

First Time Determination of the Microscopic Structure of a Stripe Phase: Low Temperature NMR in $\text{La}_2\text{NiO}_{4.17}$

I. M. Abu-Shiekah, O. Bakharev, and H. B. Brom

Kamerlingh Onnes Laboratory, Leiden University, P.O. Box 9504, 2300 RA Leiden, The Netherlands

J. Zaanen

Instituut Lorentz for Theoretical Physics, Leiden University, P.O. Box 9506, 2300 RA Leiden, The Netherlands

(Received 17 August 2001; published 15 November 2001)

The experimental observations of stripes in superconducting cuprates and insulating nickelates clearly show the modulation in charge and spin density. However, these have proven to be rather insensitive to the harmonic structure and (site or bond) ordering. Using ^{139}La NMR in $\text{La}_2\text{NiO}_{4+\delta}$ with $\delta = 0.17$, we show that in the $1/3$ hole doped nickelate below the freezing temperature the stripes are strongly solitonic and site ordered with Ni^{3+} ions carrying $S = 1/2$ in the domain walls and Ni^{2+} ions with $S = 1$ in the domains.

DOI: 10.1103/PhysRevLett.87.237201

PACS numbers: 76.60.-k, 74.72.Dn, 75.30.Ds, 75.40.Gb

Stripe phases have by now been observed in a variety of doped Mott insulators, like cuprates, nickelates, and manganites [1]. Nevertheless, there is still a remarkable lack of knowledge on the details of the charge and spin distributions in these novel electronic phases. At the same time there is a growing body of theoretical literature dealing with the microscopic origin of stripe formation, predicting stripes starting from rather different physical perspectives [1–5]. Given the potential connections of stripe formation to, for instance, the mechanism of superconductivity, it is a matter of high urgency to find out how the stripes look in detail. In this regard the various theories lead to quite different predictions. Because the spin and charge distributions are inhomogeneous, NMR with its microscopic sensitivity could in principle yield important information. However, in cuprates attempts in this direction have been frustrated due to the anomalous, glassy ordering dynamics obscuring the static signal. Neutron scattering shows that both in cuprates and nickelates the correlation length of the stripe order is relatively small. This disorderly nature comes to play in NMR [6,7] in the form of a peculiar loss of signal intensity (wipeout) [6–11] explained by a spread in very short spin-dephasing times T_2 [6,10,11]. Even at the lowest temperatures the signal recovery is only partial indicating that this dynamics is still at work at temperatures as low as 0.3 K in the cuprates, prohibiting attempts to deduce information about the static order from the NMR data [7]. We demonstrate here that in a nickelate stripe system the nature of the stripe order (established by neutron scattering) can be deduced in detail from NMR, despite the strong similarities with the cuprate NMR at higher temperatures. We find the stripe structure to be strongly solitonic, with sharply defined charge stripes with a width which is not exceeding the lattice constant by much. Surprisingly, they look quite like the site centered stripes predicted by early mean-field calculations for the nickelate system [12].

Before analyzing the low temperature line shape, which is the central issue of this Letter, we first introduce the

main features of the NMR line shape, derive the value of the hyperfine coupling, and explain the partial recovery of the signal intensity from the relaxation data.

The NMR measurements were performed on the same $\text{La}_2\text{NiO}_{4.17}$ crystals as measured before [9,13], by sweeping the external field B at various frequencies. The crystal symmetry is almost tetragonal [14], which makes the ^{139}La spectra ($I = 7/2$) strongly dependent on the direction of the magnetic field with respect to the crystallographic axes.

In $\text{La}_2\text{NiO}_{4.18}$, with hole doping very close to $\text{La}_2\text{NiO}_{4.17}$, details about the excess oxygen atoms' positions are available [14]. The interstitial oxygen sites are at $(0.183a, 0.183b, 0.217c)$ or equivalent positions. A scenario for the excess oxygen ordering is shown in Fig. 1. There is one excess oxygen atom for each six unit cells. The number of ^{139}La sites, which are far from the excess oxygen atoms, is twice as large as the number closest to the excess oxygen atoms. For $\text{La}_2\text{NiO}_{4.17}$, using simple point charge calculations, we expect the interstitial oxygen to change the electrical field gradient (EFG) around the La sites and to split the ^{139}La NMR line into two: an A line due to A sites with the main component of the EFG (V_{zz}) along the c axis and a B line due to 2 times less abundant B sites with the crystal field gradient in the (ab) plane. This prediction is in agreement with the experimental intensity ratio of line A to B [9,15].

