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Uniform infinite and Gibbs causal triangulations

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CHAPTER 5

Properties of transfer matrix for Ising model coupled to causal dynamical triangulations

We introduce a transfer matrix formalism for the (annealed) Ising model coupled to two-dimensional causal dynamical triangulations. Using the Krein-Rutman theory of positivity preserving operators we study several properties of the emerging transfer matrix. In particular, we determine regions in the quadrant of parameters $\beta, \mu > 0$ where the infinite-volume free energy converges yielding results on the convergence and asymptotic properties of the partition function and the Gibbs measure.

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5.1 | Introduction. A review of related results

In the study of two-dimensional quantum gravity and non-critical string theory, models of discrete random surfaces play an essential role.

In the 1980s, so-called *dynamical triangulations* (DT) were introduced to define a Euclidean path integral for two-dimensional quantum gravity (see [9] for an overview). In particular, the partition function has been determined as a sum over all possible

triangulations of a sphere where each configuration is weighted by a Boltzmann factor $e^{-\mu|T|}$, with $|T|$ standing for the size of the triangulation and μ being the cosmological constant. The evaluation of the partition function was reduced to a purely combinatorial problem that can be solved with the help of the early work of Tutte [4, 3]; alternatively, more powerful techniques were proposed, based on random matrix models (see, e.g., [98]) and bijections to well-labelled trees (see [99, 100]). One can then pass to a continuum limit by taking the number of triangles to infinity. An interesting property of the resulting “quantum geometry” is its fractal structure as illustrated in Figure 5.4 (a). In the physical literature such fractal structures are called “baby universes”, and they completely dominate the continuum limit leading to a fractal dimension $d = 4$.

From a probabilistic point of view there has recently been an increasing interest in DT, most notably through the work of Angel and Schramm on a uniform measure on infinite planar triangulations [14], as well as through the work of Le Gall, Miermont and collaborators on Brownian maps (see [21] for a recent review).

From a physical point of view it is interesting to study various models of matter, such as the Ising model, coupled to the DT. The calculation of the partition function in this case also reduces to a combinatorial problem. It was first solved in [101, 102] by using random matrix models and later by using a bijection to well-labelled trees [103]. It is interesting that the solution here is much simpler than in the case of a flat triangular or square lattice as given by Onsager [104]. Further, one can see that the critical exponents in the case where the model is coupled to DT differ from the Onsager values. This is related to the strong back-reaction of the Ising model with the quantum geometry. In particular, the spin clusters energetically prefer to sit within baby universes since those are connected to the main universe through a very short so-called bottleneck boundary (see Figure 5.4 (b)). The spins increase the fractal structure leading to a change in the values of the critical exponents at the critical temperature.

In a continuum framework one attempts to understand the resulting theory as Liouville theory coupled to a conformal field theory with central charge $c = 1/2$. Furthermore, this leads to a simple algebraic identity (the KPZ-relation) between the critical exponents of the Ising model on a flat lattice and the critical exponents of the Ising model coupled to DT [105].

While DT has a very rich mathematical structure which very recently has been related to the SLE (the Schramm-Loewner evolution) and level curves of a Gaussian free field [106], from the point of view of quantum gravity its fractal structure leads to

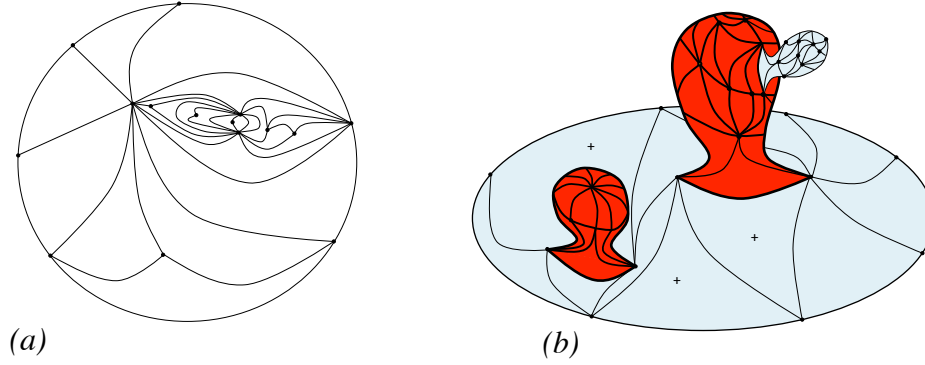


Figure 5.1: (a) A section of a typical planar triangulation of the DT ensemble illustrating the fractal structure of the quantum geometry. (b) Illustration of the baby universes and the formation of spin clusters within them. Each baby universe corresponds to a fractal structure as in part (a) drawn out of the plane.

causality-violating geometries that are arguably non-physical. This led to the development of so-called *causal dynamical triangulations* (CDT) by Ambjørn and Loll [10], to define the *Lorentzian* gravitational path integral. A causal triangulation is formed by triangulations of spatial strips as illustrated in Figure 5.2. Note that the left and right boundaries of the spatial strip are periodically identified.

A first analytical solution of two-dimensional (pure) CDT was obtained in [10] where it was shown that the resulting quantum geometry, while still random, is much more regular than in the case of a DT, leading to a fractal dimension of $d = 2$. From a probabilistic point of view, we would like to note that a uniform measure on infinite causal triangulations UIC_T has been recently introduced in [34, 88].

An interesting question is: What are the properties of the Ising model coupled to a CDT ensemble? As was said above, the CDT ensemble is more regular than that of the DT, but it is still random and allows for a back-reaction of the spin system with the quantum geometry. Monte Carlo simulations [77] (see also [107, 108]) give a strong evidence that critical exponents of the Ising model coupled to CDT are identical to the Onsager values.

While recently much progress has been made in the development of analytical tech-

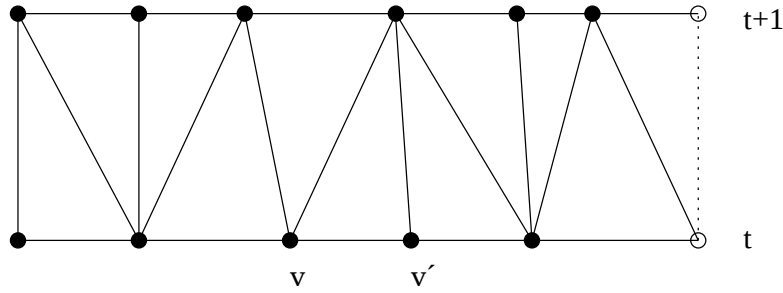


Figure 5.2: A causal triangulation.

niques for CDT [109, 110], particularly random matrix models [90, 54, 89], and their application to multi-critical CDT [93, 94, 111], the causality constraints still makes it difficult to find an analytical solution of the Ising model coupled to CDT. For the quenched Ising model coupled to two-dimensional CDT some progress has been made in proving the existence of a phase transition [52].

In this chapter we develop a transfer matrix formalism for the annealed Ising model coupled to two-dimensional CDT. Spectral properties of the transfer matrix are rigorously analysed by using the Krein-Rutman theorem [59] on operators preserving the cone of positive functions. This yields results on convergence and asymptotic properties of the partition function and the Gibbs measure and allows us to determine regions in the parameter quarter-plane where the partition function converges.

Outline: The chapter is organised as follows. Section 5.2 contains basic definitions. Section 5.3 gives a summary of the transfer-matrix formalism for CDTs. In Section 5.4 we introduce the transfer matrix for the Ising model coupled with a CDT and state and comment on our main results. This is followed with concluding remarks in Section 5.5. An Appendix contains the proofs of the stated results.

5.2 | Definitions

We will work with rooted causal dynamic triangulations of the cylinder $C_N = \mathcal{S} \times [0, N]$, $N = 1, 2, \dots$, which have N bonds (strips) $\mathcal{S} \times [j, j + 1]$. Here $\mathcal{S}$ stands for a unit circle. The definition of a causal triangulation starts by considering a connected graph G embedded in C_N with the property that all faces of G are triangles (using the convention that an edge incident to the same face on two sides counts twice, see Chapter 3 for

more details). A triangulation $\underline{t}$ of C_N is a pair formed by a graph G with the above property and the set F of all its (triangular) faces: $\underline{t} = (G, F)$.

Definition 5.2.1. A triangulation $\underline{t}$ of C_N is called a causal triangulation (CT) if the following conditions hold:

- each triangular face of $\underline{t}$ belongs to some strip $S \times [j, j + 1]$, $j = 1, \dots, N - 1$, and has all vertices and exactly one edge on the boundary $(S \times \{j\}) \cup (S \times \{j + 1\})$ of the strip $S \times [j, j + 1]$;
- if $k_j = k_j(\underline{t})$ is the number of edges on $S \times \{j\}$, then we have $0 < k_j < \infty$ for all $j = 0, 1, \dots, N - 1$.

