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ALMA RESOLVES THE TORUS OF NGC 1068: CONTINUUM AND MOLECULAR LINE EMISSION

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ABSTRACT

We used the Atacama Large Millimeter Array (ALMA) to map the emission of the CO(6–5) molecular line and the 432 μm continuum emission from the 300 pc sized circumnuclear disk (CND) of the nearby Seyfert 2 galaxy NGC 1068 with a spatial resolution of ~ 4 pc. These observations spatially resolve the CND and, for the first time, image the dust emission, the molecular gas distribution, and the kinematics from a 7–10 pc diameter disk that represents the submillimeter counterpart of the putative torus of NGC 1068. We fitted the nuclear spectral energy distribution of the torus using ALMA and near- and mid-infrared (NIR/MIR) data with CLUMPY torus models. The mass and radius of the best-fit solution for the torus are both consistent with the values derived from the ALMA data alone: $M_{\text{gas}}^{\text{torus}} = (1 \pm 0.3) \times 10^5 M_{\odot}$ and $R_{\text{torus}} = 3.5 \pm 0.5$ pc. The dynamics of the molecular gas in the torus show strong non-circular motions and enhanced turbulence superposed on a surprisingly slow rotation pattern of the disk. By contrast with the nearly edge-on orientation of the H₂O megamaser disk, we found evidence suggesting that the molecular torus is less inclined ($i = 34^{\circ}$ – 66°) at larger radii. The lopsided morphology and complex kinematics of the torus could be the signature of the Papaloizou–Pringle instability, long predicted to likely drive the dynamical evolution of active galactic nuclei tori.

Key words: galaxies: individual (NGC 1068) – galaxies: ISM – galaxies: kinematics and dynamics – galaxies: nuclei – galaxies: Seyfert

1. INTRODUCTION

NGC 1068 is the prototypical Seyfert 2 galaxy and a prime example for active galactic nuclei (AGNs) unifying schemes since the discovery of polarized optical continuum and broadline emission in this source (Miller & Antonucci 1983; Antonucci & Miller 1985). Its central engine is thought to be hidden behind a screen of obscuring material located in a dusty molecular torus of a few parsecs size. Given the distance to the galaxy ($D \sim 14$ Mpc; Bland-Hawthorn et al. 1997), which implies a scale ~ 70 pc $''$, spatially resolving the molecular/dusty torus has been challenging, as angular resolutions $< 0''$ are required.

Single-dish observations in the near- and mid-infrared (NIR and MIR) have shown extended emission around the central engine (Bock et al. 2000; Tomono et al. 2001; Rouan et al. 2004; Galliano et al. 2005; Gratadour et al. 2006,

2015). Higher-spatial resolution interferometric observations in the NIR and MIR identified two components (Jaffe et al. 2004; Weigelt et al. 2004; Raban et al. 2009; Burtscher et al. 2013; López-Gonzaga et al. 2014). First, a compact 0.5–1.4 pc sized core at the AGN with a position angle $PA \simeq 140^{\circ}$ that contains hot ($\simeq 800$ K) dust co-spatial with the H₂O megamaser disk (Greenhill et al. 1996, hereafter GR96; Greenhill & Gwinn 1997; Gallimore et al. 2001, hereafter GA01; Gallimore et al. 2004).

A second component of warm ($\simeq 300$ K) dust extends over 3 to 10 pc along the north-south axis, and may correspond to polar dust emission in the ionization cone. Due to their limited (u, v) coverage, IR interferometric data have to be processed with reconstruction techniques to provide a best-fit model of the putative torus in the (u, v) plane. This limitation can be to a large extent circumvented in the submillimeter using

