

What are the megahertz peaked-spectrum sources?

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ABSTRACT

Megahertz peaked-spectrum (MPS) sources have spectra that peak at frequencies below 1 GHz in the observer’s frame and are believed to be radio-loud active galactic nuclei (AGN). We recently presented a new method to search for high-redshift AGN by identifying unusually compact MPS sources. In this paper, we present European VLBI Network (EVN) observations of 11 MPS sources which we use to determine their sizes and investigate the nature of the sources with ~ 10 mas resolution. Of the 11 sources, we detect 9 with the EVN. Combining the EVN observations with spectral and redshift information, we show that the detected sources are all AGN with linear sizes smaller than 1.1 kpc and are likely young. This shows that low-frequency colour–colour diagrams are an easy and efficient way of selecting small AGN and explains our high detection fraction (82 per cent) in comparison to comparable surveys. Finally we argue that the detected sources are all likely compact symmetric objects and that none of the sources are blazars.

Key words: techniques: high angular resolution – techniques: interferometric – galaxies: active – galaxies: high-redshift – radio continuum: galaxies.

1 INTRODUCTION

AGN jets are powered by the accretion of material from their host galaxy on to a supermassive black hole (e.g. Blandford & Königl 1979; Falcke & Biermann 1995) and can grow to extend well beyond their host galaxy. The young jets can be distorted and even stopped by the ambient medium, while larger jets heat both the interstellar and intergalactic medium, quench star formation and expel material from the galaxy (e.g. Morganti et al. 2013). Hence to understand AGN, we need to understand galaxies and vice versa (e.g. Fabian 2012).

Young or restarted AGN can be used to study how AGN are launched and evolve from parsec-scale objects to sources of hundreds of kiloparsec such as Cygnus A and 3C175 (e.g. Snellen et al. 2000). AGN also act as beacons allowing us to observe sources out to $z > 7$ (e.g. Mortlock et al. 2011) and study how the Universe evolved from a time when it was less than 6 per cent of its current age. Specifically, the fraction of jets that are frustrated by their

host galaxy appears to increase with redshift (van Velzen, Falcke & Körding 2015). By comparing the number of young and small AGN at high redshifts to those in the modern Universe, we can therefore trace the evolution of the ambient medium which both feeds AGN and hampers or even confines their jets (e.g. Falcke, Körding & Nagar 2004). Hence, searching for both young AGN and AGN at high redshifts is critically important to understand what triggers the nuclear activity, how AGN evolve in size, how the population evolves with redshift and, ultimately, the origin of their redshift evolution.

Compact steep-spectrum (CSS), gigahertz peaked-spectrum (GPS) and high-frequency peaked (HFP) sources are all radio-loud AGN that are identified based on their spectral energy distribution (SED) in the radio. CSS, GPS and HFP sources are characterized by steep optically thin spectra that turn over and have inverted spectra (the spectral index, α , is defined as $S \propto \nu^\alpha$, where S is the flux density and ν is the frequency) above the turnover frequency. The CSS sources have typical rest-frame turnover frequencies (ν_r) smaller than 500 MHz and largest linear sizes (LLS) of 1–20 kpc (O’Dea 1998). For the GPS sources, $1 < \nu_r < 5$ GHz, while $\nu_r > 5$ GHz for the HFP sources (Dallacasa et al. 2000). Both the GPS and HFP sources have LLS < 1 kpc (O’Dea 1998).

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Morphologically GPS and HFP sources are typically classified as compact symmetric objects (CSOs) while the CSS sources are medium-size symmetric objects (MSOs; e.g. Snellen et al. 2000; Conway 2002). CSOs and MSOs are characterized by unbeamed emission from their steep-spectrum radio lobes on either side of a central position and have sizes smaller than their host galaxy (Fanti et al. 1995; Fanti 2009). To strictly classify a source as a CSO or MSO, its flat-spectrum core has to be detected (Oriente & Dallacasa 2014), which is often not the case. In addition, unlike their names suggest, CSOs and MSOs are often not symmetric around their cores. This is likely the result of the interaction of the jet with an inhomogeneous ambient medium (e.g. O’Dea 1998; Oriente & Dallacasa 2014).

Based on spectral and kinematic age estimates (determined by fitting models to the source spectra and measuring lobe expansion speeds), most of the HFP, GPS and CSS sources are believed to be young ($\lesssim 10^5$ yr) and small ($\lesssim 20$ kpc) AGN rather than being sources that are confined by the ambient medium of their host galaxy (O’Dea 1998; Conway 2002; Murgia et al. 2002; Murgia 2003; Fanti 2009). In addition, an empirical relation was found between ν_r and LLS that spans three orders of magnitude (see Section 5.1, O’Dea 1998; Snellen et al. 2000; Oriente & Dallacasa 2014). This shows that the smallest sources have the highest turnover frequencies. Based on this evidence, it is believed that the HFP sources evolve into GPS sources which in turn evolve into CSS sources (O’Dea 1998; Snellen et al. 2000; Tschager et al. 2003). There is also strong evidence from their expected luminosity evolution and the similarities between their host galaxies (Begelman 1996; Snellen et al. 2000; De Vries, O’Dea & Barthel 2002a) that the CSS sources will evolve into the FR I and FR II radio galaxies (Fanaroff & Riley 1974).

In the past, searches for high-redshift AGN in the radio have focused on ultrasteep-spectrum (USS) sources (e.g. Jarvis et al. 2001; Cruz et al. 2006; De Breuck et al. 2006). However, the reason why USS sources should be at higher redshifts than non-USS sources remains unclear (Miley & De Breuck 2008). Moreover, several recent studies have found that USS sources are not at higher redshifts than non-USS sources (Ker et al. 2012; Singh et al. 2014; Smolčić et al. 2014).

In our previous paper (Coppejans et al. 2015), we described a new method of searching for high-redshift radio-loud AGN by selecting compact megahertz peaked-spectrum (MPS; turnover frequency below 1 GHz) sources. The MPS sources are believed to be a mixture of nearby CSS sources, and smaller GPS and HFP sources whose spectral turnovers have been redshifted to lower frequencies. Hence, the most compact MPS sources should be at the highest redshifts. In Coppejans et al. (2015), we took the first steps in testing the method by making a low-frequency radio colour–colour diagram and selecting a sample of 33 MPS sources from it. Using their photometric redshifts, we concluded that there is encouraging evidence that the MPS method can be used to search for high-redshift AGN. This was the first time that a colour–colour diagram was used to select MPS sources. However, it will soon be possible to repeat the analysis over the full sky using instruments such as the Low-Frequency Array (LOFAR; Van Haarlem et al. 2013). We therefore wish to confirm that this novel selection method yields a separate class of small and likely young AGN, where the most compact sources are at high redshifts.

Here we present very long baseline interferometry (VLBI) observations of 11 MPS sources conducted with the European VLBI Network (EVN). Combining the EVN’s sub-arcsecond resolution with the spectra of the sources, we investigate the nature of the MPS

sources and test the hypothesis that these are AGN with small jets. In Section 2, we describe how we selected the sources and reduced the EVN data. Section 3 describes how the source spectra were generated, presents the source properties derived from the images and spectra, and discusses whether the radio emission is from star formation in the host galaxy or AGN activity. The individual sources are discussed in detail in Section 4. In Section 5, we discuss what the MPS sources are and present a summary in Section 6. Throughout this paper, we use the following cosmological parameters: $\Omega_m = 0.3$, $\Omega_\lambda = 0.7$, $H_0 = 72 \text{ km s}^{-1} \text{ Mpc}^{-1}$.

2 TARGET SELECTION, OBSERVATIONS AND DATA REDUCTION

2.1 Target selection

In Coppejans et al. (2015), we matched the sources in our 324.5 MHz Very Large Array (VLA) P-band image (hereafter referred to as the VLA-P image) of the National Optical Astronomy Observatory (NOAO) deep wide-field survey Boötes field to the VLA Faint Images of the Radio Sky at Twenty-Centimetres (FIRST) survey (White et al. 1997) and a 153 MHz Giant Metrewave Radio Telescope (GMRT) catalogue of the field (Williams, Intema & Röttgering 2013, hereafter referred to as the WIR catalogue). From this we generated a colour–colour diagram of the field and selected 33 MPS sources that either show a turnover in their spectra or a significant low-frequency flattening, which could indicate a turnover below 153 MHz. Sources were excluded if they were extended in the FIRST or VLA-P catalogues or had a flux density difference of more than 20 per cent between any two of the following three 1.4 GHz catalogues: FIRST, National Radio Astronomy Observatory (NRAO) VLA Sky Survey (NVSS; Condon et al. 1998) and De Vries et al. (2002b, hereafter referred to as the dVMR catalogue). Since the resolution of the FIRST, dVMR and NVSS catalogues are 5.4, 20 and 45 arcsec, respectively (see Table 3), this not only removed variable sources, but also sources with extended structures that are resolved out in one of the higher resolution catalogues.

The MPS sources presented in this paper are given in Table 1 and were observed with the EVN during two projects, EV020 and EC053. In the table the columns are (1) source name, (2) the EVN project code under which the source was observed, (3) low-frequency spectral index calculated between 153 and 325 MHz, (4) high-frequency spectral index calculated between 325 and 1400 MHz, (5,6) photometric redshifts calculated using the EAZY code from Brammer, van Dokkum & Coppi (2008) and LRT code from Assef et al. (2008) as described in Coppejans et al. (2015), and (7) the total time spent observing the source with the EVN. A description of how α_{low} and α_{high} were calculated are given in Section 3.1. Note that for the redshift values, the LRT code does not provide errors and includes an empirical AGN SED template in the fitting, and should therefore fit AGN spectra better. In the table, there are three sources (J142850+345420, J143718+364549 and J144230+355735) for which we do not have photometric redshifts. J142850+345420’s optical counterpart is too faint for us to find a photometric redshift for it, while J143718+364549 and J144230+355735 lie outside the multiwavelength coverage of the Boötes field.

The sources that were observed with the EVN were selected to be unresolved in FIRST, non-variable, have the highest possible flux density in FIRST and have not been previously observed with VLBI. Of the 11 MPS sources observed with the EVN, four are also in the selection of sources in Coppejans et al. (2015). Two

Table 1. Basic parameters of the target sources.

Source ID	Project code	α_{low}^a	α_{high}^b	z_{EAZY}	z_{LRT}	Time on source [min]
(1)	(2)	(3)	(4)	(5)	(6)	(7)
J142850+345420	EV020	0.6 ± 0.4	-0.5 ± 0.1	–	–	45
J142904+354425	EC053	0.0 ± 0.5	-0.6 ± 0.1	$0.809^{+0.084}_{-0.081}$	0.84	120
J142917+332626	EC053	0.6 ± 0.4	-0.6 ± 0.1	$1.583^{+0.322}_{-0.290}$	2.49	70
J143024+352438	EV020	0.2 ± 0.6	-0.7 ± 0.1	$1.196^{+0.116}_{-0.118}$	1.33	200
J143042+351240	EC053	0.0 ± 0.4	-0.7 ± 0.1	$1.281^{+0.202}_{-0.217}$	1.13	200
J143050+342614	EV020	0.5 ± 0.3	-0.6 ± 0.0	$2.364^{+0.535}_{-0.536}$	2.98	45
J143055+350852	EC053	-0.1 ± 0.3	-0.8 ± 0.1	$2.195^{+0.379}_{-0.397}$	0.38	60
J143213+350940	EV020	-0.1 ± 0.3	-0.3 ± 0.1	$0.978^{+0.100}_{-0.095}$	0.96	45
J143329+355042	EC053	0.2 ± 0.5	-0.5 ± 0.1	$2.821^{+2.032}_{-1.536}$	1.37	60
J143718+364549	EV020	0.2 ± 0.5^c	-0.7 ± 0.1^c	–	–	45
J144230+355735	EV020	0.1 ± 0.7^c	-1.0 ± 0.2^c	–	–	125

^a α_{low} was calculated between 153 and 325 MHz. ^b α_{high} was calculated between 325 and 1400 MHz. ^cThe spectral indices of the sources were calculated using their WENSS flux densities since they lie outside the area imaged with the VLA-P data (see Section 3.1).

