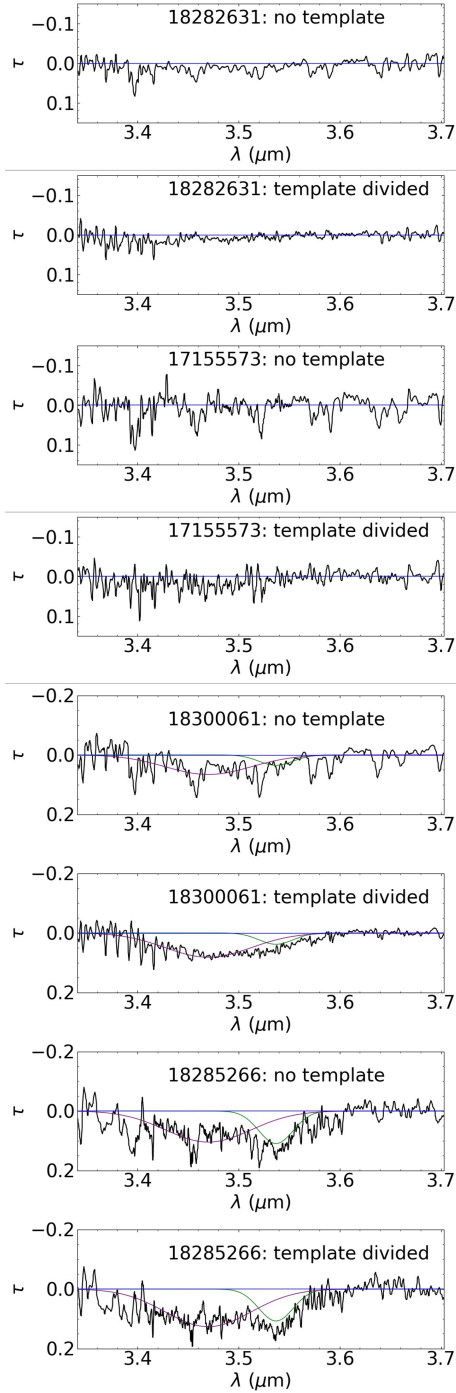


Figure 5.1: Optical depth spectra taken towards four background stars at varying visual extinctions: 2MASS J18282631 ( $A_V=5.1$  mag), 2MASS J17155573 ( $A_V=10.6$  mag), 2MASS J18300061 ( $A_V=46.0$  mag), 2MASS J18285266 ( $A_V=36.0$  mag). For each target, the two panels give the spectrum before (top) and after (bottom) template division, respectively. The blue line represents  $\tau=0$ . The green and purple lines are Gaussian fits to the (tentative) 3.537  $\mu\text{m}$  and 3.47  $\mu\text{m}$  features, respectively.

corresponding reduced  $\chi^2$  values. The photosphere-corrected flux spectra and the final optical depth spectra are found in appendix 5.7.

To demonstrate the effect of the photospheric line correction, optical depth spectra with and without the correction are shown for four sources with  $A_V$  values in the range of 5.1–46.0 mag in Figure 5.1. Clearly, the reduction of photospheric lines allows for a more sensitive search for interstellar ice features. Blended photospheric lines could mimic narrow ice features, such as those of the  $\text{CH}_3\text{OH}$  band at  $3.537 \mu\text{m}$ . Reduction of the photospheric lines also improves the fitting of the baseline. And for secure detections (e.g., 2MASS J18285266), a more reliable band profile is derived, as well as a more secure separation from the  $3.47 \mu\text{m}$  feature.

When both the S/N and the photospheric line reduction are excellent, the  $\text{CH}_3\text{OH}$  ice abundance determination is limited by the decomposition of the  $3.537 \mu\text{m}$  band from the broad ammonia hydrate feature at  $3.47 \mu\text{m}$ . Unless a distinct dip at  $3.537 \mu\text{m}$  is observed, the Gaussian decomposition (section 5.3) is uncertain.

For one of the targets in Boogert et al. (2011) we obtain different results. We do not confirm the  $\text{CH}_3\text{OH}$  ice detection toward 2MASS J18140712 ( $7 \pm 2 \times 10^{16} \text{cm}^{-2}$  versus  $< 6 \times 10^{16} \text{cm}^{-2}$  ( $3\sigma$ ) in this work). Figure 5.2 shows that significant photospheric lines are still present after the division over the template. Perhaps our disagreement with Boogert et al. (2011) is due to the baseline choice.

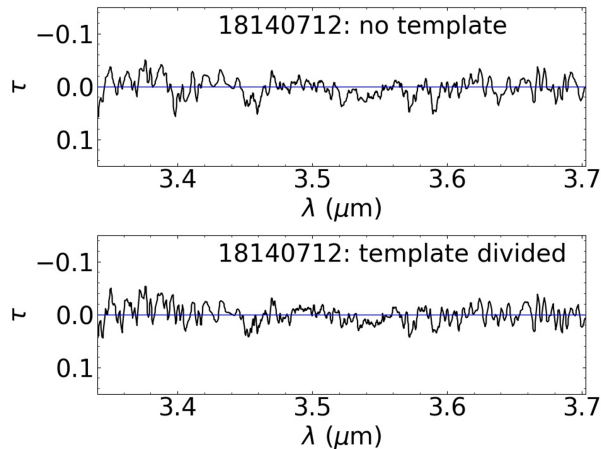


Figure 5.2: Optical depth spectra taken towards the field star, 18140712 (Boogert et al. 2011). No observable  $3.537 \mu\text{m}$   $\text{CH}_3\text{OH}$  feature is present in either panel, whereas the feature was present in Boogert et al. (2011). The choice of baseline may be the cause of such discrepancy. The blue line indicates  $\tau=0$ .

