

MEASURING PROTOPLANETARY DISK ACCRETION WITH H I PFUND β

COLETTE SALYK¹, GREGORY J. HERCZEG², JOANNA M. BROWN³, GEOFFREY A. BLAKE⁴,
KLAUS M. PONTOPPIDAN⁵, AND EWINE F. VAN DISHOCK^{6,7}

¹ National Optical Astronomy Observatory, 950 North Cherry Avenue, Tucson, AZ 85719, USA; csalyk@noao.edu

² The Kavli Institute for Astronomy and Astrophysics at Peking University, Yi He Yuan Lu 5, Hai Dian Qu, Beijing 100871, China

³ Harvard-Smithsonian Center for Astrophysics, 60 Garden Street, Cambridge, MA 02138, USA

⁴ Division of Geological & Planetary Sciences, Mail Code 150-21, California Institute of Technology, Pasadena, CA 91125, USA

⁵ Space Telescope Science Institute, 3700 San Martin Drive, Baltimore, MD 21218, USA

⁶ Leiden Observatory, Leiden University, P.O. Box 9513, 2300 RA Leiden, The Netherlands

⁷ Max-Planck-Institut für Extraterrestrische Physik, Giessenbachstrasse 1, D-85748 Garching, Germany

Received 2012 October 16; accepted 2013 March 14; published 2013 May 1

ABSTRACT

In this work, we introduce the use of H I Pfund β ($\text{P}\beta$; $4.6538 \mu\text{m}$) as a tracer of mass accretion from protoplanetary disks onto young stars. $\text{P}\beta$ was serendipitously observed in NIRSPEC and CRIRES surveys of CO fundamental emission, amounting to a sample size of 120 young stars with detected $\text{P}\beta$ emission. Using a subsample of disks with previously measured accretion luminosities, we show that $\text{P}\beta$ line luminosity is well correlated with accretion luminosity over a range of at least three orders of magnitude. We use this correlation to derive accretion luminosities for all 120 targets, 65 of which are previously unreported in the literature. The conversion from accretion luminosity to accretion rate is limited by the availability of stellar mass and radius measurements; nevertheless, we also report accretion rates for 67 targets, 16 previously unmeasured. Our large sample size and our ability to probe high extinction values allow for relatively unbiased comparisons between different types of disks. We find that the transitional disks in our sample have lower than average $\text{P}\beta$ line luminosities, and thus accretion luminosities, at a marginally significant level. We also show that high $\text{P}\beta$ equivalent width is a signature of transitional disks with high inner disk gas/dust ratios. In contrast, we find that disks with signatures of slow disk winds have $\text{P}\beta$ luminosities comparable to those of other disks in our sample. Finally, we investigate accretion rates for stage I disks, including significantly embedded targets. We find that stage I and stage II disks have statistically indistinguishable $\text{P}\beta$ line luminosities, implying similar accretion rates, and that the accretion rates of stage I disks are too low to be consistent with quiescent accretion. Our results are instead consistent with both observational and theoretical evidence that stage I objects experience episodic, rather than quiescent, accretion.

Key words: accretion, accretion disks – protoplanetary disks – stars: formation – stars: pre-main sequence

Online-only material: extended figure, machine-readable tables

1. INTRODUCTION

In the study of protoplanetary disks and protostars, much effort has been focused on the study of mass accretion rates—the rates at which mass is transferred from circumstellar disks to stars—because it is so intricately linked to processes important for star and planet formation. Mass accretion is a measure of the viscosity of the disk and determines the overall rate of mass and momentum transfer, and thus the pace of disk evolution. The rate of mass accretion will affect the disk lifetime (and thus the time available for planet formation) and the rate of planetary migration, and may in turn be a *tracer* of the presence of planets (e.g., Alexander & Armitage 2007). Accretion greatly affects the inner disk environment, with the disk truncated and material lofted onto the star at or near the stellar corotation radius (Shu et al. 1994), which could be related to the observed pile-up of giant planets at small orbital radii (Lin et al. 1996; Butler et al. 2006). Finally, accretion is a tracer of star/disk magnetic interactions and determines the early angular momentum evolution of the star (e.g., Agapitou & Papaloizou 2000).

The most accurate measurement of mass accretion is likely obtained from the spectroscopic observation and modeling of UV excess flux (e.g., Valenti et al. 1993; Gullbring et al. 1998; Herczeg & Hillenbrand 2008), as this provides a nearly direct measure of the total accretion luminosity. However, there has

long been an interest in measuring accretion rates with data that are relatively simpler to obtain and analyze, as well as data that can be obtained at longer wavelengths, where extinction is lower. This has led to a number of studies of easily observable H I emission lines believed to be produced in the accretion columns and accretion shock along with the UV continuum excess (Calvet & Gullbring 1998). Such studies have shown that H I line luminosities correlate with UV-excess-derived accretion luminosities, and thus can be used as reasonably reliable tracers of mass accretion rates (e.g., Muzerolle et al. 1998b; Natta et al. 2004, 2006; Garcia-Lopez et al. 2006; Fang et al. 2009), albeit with some systematic uncertainties and caveats (e.g., Herczeg & Hillenbrand 2008). The ease of collecting spectra of H I emission lines results in more comprehensive samples of mass accretion rates, allows for the study of accretion rates in more embedded disks, and allows for a simultaneous measure of accretion rates and other disk properties, such as veiling by the disk continuum or molecular emission line strengths. In fact, one of the best studied H I lines—H α —remains a popular means of estimating mass accretion, in spite of the fact that models suggest it saturates even at moderate accretion rates (Muzerolle et al. 1998a) and the fact that its strength and profile shape can also be affected by other parameters, including outflow rate and system inclination (Kurosawa et al. 2006). In addition, since measured accretion rates span several orders of magnitude, even measurements with error bars as large as ~ 0.5 –1 dex can

be used to broadly characterize a sample of objects and provide scientifically useful information.

In this work, we introduce a new tracer of mass accretion—specifically, H I P β —that offers several advantages over other tracers, and results in one of the most comprehensive coherent data sets of a single mass accretion tracer to date. H I P β ($n = 7 \rightarrow 5$; hereafter P β) is located at $4.6538 \mu\text{m}$ in the M band, which places it between the R(0) and R(1) lines of the CO rovibrational fundamental band. Thanks to two large campaigns designed to study CO fundamental emission in protoplanetary disks with Keck-NIRSPEC (e.g., Blake & Boogert 2004; Salyk et al. 2011) and VLT-CRIRES (e.g., Pontoppidan et al. 2011b; Herczeg et al. 2011; Brown et al. 2013), P β has been serendipitously observed in more than 100 young stars with disks. P β has similar equivalent widths to other infrared tracers, including Br γ ($n = 7 \rightarrow 4$), with which it shares the same upper level energy, and appears to be ubiquitous for disks determined by other means to be typical active accretors. It also offers several distinct advantages over other accretion tracers due to its position in the M band. First, it is an ideal tracer of accretion rates in transitional disks, as the low M -band photospheric flux results in a very high P β equivalent width when the dust continuum is low, as it is in transitional disks. Second, the high degree of veiling in the M band relative to shorter wavelengths makes contamination from stellar photospheric absorption negligible in nearly all cases. Third, tentative discoveries of young planets in disks (Huélamo et al. 2011; Kraus & Ireland 2012) have sparked interest in relating CO emission line variability to the tidal influence of protoplanets (Regály et al. 2011). An understanding of the connection between accretion and P β may help observers differentiate between accretion- and planet-induced variability observed in CO emission lines. Finally, P β is observable even in heavily extincted disks, and so can be used to study the youngest disks embedded in their natal cloud.

The ability to measure accretion rates in embedded disks is a particularly important strength of P β as an emission tracer. A star is thought to gain most of its mass while it is still embedded in a dense envelope, prompting the naive expectation that accretion rates should be higher during these stages. During this phase, however, the commonly used accretion indicators in the optical and near-IR are not usually observable and this expectation is difficult to confirm. As a substitute for direct measurements of accretion rates, the mass history of stars has been inferred from a global analysis of temperature–luminosity diagrams (Kenyon et al. 1990; Dunham et al. 2010; Zhu et al. 2010). In these analyses, the luminosity distribution is $\sim 10\%$ of the expected luminosity distribution required to build stars via steady accretion, suggesting that the star may need to grow mostly in large, short bursts. A few studies of optical and near-IR accretion tracers in more extincted disks (Muzerolle et al. 1998b; White & Hillenbrand 2004; White et al. 2007) have suggested a similar result—namely, that accretion rates in younger disks are lower than expected. However, as even near-IR accretion tracers like H α , Br γ , and Pa β are not visible in the most embedded disks, these studies may have been biased toward older systems, including possibly edge-on, evolutionarily older (yet observationally extincted) disks. Therefore, the study of P β provides an exciting new path for measuring accretion in the youngest, most embedded disks.

In this paper, we report the detection and measurement of P β emission lines observed with Keck-NIRSPEC and VLT-CRIRES. We develop a method to use P β emission to measure accretion luminosity by correlating line luminosities

to known accretion luminosities for a sample of young stars. This correlation is then applied to a large sample of targets to provide accretion luminosities for 120 young stars, including a sample of objects still embedded in their molecular envelopes. In Section 2, we briefly describe the NIRSPEC and CRIRES observations and data reduction procedures, as well as the sample selection. In Section 4, we discuss the flux extraction procedure and corresponding uncertainties. In Section 5.1, we calculate a relationship between P β line luminosity and accretion luminosity, and in Section 5.2 we apply this to our full sample to provide accretion luminosity estimates for 120 stars. In Section 6, we discuss some implications of our results.

2. OBSERVATIONS AND REDUCTION

2.1. Observing Procedures

NIRSPEC observations were obtained with the Keck II telescope as part of a large survey of CO rovibrational emission from young stars with disks (see, e.g., Salyk et al. 2011), spanning 2001–2011. Spectra were obtained in the high-resolution mode with $R \sim 25000$, in the M -wide filter. Echelle and cross-disperser positions were optimized for the observation of the P branch of CO $v = 1 \rightarrow 0$. Although the favored echelle settings evolved somewhat throughout the course of the survey, nearly all sources were observed in a setting that included the H I lines P β ($4.6538 \mu\text{m}$) and H ϵ ($4.6725 \mu\text{m}$). NIRSPEC targets were observed in ABBA nod sets, with AB pairs differenced to remove thermal emission from Earth’s atmosphere. Integration times (exposure time multiplied by coadds) in each position were limited to 1 minute to minimize atmospheric variation between frames of a pair. To correct for telluric absorption, A or B telluric standard stars with nearly blackbody spectra were observed close in time and airmass to the targets.