Table 2. Telescope participation in each project.

Radio dish	EC053 ^a	EV020 ^a
Effelsberg	Yes	Yes
Jodrell Bank	Yes	Yes
Medicina	No	No
Noto	No	Yes
Onsala	Yes	Yes
Toruń	Yes	Yes
Sheshan	No	Yes
WSRT	Yes	Yes

^a‘Yes’ indicates that the telescope provided useful data while ‘No’ indicates that it did not.

of the new sources (J143718+364549 and J144230+355735) lie outside the region that was imaged with the VLA-P data and were selected based on their Westerbork Northern Sky Survey (WENSS; Rengelink et al. 1997) flux densities (see Section 3.1). The remaining five new sources (J142850+345420, J143024+352438, J143042+351240, J143055+350852 and J143213+350940) were originally excluded in Copejans et al. (2015) because they have a flux density difference of more than 20 per cent between either FIRST and dVMR or NVSS and dVMR. In Section 3.4, we argue that the flux density difference with dVMR catalogue does not necessarily indicate that these sources are variable. We therefore believe that all five sources are genuine MPS sources.

2.2 Observing setup and data reduction

EV020 and EC053 were observed on 2014 April 15 and 2015 January 14, respectively. Since $\alpha_{\text{high}} < -0.5$ for the MPS sources, we elected to do the observations at 1.664 GHz to maximize the flux density of the sources and reduce the required observing time. During both projects we requested the targets to be observed with the radio telescopes at Effelsberg (Germany), Jodrell Bank (Mk2; UK), Medicina (Italy), Noto (Italy), Onsala (Sweden), Toruń (Poland), Sheshan (China) and the Westerbork Synthesis Radio Telescope (WSRT, the Netherlands). A list of the telescopes and whether or not they successfully participated in each project is given in Table 2. Since the baselines to Sheshan (which did not take part in EC053) form the longest baselines, the typical restoring beam size

of EC053 is 26×32 mas compared to the 3×10 mas of EV020. Both projects obtained data with 2 s integrations at 1024 Mbit s^{-1} in left and right circular polarizations with eight sub-bands per polarization and 16 MHz of bandwidth per sub-band. The technique of electronic VLBI (e-VLBI) was used, where the data are not recorded at the telescopes but streamed to the central correlator using optical fibre networks in real time. The observations of the targets were interleaved with observations of two phase calibrators, J1430+3649 and J1422+3223. The phase solutions from J1422+3223 were used to correct J142917+332626, which is $1^\circ 8'$ away from J1422+3223. The solutions of J1430+3649 were used to correct the remaining sources which are separated from it by between $0^\circ 5'$ and $2^\circ 5'$.

The data were reduced using the AIPS (Greisen 1990) software package by calibrating the visibility amplitudes using antenna gains and system temperatures measured at the telescopes. Next fringe-fitting was performed on the two phase calibrators. The phase calibrators were then imaged in the Caltech DIFMAP package (Shepherd, Pearson & Taylor 1994) by doing several iterations of CLEAN and phase self-calibration. A final round of amplitude self-calibration was done on the phase calibrators in DIFMAP to determine the global antenna gain correction factors. The gain correction factors varied between one and five per cent and were applied to the visibility amplitudes of all the sources in AIPS. Using the clean component models derived for the phase calibrators in DIFMAP, improved phase solutions were calculated for the phase calibrators in AIPS. These solutions were applied to the target sources before they were exported from AIPS for flagging and imaging in DIFMAP. To check that we did not miss any source components, or that any of the sources were significantly offset from the phase centre, we started off by making images that were at least 5×5 arcsec in size. We then cleaned the identified components in smaller images using uniform weighting to get the best possible position accuracy for the components before switching to natural weighting. Since the target sources have flux densities of only a few mJy, we did not self-calibrate them. We finally imaged all of the sources using a uv-taper with a Gaussian value of 0.1 and a Gaussian radius of 15 million wavelengths ($M\lambda$). The uv-taper downweights the visibilities on baselines to Sheshan, where the uv-plane is sampled the least. This decreases both the resolution and noise of the image, allowing for the detection of diffuse emission around the source.

Table 3. Catalogues that were matched to the target sources.

Catalogue	Frequency	Resolution	Image noise ^a	Positive matches
name	[MHz]	[arcsec]	[mJy beam ⁻¹]	
FIRST	1400	5.4 × 5.4	0.15	All
NVSS	1400	45 × 45	0.45	All except J142917+332626
dVMR	1380	13 × 27	0.03	J142850+345420, J143024+352438, J143042+351240, J143050+342614, J143055+350852 and J143213+350940
GMRT-608	608	5.0 × 5.0	0.04	J142850+345420, J142904+354425, J143024+352438, J143042+351240, J143050+342614 and J143055+350852
VLA-P	325	5.1 × 5.6	0.2*	All except J143718+364549 and J144230+355735
WENSS	325	54 × 54	3.6	J143050+342614, J143055+350852, J143213+350940, J143718+364549 and J144230+355735
WIR	153	25 × 25	2.0*	All
LOFAR-150	150	5.6 × 7.4	0.11*	All except J143718+364549 and J144230+355735
LOFAR-62	62	19 × 31	4.8	None

^aTypical catalogue noise values are quoted except for those marked with *, where the noise is measured at the centre of the image.

3 RESULTS AND GENERAL DISCUSSION

In this section we will first discuss the catalogues to which the sources were matched, before presenting the results derived from the EVN observations and the spectra of the sources. This will be followed by a discussion of the 1.4 GHz variability of the sources and the cause of the radio emission.

3.1 Matched catalogues

Table 3 contains a list of all the radio catalogues and images to which the sources were matched to constrain their radio spectra. The source matching was done using the method described in Section 3.2 of Coppejans et al. (2015). In the final column of Table 3, a list of the sources with which a positive match were found is given for each of the catalogues.

The GMRT-608 image referenced in Table 3 is a mosaic of a part of the Boötes field (at 608 MHz). A mosaic of the entire field will be published once the observations have been completed. The image was made from GMRT observations of a part of the Boötes field (project code 28_064). Four pointings covering 1.95 deg² were observed on 2015 July 24 and 26. Raw visibilities were recorded every eight seconds in two polarizations (RR and LL), using 512 frequency channels to cover 32.0 MHz of bandwidth centred on 608 MHz. The on-target time for each pointing was between 100 and 110 min. Primary flux density calibration was done with 3C286 using the wide-band low-frequency flux density standard of Scaife & Heald (2012). The data reduction follows that of De Gasperin et al. (2014) and Bonafede et al. (2014) and was done in three stages: (i) initial gain and bandpass calibration, (ii) self-calibration, and (iii) direction-dependent ionospheric phase calibration using the software package SPAM (Intema et al. 2009). The combined final mosaic reaches a root mean square (rms) noise level of $\approx 40\text{--}70 \mu\text{Jy beam}^{-1}$ with a resolution of 5×5 arcsec. We have checked the consistency of the flux density scale by interpolating between the WIR and dVMR catalogues.

The 150 MHz LOFAR-150 catalogue was made from a new LOFAR survey of the Boötes field (Williams et al., submitted). The image has a resolution of ~ 6 arcsec with half of the 19 deg² image having a local rms noise below $0.18 \text{ mJy beam}^{-1}$, both of which are better than the WIR image. The catalogue itself contains 5652 sources detected above a threshold of 5σ (Williams

Table 4. 62 MHz detection threshold for each source in the LOFAR-62 catalogue.

Source ID	Detection threshold [mJy]
J142850+345420	28.7
J142904+354425	43.8
J142917+332626	75.3
J143024+352438	46.8
J143042+351240	47.2
J143050+342614	36.2
J143055+350852	30.4
J143213+350940	32.2
J143329+355042	36.1
J143718+364549	50.9
J144230+355735	68.1

et al., submitted). After matching our sources to the catalogue we found that all but one of the LOFAR-150 flux densities were higher than the corresponding WIR flux densities, despite the LOFAR-150 catalogue having a higher resolution. Specifically, the integrated LOFAR-150 flux densities for our sources have a median difference of 27 per cent compared to those of the WIR catalogue. Since the LOFAR-150 flux density scale was checked and corrected using the flux densities of the sources in the WIR catalogue (Williams et al., submitted), we elected to use the WIR flux densities when fitting the source spectra (Section 3.2). We do however discuss (Section 4) and show (Fig. 2) the LOFAR-150 flux densities for each of the sources in their spectral plots. We finally note that J143718+364549 and J144230+355735 fall outside the area imaged by the LOFAR-150 catalogue.

Finally, the LOFAR-62 catalogue (Van Weeren et al. 2014) is a catalogue constructed from LOFAR Low Band Antenna (LBA) commissioning observations at 62 MHz of the Boötes field. While all our sources lie in the image, none of them are detected at the catalogue's 5σ detection threshold. To determine the detection threshold for each source (given in Table 4), we measured the local rms noise in a 320×320 arcsec box, as was done by Van Weeren et al. (2014), centred on each of the sources' positions and multiplied it by five. We used the local noise, rather than the typical noise of the catalogue (given in Table 3), since the noise at the position of each of our sources will likely differ from the typical noise. This helped

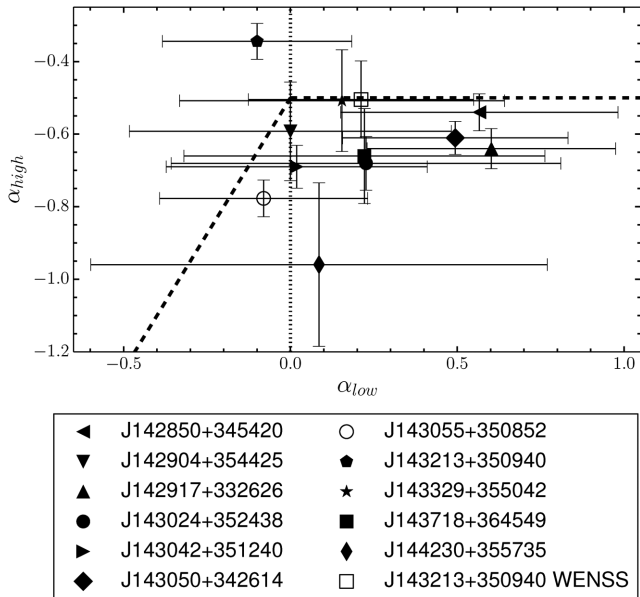


Figure 1. A colour–colour diagram of the sources. In the plot, α_{low} and α_{high} were calculated as explained in Section 3.1. The dotted vertical line indicates a spectral index of zero while the MPS sources were selected from the region below and to the right of the dashed line.

us to constrain the spectra of some of the sources (see Fig. 2 and Section 4).

In Fig. 1, the colour–colour diagram for the sources is shown. For all the sources, except J143718+364549, J144230+355735 and J143213+350940 WENSS, α_{low} was calculated between the flux densities of the WIR and VLA-P catalogues. α_{high} was calculated between those of the VLA-P and FIRST catalogues. Since J143718+364549 and J144230+355735 are not in the VLA-P image, α_{low} and α_{high} were calculated from their WENSS flux densities. Since the VLA-P and WENSS flux densities differ significantly for J143213+350940, we also plotted J143213+350940’s position in Fig. 1 if α_{low} and α_{high} are calculated using its WENSS flux density rather than the VLA-P flux density. This will be discussed in detail in Section 4.3.3. The MPS sources in Coppejans et al. (2015) were selected from the area below and to the right of the dotted lines in Fig. 1, which are defined by $\alpha_{high} < -0.5$ and $\alpha_{high} < 1.5\alpha_{low} - 0.5$. The first constraint ensures that the sources have a steep spectrum above 1 GHz. The second allows us to not only select the sources with a clear spectral peak, but also sources whose spectra flatten towards lower frequencies, which could indicate a spectral turnover below 153 MHz. We note that in Fig. 1, J143213+350940 does not satisfy the selection criteria, while J143213+350940 WENSS does (Section 4.3.3).