#### 5.4.1 $\text{CH}_3\text{OH}$ column density relative to $A_V$ and $\text{H}_2\text{O}$

Figure 5.3 shows the  $\text{CH}_3\text{OH}$  ice detections with  $1\sigma$  error bars and  $3\sigma$  upper limit values from Table 5.3 plotted as a function of the  $A_V$ . For upper limits of the same source, the lowest upper limit is plotted. Table 5.3 also shows the  $\text{CH}_3\text{OH}$  abundance relative to the  $\text{H}_2\text{O}$  ice column density. The  $A_V$  and  $N(\text{H}_2\text{O})$

values were taken from Boogert et al. (2011) and Boogert et al. (in preparation), and were determined by fitting observed 1-5  $\mu\text{m}$  broad band photometry and spectra with a model consisting of spectra from the IRTF spectral database, the infrared interstellar extinction curve, an optical reddening of  $A_V/A_K = 8$ , and an  $\text{H}_2\text{O}$  ice spectrum for small spherical grains.

A fit is only made to the detections (closed blue circles). They apparently trace a distinct environment particularly conducive to  $\text{CH}_3\text{OH}$  formation. As shown in the figure, the fit crosses the y-axis at  $A_V = 13.5 \pm 7.8$  mag (i.e., the detection threshold). Taking into account the front and back of the cloud/core, the extinction as measured from the edge to the middle of the cloud/core along the line of sight is the observed value divided by 2, which results in a  $\text{CH}_3\text{OH}$  formation threshold of  $6.8 \pm 3.9$  mag.

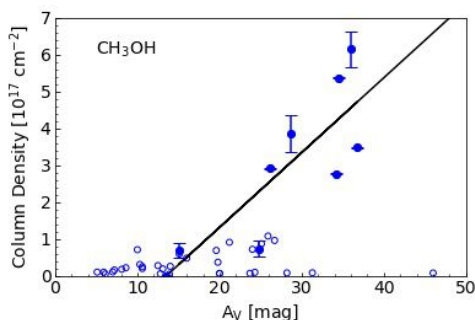


Figure 5.3:  $\text{CH}_3\text{OH}$  ice detections (solid blue circles;  $1\sigma$  error bars) and upper limits (open blue circles;  $3\sigma$ ) plotted as a function of the visual extinction. A  $\text{CH}_3\text{OH}$  detection threshold of  $13.5 \pm 7.8$  mag is found.

#### 5.4.2 Averaging spectra in $A_V$ bins

Spectra that do not show a secure  $\text{CH}_3\text{OH}$  feature at  $3.537 \mu\text{m}$  are averaged into various  $A_V$  bins to increase the sensitivity (Figure 5.4). The averaging is weighed by  $1/\sigma^2$ , with  $\sigma$  the typical value used to determine the S/N column of Table 5.2. Three bins exist: an average over the lines of sight below the  $\text{CH}_3\text{OH}$  detection threshold, at and above the  $\text{CH}_3\text{OH}$  detection threshold, and with  $\text{CH}_3\text{OH}$  detections. The noise level in the  $A_V < 18$  mag bin is 0.16, with a  $3\sigma$  upper limit of  $7.68 \times 10^{17} \text{ cm}^{-2}$  for the  $3.537 \mu\text{m}$  feature. For the  $A_V \geq 18$  mag bin, the noise level is 0.15, with a  $3\sigma$  upper limit of  $5.26 \times 10^{16} \text{ cm}^{-2}$  for the  $3.537 \mu\text{m}$  feature. A noise level of 0.056 and a  $\text{CH}_3\text{OH}$  column density of  $2.58 \times 10^{18} \text{ cm}^{-2}$  are measured for the bin containing  $\text{CH}_3\text{OH}$  detections. The  $A_V$  range of  $< 18$  mag in Figure 5.4 does not show any observable  $\text{CH}_3\text{OH}$  feature. The  $3.47 \mu\text{m}$   $\text{NH}_3 \cdot \text{H}_2\text{O}$  feature is clearly present at  $A_V \geq 18$  mag, however there is no distinct feature at  $3.537 \mu\text{m}$ .

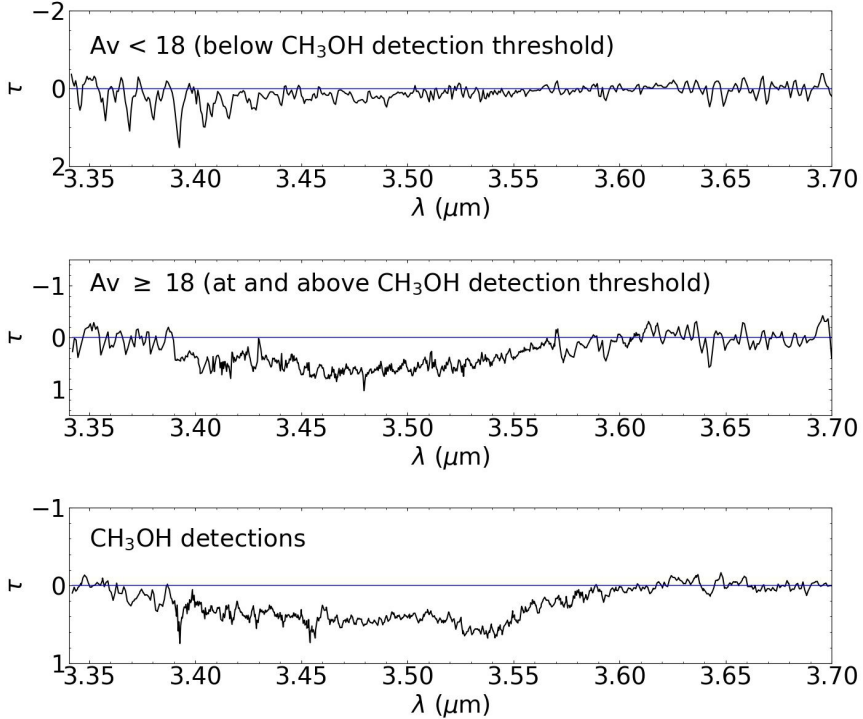


Figure 5.4: Optical depth spectra averaged across the indicated  $A_V$  bins, for targets without a  $\text{CH}_3\text{OH}$  ice detection, as probed by the  $3.537\ \mu\text{m}$  feature (see Table 5.3). For reference, a spectrum averaged over sight-lines with  $\text{CH}_3\text{OH}$  detections is shown in the bottom panel. The blue line shows  $\tau=0$ . The number of targets in each bin is 30, 18, and 7 for  $A_V < 18$ ,  $A_V \geq 18$ , and  $\text{CH}_3\text{OH}$  detections averaged, respectively. Note that the detection for sight-line 2MASS J18140712 is not included.