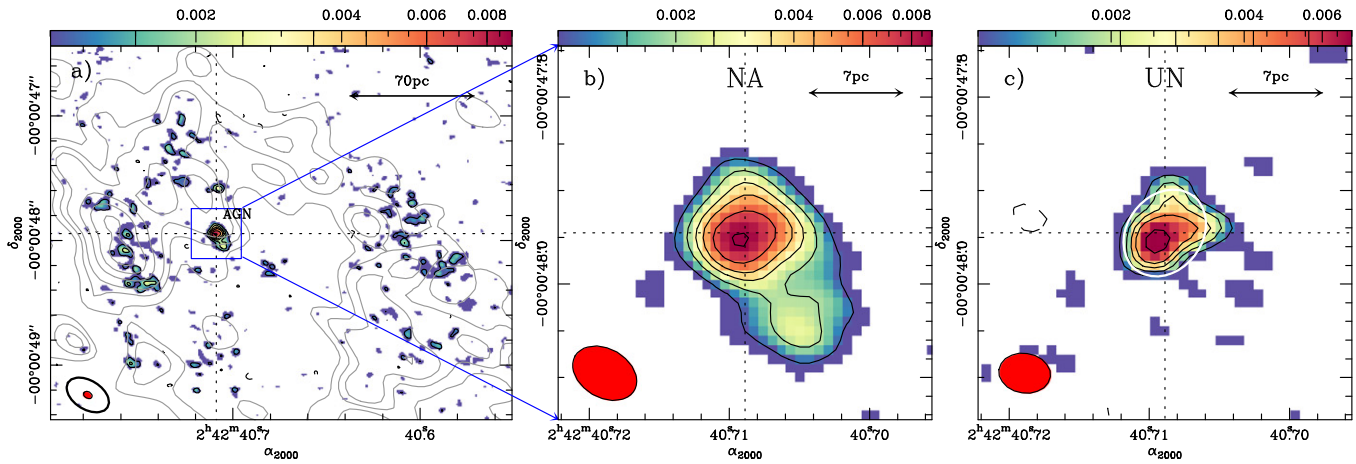


Figure 1. (a) Dust continuum emission at 694 GHz ($432 \mu\text{m}$) mapped by ALMA in the CNL of NGC 1068. The natural (NA)-weighted map is shown in color scale (in Jy beam^{-1} -units) and (black) contour levels (3σ , 5σ , 7σ , 9σ , 12σ , and 16σ where $1\sigma = 0.5 \text{ mJy beam}^{-1}$). The red-filled ellipse in the bottom left corner represents the beam size at 694 GHz ($0''.07 \times 0''.05$ at $PA = 60^\circ$). The gray contours (10%, 20%, 30% to 90% in steps of 20% of the peak value: 49 mJy beam^{-1}) identify the dust emission obtained in GB14 using ALMA with a lower resolution: $0''.4 \times 0''.2$ at $PA = 50^\circ$ (black ellipse). The dashed lines highlight the location of the AGN at $(\alpha_{2000}, \delta_{2000}) = (02^{\text{h}}42^{\text{m}}40.709^{\text{s}}, -00^{\circ}00'47''.95)$. (b) A close-up of the dust continuum emission shown in left panel. (c) The same as middle panel but using a uniform (UN)-weighted set of data with a spatial resolution: $0''.06 \times 0''.04$ at $PA = 82^\circ$. Contour levels are 3σ , 4σ , 5σ , 6σ , and 8σ , where $1\sigma = 0.7 \text{ mJy beam}^{-1}$. The white ellipse identifies the disk solution found by the task UV_FIT.

interferometers like ALMA, whose large number of antennas and baselines assure a much more complete coverage of the (u, v) -plane.

We used ALMA in Cycle 0 to image the dust continuum and CO($J = 6 - 5$) line emissions at 689 GHz in the circumnuclear disk (CND) of NGC 1068 with a spatial resolution of ~ 20 pc (García-Burillo et al. 2014, hereafter GB14). The CND appeared as a $300 \text{ pc} \times 200 \text{ pc}$ sized off-centered ellipsoidal ring with two prominent knots located east and west of the AGN. Although significant continuum and line emissions were also detected at the position of the AGN, the insufficient spatial resolution of these observations prevented us from isolating the torus. The new ALMA Cycle 2 observations of the CNL of NGC 1068 presented in this Letter have a spatial resolution of ~ 4 pc. This factor of ≥ 20 smaller beam area compared with GB14 allowed us to resolve the CND and image, and for the first time, the dust emission and the distribution and kinematics of molecular gas from the torus of NGC 1068 in the submillimeter wavelength range.