Above 50 K, the ^{139}La relaxation rates can be fitted by an Arrhenius expression $T_1^{-1} \propto \exp(E_0/T)$; see Fig. 2a [9]. Fits allow for a (full or half Gaussian) distribution in activation energies around $E_0 = 90$ K with a width of $\Delta = 25$ K. Above the magnetic freezing temperature the rates for line A are at least twice the rates of line B , and have almost the same temperature dependence [9]. To extract the amplitude of the fluctuating fields and its correlation time from the relaxation data, the value of the hyperfine constant is needed. This constant can be determined in several ways. In doped La_2CuO_4 the static internal magnetic field amounts to about 0.1 T [6], which

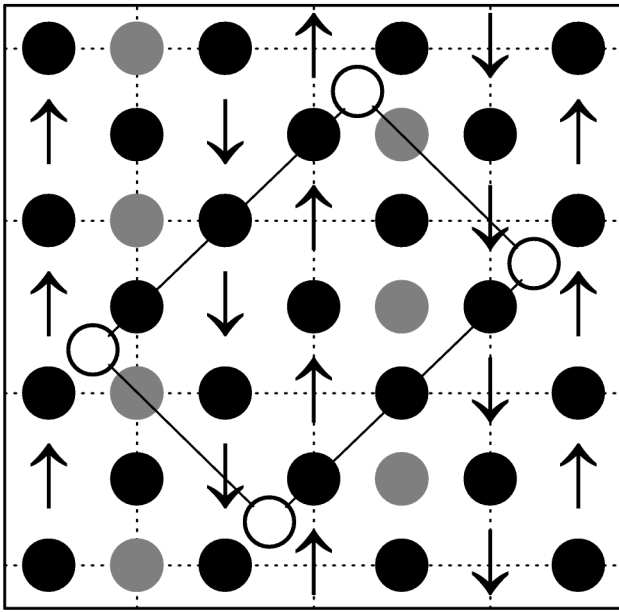


FIG. 1. Ordering of excess oxygen, holes, and spins. Open circles refer to the ordered excess oxygen sites at $z = 0.217c$, which introduce holes assumed to be Ni^{3+} with $S = 1/2$ (gray circles, site-order scenario). The arrows refer to the Ni^{2+} spins with $S = 1$. Black circles indicate the ^{139}La sites at $z = 0.14c$. The dotted lines are the in-plane boundaries for the unit cell before doping. The drawn tilted rectangle is the enlarged unit cell after doping.

is similar to the value found in the undoped case [16]. Because the hyperfine field at the lanthanum site in the nickelates is calculated to be about 18 times larger than in the cuprates [17], its value will be around 1.8 T. Yoshinari *et al.* [18] reported a value of about $2.3 \pm 0.5 \text{ T}/\mu_B$ for the hyperfine coupling in Sr-doped nickelates with an ordered moment of $\approx 0.7\mu_B$. The so-obtained internal field of 1.6 T agrees well with the internal field of about 1.78 T reported by Wadat *et al.* [19] for $\text{La}_2\text{NiO}_{4.10}$. Also by com-

paring our susceptibility and linewidth data at high temperature we can estimate the hyperfine field to be about 2.0 T, consistent with the earlier quoted values [13]. Using a hyperfine field of 1.8 T, the fit to the relaxation data gives a value of 2600 s^{-1} for $\gamma_n^2 h_0^2 \tau_\infty$, and hence $\tau_\infty = 1.8 \times 10^{-11} \text{ s}$. T_2^{-1} has almost the same T dependence as T_1^{-1} . It deviates only for $T > 200 \text{ K}$.