## 5.5 Discussion

The extinction threshold for ice formation is a direct observational indicator of the chemical evolution of molecular clouds. The observed  $A_V$  values mentioned above sample both, the front and back of the clouds/cores. And thus, the maximum depth into the clouds/cores that is traced by these observations is  $A_V/2$  (i.e., the formation threshold), which we will use in the remainder of this discussion.

The threshold value for abundant  $\text{CH}_3\text{OH}$  formation of  $A_V=6.8 \pm 3.9$  mag as derived from the linear fit is of low significance ( $1.7\sigma$ ), reflecting the large spread in the  $\text{CH}_3\text{OH}$  abundances, while the uncertainties on the individual abundances are small.

Also, for many high extinction sight-lines, tight  $\text{CH}_3\text{OH}$  abundance upper limits were derived (Table 5.3), well below the detections at similar  $A_V$  (Figure 5.3). This strengthens the trend observed in Boogert et al. (2015). We speculate that CO has not sufficiently frozen out in these sight-lines, limiting the  $\text{CH}_3\text{OH}$  formation in some of these dense clouds and dark cores. CO

freeze out could have been suppressed by lower densities at the same  $A_V$ , by higher dust temperatures or radiation fields. Alternatively,  $\text{CH}_3\text{OH}$  may have been consumed or destroyed. Other formation channels that overlap with the  $\text{CH}_3\text{OH}$  formation route, such as those of  $\text{HCOOH}$  (Qasim et al. 2019a), may also be contributing factors.

Below  $A_V=7.6$  mag, only upper limits are available, but they are tight ( $<0.39\%$   $\text{H}_2\text{O}$ ), supporting that the threshold for abundant  $\text{CH}_3\text{OH}$  formation is indeed well above that for  $\text{H}_2\text{O}$  ( $\sim 1.6$ ) and  $\text{CO}$  ( $\sim 3$ ).  $\text{CH}_3\text{OH}$  formation below these limits is expected, however. As reported in the laboratory investigations by Bergner et al. (2017) and Qasim et al. (2018),  $\text{CH}_4$  is a precursor to  $\text{CH}_3\text{OH}$  formation under low temperature molecular cloud conditions. The formation efficiency is a factor of  $\sim 20$  below that of the high density  $\text{CO}$  ice hydrogenation route (Qasim et al. 2018).  $\text{CH}_4$  is formed at low extinctions ( $A_V \sim 2$  mag), similar to  $\text{H}_2\text{O}$ . At this stage,  $\text{CO}$  formation in the gas phase is not complete, and could enhance the  $\text{CH}_4$  abundance in the first grain mantle layers. There is indeed a hint of enhanced  $\text{CH}_4$  ice abundances at low extinctions (Öberg et al. 2008). Consequently,  $\text{CH}_3\text{OH}$  is expected to be present in the ices well below the threshold for enhanced formation via the  $\text{CO}$  ice hydrogenation route discussed above.

## 5.6 Conclusions and future work

Observations targeting the  $3.537 \mu\text{m}$  feature of  $\text{CH}_3\text{OH}$  ice in dense molecular clouds and dark cores to constrain the  $\text{CH}_3\text{OH}$  ice formation threshold were pursued. A method to increase the sensitivity to weak ice absorption features and thus to the accuracy of the  $\text{CH}_3\text{OH}$  column density is presented. The following conclusions can be drawn from this work:

- 1) The method of fitting template spectra of unreddened stars to reduce photospheric lines from dense cloud and dark core background spectra increases the sensitivity to detect ice features. It also improves the accuracy of the baseline determination. With respect to the study of (Boogert et al. 2011), one new  $\text{CH}_3\text{OH}$  ice detection was made and one was not found. The total number of  $\text{CH}_3\text{OH}$  ice detections in quiescent molecular cloud and core material is now 8.

- 2) This photospheric correction method will be applicable to the analysis of future JWST spectra. Multi-object spectroscopy with JWST will make it possible to observe large sample sizes for both, background and template stars.

- 3) For high S/N spectra that are well corrected for photospheric lines, the limiting factor for accurate  $\text{CH}_3\text{OH}$  ice abundances is the decomposition of the  $3.537 \mu\text{m}$  band from the  $3.47 \mu\text{m}$  band of ammonia hydrates.

- 4) The  $\text{CH}_3\text{OH}$  formation threshold is constrained to be  $A_V=6.8 \pm 3.9$  mag. The error bars on the individual measurements are much smaller than the scatter in the column densities. No detections have been made below a formation threshold of  $A_V=7.6$  mag. The abundance upper limits obtained in many sight-lines both below and above the  $13.5$  mag detection threshold are significant.