Definition 5.2.2. A triangulation $\underline{t}$ of C_N is called rooted if it has a root. The root in the triangulation $\underline{t}$ is represented by a triangular face t of $\underline{t}$, called the root triangle, with an anticlock-wise ordering on its vertices (x, y, z) where x and y belong to $S^1 \times \{0\}$. The vertex x is identified as the root vertex and the (directed) edge from x to y as the root edge.

Definition 5.2.3. Two causal rooted triangulations of C_N , say $\underline{t} = (G, F)$ and $\underline{t}' = (G', F')$, are equivalent if there exists a self-homeomorphism of C_N which (i) transforms each slice $S^1 \times \{j\}$, $j = 0, \dots, N - 1$ to itself and preserves its direction, (ii) induces an isomorphism of the graphs G and G' and a bijection between F and F' , and (iii) takes the root of $\underline{t}$ to the root of $\underline{t}'$.

A triangulation $\underline{t}$ of C_N is identified as a consistent sequence

$$\underline{t} = (\underline{t}(0), \underline{t}(1), \dots, \underline{t}(N - 1)),$$

where $\underline{t}(i)$ is a causal triangulation of the strip $S \times [i, i + 1]$. The latter means that each $\underline{t}(i)$ is described by a partition of $S \times [i, i + 1]$ into triangles where each triangle has one vertex on one of the slices $S \times \{i\}$, $S \times \{i + 1\}$ and two on the other, together with the edge joining these two vertices. The property of consistency means that each pair $(\underline{t}(i), \underline{t}(i + 1))$ is consistent, i.e., every side of a triangle from $\underline{t}(i)$ lying in $S \times \{i + 1\}$ serves as a side of a triangle from $\underline{t}(i + 1)$, and vice versa.

The triangles forming the causal triangulation $\underline{t}(i)$ are denoted by $t(i, j)$, $1 \leq j \leq n(\underline{t}(i))$ where, $n(\underline{t}(i))$ stands for the number of triangles in the triangulation $\underline{t}(i)$. The enumeration of these triangles starts with what we call the root triangle in $\underline{t}(i)$; it is

determined recursively as follows. First, we have the root triangle $t(0, 1)$ in $\mathbf{t}(0)$ (see Definition 2.2). Take the vertex of the triangle $t(0, 1)$ which lies on the slice $\mathcal{S} \times \{1\}$ and denote it by x' . This vertex is declared the root vertex for $\mathbf{t}(1)$. Next, the root edge for $\mathbf{t}(1)$ is the one incident to x' and lying on $\mathcal{S} \times \{1\}$, so that if y' is its other end and z' is the third vertex of the corresponding triangle then x', y', z' lists the three vertices anticlock-wise. Accordingly, the triangle with the vertices x', y', z' is called the root triangle for $\mathbf{t}(1)$. This construction can be iterated, determining the root vertices, root edges and root triangles for $\mathbf{t}(i)$, $0 \leq i \leq N - 1$.

It is convenient to introduce the notion of “up” and “down” triangles (see Figure 5.2). We call a triangle $t \in \mathbf{t}(i)$ an up-triangle if it has an edge on the slice $\mathcal{S} \times \{i\}$ and a down-triangle if it has an edge on the slice $\mathcal{S} \times \{i + 1\}$. By Definition 2.1, every triangle is either of type up or down. Let $n_{\text{up}}(\mathbf{t}(i))$ and $n_{\text{do}}(\mathbf{t}(i))$ stand for the number of up- and down-triangles in the triangulation $\mathbf{t}(i)$.

Note that for any edge lying on the slice $\mathcal{S} \times \{i\}$ belongs to exactly two triangles: one up-triangle from $\mathbf{t}(i)$ and one down-triangle from $\mathbf{t}(i - 1)$. This provides the following relation: the number of triangles in the triangulation $\mathbf{t}$ is twice the total number of edges on the slices. More precisely, let n^i be the number of edges on slice $\mathcal{S} \times \{i\}$. Then, for any $i = 0, 1, \dots, N - 1$,

$$n(\mathbf{t}(i)) = n_{\text{up}}(\mathbf{t}(i)) + n_{\text{do}}(\mathbf{t}(i)) = n^i + n^{i+1}, \quad (5.1)$$

implying that

$$\sum_{i=0}^{N-1} n(\mathbf{t}(i)) = 2 \sum_{i=0}^{N-1} n^i. \quad (5.2)$$

There is another useful property regarding the counting of triangulations. Let us fix the number of edges n^i and n^{i+1} in the slices $\mathcal{S} \times \{i\}$ and $\mathcal{S} \times \{i + 1\}$. The number of possible rooted CTs of the slice $\mathcal{S} \times [i, i + 1]$ with n^i up- and n^{i+1} down-triangles is equal to

$$\binom{n^i + n^{i+1} - 1}{n^i - 1} = \binom{n(\mathbf{t}(i)) - 1}{n_{\text{up}}(\mathbf{t}(i)) - 1}. \quad (5.3)$$

5.3 | Transfer matrix formalism for pure CDT

We begin by discussing the case of pure causal dynamical triangulations, as was first introduced in [10] (see also [33] for a mathematically more rigorous account).

The partition function for rooted CTs in the cylinder C_N with periodical spatial boundary conditions (where $\mathbf{t}(0)$ is consistent with $\mathbf{t}(N-1)$) and for the value of the cosmological constant μ is given by

$$Z_N(\mu) = \sum_{\mathbf{t}} e^{-\mu n(\mathbf{t})} = \sum_{(\mathbf{t}(0), \dots, \mathbf{t}(N-1))} \exp \left\{ -\mu \sum_{i=0}^{N-1} n(\mathbf{t}(i)) \right\}. \quad (5.4)$$

Using the properties (5.2) and (5.3) we can represent the partition function (5.4) in the following way

$$Z_N(\mu) = \sum_{n^0 \geq 1, \dots, n^{N-1} \geq 1} \exp \left\{ -2\mu \sum_{i=0}^{N-1} n^i \right\} \prod_{i=0}^{N-1} \binom{n^i + n^{i+1} - 1}{n^i - 1}. \quad (5.5)$$

Moreover, $Z_N(\mu)$ admits a trace-related representation

$$Z_N(\mu) = \text{tr} (U^N). \quad (5.6)$$

This gives rise to a transfer matrix $U = \{u(n, n')\}_{n, n'=1, 2, \dots}$ describing the transition from one spatial strip to the next one. It is an infinite matrix with strictly positive entries

$$u(n, n') = \binom{n + n' - 1}{n - 1} g^{n+n'}. \quad (5.7)$$

For notational convenience we use the parameter $g = e^{-\mu}$ (a single-triangle fugacity). The entry $u(n, n')$ yields the number of possible triangulations of a single strip (say, $S \times [0, 1]$) with n lower boundary edges (on $S \times \{0\}$) and n' upper boundary edges (on $S \times \{1\}$). See Figure 5.2. The asymmetry in n and n' is due to the fact that the lower boundary is marked while the upper one is not. However, a symmetric transfer matrix $\tilde{U} = \{\tilde{u}(n, n')\}$ can be introduced, associated with a strip where both boundaries are kept unmarked:

$$\tilde{u}(n, n') = n^{-1} u(n, n'). \quad (5.8)$$

The N -strip Gibbs distribution $\mathcal{P}_N$ assigns the following probabilities to strings $(n^0, \dots, n^{N-1})$ with the number of triangles $n^i \geq 1$ for all $i = 0, \dots, N-1$:

$$\mathcal{P}_N(n^0, \dots, n^{N-1}) = \frac{1}{Z_N(\mu)} \exp \left\{ -2\mu \sum_{i=0}^{N-1} n^i \right\} \prod_{i=0}^{N-1} \binom{n^i + n^{i+1} - 1}{n^i - 1}. \quad (5.9)$$

We state two lemmas featuring properties of matrix U :

Lemma 5.3.1. *For any $g > 0$ the matrix U and its transpose U^T have an eigenvalue $\Lambda = \Lambda(g)$ given by*

$$\Lambda(g) = \left[(1 - \sqrt{1 - 4g^2}) / (2g) \right]^2. \quad (5.10)$$

The corresponding eigenvectors

$$\phi = \{\phi(n)\}_{n=1,2,\dots} \quad \text{and} \quad \phi^* = \{\phi^*(n)\}_{n=1,2,\dots}$$

have entries

$$\phi(n) = n(\Lambda(g))^n, \quad \phi^*(n) = (\Lambda(g))^n. \quad (5.11)$$

Proof. A direct verification shows that

$$\sum_{n'} u(n, n') n' \Lambda^{n'}(g) = n \Lambda^{n+1}(g) \quad \text{and} \quad \sum_n \Lambda^n(g) u(n, n') = \Lambda^{n'+1}(g).$$