CRIRES observations were obtained with the Very Large Telescope (VLT) as part of a large survey of CO rovibrational emission from young stars with circumstellar disks (Pontoppidan et al. 2011b; Brown et al. 2013), including a significant sample of embedded protostars (Herczeg et al. 2011). Only $\sim 0.02 \mu\text{m}$ are covered in each integration with the non-cross-dispersing CRIRES instrument, but the majority of targets in this survey were still observed in a setting that included P β . Observations were obtained with the $0'2$ slit, resulting in a resolution of $R \sim 90,000$. Targets were observed in ABBA nod sets, with 60 s integration times for each image. Random dithers between $0''$ and $1''$ were added to the standard $10''$ nod sequence to distribute the counts over different pixels, and ensure that the data would not always land on a bad pixel. When possible, brighter targets were observed with adaptive optics.

A log of the NIRSPEC and CRIRES observations included in this study can be found in Table 1.

2.2. Data Reduction

NIRSPEC data reduction routines were developed by our team (Blake & Boogert 2004) in IDL. NIRSPEC spectra were first linearized using the shape of telluric emission lines and the observed spectra. Fluxes were extracted with a 2.8σ aperture around the best-fit point-spread function (PSF) center, and nearby columns were subtracted to account for any residual sky emission in the A–B difference image. Extracted spectra were wavelength calibrated utilizing telluric lines. Telluric standards were processed with the same procedures as the target stars. Target spectra were then divided by the telluric standard star

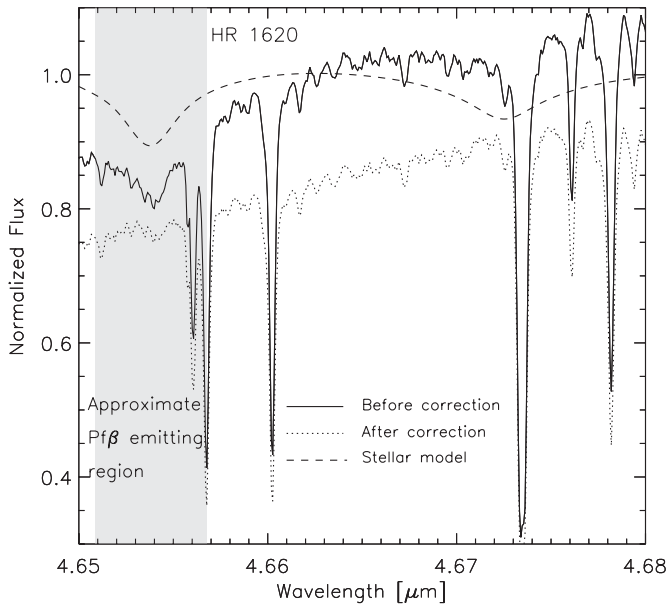


Figure 1. Solid and dotted lines show sample standard-star spectrum before and after stellar atmospheric model division. The shaded region marks an approximate location for $\text{Pf}\beta$ emission, assuming a width of 190 km s^{-1} . The actual region will depend on the line width and relative radial velocity between the source and the standard stars.

Table 1
Log of Observations

Star	Date	Exp (s)	Vdop ^a (km s^{-1})	Std
NIRSPEC				
AA Tau	2003 Nov 3	720	-15	HR 1620
	2004 Dec 27	780	13	HR 1620
AB Aur	2001 Jan 30	360	24	HR 1620
	2001 Aug 7	180	-25	HR 1620
	2002 Jan 3	240	13	HR 1620
AS 205 N	2002 Apr 21	240	-17	HR 6175
	2002 Jul 22	480	24	HR 5984
AS 205 S	2002 Apr 21	360	-17	HR 6175

Notes. ^a Earth-induced Doppler shift toward the source position.

(This table is available in its entirety in a machine-readable form in the online journal. A portion is shown here for guidance regarding its form and content.)

spectra (adjusted for small differences in airmass) to obtain spectra corrected for telluric absorption.

Although telluric standards were nearly featureless, most A stars have photospheric H I absorption, including that of $\text{Pf}\beta$. Thus, correcting for this feature is crucial for obtaining accurate $\text{Pf}\beta$ fluxes in our targets. To account for this, telluric standard spectra were first flattened utilizing Kurucz models, broadened to match observed H I absorption shapes. An example of this process is shown in Figure 1. Note that the corrected telluric spectrum should be linear (except for the narrow telluric absorption lines), and that this can easily be inspected by eye to confirm proper correction of the photospheric H I. In rare cases, the H I absorption profiles were not well accounted for by Kurucz models; in these cases, we simply allowed for an additional Gaussian-shaped absorption component to be removed from the spectrum. For all spectra, we performed an additional check on potential contamination from the standard

Table 2
Photometry

Star	M	4.5	5.8	Refs
AA Tau	6.86	6.77	6.35	12, 15
AB Aur	2.70	16
AS 205 N	4.18	17
AS 205 S	5.31	17
AS 209	5.05	5
BF Ori
CI Tau	...	6.27	5.94	23
CK 4	...	7.49	7.56	5
CrA IRS 2	4.03	9
CV Cha

Notes. “4.5” and “5.8” refer to magnitudes in the IRAC 4.5 and $5.8 \mu\text{m}$ filters.

^a Although other photometric measurements are available for this source, we use the spectroscopically derived measurements from Goto et al. (2011) as EX Lup was in outburst during the observations presented here, and thus several magnitudes brighter than in its quiescent state.

^b Corrected for observed relative fluxes of IRS 44 E and IRS 44 W from Herczeg et al. (2011).

References. (1) Bouvier & Appenzeller 1992; (2) Carpenter et al. 2008; (3) Cieza et al. 2007; (4) Ducati 2002; (5) Evans et al. 2003; (6) Evans et al. 2009; (7) Goto et al. 2011; (8) Hartmann et al. 2005; (9) Herczeg et al. 2011; (10) Hillenbrand et al. 1992; (11) Hughes et al. 1994; (12) Kenyon & Hartmann 1995; (13) Lada et al. 2006; (14) Liseau et al. 1992; (15) Luhman et al. 2006; (16) Malfait et al. 1998; (17) McCabe et al. 2006; (18) Merín et al. 2004; (19) Merín et al. 2008; (20) Padgett et al. 2006; (21) Pinte et al. 2008; (22) Rebull et al. 2010; (23) Robitaille et al. 2007; (24) Schegerer et al. 2009; (25) Wright et al. 2010.

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star: we divided a flat standard star observation by each Kurucz-model-corrected standard star observation to confirm that it did not produce a spurious H I emission line.

Since the primary observational target of these surveys, CO, can overlie telluric CO absorption lines, targets were typically observed on two or more dates, and spectra combined to create complete CO line profiles. The H I line emitting regions are relatively unaffected by telluric absorption and so the repeated observations simply increase the signal-to-noise ratio (S/N) in these lines. Alternatively, multiple epochs can also be left separate to investigate accretion variability. This aspect of the data set is left as future work, and will likely benefit from the inclusion of additional data sets from the NIRSPEC and CRIRES archives.

Approximate flux calibration was achieved by comparing telluric standard star spectra with literature photometry to derive a conversion from counts to flux. In all cases, the brightest observation of standard and target were utilized, as these would represent the best-centered observations. However, there is no guarantee that the sources were well centered, and the flux correction factor due to slit losses can be either greater than or less than one, depending on whether it is the source or the telluric standard that is observed off center. Nevertheless, we find decent agreement between absolute flux and literature photometry; the rms difference between measured and literature fluxes (see Table 2) is 50% for our sample as a whole.

Due to the inherent difficulties with absolute flux calibration for spectroscopic observations, we have chosen to correct observed fluxes with literature photometry whenever possible. The photometry used in this work can be found in Table 2. We used M -band photometry when available, or secondarily

an interpolation of other measurements, often *Spitzer* IRAC 4.5 and 5.8 μm fluxes. Although young stars with disks are known to have variable infrared emission, the variability is typically less than 10% for class 0–II disks (Luhman et al. 2010). Variability may be higher for transitional disks ($\sim 30\%$ – 50% ; Espaillat et al. 2011) and outbursting sources like EX Lup (e.g., Aspin et al. 2010), and so we note that such sources can have a corresponding systematic error in their $\text{Pf}\beta$ line flux measurement. If photometric measurements were not available in the literature, we simply used our own absolute photometry derived from the spectroscopic observations.

CRIRES reductions utilize routines written by K. Pontoppidan and described in Pontoppidan et al. (2008). Two-dimensional images for each source were added together after aligning and resampling onto a $2\times$ finer grid. Spectra were extracted using an aperture twice the size of the FWHM of the spectral profile. As with the NIRSPEC observations, nearby telluric standards were observed close in time and airmass. Wavelength calibration was performed on the telluric standards using telluric emission lines and was applied to the target spectra after small shifts to align the spectra of target and standard.

In the nominal CRIRES reduction routines, target spectra were then simply divided by the telluric spectra to remove telluric absorption features. However, this does not account for H I absorption in the standard stars, and would have produced anomalously strong H I emission in our targets. Therefore, we implemented a routine adapted from the NIRSPEC reduction routines, which uses Kurucz photospheric models to fit and correct for H I absorption features in the standard stars.

Spectra of standard stars were also compared with literature photometry to obtain a conversion from counts to pixel, which were applied to the target spectra to obtain an approximate absolute flux calibration. However, just as with the NIRSPEC data, CRIRES flux calibration suffers from an uncertain degree of slit losses for any given observation. Therefore, as with the NIRSPEC observations, we adjusted the spectra to match literature photometry whenever possible. We find a few large discrepancies (factors of ~ 2 – 3) between literature and measured photometry, including for Elias 29, IRS 44, LLN 8, LLN 17, and WL 12. Since all show higher literature fluxes than CRIRES-measured fluxes, it is possible that this is because prior observations with the *Spitzer Space Telescope* or *Infrared Space Observatory* included more than one source in their relatively large spatial PSF. For the non-outliers, we find the rms difference between observed and literature fluxes to be $\sim 52\%$.