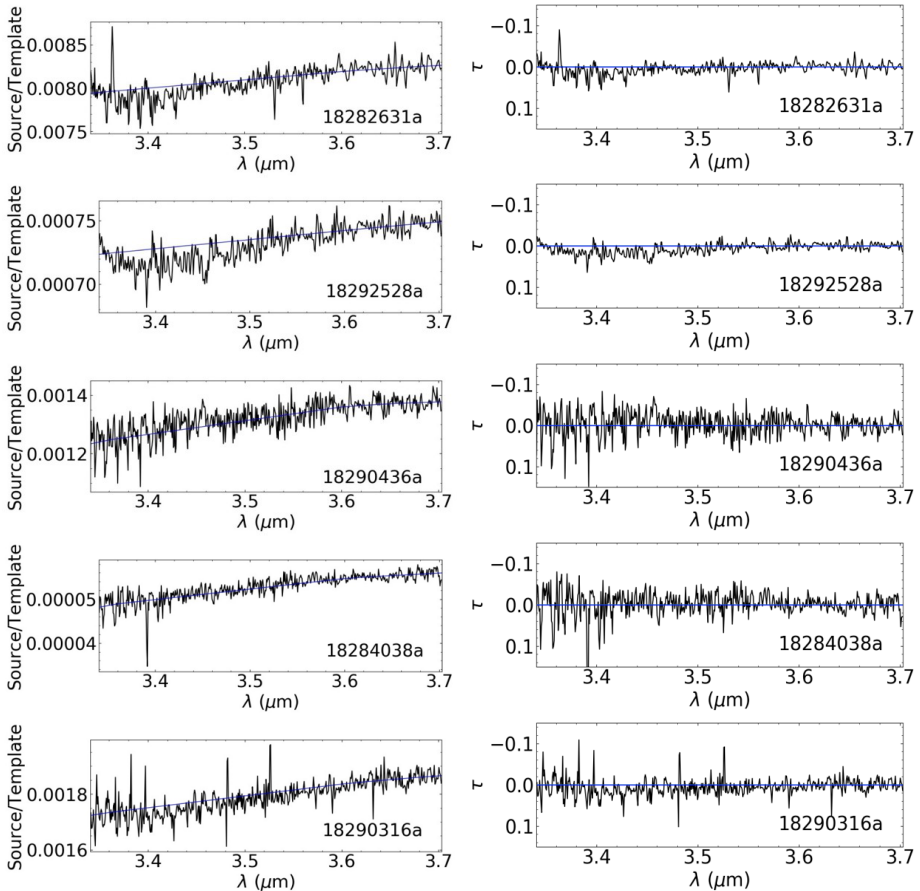
- 6) The large scatter at high  $A_V$  values is puzzling. Perhaps  $\text{CO}$  is not an abundant resource for  $\text{CH}_3\text{OH}$  formation in some of these clouds and cores. Alternatively,  $\text{CH}_3\text{OH}$  may have been effectively destroyed.

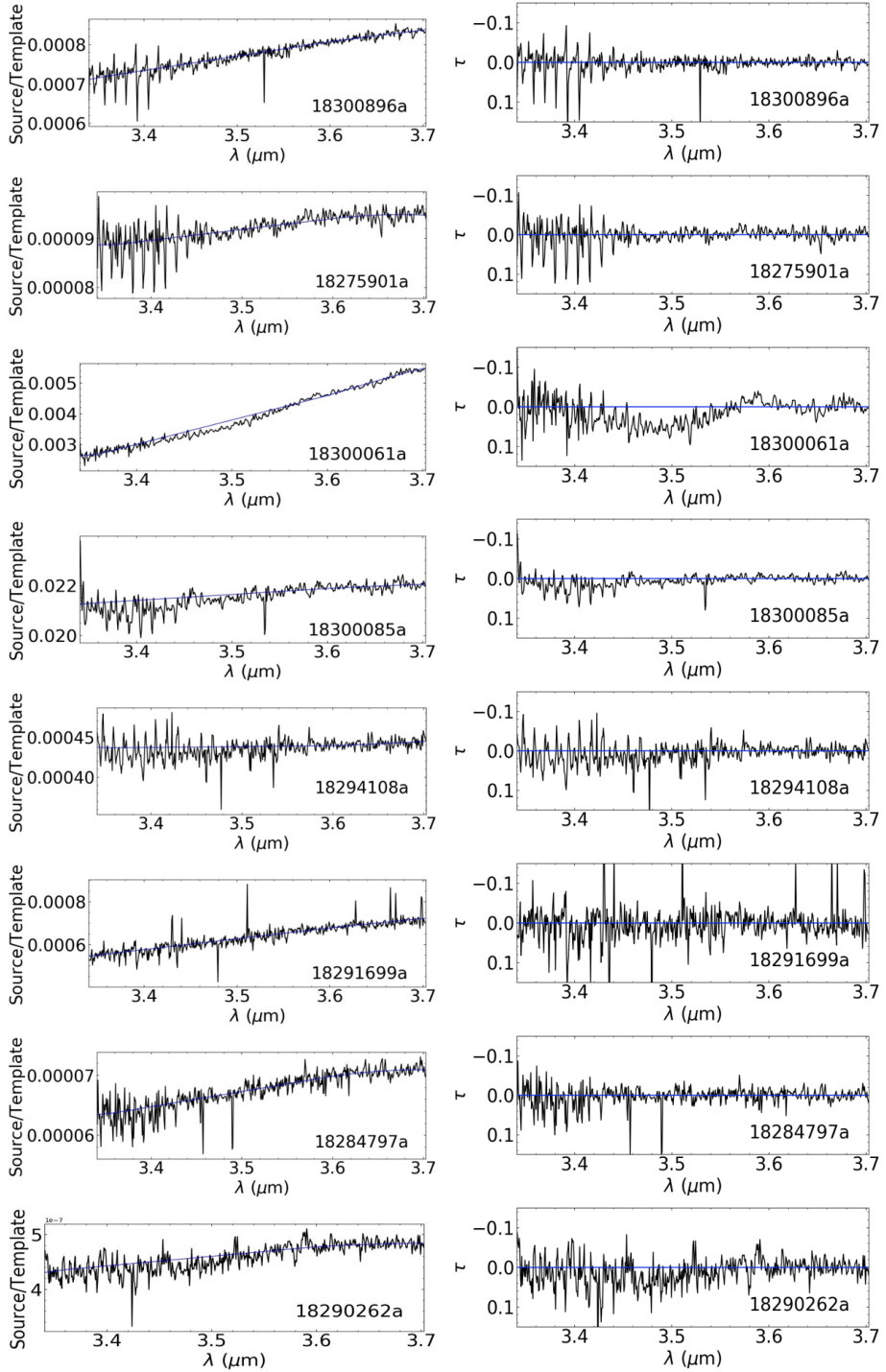
7) The threshold for the efficient formation of  $\text{CH}_3\text{OH}$  via the CO ice hydrogenation route does not preclude the formation via less efficient pathways. Such pathways may well include  $\text{CH}_4$  and/or  $\text{CH}_4$  precursors (C, CH, etc.), as supported from recent laboratory experiments.