3.2 Source properties

In Table 5, the parameters derived for the sources are presented with the sources components named as the source name followed by a letter. These components are shown in the EVN images presented in Figs 3–7.

Columns (2), (3) and (4) in Table 5 are the rms noise, the right ascension (RA) and declination (DEC), respectively, of each of the components of the source in the EVN image. The uncertainty of the RA and DEC are given in brackets after the value. The uncertainties were calculated using the equation given in Fomalont (1999), to this we added the uncertainty of the position of the phase calibrator from

the VLBA calibrator list¹ (0.14 mas and 0.11 mas for J1422+3223 and J1430+3649, respectively), in quadrature.

Column (5) gives the EVN integrated flux density at 1.7 GHz. The values were determined by fitting circular Gaussian brightness distribution models in DIFMAP to all of the sources and source components except J143050+342614, J143213+350940, J143329+355042 and J144230+355735b. These sources were fitted with elliptical Gaussians since the fit did not converge when fitting circular Gaussians or the circular fit clearly does not describe the flux density distribution of the source. Since DIFMAP does not report an error on the integrated flux density, the errors were calculated using the equations in Fomalont (1999) and adding an additional five per cent to account for the VLBI amplitude calibration uncertainty, as done by e.g. Frey et al. (2015) and An et al. (2012). The integrated flux densities for all the multicomponent sources except J143213+350940 and J144230+355735 are the sum of the individual components where the flux density errors were calculated by adding the errors of the individual components in quadrature. See Sections 4.3.3 and 4.3.5 on how the values for J143213+350940 and J144230+355735 were calculated.

Columns (6) and (7) contains the minor- and major-axis full width at half-maximum (FWHM) of the Gaussians fitted to the sources. The errors were calculated using the equations in Fomalont (1999). For the sources which were fitted with a circular Gaussian, the values in columns (6) and (7) are the same. For all of the sources, the values in column (7) were used as the source size. If the source was resolved into multiple components, the size was determined by calculating the distance between the centres of the two components that are the furthest apart, taking into account the uncertainties of the central positions. Note, however, that while J143213+350940 is composed of multiple components, we fitted both components simultaneously with a single elliptical Gaussian (see Section 4.3.3). Hence the values reported for J143213+350940 are the minor- and major-axis of the fitted Gaussian.

Column (8) gives the percentage of the predicted flux density that was recovered from the image. The value was calculated using $100S_i/S_{predicted}$, where S_i is the integrated EVN flux density of the sources given in column (5) and $S_{predicted}$ is the sources’ predicted flux density at 1.7 GHz. $S_{predicted}$ was calculated using α_{high} from Table 1 in combination with the equation $S_{predicted} = kv^\alpha$. The constant k was calculated for each source using its integrated FIRST flux density. The errors of the values in Column (8) were calculated by propagating the errors of S_i and $S_{predicted}$.

The redshift-corrected brightness temperatures of the sources in column (9) were calculated using

$$T_b = 1.22 \times 10^{12} (1+z) \frac{S_i}{\theta_1 \theta_2 \nu^2} \quad (1)$$

(Condon et al. 1982). Here, z is the redshift, S_i is the integrated flux density in Jy, θ_1 and θ_2 are the major- and minor-axis of the Gaussian fitted to the source in mas, and ν is the observing frequency in GHz. If the source component was fitted with a circular Gaussian, $\theta_1 = \theta_2$. Since we have two photometric redshifts for each source, we opted to calculate two values for each source for the relevant parameters in Table 5. Since the upper and lower uncertainties of z_{EAZY} are not symmetrical, we used the larger of the two as the uncertainty to calculate the errors reported for the relevant parameters. For the LRT code, which does not report an uncertainty on the redshift, we used an uncertainty of zero. Finally, to get robust lower limits for T_b for

¹ <http://www.vlba.nrao.edu/astro/calib/>

Table 5. Derived parameters of the sources.

Source ID	Noise [$\mu\text{Jy beam}^{-1}$]	RA ^a [J2000]	DEC ^b [J2000]	S _r ^c [mJy]	Min [mas]	Maj [mas]
(1)	(2)	(3)	(4)	(5)	(6)	(7)
J142850+345420	250	–	–	8.16 ± 0.77	–	47.6 ± 0.7 ^g
J142850+345420a	250	14:28:50.46588(0.00007)	34:54:20.8346(0.0011)	6.25 ± 0.61	29.9 ± 1.7	29.9 ± 1.7
J142850+345420b	250	14:28:50.46894(0.00015)	34:54:20.8206(0.0022)	1.91 ± 0.47	21.9 ± 1.2	21.9 ± 1.2
J142904+354425	23	14:29:04.6 ^h	35:44:25.1 ^h	–	–	–
J142917+332626	170	14:29:17.42200(0.00003)	33:26:26.6001(0.0004)	3.47 ± 0.32	13.0 ± 0.5	13.0 ± 0.5
J143024+352438	11	14:30:24.3 ^h	35:24:38.1 ^h	–	–	–
J143042+351240	18	14:30:42.57796(0.00004)	35:12:40.6253(0.0006)	0.20 ± 0.03	12.1 ± 1.5	12.1 ± 1.5
J143050+342614	198	14:30:50.90925(0.00003)	34:26:14.1890(0.0004)	3.53 ± 0.50	3.0 ± 0.4	15.4 ± 1.8
J143055+350852	88	–	–	5.81 ± 0.37	–	137.9 ± 4.4 ^g
J143055+350852a	88	14:30:55.07509(0.00001)	35:08:52.8445(0.0002)	3.78 ± 0.23	13.1 ± 0.4	13.1 ± 0.4
J143055+350852b	88	14:30:55.07953(0.00020)	35:08:52.9619(0.0030)	2.03 ± 0.29	47.0 ± 5.9	47.0 ± 5.9
J143213+350940 ⁱ	400	–	–	14.41 ± 1.64	30.1 ± 2.2	69.6 ± 5.2
J143213+350940a ⁱ	201	14:32:13.54889(0.00016)	35:09:40.8707(0.0023)	5.35 ± 0.86	31.5 ± 4.7	31.5 ± 4.7
J143213+350940b ⁱ	201	14:32:13.55036(0.00001)	35:09:40.8569(0.0002)	2.81 ± 0.32	4.1 ± 0.3	4.1 ± 0.3
J143329+355042	95	14:33:29.85779(0.00001)	35:50:42.2509(0.0001)	5.85 ± 0.33	4.0 ± 0.1	19.3 ± 0.4
J143718+364549	175	14:37:18.09615(0.00009)	36:45:49.8219(0.0013)	6.07 ± 0.53	40.2 ± 2.7	40.2 ± 2.7
J144230+355735 ⁱ	30	–	–	2.32 ± 0.17	–	134.4 ± 0.7 ^g
J144230+355735a ⁱ	36	14:42:30.69966(0.00036)	35:57:35.0093(0.0054)	0.56 ± 0.21	29.4 ± 10.8	29.4 ± 10.8
J144230+355735b ⁱ	36	14:42:30.69602(0.00019)	35:57:35.0979(0.0028)	0.63 ± 0.19	11.4 ± 3.3	33.3 ± 9.6
J144230+355735c ⁱ	36	14:42:30.69416(0.00053)	35:57:35.0776(0.0080)	0.31 ± 0.19	26.3 ± 15.8	26.3 ± 15.8

Source ID	Per cent flux density	T _b ^{d,e} [$\times 10^6$ K]	ν_0 [GHz]	ν_r^d [GHz]	LLS ^{d,f} [pc]
(1)	(8)	(9)	(10)	(11)	(12)
J142850+345420	99 ± 11	–	0.38 ± 0.08	–	<370 ± 17
J142850+345420a	–	>3.1 ± 0.4	–	–	–
J142850+345420b	–	>1.8 ± 0.5	–	–	–
J142904+354425	–	–	0.22 ± 0.13	0.38 ± 0.24 and 0.40 ± 0.25	–
J142917+332626	57 ± 6	23.4 ± 3.9 and 31.6 ± 3.4	0.39 ± 0.06	1.00 ± 0.20 and 1.35 ± 0.21	107 ± 8 and 102 ± 7
J143024+352438	–	–	0.29 ± 0.14	0.63 ± 0.32 and 0.67 ± 0.33	–
J143042+351240	4 ± 1	1.4 ± 0.4 and 1.3 ± 0.3	0.23 ± 0.11	0.51 ± 0.25 and 0.49 ± 0.23	97 ± 14 and 96 ± 12
J143050+342614	33 ± 5	113.0 ± 30.7 and 133.7 ± 29.4	0.34 ± 0.08	1.13 ± 0.31 and 1.33 ± 0.30	122 ± 20 and 115 ± 15
J143055+350852	71 ± 6	–	0.18 ± 0.09	0.56 ± 0.30 and 0.24 ± 0.13	1105 ± 99 and 698 ± 56
J143055+350852a	–	31.5 ± 4.6 and 13.6 ± 1.0	–	–	–
J143055+350852b	–	1.3 ± 0.3 and 0.6 ± 0.1	–	–	–
J143213+350940 ⁱ	94 ± 12	6.1 ± 1.0 and 6.0 ± 0.9	0.27 ± 0.13	0.54 ± 0.26 and 0.53 ± 0.25	539 ± 66 and 537 ± 53
J143213+350940a ⁱ	–	4.7 ± 1.3 and 4.7 ± 1.2	–	–	–
J143213+350940b ⁱ	–	145.7 ± 24.0 and 144.3 ± 22.6	–	–	–
J143329+355042	92 ± 9	128.1 ± 68.6 and 79.5 ± 5.0	0.35 ± 0.10	1.33 ± 0.80 and 0.83 ± 0.23	139 ± 26 and 158 ± 8
J143718+364549	68 ± 9	>1.7 ± 0.2	0.34 ± 0.10	–	<313 ± 29
J144230+355735 ⁱ	69 ± 14	–	0.30 ± 0.10	–	<1046 ± 28
J144230+355735a ⁱ	–	>0.3 ± 0.2	–	–	–
J144230+355735b ⁱ	–	>0.7 ± 0.4	–	–	–
J144230+355735c ⁱ	–	>0.2 ± 0.2	–	–	–

^aThe uncertainty, in seconds, is given in brackets after the value. ^bThe uncertainty, in arcseconds, is given in brackets after the value. ^cFor all the multicomponent sources except J143213+350940 and J144230+355735 (see note ⁱ), the value is the sum of the integrated flux densities of their components. ^dFor sources with two values, the first was calculated using z_{EAZY} and the second using z_{LRT} . ^eFor sources without redshifts, the values were calculated using a redshift of zero. ^fFor sources with upper limits, the limit was calculated using $z = 1$. ^gThe value is the source size (distance between the centroids of the two furthest components). ^hSince the source was not detected with the EVN, the RA and DEC values are taken from FIRST. ⁱThe values of the source were derived from the uv-tapered image while the values of the components were derived from the non-uv-tapered image (see Sections 4.3.3 and 4.3.5).

the sources without redshifts, we used a redshift of zero for these sources.