(In fact, each of these relations implies the other.) See Theorem 1 in [33]. $\square$

Lemma 5.3.2. *For any fixed n and any $g < 1$ (equivalently, $\mu > 0$) one has*

$$\sum_{n'} u(n, n') = \left(\frac{g}{1-g} \right)^n (1 - (1-g)^n). \quad (5.12)$$

Proof. The proof again follows from a straightforward verification. $\square$

A transfer-matrix formalism of Statistical Mechanics predicts that, as $N \rightarrow \infty$, the partition function is governed by the largest eigenvalue Λ of the transfer matrix:

$$Z_N(g) = \text{tr } U^N \sim \Lambda^N \quad (5.13)$$

We make this statement more precise in the statements of Lemma 5.3.3 and Theorem 5.3.1 below. Here the symbol ℓ^2 stands for the Hilbert space of square-summable complex sequences (infinite-dimensional vectors) $\psi = \{\psi(n)\}_{n=1,2,\dots}$ equipped with the standard scalar product $\langle \psi', \psi'' \rangle = \sum_n \psi'(n) \overline{\psi''(n)}$. Accordingly, the matrices U and U^T are treated as operators in ℓ^2 .

Lemma 5.3.3. *For any $g < 1/2$ (equivalently $\mu > \ln 2$) the following statements hold true:*

1. U and U^T are bounded operators in ℓ^2 preserving the cone of positive vectors;

2. The sum $\sum_{n,n'} u(n, n') < \infty$. Consequently, U and U^T have

$$\text{tr} (UU^T) = \text{tr} (U^T U) < \infty,$$

i.e., U and U^T are Hilbert-Schmidt operators. Therefore, $\forall N \geq 2$, U^N and $(U^T)^N$ are trace-class operators.

3. The maximal eigenvalue $\Lambda = \Lambda(g)$ of U in ℓ^2 is positive, coincides with the maximal eigenvalue of U^T and is given by Eqn (5.10). The corresponding eigenvectors $\phi, \phi^* \in \ell^2$ are unique up to multiplication by a constant factor and given in Eqn (5.11).

4. The following asymptotical formulas hold as $N \rightarrow \infty$:

$$\frac{1}{\Lambda^N} \text{tr} (U^N), \quad \frac{1}{\Lambda^N} \text{tr} ((U^T)^N) \rightarrow 1,$$

and, $\forall$ vectors $\psi', \psi'' \in \ell^2$,

$$\frac{1}{\Lambda^N} \langle \psi', U^N \psi'' \rangle = \langle \psi', \phi \rangle \langle \phi^*, \psi'' \rangle,$$

where the eigenvectors ϕ and ϕ^* are normalized so that $\langle \phi, \phi^* \rangle = 1$.

Theorem 5.3.1. For any $g < 1/2$ the following relation holds true:

$$\lim_{N \rightarrow \infty} \frac{1}{N} \log Z_N(g) = \log \Lambda \quad (5.14)$$

with $\Lambda = \Lambda(g)$ given in (5.10). Further, the N -strip Gibbs measure $\mathcal{P}_N$ converges weakly to a limiting measure $\mathcal{P}$ which is represented by a positive recurrent Markov chain on $\mathbb{Z}_+ = \{1, 2, \dots\}$, with the transition matrix $P = \{P(n, n')\}_{n=1,2,\dots}$ and the invariant distribution π . Here

$$P(n, n') = \frac{u(n, n') \phi(n')}{\Lambda \phi(n)}$$

and

$$\pi(n) = \frac{\phi^*(n) \phi(n)}{\langle \phi^*, \phi \rangle}.$$

where $\phi(n)$ and $\phi^*(n)$ are as in (5.11).

Proof. The proof is a consequence of Lemma 5.3.1 and 5.3.3 and the Krein-Rutman theory [59]. $\square$

5.4 | An Ising model coupled to the CDT: statement of results

5.4.1 | The model

With any triangle from a triangulation $\mathbf{t}$ we associate a spin taking values ± 1 . An N -strip configuration of spins is represented by a collection

$$\underline{\sigma} = (\sigma(0), \sigma(1), \dots, \sigma(N-1))$$

where $\sigma(i) = \sigma(\mathbf{t}(i))$ is a configuration of spins $\sigma(i, j)$ over triangles $t(i, j)$ forming a triangulation $\mathbf{t}(i)$, $1 \leq j \leq n(\mathbf{t}(i))$. We will say that a single-strip configuration of spins $\sigma(i)$ is supported by a triangulation $\mathbf{t}(i)$ of strip $\mathcal{S} \times [i, i+1]$. We consider a usual (ferromagnetic) Ising-type energy where two spins $\sigma(i, j)$ and $\sigma(i', j')$ interact if their supporting triangles $t(i, j)$, $t(i', j')$ share a common edge; such triangles are called nearest neighbors, and this property is reflected in the notation $\langle \sigma(i, j), \sigma(i', j') \rangle$, where we require $0 \leq i \leq i' \leq N-1$. Thus, in our model each spin has three neighbors. Moreover, a pair $\langle \sigma(i, j), \sigma(i', j') \rangle$ can only occur for $i' - i \leq 1$ or $i = 0, i' = N-1$. Formally, the Hamiltonian of the model reads:

$$\mathbb{H}(\underline{\sigma}) = - \sum_{\langle \sigma(i, j), \sigma(i', j') \rangle} \sigma(i, j) \sigma(i', j'). \quad (5.15)$$

We will use the following decomposition:

$$\mathbb{H}(\underline{\sigma}) = \sum_{i=0}^{N-1} H(\sigma(i)) + \sum_{i=0}^{N-1} V(\sigma(i), \sigma(i+1)), \quad (5.16)$$

where we assume that $\sigma(0) \equiv \sigma(N)$ (the periodic spatial boundary condition). Here $H(\sigma(i))$ represents the energy of the configuration $\sigma(i)$:

$$H(\sigma(i)) = - \sum_{\langle \sigma(i, j), \sigma(i, j') \rangle} \sigma(i, j) \sigma(i, j'). \quad (5.17)$$

Further, $V(\sigma(i), \sigma(i+1))$ is the energy of interaction between neighboring triangles belonging to the adjacent strips $\mathcal{S} \times [i, i+1]$ and $\mathcal{S} \times [i+1, i+2]$:

$$V(\sigma(i), \sigma(i+1)) = - \sum_{\langle \sigma(i, j), \sigma(i+1, j') \rangle} \sigma(i, j) \sigma(i+1, j'). \quad (5.18)$$

The partition function for the (annealed) N -strip Ising model coupled to CDT, at the inverse temperature $\beta > 0$ and for the cosmological constant μ , is given by

$$\begin{aligned} \Xi_N(\mu, \beta) &= \sum_{(\mathbf{t}(0), \dots, \mathbf{t}(N-1))} \exp \left\{ -\mu \sum_{i=0}^{N-1} n(\mathbf{t}(i)) \right\} \\ &\times \sum_{(\boldsymbol{\sigma}(0), \dots, \boldsymbol{\sigma}(N-1))} \prod_{i=0}^{N-1} \exp \left\{ -\beta H(\boldsymbol{\sigma}(i)) - \beta V(\boldsymbol{\sigma}(i), \boldsymbol{\sigma}(i+1)) \right\}. \end{aligned} \quad (5.19)$$

Here $n(\mathbf{t}(i))$ stands for the number of triangles in the triangulation $\mathbf{t}(i)$. Like before, the formula

$$\Xi_N(\mu, \beta) = \text{tr } \mathbf{K}^N \quad (5.20)$$

gives rise to a transfer matrix $\mathbf{K}$ with entries $K((\mathbf{t}, \boldsymbol{\sigma}), (\mathbf{t}', \boldsymbol{\sigma}'))$ labelled by pairs $(\mathbf{t}, \boldsymbol{\sigma}), (\mathbf{t}', \boldsymbol{\sigma}')$ representing triangulations of a single strip (say, $\mathcal{S} \times [0, 1]$) and their supported spin configurations which are positioned next to each other. Formally,

$$\begin{aligned} K((\mathbf{t}, \boldsymbol{\sigma}), (\mathbf{t}', \boldsymbol{\sigma}')) &= \mathbf{1}_{\mathbf{t} \sim \mathbf{t}'} \exp \left\{ -\frac{\mu}{2} (n(\mathbf{t}) + n(\mathbf{t}')) \right\} \\ &\times \exp \left\{ -\frac{\beta}{2} (H(\boldsymbol{\sigma}) + H(\boldsymbol{\sigma}')) - \beta V(\boldsymbol{\sigma}, \boldsymbol{\sigma}') \right\}. \end{aligned} \quad (5.21)$$

As earlier, $n(\mathbf{t})$ and $n(\mathbf{t}')$ are the numbers of triangles in the triangulations $\mathbf{t}$ and $\mathbf{t}'$. The indicator $\mathbf{1}_{\mathbf{t} \sim \mathbf{t}'}$ means that the triangulations $\mathbf{t}, \mathbf{t}'$ have to be consistent with each other in the above sense: the number of down-triangles in $\mathbf{t}$ should equal the number of up-triangles in $\mathbf{t}'$, and an upper-marked edge in $\mathbf{t}$ should coincide with a lower-marked edge in triangulation $\mathbf{t}'$. It means that the pair $(\mathbf{t}, \mathbf{t}')$ forms a CDT for the strip $\mathcal{S} \times [0, 2]$.