2. ALMA OBSERVATIONS AND DATA REDUCTION

We observed the CO($J = 6 - 5$) line and its underlying continuum emission (at $432 \mu\text{m}$) in NGC 1068 with ALMA during Cycle 2 using the Band 9 receiver (project ID: #2013.1.00055.S). The projected baselines range from 41 m to 2251 m. The data were calibrated using the ALMA reduction package CASA (McMullin et al. 2007). The calibrated UV tables were exported to GILDAS²¹ for mapping and CLEANing. One track was observed with 35 antennas during September 2015 using the extended configuration of the array, which allowed us to reach a spatial resolution of $0''.07 \times 0''.05$ using natural (NA) weighting, and a field of view of $9''$ that covers the entire CNL of the galaxy. Four spectral windows with a spectral bandwidth of 1.875 GHz were placed at sky frequencies of 687.970, 688.888, 691.701, and 706.221 GHz. The second sub-band was centered at the redshifted frequency of the CO line. To optimize the deconvolution and cleaning, we

used polygons to restrict the regions where cleaning components are identified. The polygons were defined for the CO line image for each channel based on the detection of significant ($\geq 5\sigma$) emission identified in our Cycle 0 project (GB14). The single polygon used in the cleaning of the continuum image was based on a similar detection criterion using the Cycle 0 data. The sensitivity in the line data cube is 4 mJy beam^{-1} in channels of 12.8 km s^{-1} width. We obtained images of the continuum emission by averaging in each sub-band the channels free of possible line emission. The resulting maps were combined to produce an image of the continuum emission at 694 GHz. The corresponding point source sensitivity for the continuum is $0.5 \text{ mJy beam}^{-1}$. Given the (u, v) -coverage, we expect to filter a significant amount of emission on scales $\geq 2''$.

The phase tracking center of the central field is the same as in GB14 ($\alpha_{2000}, \delta_{2000}$) = $(02^{\text{h}}42^{\text{m}}40.771^{\text{s}}, -00^{\circ}00'47''.84)$ and the reference velocity is $v_o(\text{HEL}) = 1136 \text{ km s}^{-1}$.

3. CONTINUUM AND LINE EMISSION MAPS

(1) *The CNL and the torus.* Figure 1 shows the continuum emission of the CNL imaged by the new ALMA observations. While $\sim 90\%$ of the total continuum flux is spread over a large number of faint emission clumps dispersed throughout the CNL, the strongest emission peak corresponds, within $0''.02$ ($\sim 1-2$ pc), to the position of the AGN (see Figures 1(a), (b)). The extended emission connecting the AGN with the CNL in the map of GB14 is likely resolved out in the present data. The AGN is identified as the source S1 in the radio-continuum maps of Gallimore et al. (1996) and Gallimore et al. (2004).²²

We have used the GILDAS task UV_FIT to find the best fit to the visibilities inside a radius $r \leq 0''.2$ (14 pc) from the AGN. The fit indicates that half of the total flux in this region stems from a spatially resolved ellipsoidal disk, characterized by a deconvolved diameter of 7 ± 1 pc oriented along $PA^{\text{major}} = 142^\circ \pm 23^\circ$. Hereafter, we refer to this fitted disk as the dust torus. The orientation of the dust torus is similar to that of the H₂O

²¹ <http://www.iram.fr/IRAMFR/GILDAS>

²² Imanishi et al. (2016) have also recently isolated HCN(3-2) and HCO⁺(3-2) emission from the AGN.