Column (10) contains the fitted observed turnover frequency (ν_0) for each source. Following Orienti, Dallacasa & Stanghellini (2007), Scaife & Heald (2012) and Orienti & Dallacasa (2014), we calculate ν_0 by fitting a second-order polynomial of the form $\log_{10}(S_\nu) = a(\log_{10}(\nu) - \log_{10}(\nu_0))^2 + b$ to the spectral plot of each of the

sources where a and b are constants. The spectral points to which that function was fitted are composed of the flux densities from the FIRST, GMRT-608, VLA-P and WIR catalogues and are shown in Fig. 2. For J143718+364549 and J144230+355735, which do not have VLA-P flux densities, we used the WENSS flux density. Since this involves fitting a function with three unknown parameters to three data points for the sources without GMRT-608 flux densities,

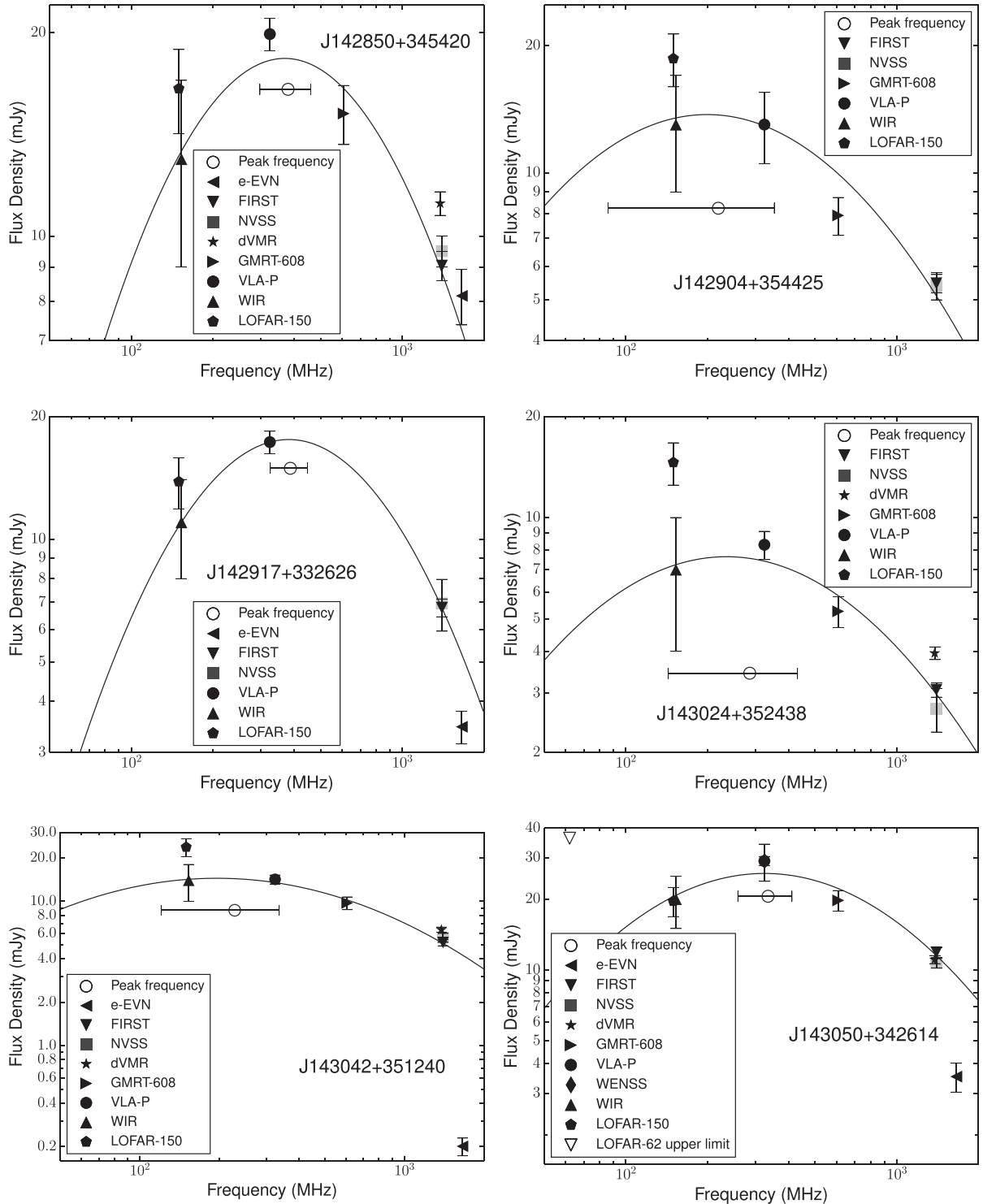


Figure 2. The spectral plots of the sources where the solid line is added to guide the eye. For J143050+342614, the VLA-P and WENSS flux densities are indistinguishable, with the larger of the two error bars being associated with the WENSS flux density. The LOFAR-62 detection limits, given in Table 4, are shown as an empty downward triangle for the sources where they help to constrain the spectrum.

the error values reported by the fitting algorithm can not be trusted for these sources. To improve the error estimates, we used a Monte Carlo method to estimate ν_0 and its error for all the sources. To do this, we used a random number generator to find new flux density values at each frequency for the source and calculated a new value of ν_0 . The flux densities returned by the random number generator are

Gaussian distributed values centred on the original flux density with a standard deviation equal to the error on the flux density. Repeating the procedure 100 000 times, we generated a histogram of all the solutions and fitted a Gaussian to it. The final value reported for ν_0 is the median of the fitted Gaussian and the error is its standard deviation. We note that the addition of the GMRT-608 flux densities

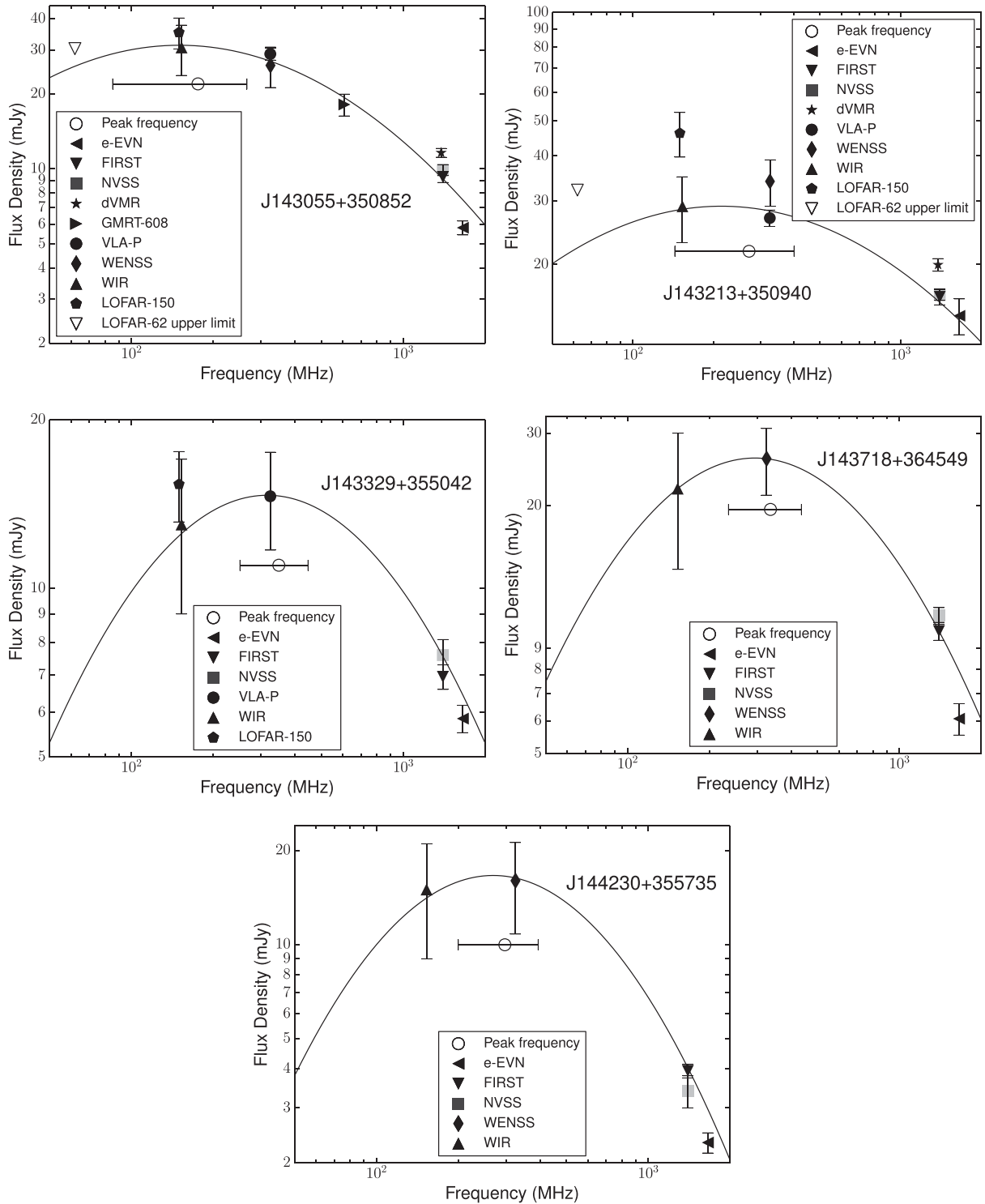


Figure 2 – continued

significantly decreased the uncertainty on the position of the peak frequency for the relevant sources.

Column (11) gives the rest-frame turnover frequency of the sources with known redshifts. The values were calculated using $\nu_r = \nu_o(1+z)$ where the errors of the z_{EAZY} values were dealt with as was described for columns (7) and (8).

Column (12) contains the LLSs of the sources calculated using their angular sizes and redshifts. For the sources which have EVN sizes but not redshifts, we calculated upper limits using $z = 1$. This can be done because, given a source with fixed linear size, its angular size will decrease as a function of redshift from $z = 0$ to $z \sim 1$. However, at $z > 1$, increasing the redshift of the source

results in a slight, but systematic, increase in its angular size (see for example fig. 10 in Falcke et al. 2004). Because of this, for a fixed angular size, a plot of the linear size to which it corresponds as a function of redshift peaks at $z = 1$.

3.3 Are they resolved?

To test if the source components are resolved in the EVN images, we calculated the minimum resolvable size of each of the Gaussian components fitted to the sources using equation (2) in Kovalev et al. (2005). Since all the fitted source sizes are larger than the minimum resolvable size, we conclude that all the sources and their components are resolved in the EVN images.

3.4 1.4 GHz variability

In Section 2.1, we pointed out that J142850+345420, J143024+352438, J143042+351240, J143055+350852 and J143213+350940 have flux density differences of more than 20 per cent between either FIRST and dVMR or NVSS and dVMR. This could potentially indicate that these sources are variable. However, it is surprising that all five sources have flux density differences of more than 20 per cent with the dVMR catalogue, while they all have a flux density difference of less than 12 per cent between FIRST and NVSS. We therefore checked the flux density offset between the dVMR catalogue and FIRST and dVMR and NVSS for the 190 VLA-P sources that we matched to all three catalogues in Coppejans et al. (2015), excluding only the sources that were flagged as extended. Doing this, we found that the median values of $S_{\text{dVMR}}/S_{\text{FIRST}}$, $S_{\text{dVMR}}/S_{\text{NVSS}}$ and $S_{\text{NVSS}}/S_{\text{FIRST}}$ are 1.16, 1.07 and 1.07, respectively. Since the dVMR flux densities were calibrated against NVSS (De Vries et al. 2002b), it is not surprising that their flux densities agree well with each other. What is interesting is that the flux densities of FIRST and NVSS, the highest and lowest resolution catalogues, agree well with each other but that the values for dVMR and FIRST do not agree as well. If anything we would expect that the dVMR values should be higher than those of FIRST since its resolution is three and a half times lower than that of FIRST. Of the 190 sources, the dVMR flux density of 124 are higher than both those of FIRST and NVSS. This seems to indicate that the dVMR flux densities are slightly higher than those of FIRST for our sub-selection of sources. Hence, the five sources that have a flux density difference of more than 20 per cent between FIRST or NVSS and dVMR are not necessarily variable.

One possible reason why the sources could be variable is that they are blazars, which are radio-loud AGN whose jet is pointed within a small angle of our line of sight (e.g. Urry 1999; Krawczynski & Treister 2013). It is however very unlikely that any of the sources (including J142904+354425 and J143024+352438 which we did not detect) are blazars since blazars have brightness temperatures above $\sim 10^{10}$ K (Readhead 1994; Homan et al. 2006), which is significantly higher than the values derived for our targets. In addition, blazars have flat spectra in the MHz regime (Massaro et al. 2013) which is not the case for any of our sources. Hence we do not believe that any of the sources are blazars and are therefore highly unlikely to be variable. It is also worth pointing out that the GPS and CSS sources are the least variable class of radio AGN (O’Dea 1998).