2.3. Sample Selection

The NIRSPEC and CRIRES target lists used in this work were compiled for prior studies of M -band CO emission and do not represent an unbiased sample of young stars. However, the large sizes of these surveys means they come close to representing a complete flux-limited sample of disks with detectable CO fundamental emission in nearby clouds. The NIRSPEC M -band survey as a whole is dominated by revealed, optically thick disks, including significant numbers of both low-mass (T Tauri) and mid-mass (Herbig Ae/Be) disks. It also includes a smaller but significant number of embedded disks. The NIRSPEC sample is strongly biased against sources with tenuous disks or no disks at all, as these disks tended not to produce observable CO rovibrational emission lines. The NIRSPEC sample is also biased toward clouds at the higher declinations observable from Mauna Kea, sampling well in Taurus, Perseus, Serpens,

and Ophiuchus, only poorly in Lupus, and not at all at lower declinations. The NIRSPEC sample was also limited to targets with M -band continuum fluxes brighter than ~ 0.01 Jy ($M \sim 9$).

The CRIRES M -band CO survey contains many of the bright Class II disks visible from the southern sky as well as a number of Class I embedded protostars and transitional disks. This includes targets in Lupus, Vela, and Chamaeleon that are difficult or impossible to observe with NIRSPEC.

A more detailed discussion of the evolutionary status of our sample as a whole is included in Section 6.4.

3. GENERAL CHARACTERISTICS OF $\text{Pf}\beta$ DETECTIONS AND NON-DETECTIONS

A sample of targets showing different strengths and types of $\text{Pf}\beta$ emission profiles is shown for NIRSPEC in Figure 2 and for CRIRES in Figure 3. Although the biases in our sample selection make any statistics difficult to interpret, $\text{Pf}\beta$ appears to be a robust tracer of accreting systems. Greater than 80% of optically thick, classical disks around T Tauri and Herbig Ae/Be stars show detectable $\text{Pf}\beta$ emission. It is interesting to note that we find a slightly lower detection fraction for embedded protostars ($\sim 60\%$). A similar fraction of embedded sources without $\text{Pf}\beta$ emission is found in a lower resolution VLT-ISAAC study of ~ 30 sources by Pontoppidan et al. (2003, their Figures 1–7). Although the difference in detection rates may not be significant, especially because of possible biases in the selection of embedded objects, the presence of a large number of embedded protostars with little or no detectable $\text{Pf}\beta$ is interesting in and of itself. Because embedded protostars have rising continua toward the infrared, one might imagine that this could decrease line/continuum ratios and introduce an observational bias. However, the embedded protostars with $\text{Pf}\beta$ detections do not appear to be systematically fainter than those without, nor do the embedded protostars appear to have systematically higher M -band fluxes than the protostars with classical disks.

The majority of targets with $\text{Pf}\beta$ non-detections represent a few distinct types of objects. Not surprisingly, debris disks and disks classified as weak-line T Tauri stars (wTTS's) or protostars with class III spectral energy distributions (SEDs) generally do not appear to show $\text{Pf}\beta$ in emission. wTTS's with weak or non-existent $\text{Pf}\beta$ emission include HD 98800 (also a circumbinary disk; Furlan et al. 2006), Hen 3-600A, LkH α 332 G1, TWA 7, TWA 8a, and WaOph 4. These additionally show no CO fundamental emission and very low levels of veiling, likely reflecting the expected relationship between warm inner disk gas, inner disk dust, and accretion. Other targets without detectable $\text{Pf}\beta$ emission are in some way straddling the evolutionary boundary between optically thick and tenuous disks. These include HD 36917 (classified as transitioning between classes II and III in Manoj et al. 2002), CoKu Tau/3 (class II in Andrews & Williams 2005; wTTS in Furlan et al. 2011), FN Tau (classified as bordering between classical T Tauri star—cTTS—and wTTS in Furlan et al. 2011), HQ Tau (wTTS by H α equivalent width definitions, but cTTS according to its H α 10% width; Furlan et al. 2011), and WSB 60 (classified as Class II in Evans et al. 2009 but shows transitional disk millimeter emission shape in Andrews et al. 2011). Other targets with no detectable $\text{Pf}\beta$ emission include the transitional disks HD 149914 (which also shows no detectable Br γ emission; Brittain et al. 2007) as well as SR 21 (Brown et al. 2007). SR 21 has detectable CO fundamental emission, although in contrast to many other transitional disks, the emission

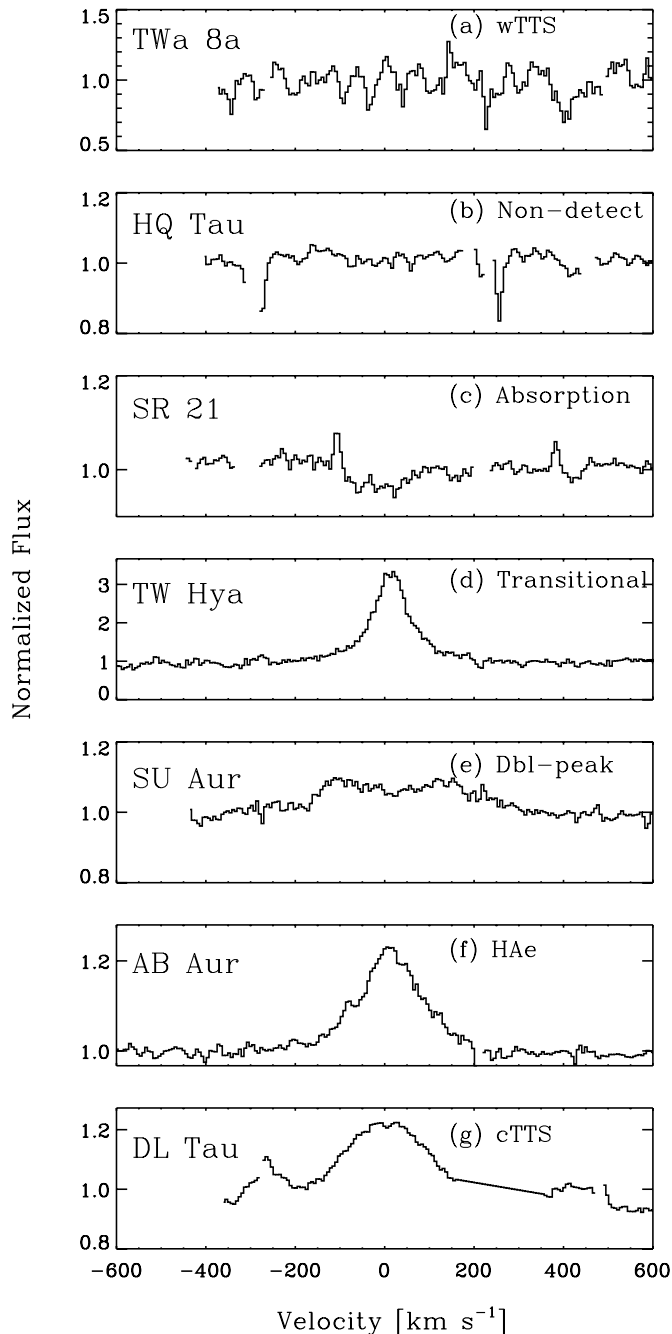


Figure 2. Sample set of $Pf\beta$ emission line profiles from full NIRSPEC sample. (Note that any overlying CO emission or absorption lines have already been removed.)

arises from moderately large disk radii (~ 5 – 8 AU; Salyk et al. 2011; Pontoppidan et al. 2011a). Two circumbinary disks, CoKu Tau/4 (7.8 AU; Ireland & Kraus 2008) and ROXs 42c (23 AU; Kraus et al. 2011), also do not show $Pf\beta$ emission, although it should be noted that this does not appear to be universally true, as close binary DP Tau (projected separation 15.5 AU; Kraus et al. 2011) has detectable $Pf\beta$ emission. A final interesting set of objects that universally shows no detectable $Pf\beta$ emission is the set of FU Orionis stars.

It is interesting to ask whether any classical, optically thick disks show no detectable $Pf\beta$ emission, either to test whether $Pf\beta$ could be unreliable as an accretion tracer, or to search for unique targets with optically thick disks but no accretion. The list of disks with class II SEDs that do not show $Pf\beta$ emission above

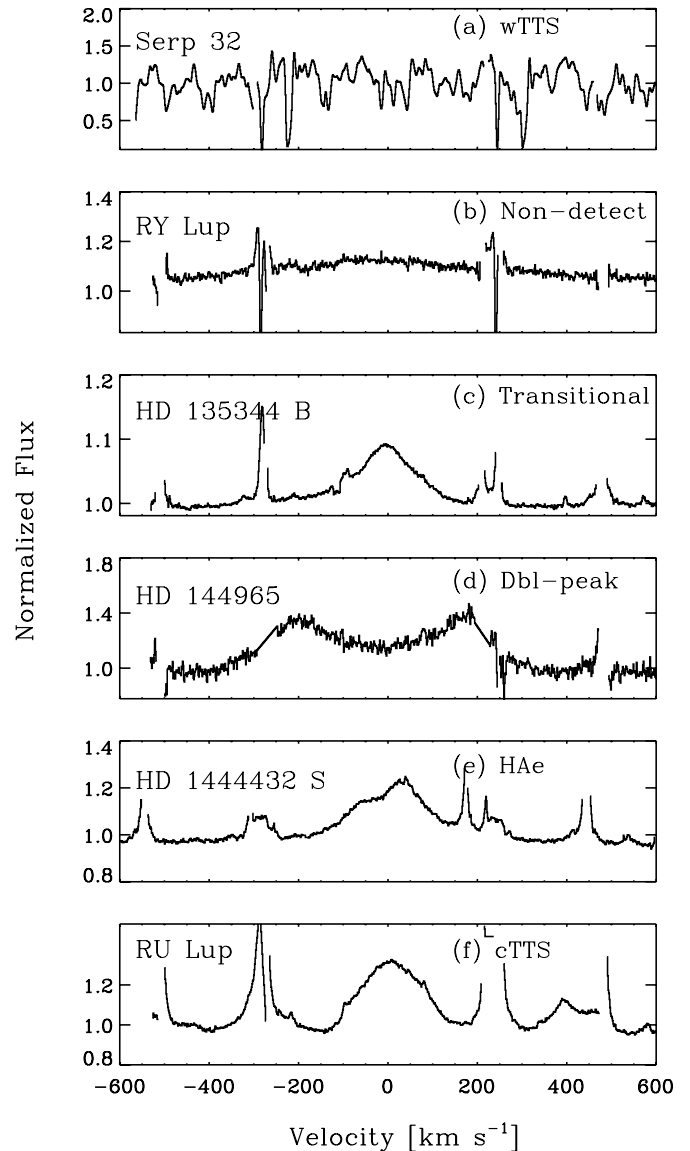


Figure 3. Sample set of profiles from full CRITES sample. (Note that any overlying CO emission lines have already been removed.)