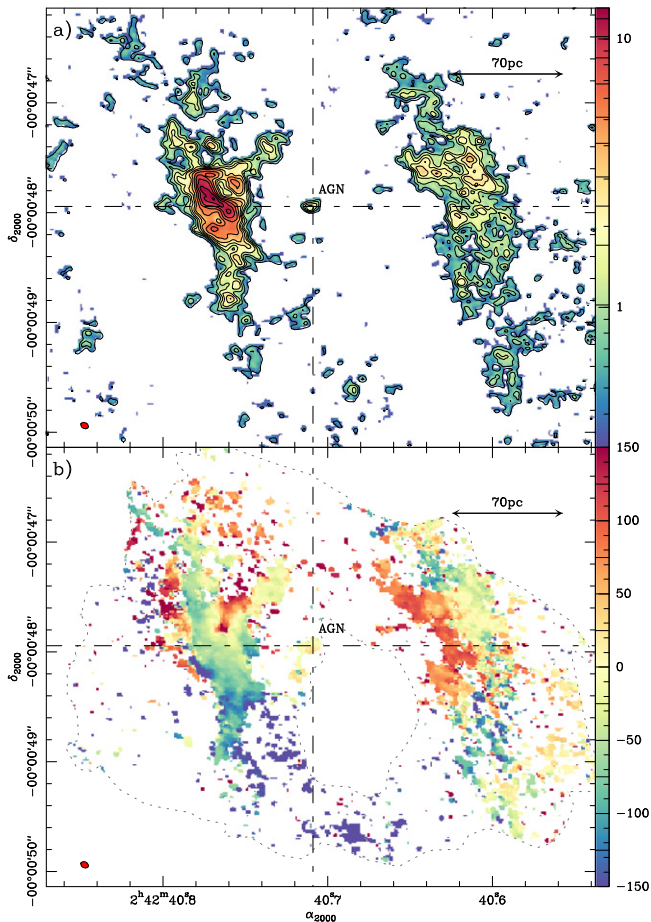


Figure 2. (a) The CO(6–5) intensity map of the CN of NGC 1068. The color scale and contours span the range: 3σ , 5σ , 9σ , 12σ , 15σ to 40σ in steps of 5σ , where $1\sigma = 0.34 \text{ Jy km s}^{-1} \text{ beam}^{-1}$. The filled ellipse at the bottom left corner represents the CO(6–5) beam size ($0''.07 \times 0''.05$ at $PA = 60^\circ$). (b) The CO(6–5) mean-velocity map (color scale). The dashed polygon identifies the region where significant ($>5\sigma$) emission was detected in GB14. Velocities refer to $v_0(\text{HEL}) = 1136 \text{ km s}^{-1}$.

maser disk ($PA^{\text{maser}} = 140^\circ \pm 5^\circ$; GR96, GA01). The minor-to-major axis ratio of the dust torus is $\simeq 0.8 \pm 0.1$, an indication that the disk is spatially resolved along its minor axis. Moreover, the residuals of the fit show polar emission that extends south to west out to a distance of ~ 10 pc relative to the torus major axis (Figure 1(b)). The enhanced spatial resolution of the uniformly weighted data set better illustrates the goodness of the fit (Figure 1(c)).

Figure 2(a) shows the CO(6–5) intensity map of the CN obtained by integrating the line emission in the velocity interval $v - v_0 = [-250, +250] \text{ km s}^{-1}$. The CO emission from the east and west knots of the CN is comparatively more prominent than in the continuum image shown in Figure 1, which likely reflects the higher percentage of flux recovered for the line on scales $\geq 2''$.

Figure 2(b) shows the mean-velocity field of the gas in the CN. The gas kinematics derived from the new observations are compatible with the picture drawn from GB14: gas motions in the CN show a superposition of outward radial flows on rotation. The molecular outflow is responsible for tilting the kinematic major axis by $\simeq 70^\circ$ from $PA \simeq 260^\circ$ in the outer galaxy disk to $PA \simeq 330^\circ$ in the CN.

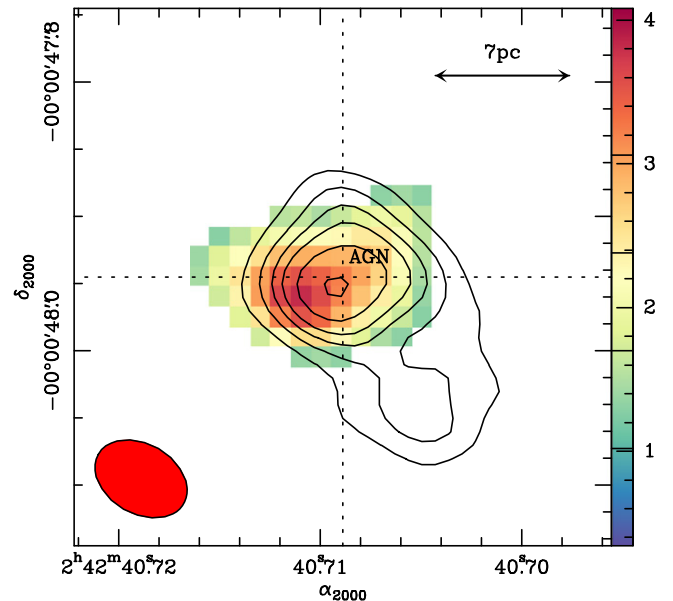


Figure 3. Overlay of the continuum emission contours of Figure 1 on the CO(6–5) emission (color scale) from the AGN torus. Units are in $\text{Jy km s}^{-1} \text{ beam}^{-1}$.