We would like to stress that the trace $\text{tr } \mathbf{K}^N$ in (5.20) is understood as the *matrix trace*, i.e., as the sum $\sum_{\mathbf{t}, \boldsymbol{\sigma}} K^{(N)}((\mathbf{t}, \boldsymbol{\sigma}), (\mathbf{t}, \boldsymbol{\sigma}))$ of the diagonal entries $K^{(N)}((\mathbf{t}, \boldsymbol{\sigma}), (\mathbf{t}, \boldsymbol{\sigma}))$ of the matrix $\mathbf{K}^N$. (Indeed, in what follows, the notation “tr” is used for the matrix trace only.) Our aim will be to verify that the matrix trace in (5.20) can be replaced with an *operator trace* invoking the eigenvalues of $\mathbf{K}$ in a suitable linear space.

As before, we can introduce the N -strip Gibbs probability distribution associated with formula (5.19):

$$\begin{aligned} \mathbb{P}_N((\mathbf{t}(0), \boldsymbol{\sigma}(0)), \dots, (\mathbf{t}(N-1), \boldsymbol{\sigma}(N-1))) \\ = \frac{1}{\Xi(\mu, \beta)} \prod_{i=0}^{N-1} \exp \left\{ -\mu n(\mathbf{t}(i)) - \beta H(\boldsymbol{\sigma}(i)) - \beta V(\boldsymbol{\sigma}(i), \boldsymbol{\sigma}(i+1)) \right\}. \end{aligned} \quad (5.22)$$

Consider several special cases of interest.

The case $\beta \approx 0$. This is the first term of the so-called high temperature expansion [77]. Here one has

$$\begin{aligned}\Xi(\mu, 0) &= \sum_{(\mathbf{t}(0), \dots, \mathbf{t}(N-1))} \exp\left\{-\mu \sum_{i=0}^{N-1} n(\mathbf{t}(i))\right\} \sum_{(\boldsymbol{\sigma}(0), \dots, \boldsymbol{\sigma}(N-1))} 1 \\ &= \sum_{n^0 \geq 1, \dots, n^{N-1} \geq 1} \exp\left\{-2(\mu - \ln 2) \sum_{i=0}^{N-1} n^i\right\} \prod_{i=0}^{N-1} \binom{n^i + n^{i+1} - 1}{n^i - 1} \\ &= Z_N(\mu - \ln 2); \text{ cf. (5.4).}\end{aligned}$$

The condition $\mu - \ln 2 > \ln 2$ which guarantees properties listed in Lemma 5.3.3 and Theorem 5.3.1 results in

$$\mu > 2 \ln 2. \quad (5.23)$$

Thus, Eqn. (5.23) yields a sub-criticality condition when $\beta = 0$.

The case $\beta \approx \infty$. Observe that for any triangulation $\mathbf{t} = (\mathbf{t}(0), \dots, \mathbf{t}(N-1))$ there are two ground states: all spins $+1$ and all spins -1 , with the overall energy equals minus three half times the total number of triangles: $-3/2 \sum_{i=0}^{N-1} n(\mathbf{t}(i))$. Discarding all other spin configurations, we obtain that

$$\Xi(\mu, \beta) > \Xi_*(\mu, \beta)$$

where

$$\begin{aligned}\Xi_*(\mu, \beta) &= \sum_{\mathbf{t}(0), \dots, \mathbf{t}(N-1)} 2 \exp\left\{\left(-\mu + \frac{3}{2}\beta\right) \sum_{i=0}^{N-1} n(\mathbf{t}(i))\right\} \\ &= 2 \sum_{n^0 \geq 1, \dots, n^{N-1} \geq 1} \exp\left\{-2\left(\mu - \frac{3}{2}\beta\right) \sum_{i=0}^{N-1} n^i\right\} \binom{n^i + n^{i+1} - 1}{n^i - 1} \\ &= 2Z_N\left(\mu - \frac{3}{2}\beta\right)\end{aligned}$$

where $\exp\left[\frac{3}{2}\beta \sum_i n(\mathbf{t}(i))\right]$ is the energy of the $(+)$ -configuration (or, equivalently, the $(-)$ -configuration). For β large, we can expect that $\Xi(\mu, \beta) \sim \Xi_*(\mu, \beta)$. Then the critical inequality

$$\mu - \frac{3}{2}\beta > \ln 2$$

yields

$$\mu > \ln 2 + \frac{3}{2}\beta. \quad (5.24)$$

Equation (5.24) gives a necessary (and probably tight) criticality condition for the Ising model under consideration for large values of β . A similar result was obtained in [77].

The case $0 < \beta < \infty$. Firstly, we note that for any fixed triangulation $\underline{t}$ the energy of any spin configuration $\underline{\sigma}$ on $\underline{t}$ will be bigger or equal than the energy of a pure configuration (all +s or all -s):

$$\begin{aligned} H(\underline{\sigma}) &= \sum_j H(\sigma(j)) + \sum_j V(\sigma(j), \sigma(j+1)) \\ &\geq -\frac{3}{2} \#(\text{of all triangles in } \underline{t}) = -3 \sum_{i=0}^{N-1} n^i, \end{aligned}$$

where n^i is the number of edges in the i th level $S \times \{i\}$, $i = 0, 1, \dots, N-1$. Thus, for any $\beta > 0$ the inequality $\Xi(\mu, \beta) < \Xi^*(\mu, \beta)$ holds true, where

$$\begin{aligned} \Xi^*(\mu, \beta) &= \sum_{(\mathbf{t}(0), \dots, \mathbf{t}(N-1))} \exp \left\{ \left(-\mu + \frac{3}{2}\beta + \ln 2 \right) \sum_{i=0}^{N-1} n(\mathbf{t}(i)) \right\} \\ &= \sum_{n^0 \geq 1, \dots, n^{N-1} \geq 1} \exp \left\{ -2 \left(\mu - \frac{3}{2}\beta - \ln 2 \right) \sum_{i=0}^{N-1} n^i \right\} \\ &= Z_N \left(\mu - \frac{3}{2}\beta - \ln 2 \right). \end{aligned}$$

Hence, the inequality

$$\mu - \frac{3}{2}\beta - \ln 2 > \ln 2 \quad \text{or} \quad \mu > 2 \ln 2 + \frac{3}{2}\beta \quad (5.25)$$

provides a sufficient condition for subcriticality of the Ising model under consideration.

5.4.2 | The transfer-matrix $\mathbf{K}$ and its powers $\mathbf{K}^N$

The main results of this chapter are summarized in Lemma 5.4.1 and Theorems 5.4.1 and 5.4.2 below. Let us start with a statement (see Proposition 5.4.1 below) which

merely re-phrases standard definitions and explains our interest in the matrices $\mathbf{K}$, $\mathbf{K}^\top$, $\mathbf{K}^\top \mathbf{K}$, $\mathbf{K} \mathbf{K}^\top$ and their powers. Cf. Definition 2.2.2 on p.83, Definition 2.4.1 on p.101, Lemma 2.3.1 on p.85 and Theorem 3.3.13 on p.139 in [112]).

We treat the transfer-matrix $\mathbf{K}$ and its transpose $\mathbf{K}^\top$ as linear operators in the Hilbert space ℓ_{T-C}^2 (the subscript T-C refers to triangulations and spin-configurations). The space ℓ_{T-C}^2 is formed by functions $\psi = \{\psi(\mathbf{t}, \sigma)\}$ with the argument $(\mathbf{t}, \sigma)$ running over single-strip triangulations and supported configurations of spins, with the scalar product $\langle \psi', \psi'' \rangle_{T-C} = \sum_{\mathbf{t}, \sigma} \psi'(\mathbf{t}, \sigma) \overline{\psi''(\mathbf{t}, \sigma)}$ and the induced norm $\|\psi\|_{T-C}$. The action of $\mathbf{K}$ in ℓ_{T-C}^2 , in the basis formed by Dirac's delta-vectors $\delta_{(\mathbf{t}, \sigma)}$, is determined by

$$(\mathbf{K}\psi)(\mathbf{t}, \sigma) = \sum_{\mathbf{t}', \sigma'} K((\mathbf{t}, \sigma), (\mathbf{t}', \sigma')) \psi(\mathbf{t}', \sigma'); \quad (5.26)$$

in following we use the notation $\mathbf{K}$, $\mathbf{K}^\top$, etc., for the matrices and the corresponding operators in ℓ_{T-C}^2 . Accordingly, the symbols $\|\mathbf{K}\|_{T-C}$, $\|\mathbf{K}^\top\|_{T-C}$ etc. refer to norms in ℓ_{T-C}^2 .