detectable limits include c2dJ033035.92+303024.4, HK Tau, IRAS 03380+3135, IRAS 04385+2550, LkH α 270, LkCa 8, LkH α 325, and WaOph 5. The first five show no apparent CO emission, while the final three targets show weak CO fundamental emission, and their spectra suggest the possible presence of $Pf\beta$ at low S/N; therefore, these suggest the presence of inner disk gas. Few disks are seen with strong CO fundamental emission and no detectable $Pf\beta$ emission, one exception being the transitional disk SR 21. Very low accretion rates are also apparently not detectable in our sample. For example, the spectroscopic binary Hen 3-600A, accreting at a rate of $\sim 5 \times 10^{-11} M_{\odot} \text{ yr}^{-1}$ (Muzerolle et al. 2000), does not have detectable $Pf\beta$ emission.

We investigated but did not determine any relationship between line shapes or widths and any other stellar or disk parameters, except that sources with double-peaked line profiles (including CI Tau, HD 141569 A, RY Tau, and SU Aur) tend to have moderately high disk inclinations (typically 50° – 60°). A full discussion of the relationship between CO vibrational emission and $Pf\beta$ emission is left as future work. Here we focus

Table 3
Stellar and Disk Parameters

Star	Dist (pc)	SpT	T_* (K)	L_* (L_\odot)	R_* (R_\odot)	M_* (M_\odot)	A_V	Refs	Stage ^a	Ref	Inst. ^b
AA Tau	140	K7	4060	0.71	1.74	0.53	0.74	38	II	...	N, C
AB Aur	140	A0	9840	48	...	2.40	0.5	32	II	...	N
AS 205 N	160	K5	4450	7.1	...	1.50	2.9	81	II	...	N, C
AS 205 S	160	M3	3450	2.19	...	0.30	2.1	81	II	...	N
AS 209	160	K5	4395	2.5	2.40	1.40	1.15	52	II	...	C
BF Ori	400	A5–F6	...	27	2.70	2.10	2.4	7, 60	II	...	N
CI Tau	140	K7	4060	0.84	...	0.77	1.8	104	II	...	N
CK 4	260	K3	4730	12.8	27, 60	II ^c	27	N
CrA IRS 2	130	K2	4900	20	44	I	44	C
CV Cha	215	G8	5451	8.0	3.20	2.10	1.67	52	II	...	C

Notes.

^a II^t refers to known transitional disks with inner disk cavities.

^b Instrument. N = NIRSPEC, C = CRIRES.

^c CK 4 (SSTc2d J182958.19+011521.8) has a weak ice absorption feature in NIRSPEC spectra, but its spectral slope, bolometric temperature and lack of extended millimeter emission suggest that it is a stage II object (Evans et al. 2009).

^d Haro 6-13 is often classified as a borderline stage I/II object (Teixeira & Emerson 1999; White & Hillenbrand 2004).

^e HL Tau is often classified as a borderline stage I/II object (Teixeira & Emerson 1999; White & Hillenbrand 2004), but is classified as stage I by Herczeg et al. (2011) due to the presence of an extended envelope, cold compact gas and molecular outflow.

^f This source is borderline stage I/II based on its “flat” spectral slope and bolometric temperature, and the presence of an extended envelope (Evans et al. 2009). However, its M -band spectrum shows no evidence for CO ice, and we adopt a stage II classification.

^g IRS 48 has a weak CO ice feature in its CRIRES spectrum, but appears to be an evolved stage II disk, perhaps transitioning to stage III (Brown et al. 2012).

^h IRS 51 is a flat spectrum source with a bolometric temperature in the Class II range, but with evidence of an extended envelope (Evans et al. 2009). We adopt a stage I classification due to its prominent M -band ice feature.

ⁱ Assuming the properties of SSTc2d +J162659.10-243503.3 (Evans et al. 2009).

References. (1) Aceke et al. 2005; (2) Andrews & Williams 2007; (3) Andrews et al. 2011; (4) Barsony et al. 2005; (5) Basri & Batalha 1990; (6) Beckwith et al. 1990; (7) Blondel & Djie 2006; (8) Boersma et al. 2008; (9) Bontemps et al. 2001; (10) Boogert et al. 2002; (11) Borges Fernandes et al. 2009; (12) Bouvier & Appenzeller 1992; (13) Briceño et al. 2002; (14) Brown et al. 1995; (15) Brown et al. 2007; (16) Brown et al. 2012; (17) Calvet et al. 2004; (18) Caratti o Garatti et al. 2004; (19) Chini 1981; (20) Cohen & Kuhl 1979; (21) Correia et al. 2006; (22) de Lara et al. 1991; (23) Doppmann et al. 2003; (24) Doppmann et al. 2005; (25) Duchêne et al. 2004; (26) Dunkin & Crawford 1998; (27) Evans et al. 2009; (28) Feigelson et al. 1993; (29) Friedemann et al. 1993; (30) Furlan et al. 2006; (31) Furlan et al. 2009; (32) Garcia-Lopez et al. 2006; (33) Ghez et al. 1993; (34) Gras-Velázquez & Ray 2005; (35) Greene & Meyer 1995; (36) Gredel et al. 1997; (37) Guenther et al. 2007; (38) Gullbring et al. 1998; (39) Hartigan & Kenyon 2003; (40) Hartmann et al. 1998; (41) Herbst & Warner 1981; (42) Herczeg et al. 2004; (43) Herczeg & Hillenbrand 2008; (44) Herczeg et al. 2011; (45) Hernández et al. 2004; (46) Hernández et al. 2005; (47) Heyer et al. 1990; (48) Hillenbrand et al. 2008; (49) Hodapp 1994; (50) Hughes et al. 1994; (51) Jensen et al. 2009; (52) Johns-Krull et al. 2000; (53) Johns-Krull & Gafford 2002; (54) Keller et al. 2008; (55) Kenyon & Hartmann 1995; (56) Kessler-Silacci et al. 2006; (57) Köhler et al. 2008; (58) Koresko et al. 1997; (59) Lada et al. 2006; (60) Lahuis et al. 2007; (61) Levreault 1985; (62) Liseau et al. 1992; (63) Loinard et al. 2008; (64) Luhman & Rieke 1999; (65) Malfait et al. 1998; (66) Mannings & Sargent 2000; (67) Manoj et al. 2006; (68) Massi et al. 2006; (69) Mendigutía et al. 2011; (70) Miroshnichenko 2007; (71) Monnier et al. 2005; (72) Montesinos et al. 2009; (73) Mora et al. 2001; (74) Myers et al. 1987; (75) Natta et al. 2006; (76) Neckel et al. 1980; (77) Nisini et al. 2005; (78) Pascucci et al. 2007; (79) Perryman et al. 1997; (80) Pontoppidan et al. 2007; (81) Prato et al. 2003; (82) Prato et al. 2009; (83) Rice et al. 2006; (84) Seperuelo Duarte et al. 2008; (85) Shevchenko & Herbst 1998; (86) Shevchenko et al. 1999; (87) Simon et al. 1992; (88) Steele et al. 1999; (89) Stelzer et al. 2009; (90) Straižys et al. 1996; (91) Strom et al. 1972; (92) Testi et al. 1998; (93) Thompson et al. 1998; (94) Valenti et al. 1993; (95) Valenti et al. 2003; (96) van den Ancker et al. 1998; (97) van Boekel et al. 2005; (98) van Leeuwen 2007; (99) Vieira et al. 2003; (100) Walter 1986; (101) Watson et al. 2009; (102) White et al. 2007; (103) White & Hillenbrand 2004; (104) White & Ghez 2001; (105) Wichmann et al. 1998; (106) Yudin 2000.

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on targets with detectable $Pf\beta$ emission, in order to measure mass accretion rates. Therefore, for this study we have selected a subset of the full NIRSPEC and CRIRES samples via visual confirmation of the presence of $Pf\beta$ emission. The complete target list used in this study, with relevant stellar parameters, is shown in Table 3.

$Pf\beta$ lines are unresolved in all natural-seeing NIRSPEC observations, constraining the emission to radii less than ~ 50 AU for typical seeing and stellar distances. $Pf\beta$ observations are also unresolved in the adaptive optics corrected CRIRES spectra. With a typical PSF core of $0''.18$, this constrains the $Pf\beta$ to radii less than ~ 13 AU at 140 pc. This is consistent with an accretion origin for $Pf\beta$, as opposed to an outflow origin.