Furthermore, ALMA detects the CO emission from a spatially resolved lopsided disk located at the AGN, as illustrated in Figure 3. The size (diameter $\simeq 10 \pm 1$ pc, minor-to-major axis ratio $\simeq 0.5 \pm 0.1$) and orientation ($PA^{\text{major}} = 112^\circ \pm 20^\circ$) of the CO disk, hereafter referred to as the CO torus, are close to those derived above for the dust torus. The orientation of the CO torus is virtually identical to that of the radio-continuum source S1 resolved in the VLBA map of Gallimore et al. (2004). Figure 3 shows no significant CO counterpart for the polar emission of the dust torus.

(2.) *The mass of the torus.* From the continuum flux of the dust torus at $432 \mu\text{m}$, $\simeq 13.8 \pm 1 \text{ mJy}$, we estimate a dust mass of $M_{\text{dust}}^{\text{torus}} \sim 1600 M_\odot$, assuming that a maximum of $\sim 18\%$ of the flux at 694 GHz may come from other mechanisms different than thermal dust emission (GB14). To estimate $M_{\text{dust}}^{\text{torus}}$ we used a modified black-body function with a dust temperature $T_{\text{dust}} \simeq T_{\text{gas}} = 150 \text{ K}$, an emissivity index $\beta = 2$, and a value for the dust emissivity $\kappa_{694 \text{ GHz}} = 0.34 \text{ m}^2 \text{ kg}^{-1}$. The values of T_{dust} and $\kappa_{694 \text{ GHz}}$ are based on the analysis of the line excitation and SED fitting done in GB14 and Viti et al. (2014), and on the adoption of a value for $\kappa_{352 \text{ GHz}} = 0.0865 \text{ m}^2 \text{ kg}^{-1}$ (Klaas et al. 2001). Assuming that the gas-to-dust ratio in the central 2 kpc of NGC 1068 is $\simeq 60$ (GB14), we derive a molecular gas mass for the torus $M_{\text{gas}}^{\text{torus}} = (1 \pm 0.3) \times 10^5 M_\odot$.

4. CLUMPY TORUS MODELS

We combined the ALMA Band 9 continuum thermal flux with the NIR/MIR continuum and MIR interferometry data obtained by several groups in apertures $\leq 0''.05$ – $0''.3$ to construct the nuclear SED of NGC 1068 from $1.65 \mu\text{m}$ to $432 \mu\text{m}$ (Marco & Alloin 2000; Tomono et al. 2001; Weigelt et al. 2004; López-Gonzaga et al. 2014; this work). We also included as an upper limit the ALMA Cycle 0 Band 7 continuum flux from GB14. The 8–13 μm spectrum is the sum of components 1 and 2 in López-Gonzaga et al. (2014). We

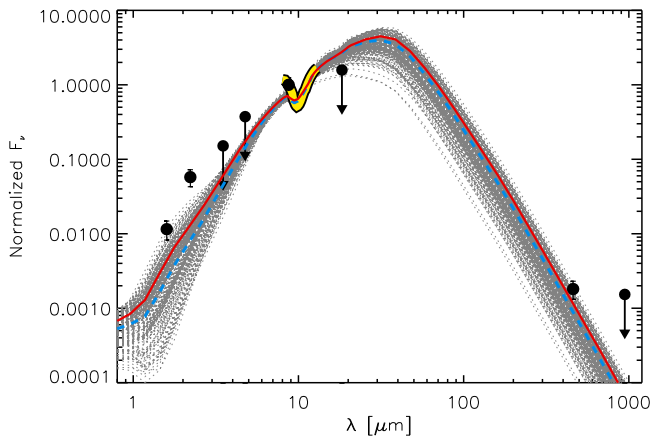


Figure 4. High spatial resolution SED of NGC 1068 (the thick yellow line represents the spectrum; the black dots represent the photometry) normalized to the $8.7 \mu\text{m}$ point. The solid red and dashed blue lines correspond to the maximum a posteriori (MAP) and median models, respectively, for the range $60^\circ < i < 90^\circ$. The upper limits correspond to the largest apertures ($\leq 0''.3$). Gray curves are the CLUMPY models sampled from the posterior distributions compatible with the data at the 68% confidence level.

note that this nuclear SED, shown in Figure 4, corresponds to the smallest spatial scales probed to date for NGC 1068.