The relaxation data below 20 K cannot be fitted by the same parameters; see Fig. 2b. The rates can be reproduced by $T_1^{-1} \propto \gamma^2 h_0^2 \exp(-E_0/T)/\omega_0^2 \tau_\infty$, with $E_0 = 2.3 \pm 0.2 \text{ K}$ and h_0 the strength of the fluctuating field. Measurements of the rates were performed for different satellites at 38.2 MHz. Using the same methods as described in Ref. [6], a magnetic origin for the fluctuating fields is found at 160, 90, and 1.6 K.

The wipeout of the ^{139}La NMR signal intensity and its partial reappearance below $\sim 11 \text{ K}$ [9] can be explained in the same terms as used before in the cuprates [6] indicating a distribution in spin-dephasing times, hence activation energies [15]. At low temperatures only nuclei relaxing with activation energies $> E_0$ lead to the measured signal, while the others will have dephasing rates higher than 10^4 Hz down to very low temperatures. Even at $T \approx 0.5 \text{ K}$ slow spin fluctuations apparently prevent the full recovery of the signal. The intensity recoveries can be fitted with $\Delta = 2.8 \text{ K}$ and $E_0 \approx 0 \text{ K}$ [15]. These results are consistent with the relaxation data because only nuclei with the higher activation energies will be visible.

What are the implications of these observations? In the renormalized classical limit below charge ordering the spin correlation rate $\tau(T)$ will be proportional to the spin correlation length $\xi(T)$ divided by the spin wave velocity c : $T_1^{-1} \sim (\xi/c)(T/2\pi\rho)^2(1 + T/2\pi\rho)^{-2}$, with ρ the spin stiffness [6,7]. In $\text{La}_2\text{NiO}_{4.17}$ according to diffraction studies the stripes have a finite and temperature independent correlation length [20,21]. If this has a dynamic origin, it will be consistent with a distribution in activation energies or τ 's. Possible sources for such a distribution in the spin dynamics are dislocations [22,23] in the domain pattern or elastic deformations [23].

We now come to the central part of this paper. We will show that although the interstitial oxygens have a direct influence on the La signal via the electric field gradients, it is still possible using ^{139}La NMR to probe the charge and spin ordering in the NiO_2 plane.

Below 15 K, when the signal starts to regain some of its intensity, the NMR lines are largely broadened. In Figs. 3 and 4 the spectra are shown for $B \parallel c$ and $B \perp c$. Simulations of the spectra taking into account the quadrupolar and Zeeman contributions [24] are shown in the lower panels. The spectra can be decomposed into 2 La sites, one with resolved satellites and the other with overlapping lines due to much stronger magnetic broadening. Magnetic freezing or ordering is also visible in the bulk magnetic susceptibility, which starts to deviate from the Curie law and peaks at $T = 17 \pm 3 \text{ K}$ [9,13,25]. The peak in $-d\chi/dT$ indicates a series of cusps at different spin freezing temperatures

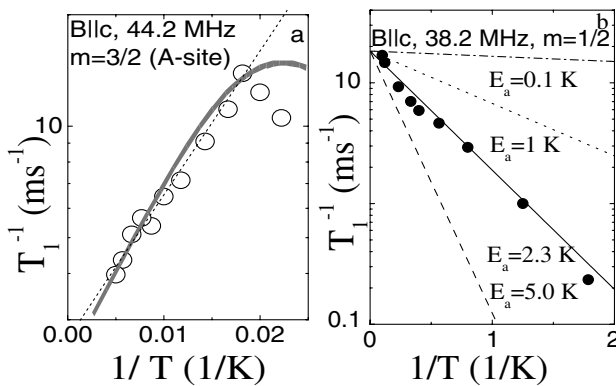


FIG. 2. (a) T_1^{-1} for line A (see text) for $m = 3/2$ at 44.196 MHz, obtained from stretched exponential fits to the recoveries. The dotted line is a fit with $E_0 = 90 \text{ K}$. Drawn line is a fit using a half Gaussian distribution in activation energy. (b) T_1^{-1} versus $1/T$ at low temperatures. An activation energy of 2.3 K fits the data well. Sites with lower activation energies will not be observable ($m = 1/2$ refers to the $+1/2, -1/2$ and $m = 3/2$ to the $3/2, 1/2$ transition).