4. DERIVATION OF $Pf\beta$ LINE LUMINOSITIES

4.1. CO Emission Corrections

The $Pf\beta$ line center is close to a number of CO vibrational lines, including $v = 1 \rightarrow 0$ R(0) ($4.6575 \mu\text{m}$) and R(1) ($4.6493 \mu\text{m}$), $v = 2 \rightarrow 1$ R(8) ($4.6523 \mu\text{m}$) and ^{13}CO R(14) ($4.6572 \mu\text{m}$) and R(15) ($4.6504 \mu\text{m}$). A majority of $Pf\beta$ profiles are affected by at least one of these CO features, which contain non-negligible amounts of flux. A number of techniques were utilized to correct for the CO emission, as demonstrated in Figure 4. If possible, CO lines were removed via subtraction of a Gaussian fit to the CO emission lines. In cases where the fits were not reliable, the line was replaced by a linear interpolation

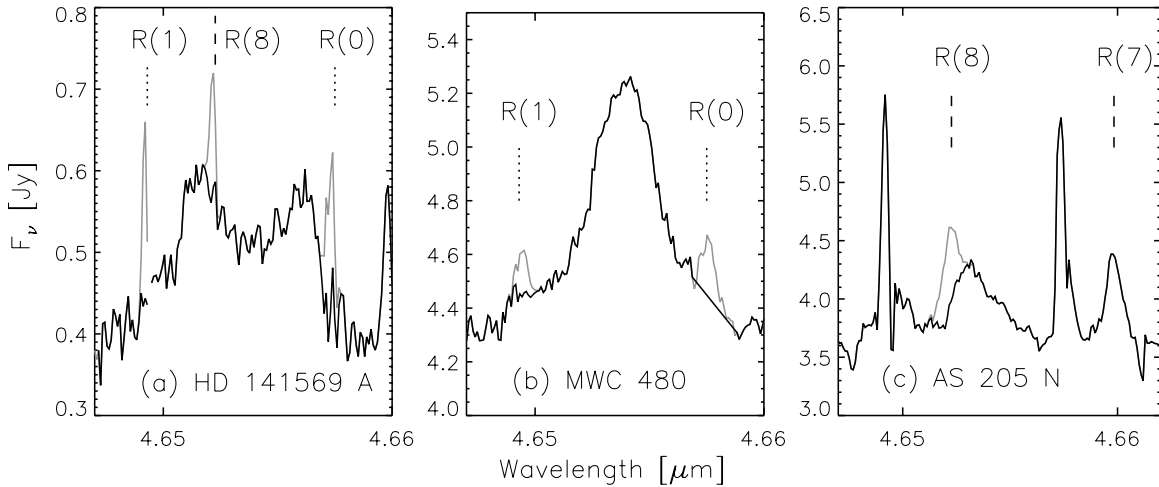


Figure 4. Gray and black lines show sample spectra before and after removal of CO emission lines, respectively. Dotted lines mark CO $1 \rightarrow 0$ lines; dashed lines mark CO $2 \rightarrow 1$ lines. (a) Lines are removed with Gaussian fits. (b) R(0) is removed with a linear fit to the underlying spectrum. (c) R(8) is removed by using a Gaussian fit to R(7).

of the surrounding flux. In some cases, the nearby $v = 2 \rightarrow 1$ or ^{13}CO lines contributed so significantly to the flux that it was difficult to separate the CO and Pf β contributions. In such cases, an adjacent CO line would be shifted to the contaminating line location and subtracted. This necessarily introduces some error (of order 10%) due to the fact that the flux in adjacent CO lines may not be identical. However, more precise modeling to derive these CO line fluxes was not pursued, due to severe blending of the $v = 2 \rightarrow 1$ and ^{13}CO rotational ladders with the much stronger $v = 1 \rightarrow 0$ emission lines.

4.2. Flux Extraction

Observed Pf β line profiles are typically not described by a Gaussian or other simple functional form, and so we chose to extract fluxes by summing all flux within a defined window, after subtraction of the continuum. The extraction window and continuum level were both determined by eye, and have associated systematic uncertainties. These can be particularly high if a number of adjacent CO lines makes the continuum hard to define, and/or if the Pf β profile has broad wings. The sample overlap between NIRSPEC and CRILES allows us to estimate this uncertainty to some degree, by comparing the fluxes derived from each data set. We find good agreement between extractions using the two data sets, with typical differences less than a factor of two. Note that since the Pf β line luminosity or the underlying continuum can be variable, some of this difference could be attributable to real changes that occurred between the two sets of observations. In our most egregious case, HL Tau, we find a factor of three discrepancy between the two data sets. This extreme case is illustrated in Figure 5, showing how uncertainty in both the continuum position and the position of the Pf β line wings results in large flux differences between the two extractions.

A few of the photometrically corrected NIRSPEC line profiles are shown in Figure 6; the remaining NIRSPEC and CRILES profiles are available in the online version of the article. Our derived Pf β luminosities, computed using the distances in Table 3, are listed in Table 4. Line luminosities have also been corrected for extinction, assuming an extinction of $0.034 \times A_V$ (Cardelli et al. 1989) at $4.6538 \mu\text{m}$, and the visual extinctions in Table 3; if the extinction is not known, we do not include any correction. Note that even the maximum extinction in our sample ($A_V = 34$) only results in a change in line flux by a factor of ~ 3 ,

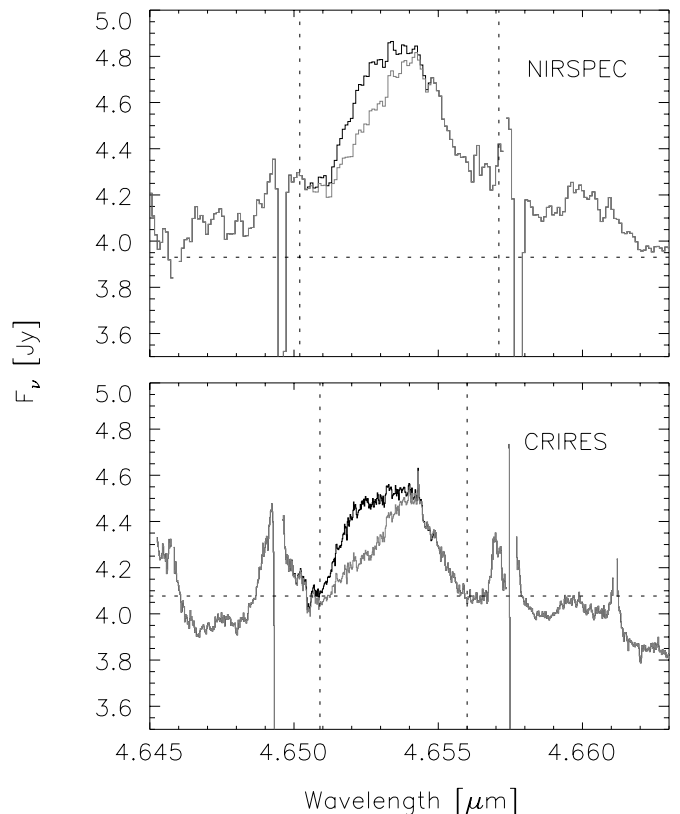


Figure 5. Pf β line profiles from HL Tau observed by NIRSPEC and CRILES before (black) and after (gray) CO emission removal. Dotted lines mark the regions chosen for flux extraction, demonstrating how uncertainty in the continuum level and line width can in extreme cases result in a factor of three difference in extracted line flux.

and thus computation of Pf β luminosity is relatively insensitive to extinction. When both NIRSPEC and CRILES data were available, we show the average of those measurements.

4.3. Importance of Photospheric H I Absorption

While accretion flows are believed to be the source of the Pf β emission, the observed spectra also include a contribution from the underlying stellar photosphere of the star/disk system.

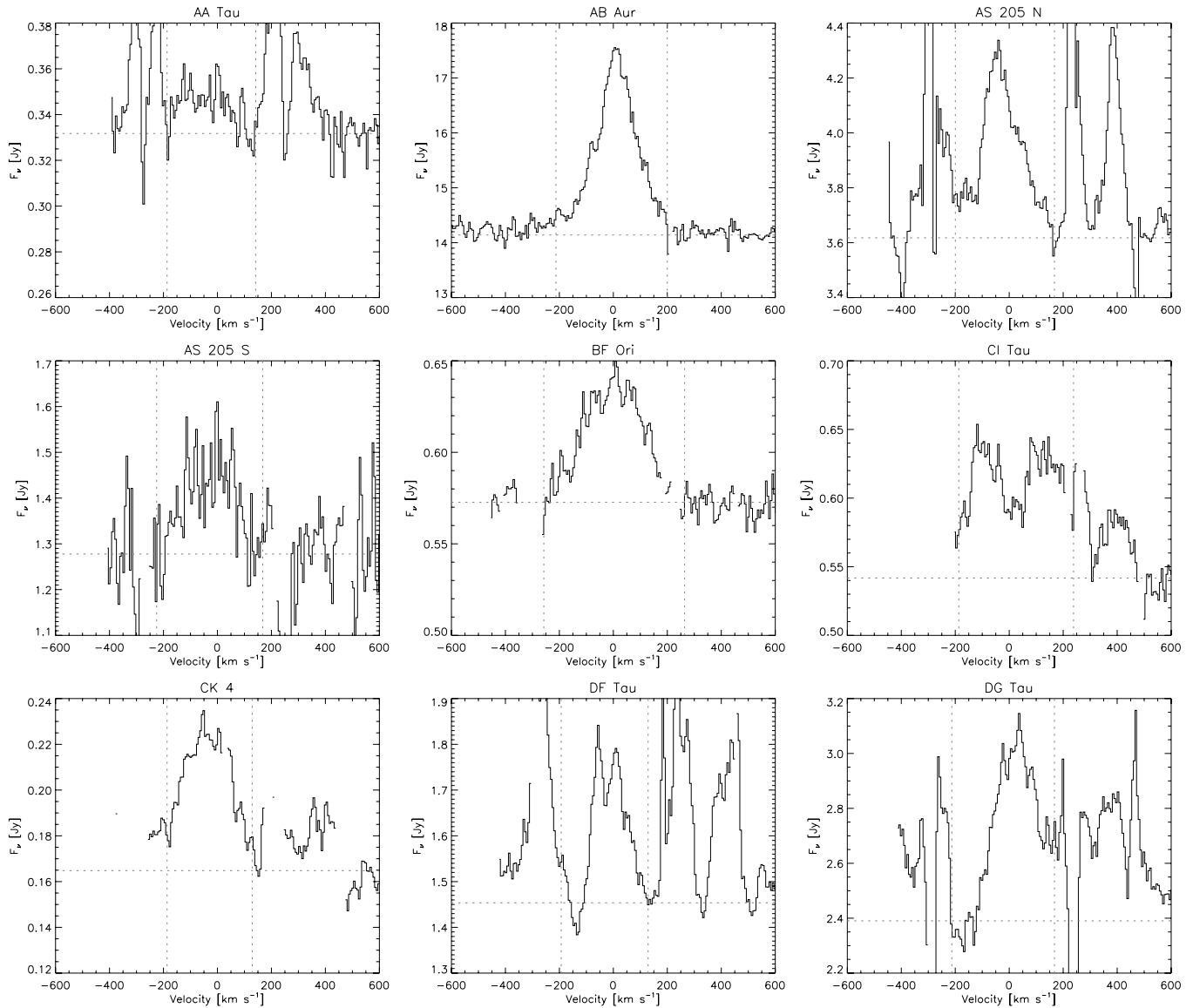


Figure 6. Pf β line profiles observed with NIRSPEC, after correction for CO emission and literature photometry, when available. Dotted lines mark the baseline and limits used for flux extraction. The full set of NIRSPEC (6a–6j) and CRILES (6k–6p) line profiles are available in an online version of this figure.