8) The sample size of  $\text{CH}_3\text{OH}$  detections is small and more observations are warranted to further investigate the origin of the large scatter in the  $\text{CH}_3\text{OH}$  abundance variations as well as the formation threshold.

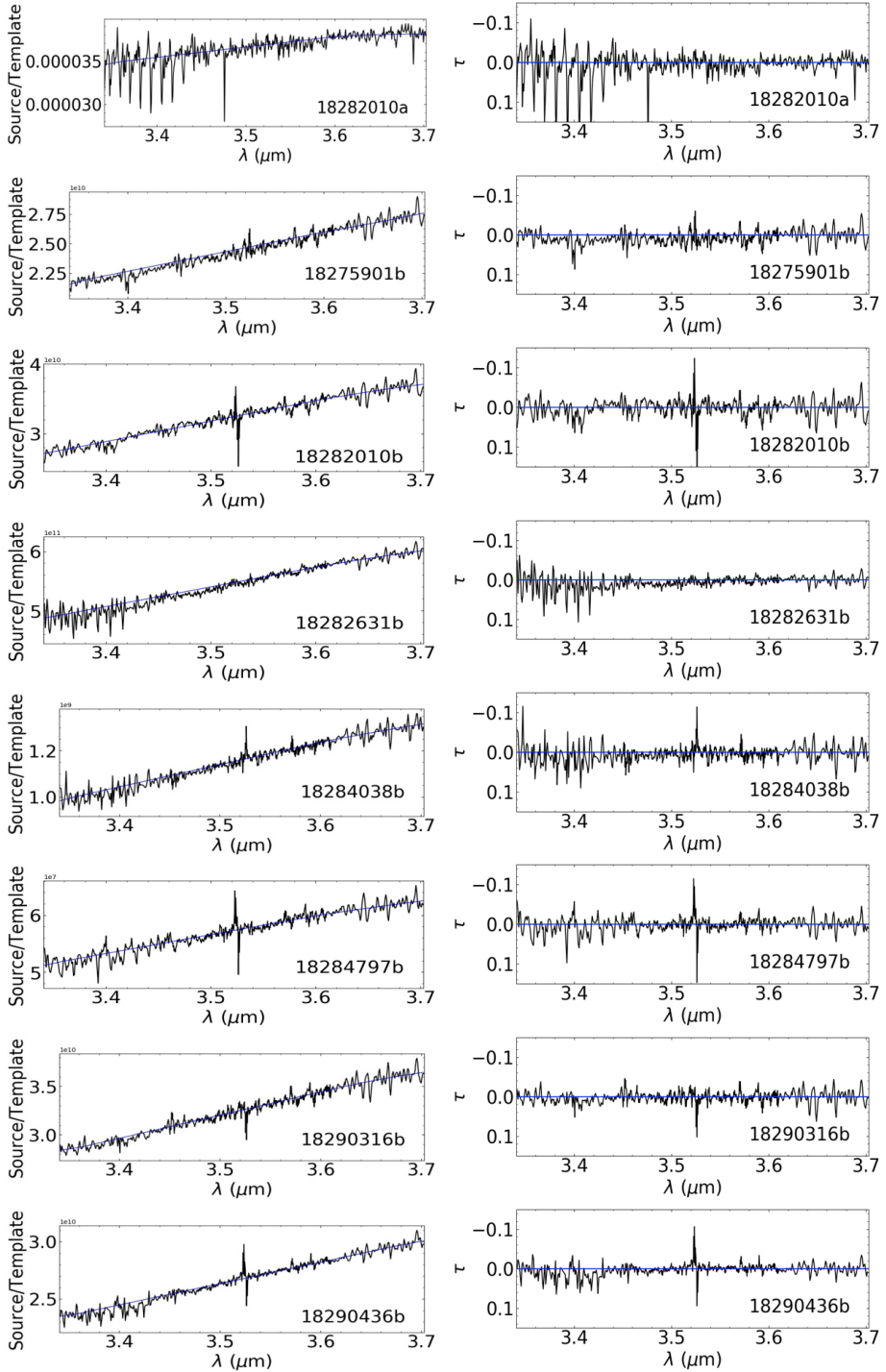
9) Tracing the columns of CO and  $\text{CH}_4$  alongside  $\text{CH}_3\text{OH}$  in the same clouds and cores will help understand how and when  $\text{CH}_3\text{OH}$  is formed. Whether CO is actually depleted in regions where  $\text{CH}_3\text{OH}$  is not detected, and if  $\text{CH}_4$  could be a precursor for less abundant  $\text{CH}_3\text{OH}$ , can be better understood by mapping these ices simultaneously. JWST is ideally suited for this, because of its large wavelength coverage, high sensitivity, and mapping capabilities.