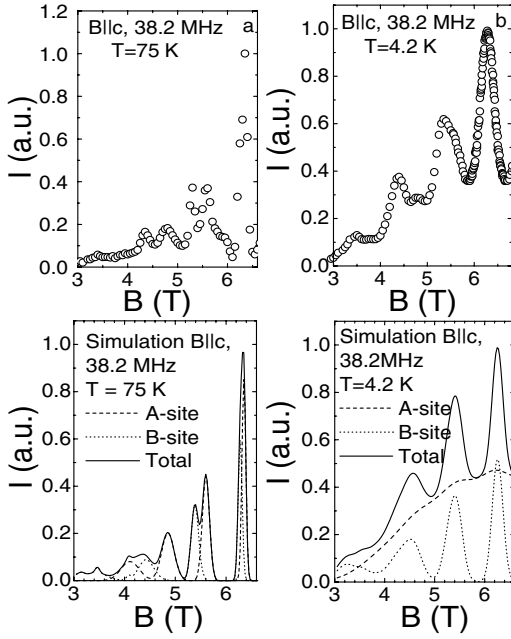


FIG. 3. High (a) and low (b) T spectra for $B \parallel c$ and their simulations (lower panels). The simulations give precise values for the EFG parameters but are not very sensitive for the strength of the internal field (at low T). Parameters are discussed in the text.

close to 17 K and can be due to spin frustration or spin clusters [25].

We first show that contrary to the observation in the cuprates [7], the low temperature line pattern in the nickelate is not due to motional narrowing. In case of two lines separated by $\delta\omega$ the condition for motional narrowing [6,24] is that the correlation time τ for the (spin)

fluctuations is much less than $1/\delta\omega$. The splitting $\delta\omega$ is given by γ_n times the value of the internal field h_0 : $\delta\omega = \gamma_n \langle h_0 \rangle$. The value of τ is coupled to the relaxation rate via $\tau \sim \gamma_n^2 \langle h_0^2 \rangle T_1 / \omega^2$. Hence the condition can be rewritten as $R_m = \gamma_n^3 \langle h_0^2 \rangle^{3/2} T_1 / \omega^2 \ll 1$. Using typical values for $h_0 = 1.5$ T, $\omega = 2\pi \times 40$ MHz, and $T_1 = 50 \mu\text{s}$, $R_m \sim 10^2$.

What is the information we might extract from the line shape data? The high T data will be sensitive for the EFG parameters and the low T data in addition for the internal or local field. Since the internal field lies in the ab plane, the spectra for $B \parallel c$ will be rather insensitive to the local field. For $B \parallel ab$ the antiferromagnetic alignment of the electron spins we expect to lead to a splitting of the lines. The simulations of the spectra are shown in the lower panels. The principal axes of the EFG with respect to the crystal axes are described by Eulerian angles (α, β, γ) [26] and the external magnetic field is described by polar angles (θ, φ) with respect to crystallographic axes.

For $B \parallel c$ ($\theta = 0$) the high T data (e.g., at 75 K) at $\nu = 38.2$ MHz are well reproduced by a quadrupolar splitting $\nu_Q = eQV_{zz}/4\hbar I(2I - 1)$ (with Q the electric quadrupole moment of the ^{139}La nucleus) of 4.5 MHz. The anisotropy parameter $\eta = (V_{xx} - V_{yy})/V_{zz} = 0$. The width of the satellites in the spectra can be simulated by introducing a spread in the EFG parameter $\Delta\nu_Q = 0.5$ MHz, while the width of the main line requires a (dipolar) field distribution of 0.05 T at the La(A) site. For the La(B) site these values are $\nu_Q = 8.5$ MHz, $\eta = 0.75$, $\Delta\nu_Q = 1.0$ MHz, and $\Delta_B = 0.05$ T. The other parameters are for the A site: $\alpha = 0, \beta = 0, \gamma = 0$; for the B site: $\alpha = 0, \beta = \pi/2, \gamma = -\pi/6$. According to the simulation of the low temperature data the main effect of the magnetic freezing is the increase of the magnetic broadening (Δ_B) to 0.5 or 0.15 T, respectively. The spectra shown use internal fields of 1.5 T at the A site and 0.5 T at the B site. These internal field values follow from the simulations for $B \perp c$, where the values are much more restricted. For the $B \perp c$ orientation simulation of the 107 K data with the same parameter set as for $B \parallel c$ reproduce the data well ($\varphi = \pi/6$). For the 4.2 K data the antiparallel alignment of the internal field leads to a splitting of the lines and is described by polar angles $(\theta = \pi/2, \varphi = \pi/3 \pm \pi/2)$ with respect to crystallographic axes.