(An extended version of this figure is available in the online journal.)

As we will discuss in this section, the correction for the underlying stellar photosphere should in nearly all cases be insignificant compared to other uncertainties, and we therefore do not apply any correction to most sources. (Note that this correction is distinct from the correction for Pf β absorption in the photosphere of the telluric standard star, which is always corrected for in the data reduction process.)

Stellar photospheres have H I absorption features that will tend to reduce the amount of observed Pf β flux, i.e., Pf β fluxes must be increased somewhat to correct for stellar absorption. The degree of correction depends sensitively on the stellar spectral type and veiling. Figures 7(a)–(e) show synthetic spectra and observed spectra for five spectral types (A0V–M0V), demonstrating that the absorption is most prominent at early spectral types. However, M -band veiling values for optically thick disks act in the opposite sense, with much higher veiling around early-type stars. Thus optically thick disks have negligible corrections to Pf β from underlying stellar absorption. Synthetic veiled spectra with veiling values typical of optically thick disks ($r = 46, 18, 10, 7,$ and 3 for spectral types A–M, derived from disk

SEDs produced with RADMC; Dullemond & Dominik 2004) are shown as dotted lines, and are indistinguishable from straight lines.

Stars with G–M spectral types have M -band spectra that are dominated by a molecular pseudo-continuum, rather than simply H I absorption (see panels (c), (d), and (e)), which can result in some change to Pf β flux due to coincident molecular absorption features near Pf β . The equivalent width of the underlying photospheric features drop by a factor of $1 + r$ where r is the veiling (defined as the dust continuum flux density divided by the stellar photospheric flux density). Considering the spectrum of HD 79210 shown in Figure 7, a reasonable M -band veiling value of ~ 3 for an optically thick disk reduces the strongest photospheric features to the $\sim 7\%$ level; this and their narrow width results in a negligible correction for optically thick disks.

Transitional disks—disks with inner regions depleted of small dust grains—have low continuum fluxes, and so could potentially have large relative contributions from the underlying photosphere. However, we find that such targets (see Figures 7(f), (h), and (j)) often have high line/continuum ratios, likely

Table 4
Derived Pf β Accretion Parameters; EW's of Other Accretion Tracers

Star	EW _{Pfβ} (Å)	log $L_{\text{Pf}\beta}$ ^a (L_{\odot})	log L_{acc} (L_{\odot})	log \dot{M} ^b ($M_{\odot} \text{ yr}^{-1}$)	EW (Br γ) (Å)	EW (Pa β) (Å)	Refs
AA Tau	-2.21	-5.21	-1.43	-8.31	-0.59	-2.58	2
AB Aur	-6.75	-3.09	0.50	-6.90	-5.50	...	3
AS 205 N	-3.59	-3.83	-0.18	-7.10
AS 205 S	-5.29	-4.10	-0.42	-6.68
AS 209	-4.41	-4.02	-0.35	-7.52
BF Ori	-4.74	-3.69	-0.05	-7.34	-3.60
CI Tau	-8.88	-4.37	-0.66	-7.68	-4.20	-7.00	5
CK 4	-12.67	-4.04	-0.37
CrA IRS 2	-7.56	-3.33	0.28
CV Cha	-11.17	-3.18	0.41	-6.81

Notes.

^a Adjusted for photometry and extinction, as described in text. Assumes distances in Table 3. If both NIRSPEC and CRIRES data were available, the average value is shown.

^b Calculated assuming $L_{\text{acc}} = 0.8 GM_{\star} \dot{M} / R_{\star}$ (Gullbring et al. 1998). If not shown in Table 3, R_{\star} is derived from L_{\star} and T_{\star} assuming $L_{\star} = 4\pi R_{\star}^2 \sigma T_{\star}^4$. Missing data indicate that either the stellar radius or mass is not available.

References. (1) Donehew & Brittain 2011; (2) Folha & Emerson 2001; (3) Garcia-Lopez et al. 2006; (4) Muzerolle et al. 1998b; (5) Natta et al. 2006.

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because in spite of their lowered near-IR continuum fluxes, transitional disks have only slightly lowered accretion rates (Najita et al. 2007). (See the extended discussion in Section 6.) As examples, corrections to the Pf β flux due to photospheric H I from transitional disks HD 141569 A, LkH α 330, and GM Aur (panels (f), (h), and (j)) would be only $\sim 20\%$, 45% , and 3% , respectively, even in the worst case scenario (no veiling). Actual corrections for these sources should be a factor of a few lower, due to non-zero M -band veiling (Salyk et al. 2009). The molecular photospheric absorption features might be more problematic, but luckily the spectra can be examined at other wavelengths to determine whether there is likely significant contamination from the underlying photosphere. With visual examination, we found one target—DoAr 21—which shows both low Pf β equivalent width and low veiling, and in this case the molecular absorption features in its stellar photosphere have a non-trivial effect on the determination of the disk Pf β line flux. We describe this example in detail in the Appendix.

The largest possible correction would be required for a naked A or F star with a low Pf β line/continuum ratio (see, e.g., panel (g)), perhaps doubling the ratio of true to observed Pf β emission flux. Such a target would be difficult to identify because its relatively featureless M -band spectrum looks the same whether it is naked or veiled. To our knowledge, there are no known diskless targets in our sample that also show detectable Pf β emission. Of course, stellar Pf β absorption could erase signatures of weak Pf β emission from weakly accreting disks; however, we have no good way to pinpoint such targets. A possible example of such a target could be SR 21 (see Figure 2) a G2.5 star with a transitional disk, which has CO gas at ~ 7 AU but no measurable accretion (Pontoppidan et al. 2008).

5. Pf β LUMINOSITY AND ACCRETION RATES

5.1. Correlation between Pf β Luminosity and Accretion Luminosity

In Figure 8, we show the correlation between our measured Pf β line luminosities and previously measured accretion luminosities, listed in Table 5. We chose to maximize the number of

Table 5
Literature Accretion Luminosities and Mass Accretion Rates

Name	L_{acc} (L_{\odot})	Ref	Method ^a	log \dot{M} ($M_{\odot} \text{ yr}^{-1}$)	Ref
AA Tau	0.025	10	UV Exc	-8.48	10
AB Aur	4.3	9	Br γ	-6.85	9
AS 205 N	-6.14	7
AS 205 S
AS 209	0.452	18	UV Exc	-7.29	18
BF Ori	-6.96	4
CI Tau	0.112 ^b	18	UV Exc	-7.83	18
CK 4
CrA IRS 2
CV Cha

Notes.

^a Description of abbreviations are as follows. UV Exc. = measurement of UV excess emission spectrum; Br γ = use of Br γ -accretion luminosity relationship; Pa β = use of Pa β -accretion luminosity relationship; U -Band Phot. = use of optical U -band photometric data; Balmer depth = based on relationship between the mass accretion rate and the excess in the Balmer discontinuity derived by Muzerolle et al. (2004).

^b This and all other Taurus accretion luminosities from Valenti et al. (1993) scaled to 140 pc distance.

References. (1) Andrews & Williams 2007; (2) Andrews et al. 2011; (3) Aspin et al. 2010; (4) Blondel & Djie 2006; (5) Calvet et al. 2004; (6) Donehew & Brittain 2011; (7) Eisner et al. 2005; (8) Fernandez et al. 1995; (9) Garcia-Lopez et al. 2006; (10) Gullbring et al. 1998; (11) Hartmann et al. 1998; (12) Herczeg & Hillenbrand 2008; (13) Muzerolle et al. 1998b; (14) Mendigutia et al. 2011; (15) Natta et al. 2006; (16) Nisini et al. 2005; (17) Prato et al. 2009; (18) Valenti et al. 1993; (19) White & Ghez 2001; (20) White & Hillenbrand 2004.

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sources in this correlation and, therefore, the literature accretion luminosities are derived from a variety of observational methods. The majority are derived from models and observations of the UV excess continuum, while others are derived somewhat less directly from UV data, utilizing a relationship between mass accretion and the Balmer discontinuity (Muzerolle et al. 2004)