We fitted the nuclear SED with the CLUMPY torus models of Nenkova et al. (2008a, 2008b), adopting a Bayesian approach using the `BayesClumpy` tool (Asensio Ramos & Ramos Almeida 2009). This scheme was applied to model the SED of other Seyferts (Alonso-Herrero et al. 2011; Ramos Almeida et al. 2011). Using this approach we can specify a priori information about the six model parameters. We restricted the torus angular width σ to be $40^\circ < \sigma < 60^\circ$ to match the opening angle of the ionization cones (80° ; Das et al. 2006). The torus radial extent, defined by the outer-to-inner radius ratio Y , is in the range $5 < Y < 30$ based on the torus size measured by ALMA at $432 \mu\text{m}$. In this first fit, we restricted the inclination angle of the torus to $60^\circ < i < 90^\circ$ based on the H_2O maser detection for this galaxy (GR96). We translate the results into the corresponding SED shown in Figure 4. The NIR fluxes are underpredicted, an indication that we need to include some contribution from starlight or a distinct component to reproduce the NIR bump (e.g., Kishimoto et al. 2012; Lira et al. 2013).

To compare our torus model with direct measurements from the ALMA observations presented here, we estimated from the fitted parameters the gas mass $M_{\text{gas}}^{\text{torus}}$, the outer radius R_{torus} , and the inclination i of the torus. Using the median values of the fitted parameters, we derive $M_{\text{gas}}^{\text{torus}} = 0.9 \pm_{0.4}^{0.2} \times 10^5 M_\odot$, $R_{\text{torus}} = 4 \pm 1 \text{ pc}$, and $i = 66^\circ \pm_4^9$. These values are consistent with those derived in Section 3 from the Band 9 observations alone. Nevertheless, if we allow the viewing angle of the torus to have smaller values in the range $30^\circ < i < 90^\circ$ (see Section 5), the best-fit median values are $M_{\text{gas}}^{\text{torus}} = 0.6 \pm_{0.2}^{0.4} \times 10^5 M_\odot$, $R_{\text{torus}} = 2 \pm 1 \text{ pc}$, and $i = 34^\circ \pm_5^5$.

5. KINEMATICS OF THE TORUS

Figure 5 shows that ALMA has spatially resolved the molecular gas kinematics in the torus. Although the mean velocities derived from CO, shown in Figure 5(a), span a moderate range ($v - v_0 = [-30, +30] \text{ km s}^{-1}$), we identify spatially resolved velocity gradients within the torus. Overall,

the line emission appears redshifted to the south and blueshifted to the north. Therefore, the apparent *kinematic* major axis of the CO disk, at $PA^{\text{kin}} \simeq 180^\circ$, is tilted by a significant angle ($\simeq 53^\circ \pm 15^\circ$) relative to the morphological major axis of the CO/dust torus, oriented at a weighted average $PA^{\text{major}} \simeq 127^\circ \pm 15^\circ$ (see Section 3). Taken at face value, the tilt between both axes indicates that strong non-circular motions are superposed to the slow rotation pattern of the CO torus.

Figure 5(b) shows the position–velocity (p–v) diagram taken along the morphological major axis of the CO torus ($PA^M = 112^\circ \pm 20^\circ$). The bulk of the CO emission in the torus spans a wide range of velocities: $v - v_0 \simeq [-80, 80] \text{ km s}^{-1}$. Figure 5(b) compares the observed kinematics with the velocity range expected from the H_2O sub-Keplerian rotation curve of Gallimore et al. (1996) projected along the CO major axis ($v_{\text{rot}} \propto v^{-\alpha}$, with $\alpha = 0.31$; see also Lodato & Bertin 2003). Even in the most favorable case of a purely Keplerian extrapolation ($\alpha = 0.5$), the CO velocities are still a factor ≥ 2 lower compared with the H_2O -based kinematics for an edge-on disk ($i = 90^\circ$). To reach a better agreement, we would need to incline the CO torus down to $i = 30^\circ\text{--}40^\circ$, i.e., close to the inclination of the galaxy disk (GB14). Although these small values for i can be discarded for the mostly edge-on maser disk (at $r \leq 1 \text{ pc}$), we cannot rule out that the CO torus is tilted toward a more face-on orientation (at $r \geq 2\text{--}4 \text{ pc}$). Furthermore, this scenario is still within the range of solutions found in Section 4 for the clumpy torus. However, as argued below, the intricate kinematics of the CO torus cannot be attributed alone to the changing orientation of a purely rotating disk.