Given $n = 1, 2, \dots$, suppose that the operator $\mathbf{K}^n$ (respectively, $(\mathbf{K}^\top)^n$) is of trace class. Then the following series absolutely converges:

$$\sum_j \Lambda_j^{(n)} \left(\text{respectively, } \sum_j \Lambda_j^{*(n)} \right), \quad (5.27)$$

where $\Lambda_j^{(n)}$ ($\Lambda_j^{*(n)}$) runs through the eigenvalues of $\mathbf{K}^n$ ($(\mathbf{K}^\top)^n$), counted with their multiplicities. In this case the sum (5.27) is called the operator trace of $\mathbf{K}^n$ (respectively, $(\mathbf{K}^\top)^n$) in ℓ_{T-C}^2 . We adopt an agreement that the eigenvalues in (5.27) are listed in the decreasing order of their moduli, beginning with $\Lambda_0^{(n)}$ ($\Lambda_0^{*(n)}$).

Set $|\mathbf{K}^n| = \sqrt{(\mathbf{K}^\top)^n \mathbf{K}^n}$ and $|(\mathbf{K}^\top)^n| = \sqrt{\mathbf{K}^n (\mathbf{K}^\top)^n}$.

Proposition 5.4.1. *For any positive integer r , the following inequalities are equivalent:*

$$\begin{aligned} \text{tr}((\mathbf{K}^r (\mathbf{K}^\top)^r)) &= \text{tr}((\mathbf{K}^\top)^r \mathbf{K}^r) < \infty \text{ and} \\ \text{tr}|\mathbf{K}^{2r}| &= \text{tr}|(\mathbf{K}^\top)^{2r}| < \infty. \end{aligned} \quad (5.28)$$

Moreover, each of the inequalities in (5.28) implies that $\forall N \geq 2r$, the operators $\mathbf{K}^N$ and $(\mathbf{K}^\top)^N$ are of trace class in ℓ_{T-C}^2 . Hence, for $N \geq 2r$, the matrix traces $\text{tr}(\mathbf{K}^N)$ and $\text{tr}((\mathbf{K}^\top)^N)$ are finite and coincide with the corresponding operator traces in ℓ_{T-C}^2 .

Theorem 5.4.1. *Suppose that the condition (5.28) is satisfied with $r = 1$. Then the following properties of transfer matrix $\mathbf{K}$ are fulfilled.*

1. The square $\mathbf{K}^2$ and its transpose $(\mathbf{K}^\top)^2$ are trace-class operators in $\ell_{\text{T-C}}^2$.
2. $\mathbf{K}$ and $\mathbf{K}^\top$ have a common eigenvalue, $\Lambda = \Lambda_0(\beta, \mu) > 0$ such that the norms $\|\mathbf{K}\|_{\text{T-C}} = \|\mathbf{K}^\top\|_{\text{T-C}} = \Lambda$. Furthermore, $\mathbf{K}^2$ and $(\mathbf{K}^\top)^2$ have the common eigenvalue $\Lambda^2 = \Lambda_0^{(2)} = \Lambda_0^{*(2)}$ such that the norms $\|\mathbf{K}^2\|_{\text{T-C}} = \|(\mathbf{K}^\top)^2\|_{\text{T-C}} = \Lambda^2$.
3. Λ is a simple eigenvalue of $\mathbf{K}$ and $\mathbf{K}^\top$, i.e., the corresponding eigenvectors $\phi = \{\phi(\mathbf{t}, \sigma)\}$ and $\phi^* = \{\phi^*(\mathbf{t}, \sigma)\}$ are unique up to multiplicative constants. Moreover, ϕ and $\phi^\top$ can be made strictly positive: $\phi(\mathbf{t}, \sigma), \phi^\top(\mathbf{t}, \sigma) > 0 \forall (\mathbf{t}, \sigma)$. Furthermore, Λ is separated from the remaining singular values and the remaining eigenvalues of $\mathbf{K}$ and $\mathbf{K}^\top$ by a positive gap. The same is true for Λ^2 and $\mathbf{K}^2$ and $(\mathbf{K}^\top)^2$.

Proof of Theorem 5.4.1. Because the entries $K((\mathbf{t}, \sigma), (\mathbf{t}', \sigma'))$ are non-negative, the condition (5.28) with $r = 1$ means that

$$\sum_{(\mathbf{t}, \sigma), (\mathbf{t}', \sigma')} K^2((\mathbf{t}, \sigma), (\mathbf{t}', \sigma')) < \infty, \quad (5.29)$$

that is, $\mathbf{K}$ and $\mathbf{K}^\top$ are Hilbert-Schmidt operators. It means that the operator $\mathbf{K}\mathbf{K}^\top$ has an orthonormal basis of eigenvectors and the series of squares of its eigenvalues (counted with multiplicities) converges and gives the trace $\text{tr}_{\text{T-C}}(\mathbf{K}\mathbf{K}^\top)$. Consequently, the operators $\mathbf{K}$ and $\mathbf{K}^\top$ are bounded (and even completely bounded) and $\mathbf{K}^2$ and $(\mathbf{K}^\top)^2$ are of trace class. The latter fact means that the matrix trace of the operator $\mathbf{K}^2$ coincides with its operator trace in $\ell_{\text{T-C}}^2$, and the same is true of $(\mathbf{K}^\top)^2$. In addition, the operator $\mathbf{K}^2$ has the property that its matrix entries $K^{(2)}((\mathbf{t}, \sigma), (\mathbf{t}', \sigma'))$ are strictly positive:

$$K^{(2)}((\mathbf{t}, \sigma), (\mathbf{t}', \sigma')) = \sum_{(\tilde{\mathbf{t}}, \tilde{\sigma})} K((\mathbf{t}, \sigma), (\tilde{\mathbf{t}}, \tilde{\sigma})) K((\tilde{\mathbf{t}}, \tilde{\sigma}), (\mathbf{t}', \sigma')) > 0. \quad (5.30)$$

The Krein–Rutman theory (see [59], Proposition VII') guarantees that both $\mathbf{K}$ and $\mathbf{K}^\top$ have a maximal eigenvalue Λ that is positive and non-degenerate, or simple. That is, the eigenvector ϕ of $\mathbf{K}$ and the eigenvector ϕ^* of $\mathbf{K}^\top$ corresponding with Λ are unique up to multiplication by a constant, and all entries $\phi(\mathbf{t}, \sigma)$ and $\phi^*(\mathbf{t}, \sigma)$ are non-zero and have the same sign. In other words, the entries $\phi(\mathbf{t}, \sigma)$ and $\phi^*(\mathbf{t}, \sigma)$ can be made positive. The spectral gaps are also consequences of the above properties. $\square$

Set:

$$\lambda(\mu, \beta) = c^2 (m^2 + 1) (\cosh 2\beta) \left(1 + \sqrt{1 - \frac{1}{(\cosh 2\beta)^2} \frac{(m^2 - 1)^2}{(m^2 + 1)^2}} \right) \quad (5.31)$$

where c and m are determined by

$$c = \frac{\exp(\beta - \mu)}{e^{2\beta}(1 - \exp(\beta - \mu))^2 - e^{-2\mu}} \quad (5.32)$$

$$m = e^{2\beta} + (1 - e^{4\beta}) \exp(-(\beta + \mu)). \quad (5.33)$$

Lemma 5.4.1. *For any $\beta, \mu > 0$ such that*

$$\lambda(\mu, \beta) < 1, \quad (5.34)$$

the condition (5.28) is satisfied for $r = 1$:

$$\text{tr}(\mathbf{K}\mathbf{K}^\top) = \text{tr}(\mathbf{K}^\top\mathbf{K}) < \infty \text{ and } \text{tr}|\mathbf{K}^2| = \text{tr}|(\mathbf{K}^\top)^2| < \infty, \quad (5.35)$$

implying the assertions of Proposition 5.4.1 and Theorem 5.4.1. Moreover, the condition (5.28) implies (5.34)

The proof of Lemma 5.4.1 is given in the next section. Here we only remark that the proof is based on the following representation of the trace (5.35): there exists matrices Q, Q_m (see formulas (5.47) and (5.48)) with positive entries and of size 4×4 such that (see formula (5.51))

$$\text{tr}(\mathbf{K}\mathbf{K}^\top) = \text{tr}\left(\left(\sum_{k \geq 1} Q^k\right) Q_m\right) + \dots$$

The convergence of the matrix series $\sum_{k \geq 1} Q^k$ is equivalent to the condition that the maximal eigenvalue of the matrix Q is less than 1. This is exactly the condition (5.34).