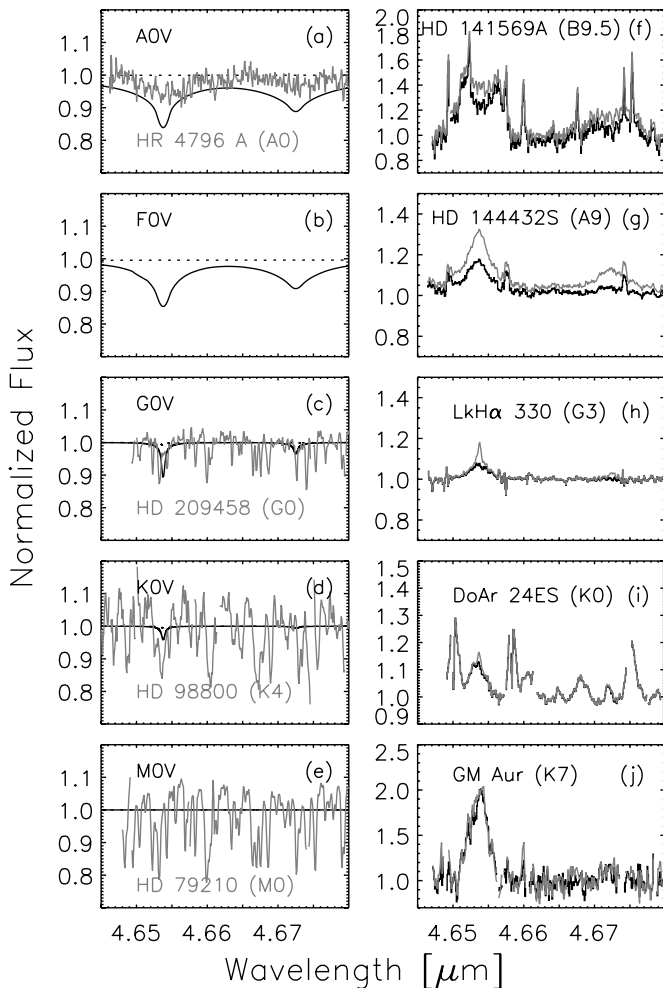


Figure 7. Panels (a)–(e) show synthetic stellar H I absorption spectra without (solid curves) and with (dotted curves) the addition of typical veiling from an optically thick disk. (A0V and F0V spectra are produced from ATLAS models using the ATLAS, WIDTH, and SYNTHÉ Linux port (Kurucz 1993; Sbordone et al. 2004, 2005). G0V, K0V, and M0V models are produced with the MOOG stellar synthesis code (Snedden 1973) using MARCS atmospheric models (Gustafsson et al. 2008), but show only the contribution from H I lines. The actual spectra are dominated by molecular absorption, and stellar models in this wavelength range have not been well benchmarked. All spectra assume a rotational broadening of 10 km s^{-1} . Representative veiling values were derived using RADMC (Dullemond & Dominik 2004), assuming optically thick disks that extend to $T_{\text{in}} = 1500 \text{ K}$.) Panels (a), (c), (d), and (e) additionally show observed spectra of unveiled stars, with the K4 and M0 spectra demonstrating that the spectra of cool stars are dominated by a molecular pseudo-continuum. Panels (f)–(j) show example spectra (black) and *worst-case* Pf β flux corrections (i.e., assuming no veiling; gray) for the following targets: (f) transitional disk HD 141569 A, (g) Herbig Ae/Be star HD 144432 S, (h) transitional disk LkH α 330, (i) cTTS DoAr 24 ES, and (j) transitional disk GM Aur.

or approximating the UV excess using photometry. About a third of the accretion rates are derived from empirical relationships between H I emission lines (Br γ or Pa β) and accretion luminosity. Although the Br γ and Pa β -derived accretion rates are themselves an indirect measure of accretion rate, we find no significant biases between these and the UV-derived accretion luminosities in Figure 8.

Because the accretion rates of high-inclination sources can be unreliable, we remove these from our correlation. We also omit HL Tau from the correlation; although its disk appears to be viewed at moderate inclination (Kwon et al. 2011), its high level of veiling makes literature accretion rates very uncertain (e.g., White & Ghez 2001).

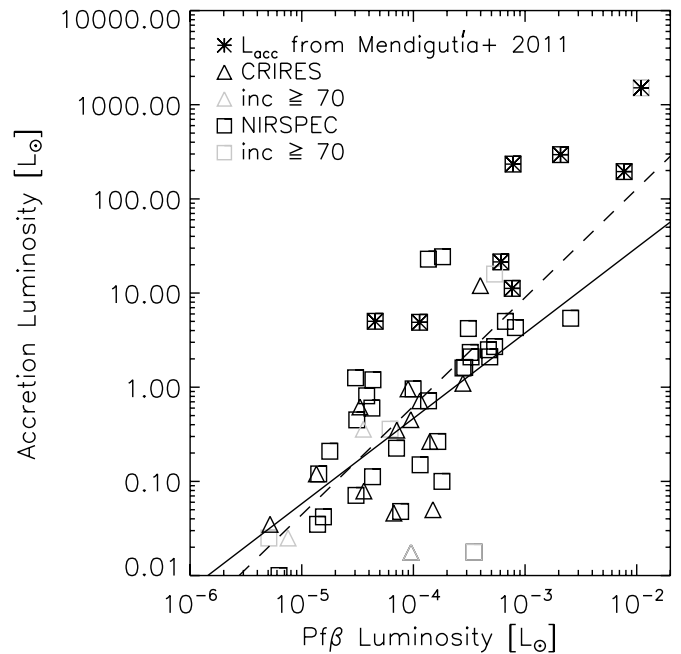


Figure 8. Correlation between literature accretion luminosity (see Table 5) and Pf β luminosity (this work). If Pf β is measured by both NIRSPEC and CRIRES, we show two symbols (squares and triangles, respectively) on the plot. Sources with high inclination or otherwise unreliable accretion rates (in gray) are excluded from the correlation analysis. The solid line shows the best linear fit to the data excluding points from Mendigutía et al. (2011) (Equation (1)). The dashed line shows the best linear fit including points where L_{acc} is derived from Mendigutía et al. (2011).

We find that the correlation between Pf β line luminosity and accretion luminosity is heavily influenced by targets with high accretion luminosities derived by Mendigutía et al. (2011). Mendigutía et al. (2011) derived accretion rates in 38 HAeBe stars by estimating the Balmer excess using U - and B -band photometry compared with accretion shock models. Since accretion rates correlate with stellar mass, these accretion rates populate the upper right corner of Figure 8. If these points are included in the fit, we find that our correlation then poorly reproduces previously measured accretion rates for low-mass stars. Indeed, Mendigutía et al. (2011) note a change in the slope of the correlation between accretion rate and stellar mass between low-mass and high-mass stars, and find larger accretion rates than those determined from Br γ emission lines (García-López et al. 2006), implying that correlations between accretion tracers and accretion rates may not extend linearly to higher mass stars. However, UV excesses are significantly more difficult to measure around high-mass stars, and additional study is warranted to understand whether the emission line-accretion luminosity relationship needs to be modified at higher stellar masses. Therefore, in this study, we omit the accretion luminosities derived by Mendigutía et al. (2011) from our analysis. If the results of Mendigutía et al. (2011) are correct, our study will tend to underestimate accretion luminosities for $M_{\star} \gtrsim 3 M_{\odot}$.

Excluding the accretion luminosities in Mendigutía et al. (2011), we find the following relationship between Pf β luminosity and previously measured accretion luminosities:

$$\log L_{\text{acc}}[L_{\odot}] = (0.91 \pm 0.16) \times \log L_{\text{Pf}\beta}[L_{\odot}] + (3.29 \pm 0.67). \quad (1)$$

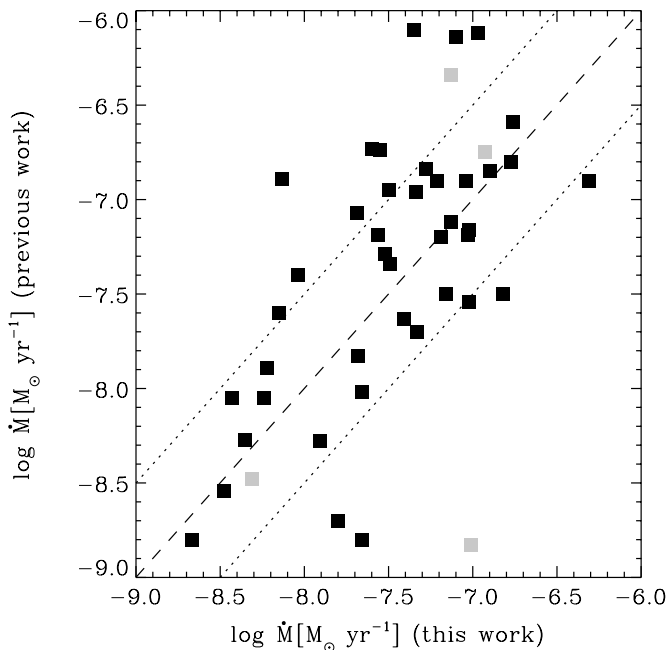


Figure 9. Mass accretion rates derived from this work compared with values from the literature. High-inclination disks are shown plotted in gray. The dashed line shows a 1:1 correlation and dotted lines show a difference of 0.5 dex.

5.2. New Accretion Rates

We use Equation (1) to compute accretion luminosities for our entire sample. These are listed in Table 4. An estimate of \dot{M} is derived using the standard relationship between L_{acc} and \dot{M} :

$$L_{\text{acc}} = 0.8 GM_{\star} \dot{M} / R_{\star} \quad (2)$$

from Gullbring et al. (1998). M_{\star} and R_{\star} are taken from Table 3. If not shown in Table 3, R_{\star} is derived from L_{\star} and T_{\star} assuming $L_{\star} = 4\pi R_{\star}^2 \sigma T_{\star}^4$. We do not compute \dot{M} if these parameters are not available. However, \dot{M} could easily be computed at a later time when these parameters are measured.

In Figure 9, we show accretion rates derived here compared to existing values from the literature (Table 5). The accretion rates derived here show no systematic bias, and most accretion rates are within 0.5 dex of previous measurements. The standard deviation of the current and previous measurements is 0.77 dex. These results are similar to those for other emission line accretion tracers, albeit over a somewhat more limited range of accretion rates. For example, Herczeg & Hillenbrand (2008) quote standard deviations of 1.0, 0.71 and 1.1 dex between UV continuum excess and Ca II $\lambda 8542$, Ca II $\lambda 8662$, and He I $\lambda 8576$, probing accretion rates as low as $10^{-12} M_{\odot} \text{ yr}^{-1}$. In addition, Herczeg & Hillenbrand (2008) estimate a factor of four (0.6 dex) random error in UV excess measures alone, due to errors in extinction and distance. Also, accretion rates are known to be variable, with Nguyen et al. (2009) finding typical accretion rate variations of 0.35 dex, but variations commonly as high as 0.5 dex. Therefore, Pf β appears to be an accretion tracer on par with existing tracers, with scatter not much larger than what would be expected from accretion variability and errors related to the derivation of UV-based accretion luminosities.

5.3. About Upper Limits

Computing upper limits on the Pf β line flux, and thus accretion rate, for any given source is not straightforward,

and because these are not crucial for the work presented here, we do not make an attempt to compute them. Complications include the fact that the Pf β line shapes are complex, and their widths varied, so the correct choice of assumed line shape and width is not obvious. In addition, sources with no detectable Pf β emission are often less-heavily veiled sources, whose underlying photospheres must be modeled to derive an accurate Pf β flux. We encourage those interested in an upper limit for a particular source in the NIRSPEC or CRIRRES archives to contact the authors of this work to discuss whether reasonable assumptions can be made to derive limits for this source.