Figures 5(a), (b) show that the CO emission is redshifted on the southeast side of the major axis at $\Delta x \simeq -0''.06$. The emission is nevertheless blueshifted northwest at $\Delta x \simeq +0''.06$. This velocity gradient, measured on scales $\Delta x \simeq \pm 4 \text{ pc}$ from the AGN, is inverted relative to that measured for the H_2O maser disk at $\Delta x \simeq \pm 1 \text{ pc}$. The apparent counterrotation of the outer disk relative to the inner disk could be caused by an enhancement of the gas turbulence in the torus. A rough estimate of the ratio of the deprojected mean velocity (v_{mean}) to the velocity dispersion (σ) of the gas along the major axis yields very low values: $v_{\text{mean}}/\sigma < 1$. However, Figure 5(c) shows that a comparable spatially resolved velocity gradient is measured along the minor axis p–v diagram. The tilted disk scenario or the simple addition of turbulence to rotation cannot alone explain the observed pattern along the minor axis, an indication that non-circular motions in the torus need to be invoked.

Papaloizou & Pringle (1984) first predicted that non self-gravitating AGN tori should likely develop low-order non-axisymmetric instabilities on a dynamical timescale. This type of instability, known as the Papaloizou–Pringle instability (PPI), has been studied in isolated self-gravitating tori (Kiuchi et al. 2011; Korobkin et al. 2013) and also in tori perturbed by the accretion of matter coming from the outer boundary of the system (Dönmez 2014). These new simulations show the growth of a long-lasting $m = 1$ mode related to the onset of the PPI. The $m = 1$ mode generates a lopsided distribution, as well as enhanced turbulence and strong non-circular motions in the torus gas. Evidence of inflowing gas connecting the CND with the torus has been found in the $2.12 \mu\text{m}$ H_2 lines by Müller Sánchez et al. (2009). The remarkably lopsided distribution and