3.5 The origin of the radio emission

To test if the observed radio emission could be the result of star formation in the host galaxies, we used the same method as Magliocchetti et al. (2014) to differentiate between star-forming

galaxies and AGN using only radio luminosity. The method presented by Magliocchetti et al. (2014) is based on the results of McAlpine, Jarvis & Bonfield (2013) who used the optical and near-infrared SEDs of 942 1.4 GHz radio sources to calculate luminosity functions and redshifts for star-forming and AGN-dominated radio galaxies. Using these results, Magliocchetti et al. (2014) calculated that the radio power beyond which AGN-powered sources are dominant over star-forming sources scale with redshift as

$$\log_{10} P_{\text{cross}}(z) = \log_{10} P_{0,\text{cross}} + z \quad (2)$$

up to $z \sim 1.8$. Where $P_{0,\text{cross}} = 10^{21.7} \text{ W Hz}^{-1} \text{ sr}^{-1}$ is the value at $z = 0$. Magliocchetti et al. (2014) notes that for a given redshift, the amount of star-forming galaxies with radio powers larger than this value drops steeply and that at all redshifts, the radio luminosity function of star-forming galaxies drops off much steeper than that of AGN. Hence, the authors expect there to be very little contamination between the selections of star-forming and AGN galaxies (Magliocchetti et al. 2014).

To calculate the radio power of our sources, we used the same relation as Magliocchetti et al. (2014):

$$P_{1.4\text{GHz}} = S_{1.4\text{GHz}} D^2 (1+z)^{3-\alpha} \quad (3)$$

Here $P_{1.4\text{GHz}}$ is in units of $[\text{W Hz}^{-1} \text{ sr}^{-1}]$, $S_{1.4\text{GHz}}$ is the FIRST flux density in Jy converted to units of $[\text{W m}^{-2} \text{ Hz}^{-1} \text{ sr}^{-1}]$, D is the angular diameter distance in metres and α is the high-frequency spectral index of the source reported in Table 1. Note that we use an exponent of $3 - \alpha$ as opposed to $3 + \alpha$ in equation (3) since Magliocchetti et al. (2014) define the spectral index as $S \propto \nu^{-\alpha}$.

Following Magliocchetti et al. (2014), sources at $z \leq 1.8$ are classified as AGN powered if $P_{1.4\text{GHz}} > P_{\text{cross}}(z)$, and as star formation powered if $P_{1.4\text{GHz}} < P_{\text{cross}}(z)$. For sources at $z > 1.8$, the classification is done in the same way except that $P_{\text{cross}}(z)$ is always equal to $10^{23.5} \text{ W Hz}^{-1} \text{ sr}^{-1}$. Using both the minimum EAZY redshift allowed inside the errors and the LRT redshift, we found that in all the sources for which we have redshifts, the radio emission is from an AGN not star formation. While we do not have redshifts for J142850+345420, J143718+364549 and J144230+355735, the minimum redshift at which the emission would be the result of star formation is 0.06 for both J142850+345420 and J143718+364549 and 0.1 for J144230+355735.

Brightness temperatures can also be used to differentiate between star formation and AGN activity. Thermal radio emission caused by star formation typically has $T_b < 10^5$ K (Sramek & Weedman 1986; Condon et al. 1991; Kewley et al. 2000) while $T_b \geq 10^6$ K can be used as an indicator of non-thermal emission from AGN (e.g. Kewley et al. 2000; Middelberg et al. 2011). Since all the sources have $T_b \geq 10^6$ K inside their errors, this confirms that the radio emission is from AGN activity. This is further supported by the morphologies of the sources that are resolved into multiple components (Section 4). We note that non-thermal emission could also originate from a supernova remnant or a nuclear supernova remnant complex (e.g. Alexandroff et al. 2012). However, this possibility is excluded by the power cut described above.

² We note that in Magliocchetti et al. (2014) the units for $P_{0,\text{cross}}$ are incorrectly shown as $[\text{W Hz sr}^{-1}]$.

4 COMMENTS ON INDIVIDUAL SOURCES

4.1 Non-detections

The following two sources were not detected in the EVN images. Tapering the data also did not result in a detection, and from Section 3.5 we do not expect that the radio emission is caused by star formation. Below we conclude that the non-detections are because the sources have extended lobes which were resolved out by the high-resolution EVN observations while the faint core remains undetected. We note that this conclusion could be tested by observing the sources with an instrument that probes angular sizes between 0.15 and 5 arcsec.

4.1.1 J142904+354425

J142904+354425 is not detected in the naturally weighted, 5×5 arcsec EVN image with $23 \mu\text{Jy beam}^{-1}$ noise and a 30×34 mas restoring beam. J142904+354425 was included in our MPS selection in Coppejans et al. (2015) and was not flagged as being variable based on its FIRST and NVSS flux densities. No match could be found for it in the dVMR catalogue because it lies outside the area that was imaged. J142904+354425 was detected by Garrett, Wrobel & Morganti (2005) using the WSRT at 1.4 GHz and found to be smaller than 5 arcsec, having an integrated flux density of $3.026 \pm 0.026 \text{ mJy}^3$ which is only 45 per cent of the FIRST value. Hence we cannot rule out the possibility that J142904+354425 is variable, but based on the arguments at the end of Section 3.4, we consider this to be very unlikely.

To constrain J142904+354425 size we note that it has a deconvolved major and minor axis FWHM of 2.8 ± 0.5 and 1.7 ± 0.5 arcsec, respectively, in FIRST at a position angle of $41^\circ 6'$ east of north. We can estimate its minimum size by assuming it consists of a single component that has a surface brightness below the EVN's detection threshold. Taking the detection threshold to be five times the rms noise of the image ($0.115 \text{ mJy beam}^{-1}$), the flux density has to be spread out over at least $3.026/0.115 = 26.3$ beams which translates to a minimum angular size of approximately 0.9 arcsec.

Considering that J142904+354425's predicted flux density at 1.7 GHz is $4.9 \pm 0.4 \text{ mJy}$, that it is an AGN (Section 3.5) and non-variable, we conclude that the non-detection is because J142904+354425 was resolved out by the EVN observations. Finally, as is evident from the Fig. 2, we cannot say for certain that J142904+354425's spectrum turns over.

4.1.2 J143024+352438

J143024+352438 is not detected in the naturally weighted, 5×5 arcsec EVN image with $11 \mu\text{Jy beam}^{-1}$ noise and a restoring beam of 6×18 mas. Considering that J143024+352438 has a predicted 1.7 GHz flux density of $2.7 \pm 0.2 \text{ mJy}$, we should have easily detected it if it is a compact source. From Fig. 2 it is clear that the flux density difference between the WIR and LOFAR-150 catalogues (7.0 ± 3.0 and $14.6 \pm 1.5 \text{ mJy}$, respectively) can have a significant impact on the shape of J143024+352438's spectrum. It is worth pointing out that J143024+352438 has a signal-to-noise ratio (SNR; defined as the peak brightness divided by the local rms noise) of 49 and 5 in the LOFAR-150 and WIR images, respectively.

³ We note that Garrett et al. (2005) did not include the absolute calibration error (which the authors estimate to be less than 2 per cent) in their flux density error.

It is therefore possible that the spectrum can be described by a single power law between 150 and 1400 MHz.

Since we do not expect that the radio emission is caused by star formation⁴, we conclude that J143024+352438 has lobes which were resolved out by the EVN observations while the faint core remains undetected. J143024+352438's maximum size is set by its deconvolved major axis FWHM of 4.2 ± 0.8 arcsec in FIRST, which is at a position angle of $164^\circ 7'$. J142904+354425's minor axis is unresolved in FIRST. Using the same argument presented for J142904+354425, we calculated a minimum angular size of approximately 1 arcsec for J143024+352438.

4.2 Marginally resolved sources

The following four sources appear as single components in the EVN images that do not have a discernible structure. Note that from Section 3.3, all of the sources are resolved. Below we discuss each of the sources individually, focusing on, among other things, the percentage of the predicted flux density that was recovered from the image and the source variability. If the percentage of the predicted flux density recovered from the image is low, it could indicate that the measured source size is not a good estimate of the true source size. If, however, the source is variable, it would mean that the predicted flux density is unreliable. Hence the source could be more extended or has an additional component that could have been missed in the EVN image. While we argue that none of the sources are variable, based on the percentage of the predicted flux densities recovered from the images, we conclude that J143042+351240 and J143050+342614 likely have undetected structure, and are therefore larger than indicated in Table 5.

4.2.1 J142917+332626

J142917+332626 is present in both the NVSS and dVMR image cutouts, but not in the catalogues because there is another source ~ 35 arcsec away which blends with it. To determine its flux density in NVSS we simultaneously fit it and the nearby source using the PYBDSM source detection package⁵. From this we found that it has an integrated flux density of $6.96 \pm 1.01 \text{ mJy}$, which differs by 2.5 per cent from the FIRST value. J142917+332626 was also observed by Ciliegi et al. (1999) at 1.4 GHz with the VLA in C configuration. The authors found a flux density of $6.27 \pm 0.04 \text{ mJy}$ for J142917+332626 which differs by 11 per cent from the FIRST value. Considering that we recovered between 51 and 63 per cent of the predicted flux density, we expect that the reported size for J142917+332626 is a good estimate of the true size.

Matching the VLA-P sources to the *Chandra* XBoötes X-ray survey of the Boötes field (Murray et al. 2005) with a search radius of 10 arcsec, J142917+332626 was the only source from our 11 sources for which we found a counterpart. The centroid of the matched source, CXOXB J142917+332626.4, is 0.32 arcsec away

⁴ We note that matching J143024+352438 to the SDSS Data Release 10 catalogue (<https://www.sdss3.org/dr10/>; Ahn et al. 2014), we found a source (SDSS J143025.19+352441.3) 3.3 arcsec away from J143024+352438 with a photometric redshift of 0.342 ± 0.1167 . This could influence whether J143024+352438 is classified as being dominated by star formation or an AGN. However, using the analysis in Section 3.5, J143024+352438 would only be classified as being dominated by star formation if it is at a redshift below 0.11.

⁵ <http://dl.dropboxusercontent.com/u/1948170/html/index.html>

from the VLA-P source and was detected in all three of *Chandra*'s bands with seven, one and six counts in the 0.5–7, 0.5–2 and 2–7 keV bands, respectively. We note that, since J142917+332626 can unambiguously be classified as an MPS source (Fig. 1), it is one of the few turnover sources with a X-ray counterpart.

Matching the FIRST sources to those in XBoötes, El Boucheffy (2009) associated CXOXB J142917+332626.4 with the same FIRST source to which we matched the VLA-P source. In addition, CXOXB J142917+332626.4 was matched to an optical source in the NOAO Deep Wide-Field Survey (NDWFS) survey of the field by Brand et al. (2006). Using the publicly available HYPERZ⁶ code along with the optical information, El Boucheffy (2009) determined a photometric redshift of $0.960_{-0.806}^{+1.227}$ for J142917+332626. This value is consistent with the EAZY redshift of J142917+332626 and is lower than the LRT value.

Based on its high X-ray to optical flux ratio, $\log_{10}(f_{(2-7)\text{keV}}/f_{\text{opt}}) = 1.37$, J142917+332626 is expected to either be a high-redshift source and/or dust obscured (El Boucheffy 2009, and references therein). The author classifies it as an obscured AGN (AGN-2) using its hardness ratio and X-ray luminosity.