To explain these remarkable differences, let us assume that we deal with site centered stripes. These stripes naturally lead to two kinds of Ni ions: Ni^{2+} (A sites) and Ni^{3+} (B sites) with two kinds of spins $S = 1$ and $S = 1/2$, respectively. As a consequence, the La sites will experience different hyperfine fields. The internal dipolar magnetic field arising from the Ni spins at the La sites will be of the order of 0.2 T. However, the La nuclei just above and below Ni^{2+} in addition will have an exchange coupling via the oxygens. This hyperfine field due to the overlap of the Ni $3d_{z^2}$ and La $6s$ orbitals through the $2p_z$ orbital of apical oxygen, is about 1.8 T in the undoped samples and

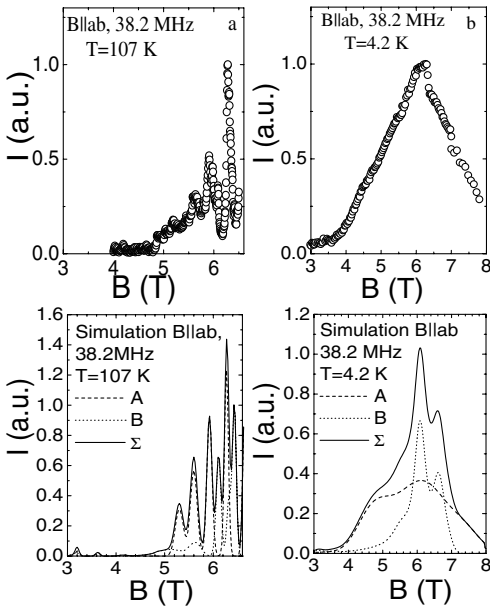


FIG. 4. High (a) and low (b) T spectra for $B \parallel ab$ and their simulations (lower panels). Spectra are simulated with the same parameters as in Fig. 3. The internal field leads to a splitting of the lines and hence can be determined precisely.

is almost doping independent [27]. Because of the different occupation of the $3d_{z^2}$ orbital of Ni^{3+} the exchange coupling between the $S = 1/2$ Ni^{3+} spins and the La sites will be weaker [28,29]. The ratio between ^{139}La sites that feel the exchange field of Ni^{3+} and Ni^{2+} will be close to 1:2. This difference in intensity ratio and hyperfine field are indeed the main characteristics of the line shape and hence are well accounted for by this scenario. Can bond centered stripes explain the observations as well? The hyperfine fields will have the same maximum, but the distribution will be different. The line shapes for $B \perp c$ puts a limit to the field on the B sites of at most 0.3 T, which rules out this possibility.

The experiments show that apart from the internal field we need to introduce an appreciable magnetic broadening. Part of the broadening might be due to canting of the spins in the ordered phase in the NiO_2 plane away from the charge and spin stripe direction [21], which effect we have not included in this calculation. Another reason for the extra broadening might be found in the finite size of the correlated magnetic regions.

To summarize, in $\text{La}_2\text{NiO}_{4.17}$ interstitial oxygens determine the line profile above the wipeout temperature. The NMR intensity loss above the spin freezing or ordering temperature around 20 K is linked to a spread in spin dephasing as in the cuprates. From the angular and temperature dependence of the La line profiles, we show that the distribution of the internal fields is in agreement with two kinds of Ni sites with different ionicity and hence different hyperfine interaction with the visible La sites. Site centered stripes of the kind predicted by mean-field theory [12] fit the low temperature data of the visible La nuclei remarkably well.