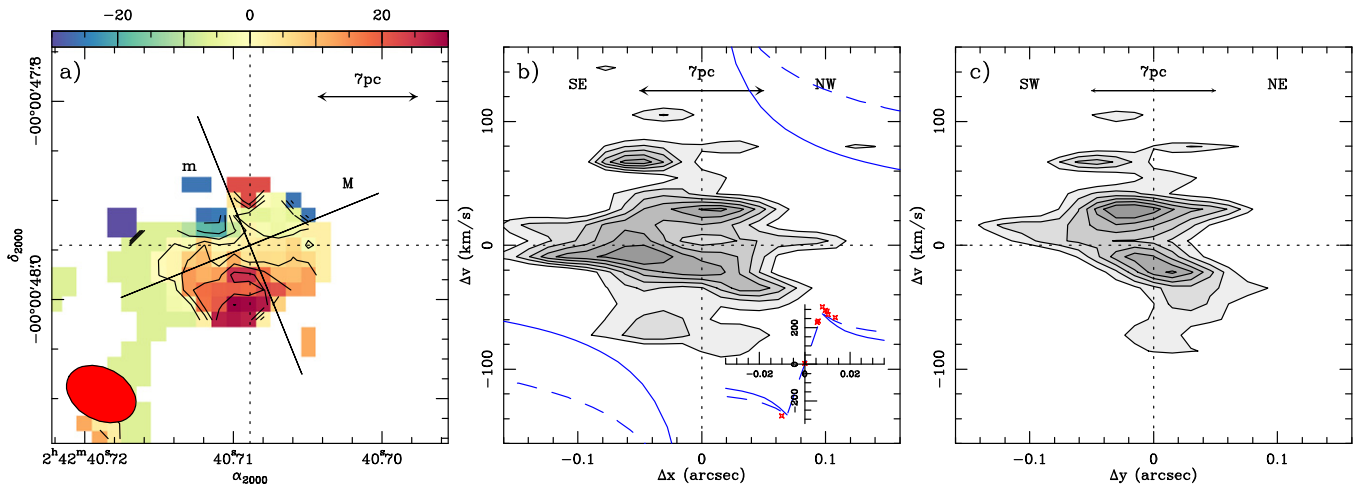


Figure 5. (a) CO(6–5) mean-velocity map of the torus. The contours span the range (-30 km s^{-1} to 30 km s^{-1}) in steps of 10 km s^{-1} . The black lines labeled as M and m show, respectively, the orientations of the major and minor axes of the CO torus ($PA^M = 112^\circ \pm 20^\circ$, $PA^m = 22^\circ \pm 5^\circ$). (b) CO(6–5) position–velocity diagram along the M axis. Contours go from 2σ to 7σ in steps of 1σ , where $1\sigma = 4 \text{ mJy}^{-1} \text{ beam}^{-1}$. The inset shows the velocities relative to v_{sys} as a function of radius (in arcseconds) as derived for the H_2O megamaser spots (red markers) detected along $PA^{\text{maser}} = 140^\circ \pm 5^\circ$ (GR96, GA01). The dashed (solid) blue curve shows the best-fit sub-Keplerian (Keplerian) rotation curve $v_{\text{rot}} \propto r^{-\alpha}$ of GR96 with $\alpha = 0.31$ (0.50). (c) Same as middle panel, but for the m axis. In all panels, velocities refer to $v_{\text{o}}(\text{HEL}) = 1136 \text{ km s}^{-1}$ and $(\Delta x, \Delta y)$ —offsets are relative to the AGN position.

complex kinematics of the ~ 7 – 10 pc CO torus of NGC 1068 seen by ALMA could be the footprint of the PPI in this Seyfert. Alternatively, the velocity gradient observed along the minor axis could be explained by gas being entrained in the outflow, as similarly detected by Cecil et al. (2002) in high-ionization lines close to the AGN. This radial shift is noticeably reversed farther out at the CND where the CO outflow follows the pattern of low-ionization lines (GB14).

6. SUMMARY AND CONCLUSIONS

We mapped with ALMA the dust continuum at $432 \mu\text{m}$ and the CO(6–5) molecular line emission in the CND of NGC 1068 with a spatial resolution of $\sim 4 \text{ pc}$. These observations allowed us to resolve the CND and image the dust emission and also the distribution and kinematics of molecular gas from a 7 – 10 pc diameter disk. This is the first detection of a submillimeter counterpart of the putative torus of NGC 1068. As expected for the contribution of a comparatively cooler component, the CO/dust torus extends on spatial scales that are twice larger compared with the brightest MIR sources detected by the VLTI at the position of the central engine of NGC 1068 (Burtscher et al. 2013; López-Gonzaga et al. 2014).

The ALMA measurements have been combined with the NIR/MIR continuum and MIR interferometry data available from $1.65 \mu\text{m}$ to $432 \mu\text{m}$. We fitted the nuclear SED with the CLUMPY torus models of Nenkova et al. (2008a, 2008b). The mass and the radius of the best-fit torus are consistent with the values derived from the ALMA data alone: $M_{\text{gas}}^{\text{torus}} = (1 \pm 0.3) \times 10^5 M_{\odot}$ and $R_{\text{torus}} = 3.5 \pm 0.5 \text{ pc}$.

We spatially resolved the kinematics of molecular gas in the torus. The overall slow rotation pattern of the disk is perturbed by strong non-circular motions and enhanced turbulence. Furthermore, while the maser condition of the H_2O line suggests that the inner molecular disk should be oriented mostly edge-on at $r \leq 1 \text{ pc}$, we found evidence that the molecular torus could be tilted toward lower inclination angles ($i = 34^\circ$ – 66°) at $r \geq 2$ – 4 pc . The lopsided morphology and complex kinematics of the NGC 1068 torus could be the signature of the PPI instability predicted by Papaloizou &

Pringle (1984). This instability was foreseen to rule the dynamical evolution of AGN tori.

The results of this work are relevant to caution against the use of CO lines as a direct tracer of the supermassive black hole mass (M_{BH}) in galaxies (e.g., Davis et al. 2013). In cases like NGC 1068 where even mid-J CO lines trace the outer boundary of a kinematically distorted tilted torus, careful modeling would be mandatory to reliably estimate M_{BH} .

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