4.2.2 J143042+351240

J143042+351240's EVN flux density of 0.2 mJy is only 4 ± 1 per cent of the predicted value. Tapering the data did not increase J143024+352438's flux density significantly. Since the emission is related to AGN activity and J143042+351240 is not variable (Sections 3.5 and 3.4), this leads us to conclude that the true size of J143042+351240 is larger than indicated in columns (7) and (12) in Table 5. If this is the case, the remaining flux density could be in low surface brightness emission surrounding the source. In this case the true source size will be larger than the measured value, but likely not significantly. If however the missing flux density is in a second component, the size could be a significant underestimate of the true source size.

We were able to match J143042+351240 to a source in the AGN and Galaxy Evolution Survey (AGES; Kochanek et al. 2012) catalogue which measured redshifts for 23 745 galaxies and AGN in the Boötes field. The photometric redshift of the AGES source (0.96) agrees very well with the values reported in Table 1. The flux density difference between the WIR and LOFAR-150 catalogues (14.0 ± 4.0 and 23.8 ± 2.4 mJy, respectively) can have a significant impact on the shape of J143042+351240's spectrum (Fig. 2). The two catalogues have SNRs of 9 and 98, respectively. It is therefore possible that the spectrum (Fig. 2) can be described by a single power law between 150 and 1400 MHz.

4.2.3 J143050+342614

We only recovered 33 per cent of the predicted flux density for J143050+342614. This percentage did not increase when we applied an uv-taper to the image. The emission is related to AGN activity and the source is not variable (Sections 3.5 and 3.4). As is the case with J143042+351240, this leads us to conclude that J143050+342614 is larger than indicated in columns (7) and (12) in Table 5. We note that J143050+342614 can unambiguously be classified as an MPS source when including the errors on its spectral indices. It is also matched to a WENSS source with the WENSS and VLA-P flux densities differing by less than one per cent. This re-

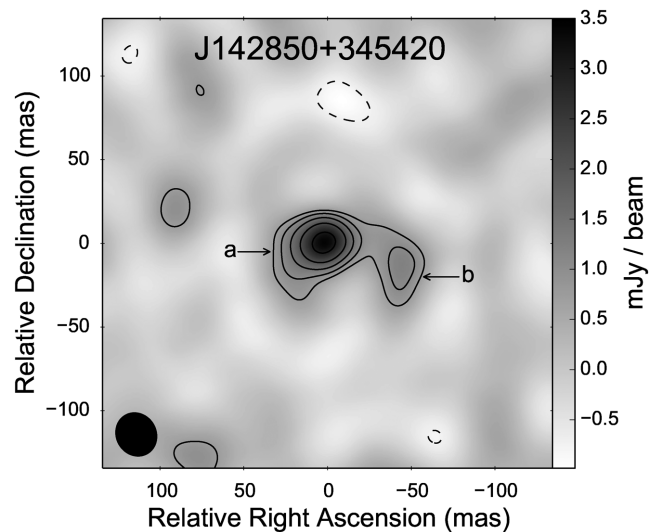


Figure 3. Naturally weighted EVN image of J142850+345420 that was made using an uv-taper. The restoring beam is shown in the bottom left corner and has a size of 24.1×26.8 mas at a major axis position angle of $29^\circ.4$. The contours are drawn at -3 and 3 times the image noise, increasing in factors of $\sqrt{2}$ thereafter.

sults in the two points being indistinguishable in J143050+342614's spectral plot (Fig. 2) with the larger of the two error bars being associated with the WENSS flux density.

4.2.4 J143329+355042

J143329+355042 is composed of a single component in the EVN image and lies outside the survey area covered by the dVMR catalogue. We recovered 92 ± 9 per cent of J143329+355042's predicted flux density. This indicates that its measured size is a good estimate of its true size. J143329+355042's spectrum (shown in Fig. 2) flattens towards lower frequencies and could turn over (Table 1).

4.3 Resolved sources

The following five sources are resolved in the EVN images and have a discernible structure. Below we discuss each of the sources individually and argue that the observed structures are lobes and/or hotspots in their jets.

4.3.1 J142850+345420

J142850+345420, shown in Fig. 3, has a double structure in the EVN image with J142850+345420a and J142850+345420b being detected at a 14σ and 5σ level, respectively. Fig. 3 was made using natural weighting and applying an uv-taper. The full-resolution, uniformly weighted image has a typical resolution of 3×10 mas, but we could not confirm the detection of J142850+345420b. Applying the uv-taper increased the beam size to 24.1×26.8 mas and resulted in the flux density of both components increasing. Specifically, the flux density of J142850+345420a increased by a factor of 1.7 and the percentage of recovered flux density of the source as a whole increases from 44 to 99 per cent. This is a clear indication that both components were resolved in the full resolution image.

For J142850+345420, we recovered 99 per cent of the predicted flux density, this indicates that its measured size is a good estimate

⁶ <http://webast.ast.obs-mip.fr/hyperz/>

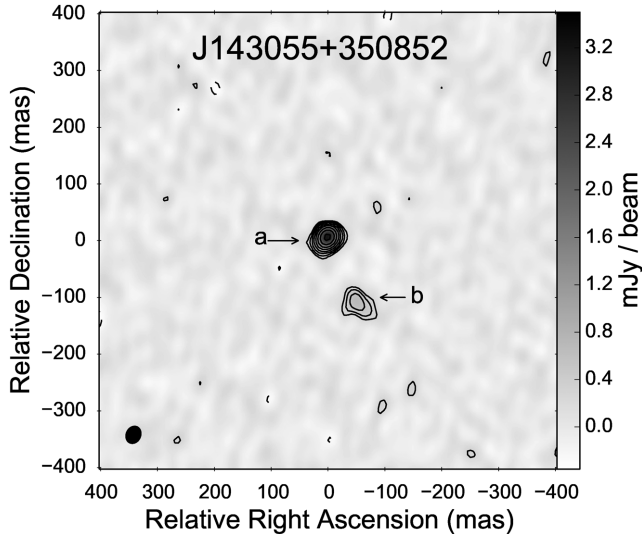


Figure 4. Uniformly weighted EVN image of J143055+350852. The restoring beam is shown in the bottom left corner and has a size of 26.4×32.1 mas at a major axis position angle of $-24^\circ.5$. The contours are drawn at -3 and 3 times the image noise, increasing in factors of $\sqrt{2}$ thereafter.

of its true size. J142850+345420 was however flagged as being variable because of a 26 per cent flux density difference between the FIRST and dVMR catalogues. Considering that the NVSS and FIRST flux densities only differ by 7 per cent, and that, as discussed in Section 2.1, the dVMR flux densities seem to be systematically higher than the FIRST and NVSS values for our sample of sources, it is unlikely that J142850+345420 is variable.

J142850+345420 lies inside the multiwavelength optical coverage of the Boötes field, but no match could be found for the source. Consequently, we could not determine a photometric redshift for it. This likely indicates that J142850+345420 is at $z > 2$. If this is the case, J142850+345420 is AGN dominated since we calculated in Section 3.5 that the radio emission would be the result of AGN activity if the source is at $z > 0.06$. Since we do not have a redshift for J142850+345420, we calculated the brightness temperatures for J142850+345420a,b using a redshift of zero. The values are therefore lower limits and allow us to conclude that J142850+345420 is consistent with being an AGN with a core-jet structure.

Since we do not have a redshift for J142850+345420, we calculated the brightness temperatures for J142850+345420a,b using a redshift of zero. The values are therefore lower limits and allow us to conclude that J142850+345420 is consistent with being an AGN with a core-jet structure. We also note that, J142850+345420 lies inside the multiwavelength optical coverage of the Boötes field, but no match could be found for the source. Consequently, we could not determine a photometric redshift for it. This likely indicates that J142850+345420 is at $z > 2$. It also indicates that J142850+345420 is AGN dominated since we calculated in Section 3.5 that the radio emission would be the result of AGN activity if the source is at $z > 0.06$.

4.3.2 J143055+350852

J143055+350852, shown in Fig. 4, is composed of two components. The EAZY and LRT redshift values for J143055+350852 differ significantly. Despite this, the brightness temperature of J143055+350852a is consistent with it being an AGN. Considering that J143055+350852b is extended in the direction

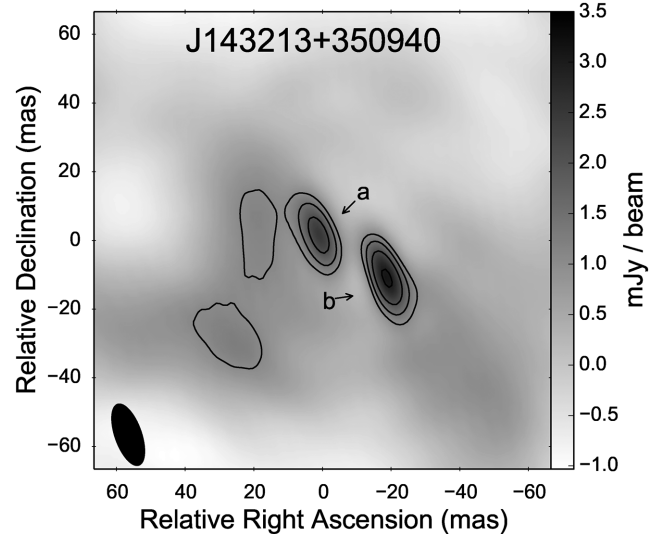


Figure 5. Uniformly weighted EVN image of J143213+350940 without uv-tapering showing a likely CSO structure. The restoring beam is shown in the bottom left corner and has a size of 7.8×18.9 mas at a major axis position angle of $19^\circ.4$. The contours are drawn at -3 and 3 times the image noise, increasing in factors of $\sqrt{2}$ thereafter.

parallel to the line between the components and that it is fainter than J143055+350852a, we suggest that J143055+350852 has a core-jet structure. Since we recovered most of the predicted flux density (71 ± 6 per cent), and the remaining flux density is likely located either between the two components or very near them, we do not expect that J143055+350852 is significantly more extended than indicated in Table 5. From Fig. 2, it is clear that the upper limit on the LOFAR-62 flux density indicates that the spectrum is turning over.

4.3.3 J143213+350940

The uniformly weighted EVN image of J143213+350940 (Fig. 5) has a likely symmetric double structure. Both components have brightness temperatures consistent with emission from AGN activity. After imaging J143213+350940 with uniform weighting, we also imaged it using natural weighting with an uv-taper. Note that the information given for J143213+350940 in Table 5 was derived from the naturally weighted, uv-tapered image, while the information for J143213+350940a,b was derived using the uniformly weighted image. Because the beam size of the uv-tapered image is 24×29 mas and J143213+350940a,b are only separated by ~ 31 mas, their flux densities merge in the uv-tapered image. Fitting the resulting single source with a single elliptical Gaussian, we found it to have an FWHM of nearly 70 mas and flux density of 14.41 mJy. This is ~ 6 mJy more than the sum of the flux densities of the individual components. Hence, while the source is resolved, it is clear that there is a significant amount of extended emission between and surrounding the components. From the uv-tapered image, we recover 94 ± 12 per cent of the predicted flux density. Hence, the size measured from the uv-tapered image is a good estimate of J143213+350940's true size.

From Fig. 2 it is clear that the VLA-P and WENSS flux densities differ (26.9 ± 1.4 and 34.0 ± 5.0 , respectively), but are nearly consistent within their 1σ errors. If we take the VLA-P flux density to be the correct value at 325 MHz, J143213+350940 is not an MPS source (Fig. 1) whose spectrum either flattens towards lower

frequencies or can be described by a single power law between 153 and 1400 MHz (Table 1). This means that the turnover frequencies and linear sizes reported in Table 5 will not be good estimates of the true values. However, if we take the WENSS flux density to be the correct value, α_{low} and α_{high} changes to 0.2 ± 0.3 and -0.5 ± 0.1 , respectively. This means that J143213+350940 is an MPS source as is shown by the point labelled ‘J143213+350940 WENSS’ in Fig. 1. The percentage of the predicted flux density that was recovered from the image changes to 97 ± 13 per cent while $\nu_o = 337 \pm 75$ MHz and $\nu_r = 666 \pm 152$ and 660 ± 146 MHz in Table 5.