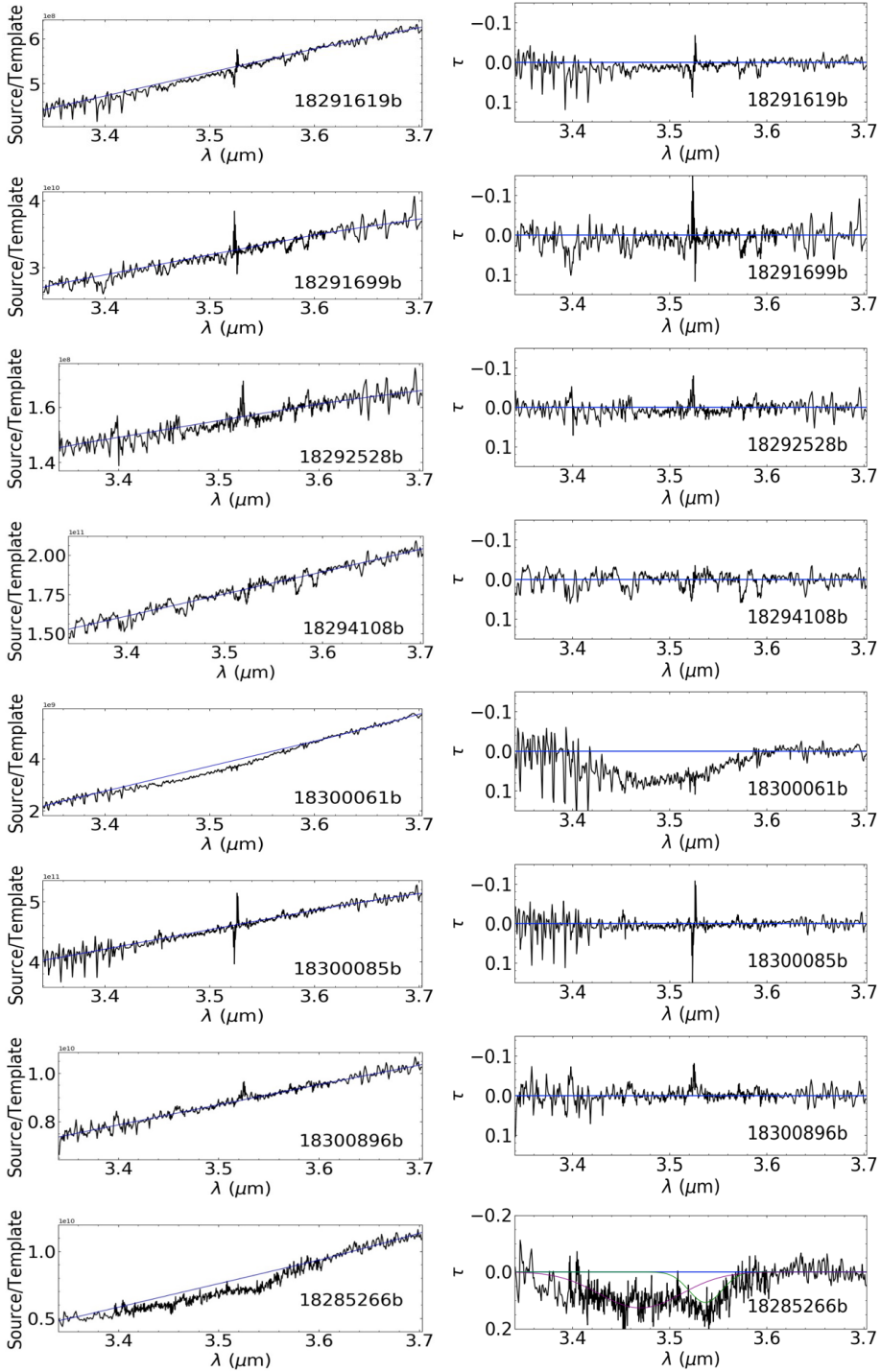
## 5.7 Appendix information

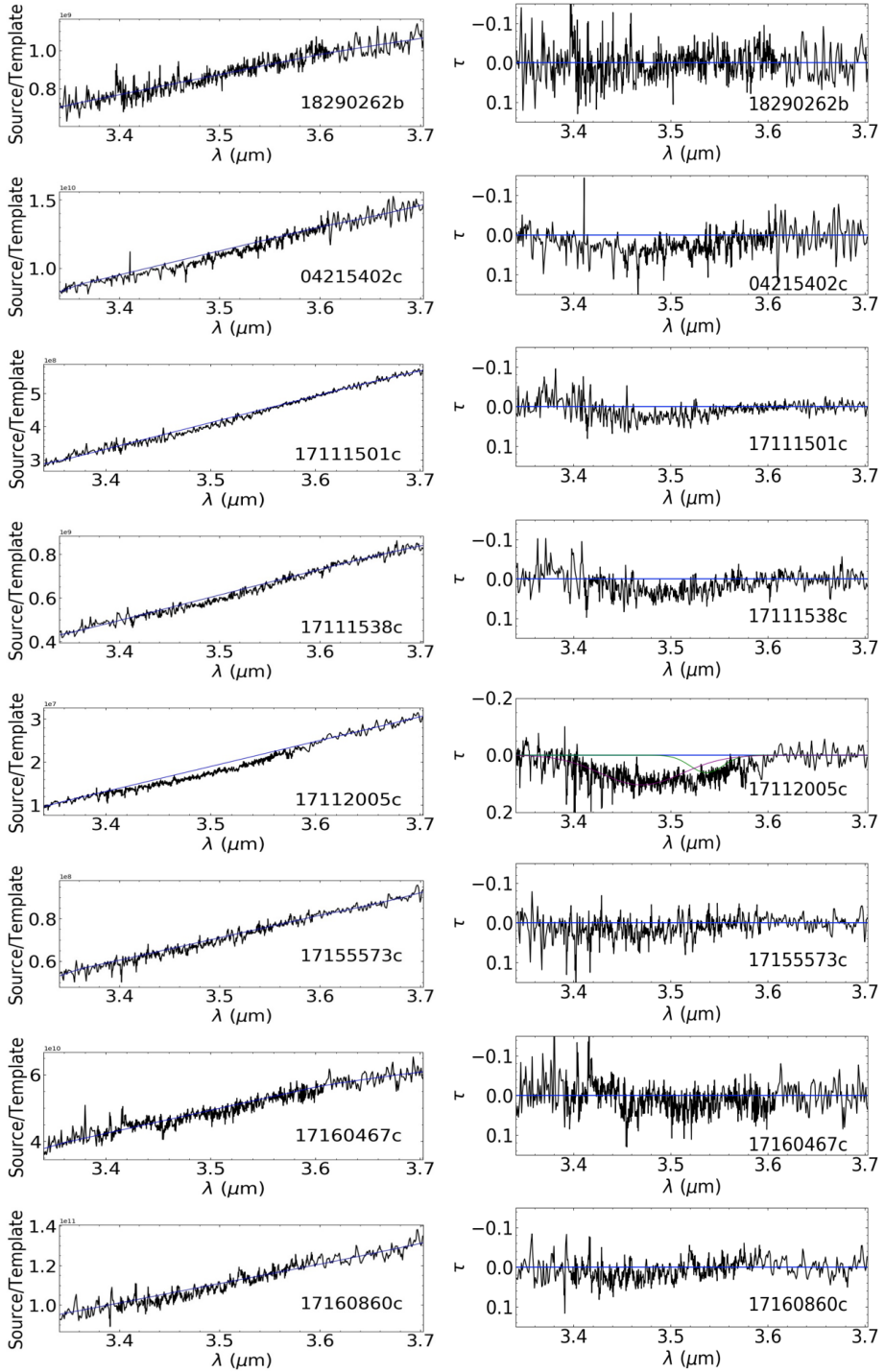


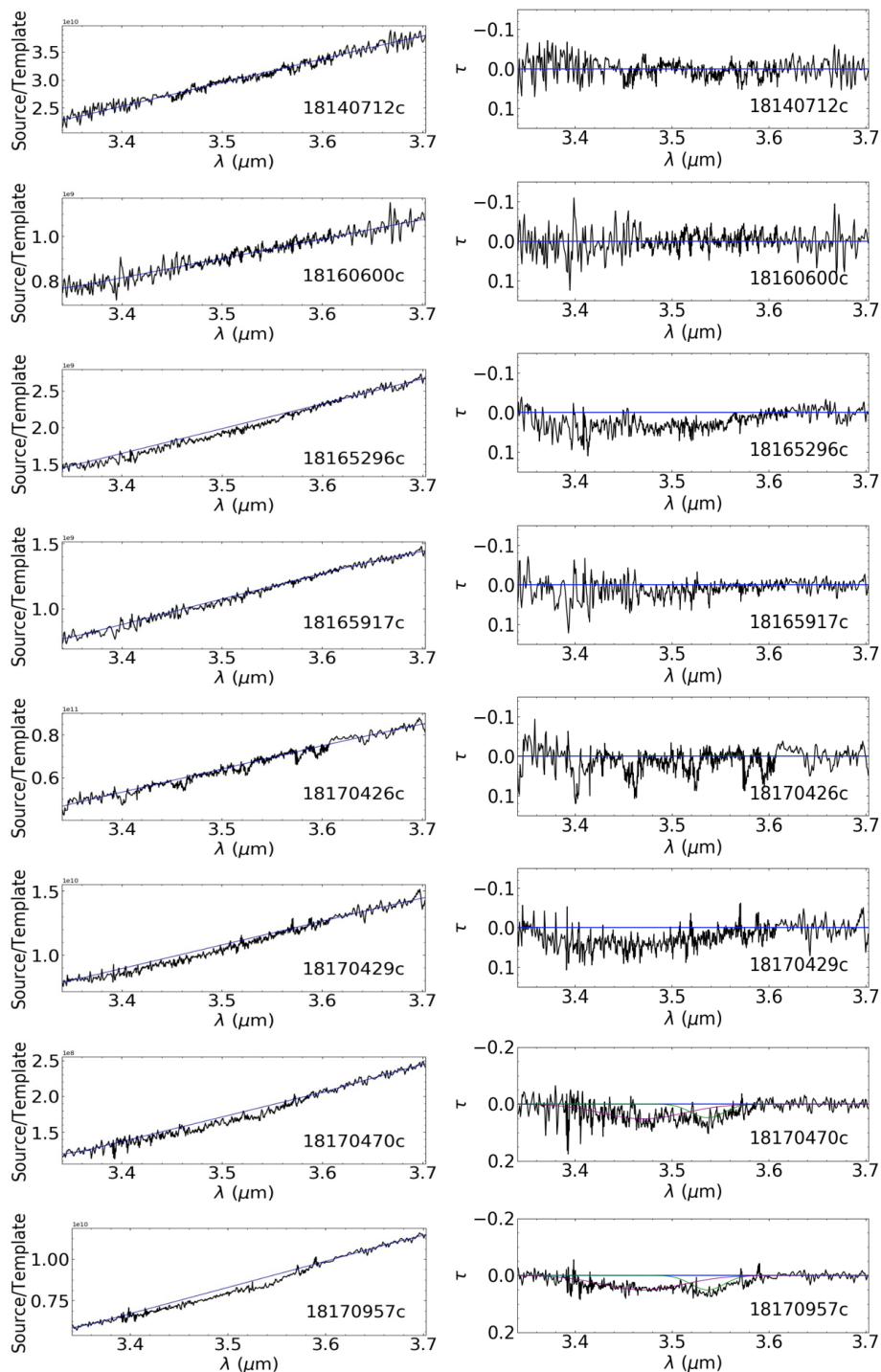


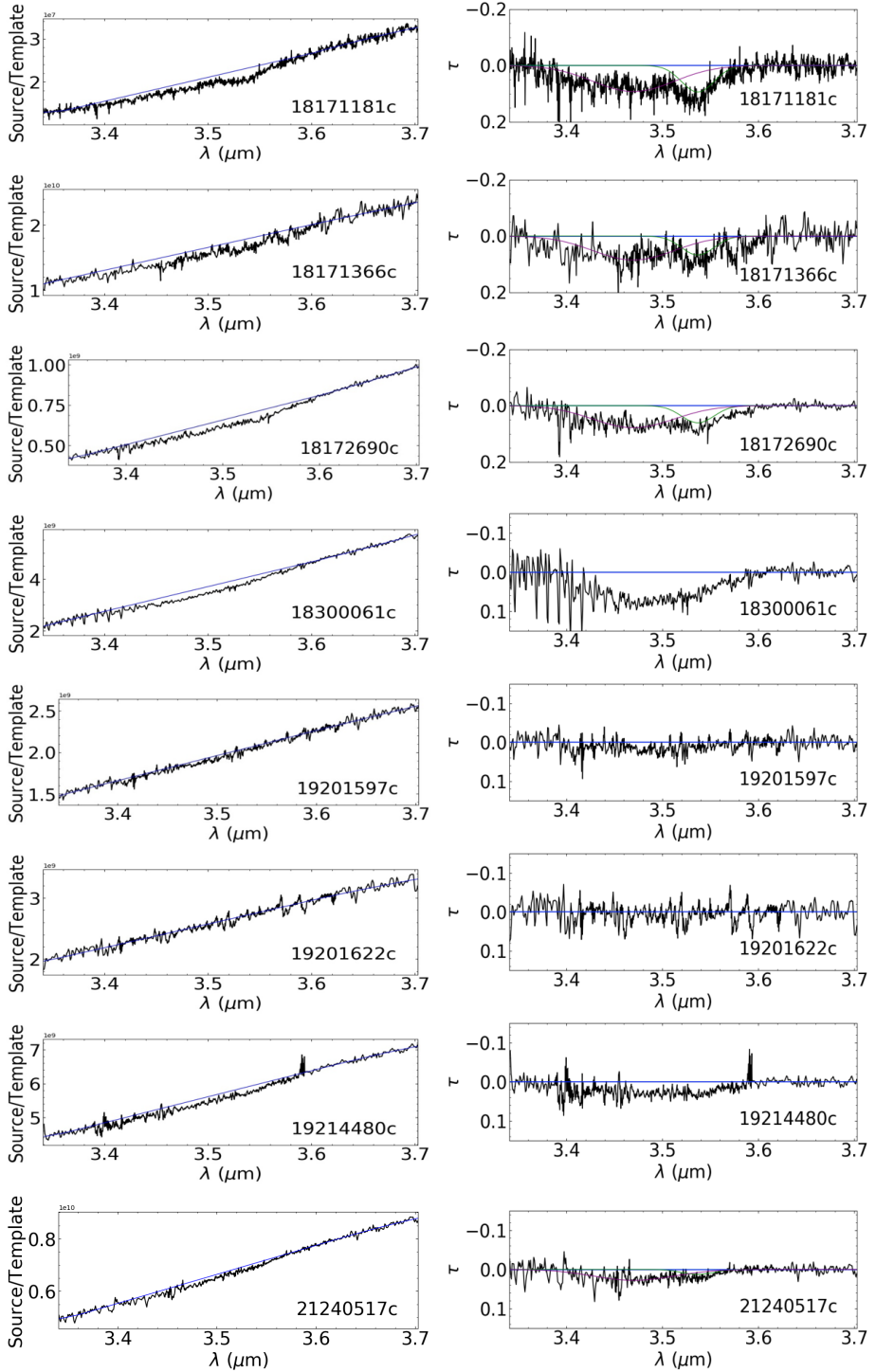