We acknowledge A.A. Menovsky, A.A. Nugroho, and Y. Mukosvkii for providing the nickelate samples, O.O. Bernal and P.M. Paulus for useful discussions, and O. Berfelo for performing the TGA. This work was supported by FOM-NWO.

-
- [1] V.J. Emery, S. A. Kivelson, and J. Tranquada, *Proc. Natl. Acad. Sci. U.S.A.* **96**, 15380 (1999), and references therein.

- [2] J. Zaanen, *J. Phys. Chem. Solids* **59**, 1769 (1998), and references therein.
- [3] S.R. White and D.J. Scalapino, *Phys. Rev. Lett.* **81**, 3227 (1998).
- [4] M. Vojta, Y. Zhang, and S. Sachdev, *Phys. Rev. B* **62**, 6721 (2000).
- [5] *Proceedings of the Third International Conference on Stripes and High T_c Superconductivity, Rome, 2001*, edited by N.L. Saini and A. Bianconi [*Int. J. Mod. Phys. B* **14**, No. 29–31 (2001)].
- [6] G.B. Teitelbaum *et al.*, *Phys. Rev. B* **63**, 020507(R) (2001).
- [7] A. W. Hunt, P.M. Singer, A.F. Cederström, and T. Imai, *Phys. Rev. B* **64**, 134525 (2001).
- [8] A. W. Hunt, P.M. Singer, K.R. Thurber, and T. Imai, *Phys. Rev. Lett.* **82**, 4300 (1999); P.M. Singer, A. W. Hunt, A.F. Cederström, and T. Imai, *Phys. Rev. B* **60**, 15345 (1999).
- [9] I. M. Abu Shiekh *et al.*, *Phys. Rev. Lett.* **83**, 3309 (1999).
- [10] N. J. Curro *et al.*, *Phys. Rev. Lett.* **85**, 642 (2000).
- [11] B. J. Suh *et al.*, *Phys. Rev. B* **61**, R9265 (2000).
- [12] J. Zaanen and P. B. Littlewood, *Phys. Rev. B* **50**, 7222 (1994).
- [13] O. O. Bernal *et al.*, *Physica (Amsterdam)* **282–287C**, 1393 (1997).
- [14] A. Mehta and P. J. Heaney, *Phys. Rev. B* **49**, 563 (1994).
- [15] I. M. Abu-Shiekh, Ph.D. thesis, Leiden University, 2001.
- [16] D. E. MacLaughlin *et al.*, *Phys. Rev. Lett.* **72**, 760 (1994).
- [17] M. Takahashi, T. Nishino, and J. Kanamori, *J. Phys. Soc. Jpn.* **60**, 1365 (1991).
- [18] Y. Yoshinari, P. C. Hammel, and S.-W. Cheong, *Phys. Rev. Lett.* **82**, 3536 (1999).
- [19] S. Wadat *et al.*, *J. Phys. Condens. Matter* **5**, 765 (1993).
- [20] C.-H. Du *et al.*, *Phys. Rev. Lett.* **84**, 3911 (2000).
- [21] S.-H. Lee, S.-W. Cheong, K. Yamada, and C. F. Majkrzak, *Phys. Rev. B* **63**, 060405(R) (2001).
- [22] J. Zaanen, *Phys. Rev. Lett.* **84**, 753 (2000).
- [23] O. Zachar, *Phys. Rev. B* **62**, 13836 (2000).
- [24] C. P. Slichter, *Principles of Magnetic Resonance* (Springer, Berlin, 1991), 3rd ed.
- [25] P. Odier, N. J. Poirot, P. Simon, and D. Desousa Meneses, *Eur. Phys. J. AP* **5**, 123 (1999).
- [26] H. Goldstein, *Classical Mechanics* (Addison-Wesley, London, 1950).
- [27] Y. Furukawa and S. Wada, *J. Phys. Condens. Matter* **6**, 8023 (1994).
- [28] E. Pellegrin *et al.*, *Phys. Rev. B* **53**, 10667 (1996).
- [29] V. I. Anisimov *et al.*, *Phys. Rev. B* **59**, 7901 (1999).