To try to find a reason for the flux density difference between the VLA-P and WENSS catalogues, we double checked both catalogues and images of the source. Since the nearest neighbouring source is more than 3 arcmin away, J143213+350940 is unresolved in the WENSS, VLA-P, FIRST and NVSS images, and there is nothing in the images that can explain the difference. One possibility is that J143213+350940 is intrinsically variable. However, as discussed in Section 3.4, we consider this to be very unlikely. Another possibility is that the difference could be caused by (as a whole or in part) interplanetary or interstellar scintillation. Typically, interstellar scintillation results in flux density variations of a few per cent (e.g. Stinebring et al. 2000; Gelfand, Lazio & Gaensler 2005), which is much less than the 21 per cent difference observed here. In addition, the VLA-P observations were carried out over four days during a period of more than a month (Copejans et al. 2015), while each field in WENSS was observed in 18 snapshots spread over a 12-h period (Rengelink et al. 1997). This averaged out the variations, which have maximum typical time-scales of about an hour (Cordes & Rickett 1998; Gelfand et al. 2005).

The WIR and LOFAR-150 flux densities for J143213+350940 are 29.0 ± 6.0 and 46.3 ± 4.6 mJy, respectively. If the LOFAR-150 flux density is correct, it would indicate that the spectrum is steep between 150 and 1400 MHz, irrespective of whether the VLA-P or WENSS flux densities are correct. However, it is clear from Fig. 2 that the LOFAR-62 flux density upper limit shows that J143213+350940’s spectrum either flattens towards lower frequencies or turns over irrespective of which of the four catalogues are correct at 153 and 325 MHz.

Finally, we note that matching J143213+350940 to the AGES catalogue, we found a counterpart within 1.1 arcsec from our central position that has a photometric redshift of 1.02. The value agrees very well with those in Table 1.

4.3.4 J143718+364549

We selected J143718+364549 for observation based on its WENSS flux density. Unfortunately, because it is 2.7 from the phase centre of the VLA-P image, we do not have a VLA-P flux density for it. The naturally weighted EVN image of J143718+364549 (Fig. 6) shows that the source is composed of a single resolved component that is significantly extended towards the north. To check if any flux density is resolved out in the naturally weighted image, we made an image in which we applied an uv-taper. The values derived from this image is reported in Table 5. Tapering increased the recovered flux density from 2.9 mJy to 6.1 mJy, resulting in 68 ± 9 per cent of the predicted flux density being recovered.

From J143718+364549’s spectrum (Fig. 2), it is clear that it either turns over or flattens towards lower frequencies. Finally, we point out that because we do not have a redshift for J143718+364549, the brightness temperature lower limit in Table 5 was derived using a redshift of zero.

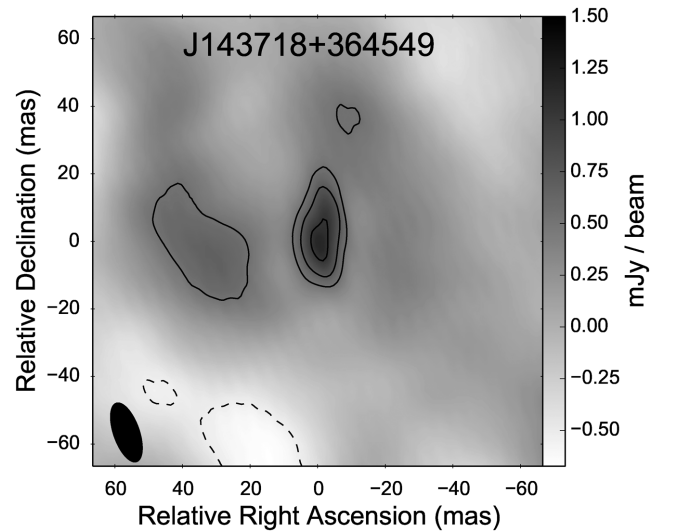


Figure 6. Naturally weighted EVN image of J143718+364549 without uv-tapering. J143718+364549 is significantly extended towards the north. The restoring beam is shown in the bottom left corner and has a size of 7.7×18.4 mas at a major axis position angle of 19.3° . The contours are drawn at -3 and 3 times the image noise, increasing in factors of $\sqrt{2}$ thereafter.

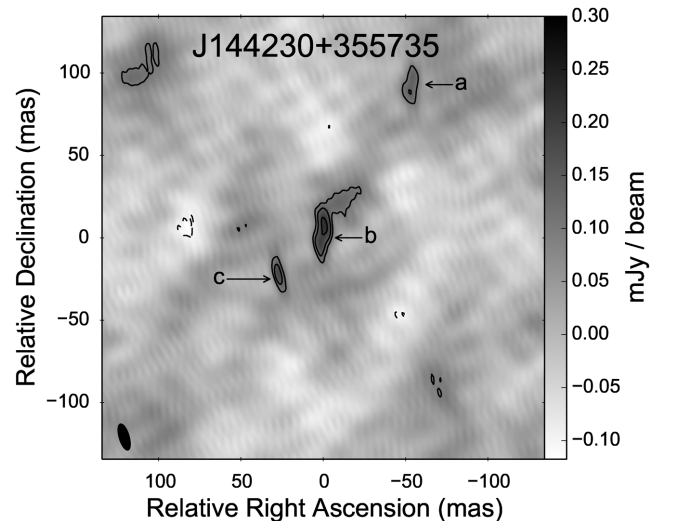


Figure 7. Naturally weighted EVN image of J144230+355735 without uv-tapering showing a CSO structure. The restoring beam is shown in the bottom left corner and has a size of 6.3×16.7 mas at a major axis position angle of 16.6° . The contours are drawn at -3 and 3 times the image noise, increasing in factors of $\sqrt{2}$ thereafter.

4.3.5 J144230+355735

Like J143718+364549, J144230+355735 lies outside the area imaged in the VLA-P image and was selected for observations with the EVN based on its WENSS flux density. From its spectrum (Fig. 2), the source can either be turning over or flattening towards lower frequencies.

The EVN image of J144230+355735 (Fig. 7) is composed of three resolved components and has a CSO structure. J144230+355735’s overall structure is that of a core with two lobes. From J144230+355735b a jet-like structure extends towards J144230+355735c.

To check if any of the flux density is resolved out in the naturally weighted image, we made an image in which we applied an uv-taper. The values derived from the uv-tapered image are reported for J144230+355735 in Table 5 while the values for J144230+355735a,b,c are derived from the non-uv-tapered image. Since the flux densities of J144230+355735b and J144230+355735c merge in the uv-tapered image, the component was fitted using a single elliptical Gaussian. In this image, J144230+355735a has a flux density of 1.07 ± 0.13 mJy and a size of 46.7 ± 3.4 mas. The combined component from J144230+355735b and J144230+355735c has a minor and major axis of 21.9 ± 1.3 and 64.0 ± 3.7 mas, respectively, and a flux density of 1.26 ± 0.11 mJy. Hence, after tapering, the total flux density of J144230+355735 increased by 0.6 mJy resulting in 69 ± 14 per cent of the predicted flux density being recovered. The missing flux density likely originates from low surface brightness emission along the jet axis between the central and outer components. It is however possible that at least some of the missing flux density is located beyond the outer components and that the jet is larger than what we measured. We note that while J144230+355735a is only detected at a 4.2σ level in Fig. 7, it is detected above 6σ in the uv-tapered image. Similarly the central position of J144230+355735c is encircled by a 6σ contour in the uv-tapered image.

The brightness temperature of the central component is consistent with the emission being non-thermal within its uncertainty. Furthermore, the brightness temperatures of all three components are lower limits as they were calculated using a redshift of zero

since we could not determine a redshift for the source. Hence, if J144230+355735 is at a redshift above 0.45, the central component would clearly indicate AGN activity. Combining this with the structure of the source and the results in Section 3.5, we conclude that J144230+355735 is an AGN with a central core-jet structure and two lobes and/or hotspots.

5 DISCUSSION

5.1 What are the MPS sources?

From Sections 3.5 and 4, we know that all the sources detected with the EVN are AGN. For the sources with redshifts, columns (11) and (12) in Table 5 show that they are GPS and CSS sources whose spectral turnovers have been redshifted to lower frequencies. Column (12) shows that J142850+345420 and J143718+364549, for which we do not have redshifts, are both smaller than ~ 1 kpc and are therefore either GPS or HFP sources. Hence, since all the sources that were detected with the EVN are either CSS, GPS or HFP sources, they are likely all young AGN (O’Dea 1998; Conway 2002; Murgia et al. 2002; Murgia 2003; Fanti 2009).

In Fig. 8, we used the information in table 5 of Orienti & Dallacasa (2014) to generate the plot that the authors used to derive the turnover frequency–linear size relation. To determine the empirical relation, the authors compiled a list of CSOs for which the core components have been detected. The authors argue that this allows them to select genuine young radio galaxies and exclude other

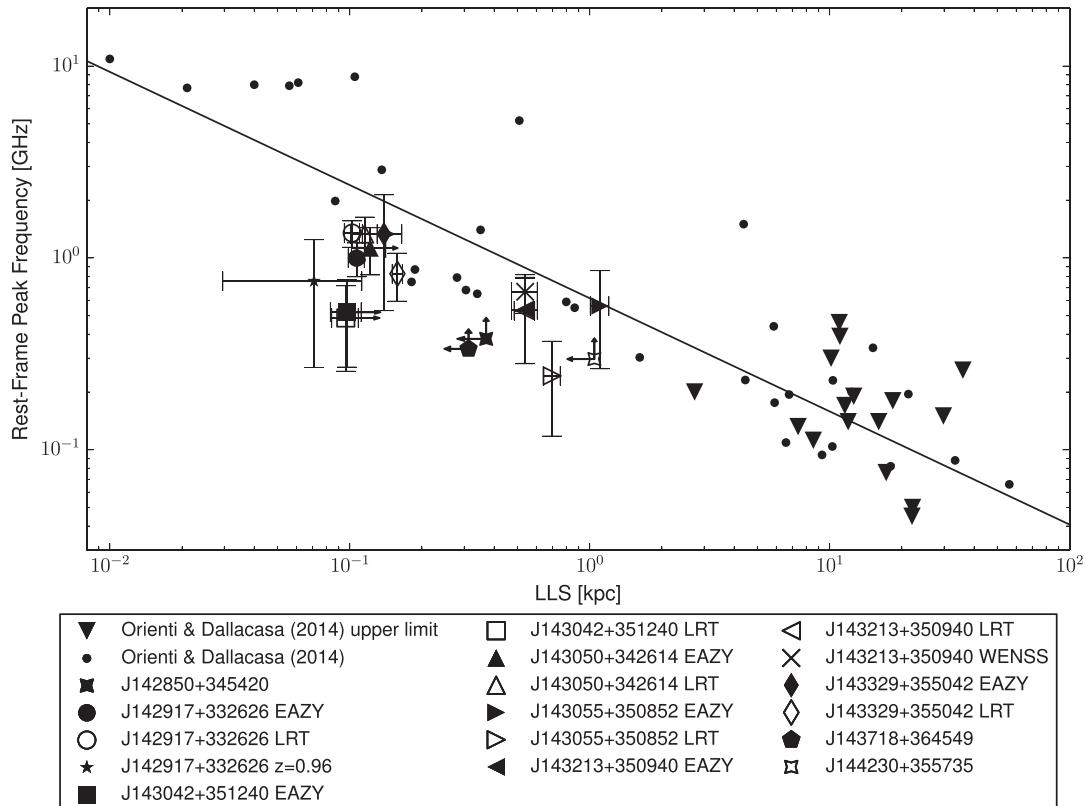


Figure 8. The figure used by Orienti & Dallacasa (2014) to derive the turnover frequency–linear size relation (shown as the solid line) with our sources included. The points and downward triangles show the sources from which Orienti & Dallacasa (2014) derived the relation. The downward triangles show sources for which the authors could only determine an upper limit for the turnover frequency. For each of our sources, the position is plotted using the redshifts from both the EAZY and LRT codes. For J142850+345420, J143718+364549 and J144230+355735, for which we cannot calculate ν_r , we plotted ν_o which is a lower limit for ν_r . Lower and upper limits on the LLS are indicated by arrows. A zoom-in of the region around our sources is shown in Fig. 9.