Theorem 5.4.2. *Under condition (5.34), the following limit holds:*

$$\lim_{N \rightarrow \infty} \frac{1}{N} \log \Xi_N(\beta, \mu) = \log \mathbf{\Lambda}. \quad (5.36)$$

Moreover, as $N \rightarrow \infty$, the N -strip Gibbs measure $\mathbb{P}_N$ (see Eqn (5.22)) converges weakly to a limiting probability distribution $\mathbb{P}$ that is represented by a positive recurrent Markov chain with states $(\mathbf{t}, \sigma)$, the transition matrix

$\mathbf{P} = \{P((\mathbf{t}, \sigma), (\mathbf{t}', \sigma'))\}$ and the invariant distribution $\pi = \{\pi(\mathbf{t}, \sigma)\}$ where

$$\begin{aligned} P((\mathbf{t}, \sigma), (\mathbf{t}', \sigma')) &= \frac{K((\mathbf{t}, \sigma), (\mathbf{t}', \sigma'))\phi(\mathbf{t}', \sigma')}{\Lambda\phi(\mathbf{t}, \sigma)} \\ \pi(\mathbf{t}, \sigma) &= \phi(\mathbf{t}, \sigma)\phi^\top(\mathbf{t}, \sigma) / \langle \phi, \phi^\top \rangle_{\mathbf{T}-\mathbf{C}} \end{aligned}$$

with the norm $\|\phi\|_{\mathbf{T}-\mathbf{C}}^2 = \sum_{\mathbf{t}, \sigma} \phi(\mathbf{t}, \sigma)^2$.

Proof of Theorem 5.4.2. The spectral gap for $\mathbf{K}$ implies that $\forall \psi \in \ell_{\mathbf{T}-\mathbf{C}}^2$, we have the convergence

$$\lim_{N \rightarrow \infty} \frac{1}{\Lambda^N} \mathbf{K}^N \psi = (\langle \psi, \phi \rangle_{\mathbf{T}-\mathbf{C}}) \phi$$

in the norm of space $\ell_{\mathbf{T}-\mathbf{C}}^2$. Moreover, let Π denote the operator of projection to the subspace spanned by the eigenvectors of $\mathbf{K}$ different from ϕ . Then

$$\frac{1}{\Lambda} \|\Pi \mathbf{K} \Pi\|_{\mathbf{T}-\mathbf{C}} < 1 \implies \lim_{N \rightarrow \infty} \frac{1}{\Lambda^N} \|(\Pi \mathbf{K} \Pi)^N\|_{\mathbf{T}-\mathbf{C}} = 0.$$

In turn, this implies that

$$\frac{1}{N} \log \Xi_N(\mu, \beta) = \frac{1}{N} \log \text{tr}_{\mathbf{T}-\mathbf{C}} \mathbf{K}^N \rightarrow \log \Lambda.$$

Convergence of the Gibbs measure $\mathbb{P}_N$ follows as a corollary. $\square$

5.5 | Concluding remarks

This chapter makes a step towards determining the subcriticality domain for an Ising-type model coupled to two-dimensional causal dynamical triangulations (CDT). In doing so we employ transfer-matrix techniques and in particular the Krein-Rutman theorem. We complement the discussion of the previous sections with the following two concluding remarks:

Remark 5.5.1. *It is instructive to summarise the logical structure of the argument establishing Lemma 5.4.1 and Theorems 5.4.1 and 5.4.2:*

- First, (5.35) holds iff condition (5.34) holds: see the proof of Lemma 5.4.1.

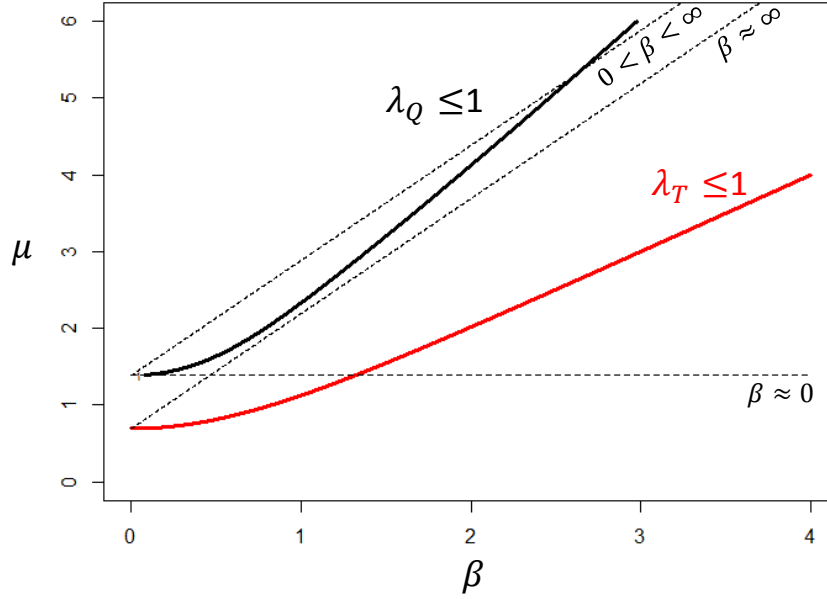


Figure 5.3: $\lambda_Q = \lambda$ and λ_T are the maximal eigenvalues of the matrix Q and a related matrix T respectively (see Appendix 5.A). The area above the black curve is where the condition (5.34) holds true.

- Next, (5.35) implies that $\mathbf{K}$ is a Hilbert–Schmidt operator and $\mathbf{K}^2$ is a trace class operator in $\ell^2_{T-\mathbb{C}}$.
- The last fact, together with the property of positivity (5.30), allow us to use the Krein–Rutman theory, deriving all assertions of Theorems 5.4.1 and 5.4.2.

On the other hand, if (5.34) fails (and therefore (5.35) fails), it does not necessarily mean that the assertions Theorems 5.4.1 and 5.4.2 fail. In other words, we do not claim that the boundary of the domain of parameters β and μ where the model exhibits uncritical behavior is given by Eqn. (5.34). Moreover, Figure 5.3 shows the result of a numerical calculation indicating that the condition (5.34) is worse than (5.25) for (moderately) large values of β .

An apparent condition closer to necessity is the pair of inequalities (5.28) for some (possibly) large r . This issue needs a further study.

Remark 5.5.2. Physical considerations suggest that the critical curve in the (β, g) quarter-plane would have some predictable patterns of behavior: as a function of β , it would decay and exhibit a first-order singularity at a unique point $\beta = \beta_{\text{cr}} \in (0, \infty)$.

A plausible conjecture is that the boundary of the critical domain coincides with the locus of points (β, μ) where Λ loses either the property of positivity or the property of being a simple eigenvalue. This direction also requires further research.

5.A | Proof of Lemma 5.4.1.

By definition the trace (5.35) we need to calculate the series

$$\begin{aligned}
 \text{tr}(\mathbf{K}^T \mathbf{K}) &= \sum_{(\mathbf{t}, \sigma)} \mathbf{K}^T \mathbf{K}((\mathbf{t}, \sigma), (\mathbf{t}, \sigma)) \\
 &= \sum_{(\mathbf{t}, \sigma), (\mathbf{t}', \sigma')} K((\mathbf{t}, \sigma), (\mathbf{t}', \sigma')) K((\mathbf{t}, \sigma), (\mathbf{t}', \sigma')) \\
 &= \sum_{(\mathbf{t}, \sigma), (\mathbf{t}', \sigma')} K^2((\mathbf{t}, \sigma), (\mathbf{t}', \sigma')). \tag{5.37}
 \end{aligned}$$

A single-strip triangulation $\mathbf{t}$ consists of up- and down-triangles. Accordingly, it is convenient to employ new labels for spins: if a triangle $t(l)$ is an l th up-triangle then we denote it by t_{up}^l ; the corresponding spin $\sigma(j)$ will be denoted by σ_{up}^l . Similarly, if $t(j)$ is an l th down-triangle then we denote it by t_{do}^l ; the spin $\sigma(j)$ will be denoted by σ_{do}^l . Consequently, the triangulation $\mathbf{t}$ and its supported spin-configuration σ are represented as

$$\mathbf{t} := (\mathbf{t}_{\text{up}}, \mathbf{t}_{\text{do}}) \text{ and } \sigma := (\sigma_{\text{up}}, \sigma_{\text{do}}).$$