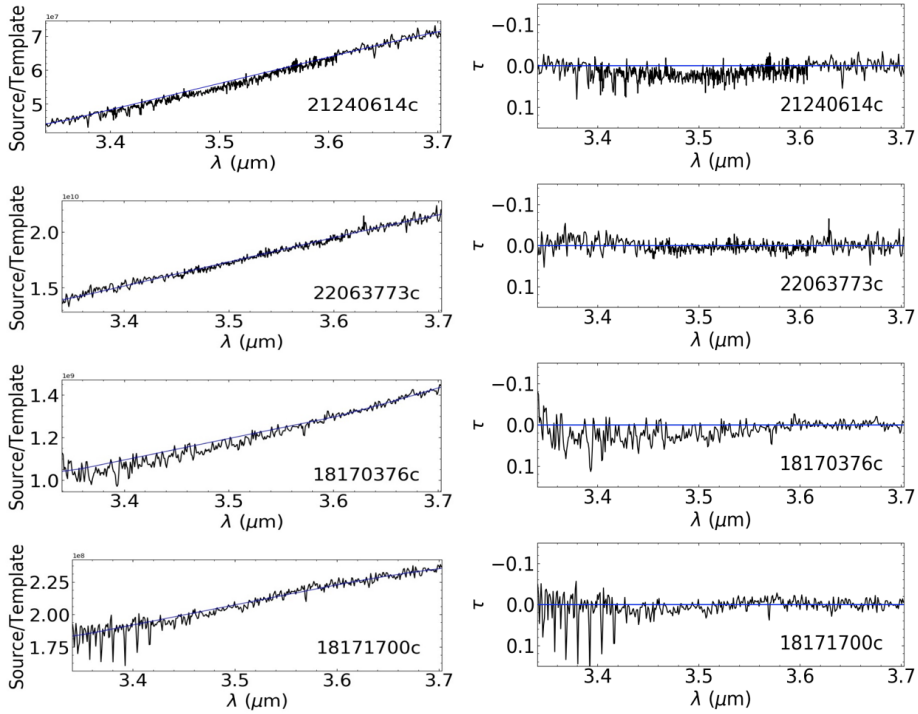


Figure 5.5: Flux and optical depth spectra of the 3.3-3.7  $\mu\text{m}$  wavelength region towards the entire sample. The blue line in each left panel is the local continuum, and in the right panel is  $\tau=0$ . The green and purple lines are Gaussian fits for the 3.537  $\mu\text{m}$  and 3.47  $\mu\text{m}$  absorption features, respectively.

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