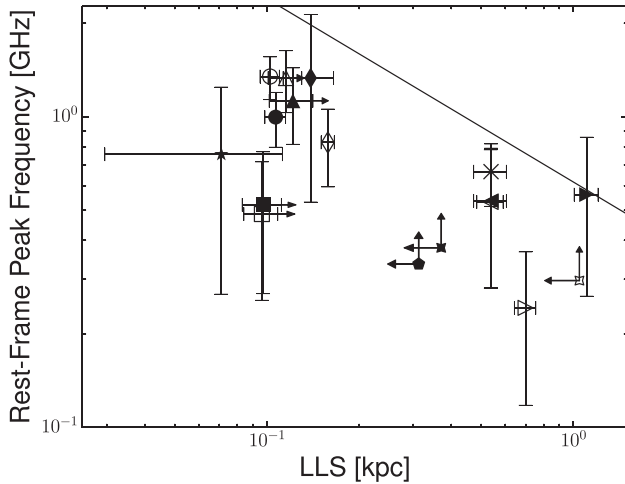


Figure 9. A zoomed-in image of the region around our sources in Fig. 8 in which the sources from which Orienti & Dallacasa (2014) derived the relation was omitted for clarity. The figure legend is the same as that of Fig. 8 and is not shown here for clarity. For J142850+345420, J143718+364549 and J144230+355735, for which we cannot calculate ν_r , we plotted ν_o which is a lower limit for ν_r . Lower and upper limits are indicated by arrows.

populations of sources such as blazars that could contaminate the sample. The resulting list of sources span redshifts between 0.08 and 2.93, linear sizes between 10 pc and 56 kpc and rest-frame turnover frequencies from 10.9 GHz to below 45 MHz. By minimizing the chi-squared error statistics, Orienti & Dallacasa (2014) found the relation

$$\log_{10} \nu_r = (-0.21 \pm 0.04) - (0.59 \pm 0.05) \log_{10}(\text{LLS}), \quad (4)$$

where ν_r is in GHz and the LLS is in kpc. This relation agrees well with those found by both O’Dea (1998) and Falcke et al. (2004). Fig. 8 shows the data used to derive the relation, with Fig. 9 zooming into the region around our sources.

To check if our sources lie within the scatter of points around the relation, we calculated the vertical distance between the relation and the sources used to derive it. From this we found that the median distance between the points and the line is 0.06 and has a 1σ deviation of 0.29. J143042+351240 EAZY and J143042+351240 LRT are 2.5σ and 2.7σ from the relation, with the remaining sources being within 2σ and J143050+342614 LRT, J143055+350852 EAZY, J143213+350940 EAZY, J143213+350940 LRT, J143213+350940 WENSS and J143329+355042 EAZY being within 1σ . The maximum distance between the relation and the points used to derive it is 5.8 GHz, which is more than the maximum distance between the relation and our sources. We note that since equation (4) has a constant slope, the vertical, horizontal and diagonal distances between each point and the relation are related to each other by constant factors. Hence, the number of σ that each point is from the relation is independent of whether the vertical, horizontal or diagonal distance is used.

As discussed in Section 4.3.3, the VLA-P and WENSS flux densities for J143213+350940 differ, resulting in a large uncertainty on the shape of its spectrum. Consequently, we plotted its position in Figs 8 and 9 for both the cases where ν_r was derived using the VLA-P and WENSS flux densities. Since the EAZY and LRT redshifts are so close to each other (Table 1), the points are indistinguishable in the plots and we only plot the position of the EAZY redshift point for the case where ν_r was derived from the WENSS flux density.

Since the redshift derived for J142917+332626 by El Boucheffry (2009) differs from that of the EAZY and LRT values, we added an additional point (labelled as ‘J142917+332626 $z = 0.96$ arcsec) showing its position using their redshift. J142917+332626 and J143055+350852 (for which the EAZY and LRT redshift values differ significantly) illustrate the effect of redshift uncertainty on the position of the sources. Typically, an uncertainty in the redshift will have a larger effect on the rest-frame turnover frequency than on the linear size of the source, resulting in the sources primarily being displaced vertically. The reason for this is, for a source with a fixed linear size, its angular size does not simply decrease as a function of redshift [see the discussion on Column (12) of Table 5 in Section 3], as is the case for its observed turnover frequency. Hence for sources around $z = 1$ and above, an uncertainty in the redshift primarily has an effect on the rest-frame turnover frequency. However, for sources at $z < 1$, a redshift uncertainty can have a significant effect on both ν_r and the linear size.

Since we do not have redshifts for J142850+345420, J143718+364549 and J144230+355735, we used their observed turnover frequencies and size upper limits to plot their positions in Figs 8 and 9. Since $\nu_r = \nu_o(1+z)$, ν_o is a lower limit for ν_r . From the arguments in the previous paragraph, we expect that the true positions of the sources will primarily be vertical displaced compared to their current positions. As plotted, the sources are between 1.2σ and 2.1σ from the relation.

Looking at Fig. 8, it is striking that all but two of our sources unambiguously lie below the correlation. There are several possible reasons for this, the first of which being that the sources were resolved out by the EVN, or that we missed components because they are too faint to be detected in the EVN images. In this case, the LLS will be underestimated and the sources should be towards the right of where they are plotted, closer to the relation. This is specifically the case for J143042+351240 and J143050+342614, for which, based on the percentage of the predicted flux density that was recovered from the EVN image, we concluded in Sections 4.2.2 and 4.2.3 that they are larger than their measured sizes. We also cannot rule out this possibility for the remainder of the detected sources (for which we recovered a larger percentage of the predicted flux density). For J143042+351240 and J143050+342614, the limits are indicated by arrows in Figs 8 and 9. The second possibility why the sources lie below the correlation is selection effects. The first selection effect is that we can only select sources that have $\nu_o \lesssim 500$ MHz. Assuming that all of the sources that we can detect are at $z < 10$, this excludes sources with $\nu_r \gtrsim 5.5$ GHz. The second selection effect is that the EVN’s lack of short spacings prevents us from detecting extended emission on scales larger than ~ 35 mas. Hence, we cannot detect a source if it has a single, or multiple, components larger than ~ 35 mas, in which the flux density is uniformly distributed over the component or components. This is also the most likely explanation for why we did not detect J142904+354425 and J143024+352438. Since we cannot place an upper limit on the distance between two $\gtrsim 35$ mas components, it is impossible to determine the true limitations imposed by this selection effect. Another selection effect is that the MPS sources were selected to be compact in FIRST (see section 3.2 in Coppejans et al. 2015). However since this only sets the constraint that $\text{LLS} \lesssim 54.5$ kpc, this effect is irrelevant. In conclusion, we cannot therefore say whether the sources lying below the correlation is because of selection effects or not.

From the above it appears that the sources plotted in Figs 8 and 9 lie within the scatter of the points around the relation. We do note that the uncertainties on LLS and the peak frequency could influence this result. Since all the sources are AGN (Sections 3.5

and 4), we therefore expect that at least some, if not all, of the sources will reveal themselves as CSOs if observed at high enough resolution. This can be said with certainty for, and is confirmed by, J143213+350940 and J144230+355735 which have CSO structures. We also point out again that, as discussed in Section 3.4, it is very unlikely that any of the sources are blazars.

As a final point, we note that the redshift range of the sources is similar to those of the sources from which Orienti & Dallacasa (2014) constructed the relation. Hence, if we assume that the sources are below the correlation purely because of their LLS's, and that they were not resolved out by the EVN, it is possible that they are less powerful than the sources in the relation and, therefore, get frustrated at smaller distances. Since AGN activity can be discontinuous (e.g. Schoenmakers et al. 2000; Jamrozy et al. 2007), it is also possible that their central engines have (temporarily) switched off. This would result in the core being undetected while the low surface density radio lobes continue to expand and fade. There are, however, two arguments against this. First, switched off AGN typically have spectra with $\alpha < -1.5$ (Komissarov & Gubanov 1994; Shulevski et al. 2015), which is steeper than the spectra of our sources. Secondly, the core accounts for a relatively small fraction of the total flux density (e.g. Orienti & Dallacasa 2014), and therefore often remains undetected.

5.2 Detection fraction

In this section we compare our detection fraction of 0.818 to that of the mJIVE VLBA Exploration at 20 cm (mJIVE-20; Deller & Middelberg 2014) survey. mJIVE-20 is a VLBA survey to study the VLBI structure of FIRST sources. The survey has an angular resolution of 5 mas, median noise of $0.178 \text{ mJy beam}^{-1}$ and targets sources with FIRST peak intensities (I_{FIRST}) between 1 and $200 \text{ mJy beam}^{-1}$. To date, nearly 20 000 sources have been observed of which 4336 were detected at or above a detection threshold of 6.75σ , giving the survey an overall detection fraction of 0.217. Since our sources have $2.5 < I_{\text{FIRST}} < 15.7 \text{ mJy beam}^{-1}$, we used the information in table 2 of Deller & Middelberg (2014) to calculate that mJIVE-20 has a detection fraction of 0.151 ± 0.007 for sources with $2 < I_{\text{FIRST}} < 16 \text{ mJy beam}^{-1}$. The fact that we have a higher detection fraction is in part due to our lower detection threshold (6σ) and median noise ($0.095 \text{ mJy beam}^{-1}$). To correct for this, we calculated that with mJIVE-20's median noise and detection threshold, our survey would detect 7 of our 11 sources. This means a detection fraction of 0.636, which is 3 times that of the mJIVE-20 survey as a whole and 4 times that achieved for the sources in our peak intensity range.

As discussed in Section 5.1, MPS sources are some of the smallest AGN, which explains our high detection fraction in comparison to mJIVE-20. mJIVE-20 targets all sources with FIRST peak intensities above 1 mJy beam^{-1} , while all our sources should be intrinsically small and thus good targets for VLBI.

6 SUMMARY AND FINAL POINTS

In this paper, we presented high-resolution 1.7 GHz EVN observations of 11 MPS sources. Of the 11 sources we detected 9 with the EVN, recovering more than 50 per cent of the predicted flux density for seven. Based on their radio luminosities, brightness temperatures and morphologies, we conclude that the radio emission of the sources detected with the EVN is from AGN activity (Sections 3.5 and 4). In Section 5.1, we used their peak frequencies and linear sizes to show that the MPS sources are a mixture of CSS, GPS

and HFP sources whose spectral turnovers have been redshifted to lower frequencies. Hence, the MPS sources are likely all young AGN. We also argue that the detected sources are all likely CSOs. We do, however, emphasize that the uncertainties on the rest-frame turnover frequencies and the LSSs could influence this conclusion. From their steep high-frequency spectra and brightness temperatures, we argue that none of the sources are blazars. We conclude that the radio emission from the two sources that were not detected with the EVN are likely from AGN activity, and that the reason for the non-detections is that they were resolved out by the EVN. Finally we conclude that our high detection fraction of 82 per cent is due to the MPS sources being intrinsically small AGN, which makes them good targets for VLBI observations.

Because of small number statistics, we cannot say whether the sources at the highest redshifts are more compact than those at lower redshifts. However, all the detected sources, with the possible exceptions of J143042+351240 and J143050+342614, have linear and angular sizes smaller than 1.1 kpc and 0.15 arcsec, respectively. This shows that low-frequency colour-colour diagrams are an easy and efficient way of selecting small and likely young AGN. This also seems to indicate that when selecting a group of MPS sources that are compact on a scale of 5 arcsec, most of the sources are compact on a scale below 0.5 arcsec. Hence it is possible that to fulfill the criteria of selecting compact sources for the MPS method of searching for high-redshift AGN, all that is required is to select sources that are unresolved at 5 arcsec.

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