Here

$$\mathbf{t}_{\text{up}} = (t_{\text{up}}^1, \dots, t_{\text{up}}^n), \quad \mathbf{t}_{\text{do}} = (t_{\text{do}}^1, \dots, t_{\text{do}}^m),$$

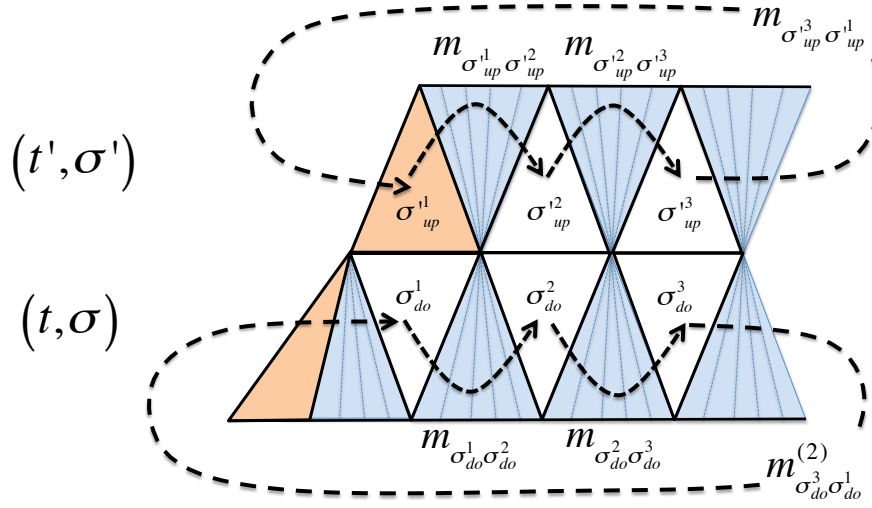
and

$$\sigma_{\text{up}} = (\sigma_{\text{up}}^1, \dots, \sigma_{\text{up}}^n), \quad \sigma_{\text{do}} = (\sigma_{\text{do}}^1, \dots, \sigma_{\text{do}}^m),$$

assuming that the supporting single-strip triangulation $\mathbf{t}$ contains n up-triangles and m down-triangles. (The actual order of up- and down-triangles and supported spins does not matter.)

The same can be done for the pair $(\mathbf{t}', \sigma')$ as illustrated in (5.37). Let recall that the triangulations $\mathbf{t}$ and $\mathbf{t}'$ are consistent ($\mathbf{t} \sim \mathbf{t}'$) iff number of the down-triangles in $\mathbf{t}$ equals that of up-triangles in $\mathbf{t}'$.

To calculate the sum (5.37) we divide the summation over $(\mathbf{t}', \sigma')$ into a summation over $(\mathbf{t}'_{\text{up}}, \sigma'_{\text{up}})$ and $(\mathbf{t}'_{\text{do}}, \sigma'_{\text{do}})$. Firstly, fix a pair $(\mathbf{t}'_{\text{up}}, \sigma'_{\text{up}})$ and make the sum over

Figure 5.4: Interaction between the elements z_i, η_i

$(\mathbf{t}'_{do}, \sigma'_{do})$. Note that the term $V((\mathbf{t}, \sigma), (\mathbf{t}', \sigma'))$ depends only on σ_{do} and σ'_{up} . Consequently,

$$\sum_{\mathbf{t}'_{do}, \sigma'_{do}} K^2((\mathbf{t}, \sigma), (\mathbf{t}', \sigma')) \quad (5.38)$$

$$= e^{-\beta H(\sigma)} e^{-2\beta V((\mathbf{t}, \sigma), (\mathbf{t}', \sigma'))} e^{-\mu n(\mathbf{t})} \sum_{(\mathbf{t}'_{do}, \sigma'_{do})} e^{-\beta H(\sigma')} e^{-\mu n(\mathbf{t}')}.$$

The sum in the right-hand side of (5.38) can be represented in a matrix form. Denote by $e_{\pm 1}$ the standard spin-1/2 unit vectors in $\mathbb{R}^2$:

$$e_{+1} = \begin{pmatrix} 1 \\ 0 \end{pmatrix} \quad \text{and} \quad e_{-1} = \begin{pmatrix} 0 \\ 1 \end{pmatrix}.$$

Next, let us introduce a 2×2 matrix T where

$$T = e^{-\mu} \begin{pmatrix} e^{\beta} & e^{-\beta} \\ e^{-\beta} & e^{\beta} \end{pmatrix} := \begin{pmatrix} t_{++} & t_{+-} \\ t_{-+} & t_{--} \end{pmatrix}. \quad (5.39)$$

Denote by $n(i), i = 1, \dots, n_{up}(\mathbf{t}')$ the number of down-triangles in $\mathbf{t}'$ which are between

the i th and $(i + 1)$ th up-triangles in $\mathbf{t}'$. Let $n_{up}(\mathbf{t}') = k$ then

$$\begin{aligned} \sum_{\mathbf{t}'_{do}, \boldsymbol{\sigma}'_{do}} e^{-\beta H(\boldsymbol{\sigma}')} e^{-\mu n(\mathbf{t}')} &= \sum_{n(i) \geq 0: \sum_i n(i) \geq 1} \prod_{l=1}^k \left(e_{\sigma'_{up}^l}^T T^{n(l)+1} e_{\sigma'_{up}^{l+1}} \right) \\ &= \prod_{l=1}^k \left(e_{\sigma'_{up}^l}^T M e_{\sigma'_{up}^{l+1}} \right) - \prod_{l=1}^k \left(e_{\sigma'_{up}^l}^T T e_{\sigma'_{up}^{l+1}} \right) \end{aligned} \quad (5.40)$$

where the matrix M is the sum of the geometric progression

$$M = \sum_{n=1}^{\infty} T^n := \begin{pmatrix} m_{++} & m_{+-} \\ m_{-+} & m_{--} \end{pmatrix}. \quad (5.41)$$

Using the same procedure we can obtain the sum over all up-triangles into the triangulation $\mathbf{t}$. The only difference is the existence of marked up-triangle in the strip: let as before $n_{up}(\mathbf{t}') = n_{do}(\mathbf{t}) = k$ then

$$\sum_{\mathbf{t}_{up}, \boldsymbol{\sigma}_{up}} e^{-\beta H(\boldsymbol{\sigma})} e^{-\mu n(\mathbf{t})} = \prod_{l=1}^{k-1} \left(e_{\sigma_{up}^l}^T M e_{\sigma_{up}^{l+1}} \right) \left(e_{\sigma_{up}^k}^T M^2 e_{\sigma_{up}^1} \right) \quad (5.42)$$

See Figure 5.4 for illustration of these calculations (5.40) and (5.42). Further, supposing the existence of the matrix M and using (5.40) and (5.42) we obtain the following:

$$\begin{aligned} \sum_{\mathbf{t}_{up}, \boldsymbol{\sigma}_{up}} \sum_{\mathbf{t}'_{do}, \boldsymbol{\sigma}'_{do}} K^2((\mathbf{t}, \boldsymbol{\sigma}), (\mathbf{t}', \boldsymbol{\sigma}')) &= e^{-2\beta V((\mathbf{t}_{do}, \boldsymbol{\sigma}_{do}), (\mathbf{t}'_{up}, \boldsymbol{\sigma}'_{up}))} \\ &\times \sum_{\mathbf{t}_{up}, \boldsymbol{\sigma}_{up}} e^{-\beta H(\boldsymbol{\sigma})} e^{-\mu n(\mathbf{t})} \sum_{(\mathbf{t}'_{do}, \boldsymbol{\sigma}'_{do})} e^{-\beta H(\boldsymbol{\sigma}')} e^{-\mu n(\mathbf{t}')} \\ &= e^{-2\beta V((\mathbf{t}_{do}, \boldsymbol{\sigma}_{do}), (\mathbf{t}'_{up}, \boldsymbol{\sigma}'_{up}))} \\ &\times \left[\prod_{l=1}^k \left(e_{\sigma'_{up}^l}^T M e_{\sigma'_{up}^{l+1}} \right) \prod_{l=1}^{k-1} \left(e_{\sigma_{do}^l}^T M e_{\sigma_{do}^{l+1}} \right) \left(e_{\sigma_{up}^k}^T M^2 e_{\sigma_{up}^1} \right) \right. \\ &\quad \left. - \prod_{l=1}^k \left(e_{\sigma'_{up}^l}^T T e_{\sigma'_{up}^{l+1}} \right) \prod_{l=1}^{k-1} \left(e_{\sigma_{do}^l}^T M e_{\sigma_{do}^{l+1}} \right) \left(e_{\sigma_{up}^k}^T M^2 e_{\sigma_{up}^1} \right) \right]. \end{aligned} \quad (5.43)$$

Necessary and sufficient condition for the convergence of the matrix series for M is that the maximal eigenvalue of matrix T is less than 1. The eigenvalues of T are

$$\lambda_{\pm} = e^{(\beta-\mu)} \pm e^{-(\beta+\mu)}, \quad (5.44)$$

and the above condition means that $\lambda_+ < 1$ or, equivalently,

$$\mu > \ln(2\cosh(\beta)). \quad (5.45)$$

Under this condition (5.45), the matrix M is calculated explicitly:

$$M = \frac{e^{(\beta-\mu)}}{e^{2\beta}(1 - e^{(\beta-\mu)})^2 - e^{-2\mu}} \times \begin{pmatrix} e^{2\beta} + (1 - e^{4\beta})e^{-(\beta+\mu)} & 1 \\ 1 & e^{2\beta} + (1 - e^{4\beta})e^{-(\beta+\mu)} \end{pmatrix}. \quad (5.46)$$

We are now in a position to calculate the sum in (5.37). To this end, we again represent it through the product of transfer matrices. Pictorially, we express the above sum as the partition function of a one-dimensional Ising-type model where states are pairs of spins $(\sigma_{\text{do}}^I, \sigma_{\text{up}}^I)$ and the interaction is via the matrix T between the members of the pair and via matrix M between neighboring pairs. More precisely, define the following 4×4 matrices:

$$Q = \begin{pmatrix} e^{2\beta} m_{++} m_{++} & m_{++} m_{+-} & m_{+-} m_{++} & e^{2\beta} m_{+-} m_{+-} \\ m_{++} m_{-+} & e^{-2\beta} m_{++} m_{--} & e^{-2\beta} m_{+-} m_{-+} & m_{+-} m_{--} \\ m_{-+} m_{++} & e^{-2\beta} m_{-+} m_{+-} & e^{-2\beta} m_{--} m_{++} & m_{--} m_{+-} \\ e^{2\beta} m_{-+} m_{-+} & m_{-+} m_{--} & m_{--} m_{-+} & e^{2\beta} m_{--} m_{--} \end{pmatrix} \quad (5.47)$$

$$Q_m = \begin{pmatrix} e^{2\beta} m_{++} m_{++}^{(2)} & m_{++} m_{+-}^{(2)} & m_{+-} m_{++}^{(2)} & e^{2\beta} m_{+-} m_{+-}^{(2)} \\ m_{++} m_{-+}^{(2)} & e^{-2\beta} m_{++} m_{++}^{(2)} & e^{-2\beta} m_{+-} m_{-+}^{(2)} & m_{+-} m_{--}^{(2)} \\ m_{-+} m_{++}^{(2)} & e^{-2\beta} m_{-+} m_{+-}^{(2)} & e^{-2\beta} m_{--} m_{++}^{(2)} & m_{--} m_{+-}^{(2)} \\ e^{2\beta} m_{-+} m_{-+}^{(2)} & m_{-+} m_{--}^{(2)} & m_{--} m_{-+}^{(2)} & e^{2\beta} m_{--} m_{--}^{(2)} \end{pmatrix} \quad (5.48)$$

$$Q_t = \begin{pmatrix} e^{2\beta} t_{++} m_{++} & t_{++} m_{+-} & t_{+-} m_{++} & e^{2\beta} t_{+-} m_{+-} \\ t_{++} m_{-+} & e^{-2\beta} t_{++} m_{--} & e^{-2\beta} t_{+-} m_{-+} & t_{+-} m_{--} \\ t_{-+} m_{++} & e^{-2\beta} t_{-+} m_{+-} & e^{-2\beta} t_{--} m_{++} & t_{--} m_{+-} \\ e^{2\beta} t_{-+} m_{-+} & t_{-+} m_{--} & t_{--} m_{-+} & e^{2\beta} t_{--} m_{--} \end{pmatrix} \quad (5.49)$$

$$Q_{tm} = \begin{pmatrix} e^{2\beta} t_{++} m_{++}^{(2)} & t_{++} m_{+-}^{(2)} & t_{+-} m_{++}^{(2)} & e^{2\beta} t_{+-} m_{+-}^{(2)} \\ t_{++} m_{-+}^{(2)} & e^{-2\beta} t_{++} m_{++}^{(2)} & e^{-2\beta} t_{+-} m_{-+}^{(2)} & t_{+-} m_{--}^{(2)} \\ t_{-+} m_{++}^{(2)} & e^{-2\beta} t_{-+} m_{+-}^{(2)} & e^{-2\beta} t_{--} m_{++}^{(2)} & t_{--} m_{+-}^{(2)} \\ e^{2\beta} t_{-+} m_{-+}^{(2)} & t_{-+} m_{--}^{(2)} & t_{--} m_{-+}^{(2)} & e^{2\beta} t_{--} m_{--}^{(2)} \end{pmatrix} \quad (5.50)$$

where $m_{ij}, m_{ij}^{(2)}$ and $t_{i,j}$ ($i, j \in \{-, +\}$) are elements of the matrices M, M^2 , and T respectively.

Now for the sum under consideration (5.37) we obtain using representation (5.43)

$$\begin{aligned}
\sum_{(\mathbf{t}, \boldsymbol{\sigma}), (\mathbf{t}', \boldsymbol{\sigma}')} K^2((\mathbf{t}, \boldsymbol{\sigma}), (\mathbf{t}', \boldsymbol{\sigma}')) &= \sum_{(\mathbf{t}_{\text{do}}, \boldsymbol{\sigma}_{\text{do}}), (\mathbf{t}'_{\text{up}}, \boldsymbol{\sigma}'_{\text{up}})} e^{-2\beta V((\mathbf{t}_{\text{do}}, \boldsymbol{\sigma}_{\text{do}}), (\mathbf{t}'_{\text{up}}, \boldsymbol{\sigma}'_{\text{up}}))} \\
&\times \left[\prod_{l=1}^k \left(e_{\sigma'_{\text{up}}^l}^T M e_{\sigma_{\text{up}}^{l+1}} \right) \prod_{l=1}^{k-1} \left(e_{\sigma_{\text{do}}^l}^T M e_{\sigma_{\text{do}}^{l+1}} \right) \left(e_{\sigma_{\text{up}}^k}^T M^2 e_{\sigma_{\text{up}}^1} \right) \right. \\
&\quad \left. - \prod_{l=1}^k \left(e_{\sigma'_{\text{up}}^l}^T T e_{\sigma_{\text{up}}^{l+1}} \right) \prod_{l=1}^{k-1} \left(e_{\sigma_{\text{do}}^l}^T M e_{\sigma_{\text{do}}^{l+1}} \right) \left(e_{\sigma_{\text{up}}^k}^T M^2 e_{\sigma_{\text{up}}^1} \right) \right] \\
&= \text{tr} \left(\left(\sum_{k=0}^{\infty} Q^k \right) Q_m \right) - \text{tr} \left(\left(\sum_{k=1}^{\infty} Q_t^k \right) Q_{tm} \right). \tag{5.51}
\end{aligned}$$

By the construction the matrix Q is greater then Q_t elementwise. Thus the eigenvalue of matrix Q is greater than the eigenvalue of the matrix Q_t (it follows from the Perron-Frobenius theorem). Therefore the necessary and sufficient condition for the convergence in (5.37) is that the largest eigenvalue of Q is less than 1. It is possible to calculate its eigenvalue analytically. In order to express the eigenvalues of Q it is convinient to use notations (5.32) and (5.33). In this notations the matrix M , i.e. (5.46), is represented as following

$$M = c \begin{pmatrix} m & 1 \\ 1 & m \end{pmatrix}.$$

The equations for the eigenvalues of Q are:

$$\begin{aligned}
\lambda_1 &= c^2 e^{\beta} (m^2 - 1) \\
\lambda_2 &= c^2 e^{-\beta} (m^2 - 1) \\
\lambda_3 &= c^2 (m^2 + 1) (\cosh \beta) \left(1 - \sqrt{1 - \frac{(m^2 - 1)^2}{(\cosh \beta)^2 (m^2 + 1)^2}} \right) \\
\lambda_4 &= c^2 (m^2 + 1) (\cosh \beta) \left(1 + \sqrt{1 - \frac{(m^2 - 1)^2}{(\cosh \beta)^2 (m^2 + 1)^2}} \right)
\end{aligned}$$

A straightforward inspection confirms that the largest eigenvalue is given by λ_4 . The condition $\lambda_4 < 1$ coincides with (5.37). This completes the proof of Lemma 5.4.1. $\square$