Nevertheless, we can make some general observations about how sensitive Pf β measurements are to accretion rates. Assuming the Pf β line is a Gaussian with $\sigma_{\text{width}} = 0.001 \mu\text{m}$ (derived from a fit to AB Aur’s Pf β emission line), then with a continuum flux density of $F_{\nu} = 1 \text{ Jy}$, a S/N of 50, a 3σ line peak, and for a 1 solar mass, 1 solar radius star, the minimum measurable accretion rate is $\sim 4 \times 10^{-9} M_{\odot} \text{ yr}^{-1}$. The minimum measurable accretion rate then scales roughly proportional to $F_{\nu}/(\text{S/N})$. For the integration times used for the acquisition of the data used in this work, Pf β -derived accretion rates do not appear to be as sensitive as previous studies using Bry γ or Pa β (e.g., Natta et al. 2006).

The lowest detected accretion rates in our sample are for the transitional disks TW Hya, and DM Tau, with rates of 2.6×10^{-9} and $4.3 \times 10^{-9} M_{\odot} \text{ yr}^{-1}$, respectively. Pf β is especially well-suited for measuring small accretion rates from transitional disks, as the dust-cleared inner holes decrease F_{ν} and increase the line/continuum ratio (see Section 6.2 for more detail). However, as the S/N per unit exposure time also decreases as F_{ν} decreases, sufficient time must be expended to reach a reasonable S/N.

6. ANALYSIS AND DISCUSSION

6.1. Utility of Pf β and Comparison to Other Accretion Tracers

In this work, we have introduced the use of H I Pf β to measure mass accretion from protoplanetary disks. Although numerous accretion tracers have been developed prior to this work, the use of Pf β offers several advantages, which we outline here.

One major advantage of using Pf β as an accretion tracer is that it comes “for free” and contemporaneous with observations of CO fundamental emission. The NIRSPEC and CRIRRES surveys from which we extract the data in this work provide a self-consistent sample of accretion luminosity estimates for 120 targets, and the NIRSPEC and CRIRRES archives likely include many additional protoplanetary disk targets for which accretion luminosities could be estimated. To our knowledge, this is the first work to make extensive use of M -band spectra of protostars from both NIRSPEC and CRIRRES—data which are all currently available in their respective archives. The varied and large set of targets resulting from this combined data set is a great demonstration of the utility of these large archives. As future work, we plan to extend our analysis to investigate long-term (\sim year timescales) accretion variability for different classes of disks.

The simultaneous observation of CO and Pf β also allows for a comparison between the accretion flow and molecular disk gas at distances of ~ 0.1 –1 AU from the star (e.g., Pontoppidan et al. 2011a; Salyk et al. 2011), which is traced by CO fundamental emission. The relationship between these two disk components is being investigated, e.g., by Brown et al. (2013). The contemporaneous nature of these measurements also allows

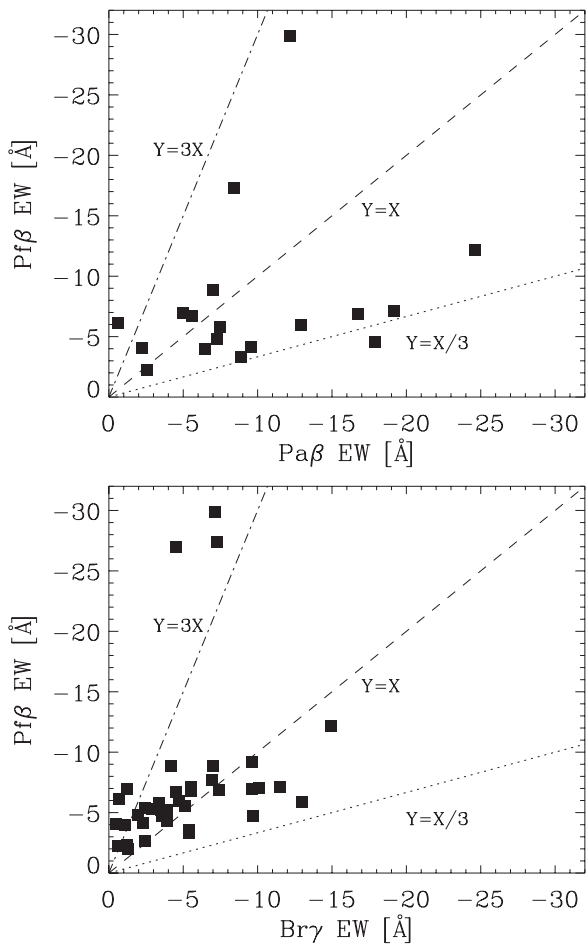


Figure 10. EW of $Pf\beta$ compared to those of $Br\gamma$ (bottom) and $Pa\beta$ (top). The outliers with highest $Pf\beta$ EW in the bottom plot are transitional disks HD 141569 A, GM Aur, and TW Hya.

for simultaneous measurements of disk gas and accretion rate as a function of time, although it should be cautioned that there can be changes in emission line fluxes without corresponding changes in accretion rate (e.g., Gahm et al. 2008). The tidal influence of protoplanets is predicted to cause CO emission line variability (Regály et al. 2011), but understanding the influence of protoplanets requires ruling out other possible sources of line variability, such as changes in accretion rate or accretion geometry.

$Pf\beta$ ($n = 7 \rightarrow 5$) is in many ways similar to $Br\gamma$ ($n = 7 \rightarrow 4$; $2.160 \mu\text{m}$), originating from the same upper level energy and appearing in the infrared. In Figure 10, we compare $Pf\beta$ EW with literature values of $EW(Br\gamma)$ as well as $EW(Pa\beta)$ (see Table 4). Typical values of $Pf\beta$ equivalent widths are between $1/3$ and $3 \times EW(Br\gamma)$ or $EW(Pa\beta)$, with $EW(Pf\beta)$ typically being similar to $EW(Br\gamma)$ but somewhat smaller than $EW(Pa\beta)$. The lack of a strong correlation between tracers likely reflects the varied temperatures of the accretion column, the underlying photosphere, and the continuum veiling. Muzerolle et al. (1998b), Calvet et al. (2004), and Donehew & Brittain (2011) all find correlations between accretion luminosity and $Br\gamma$ luminosity very similar to our Equation (1). Equation (1) combined with the relationships in Muzerolle et al. (1998b) and Donehew & Brittain (2011) predict somewhat higher $EW(Pf\beta)$ than $EW(Br\gamma)$, while the relationship in Calvet et al. (2004) predicts somewhat lower $EW(Pf\beta)$ as compared to $EW(Br\gamma)$ (with differences less than ~ 0.2 dex). Line shapes of the two

accretion tracers appear to be quite similar—typically single-peaked with line widths near $\sim 200 \text{ km s}^{-1}$ (Garcia-Lopez et al. 2006). HD 141569 A, on the other hand, is an example of a disk with both double-peaked $Br\gamma$ and $Pf\beta$ emission lines. Interestingly, however, several targets for which $Pf\beta$ is clearly in emission show no $Br\gamma$, or $Br\gamma$ in absorption—for example, HD 142527, HD 142666, T CrA, and TY CrA. Thus, $Pf\beta$ may be a more robust accretion tracer in targets with low near-IR veiling.

Although equivalent widths are similar, $Pf\beta$ offers some distinct advantages over $Br\gamma$ and $Pa\beta$. One advantage is the relatively higher continuum veiling in the M band as compared to shorter wavelengths. Since the photospheric emission is a smaller fraction of the total flux at $5 \mu\text{m}$ than it is at shorter wavelengths, $Pf\beta$ suffers less contamination and uncertainty from the underlying photospheric H I absorption (as discussed in detail in Section 4.3). Being at longer wavelengths also makes $Pf\beta$ less sensitive to extinction. Using the reddening law of Cardelli et al. (1989), $A_{Br\gamma}/A_V = 0.12$ while $A_{Pf\beta}/A_V = 0.03$, a factor of four difference in *magnitudes* of correction. At $A_V \sim 30$, for example, this results in a flux correction factor of ~ 28 for $Br\gamma$ but only ~ 3 for $Pf\beta$. Finally, $Pf\beta$ equivalent widths are quite high for accreting disks with low continuum veiling; thus, $Pf\beta$ is an excellent tracer for low accretion rates in disks with inner regions depleted in small dust grains (discussed further in Section 6.2). One disadvantage of $Pf\beta$ as an accretion tracer is the systematic errors that are introduced by complex CO emission/absorption spectra, which can result in up to a factor of a few uncertainty in flux in extreme cases (as in Figure 5). Another is that the veiling is difficult to determine empirically in earlier-type stars with few photospheric features, and so the importance of H I stellar photospheric absorption is difficult to assess in these stars (see Section 4.3).

6.2. Accretion in Transitional Disks

6.2.1. Equivalent Widths and the Identification of Transitional Disks

In Figure 11 we show $Pf\beta$ equivalent width (EW) versus $Pf\beta$ luminosity. There is no significant correlation between these two variables and thus $Pf\beta$ EW is not a good predictor of accretion luminosity. However, we note that several transitional disks have notably large $Pf\beta$ EW's (between -20 and -30 \AA). Since the $Pf\beta$ luminosities of these disks are not higher than average, the high EW's instead reflect the fact that these disks have reduced continuum flux levels in the near-IR. This suggests that $Pf\beta$ may be a good tracer for some accreting transitional disks, and a means to detect these disks with a single spectrum—one simple, robust measurement that requires no knowledge of the stellar parameters or absolute flux level. We arbitrarily label all disks with $Pf\beta$ EW < -15 and suggest that these targets may have depletions of dust in their inner regions. Aside from the transitional disks, which are known to have inner disk dust depletions, we also note that TY CrA is a tertiary system with an eclipsing binary (e.g., Vaz 2001) and that Haro 1-1 is an anomalously fast rotator (Shevchenko & Herbst 1998). To our knowledge, no unique properties have been discussed for the other disks with high $Pf\beta$ EW.

It is interesting to note that $Pf\beta$ EW's are not high for all transitional disks. In our sample, LkH α 330, HD 135344 B, IRS 48, and DoAr 44 have typical or even slightly low $Pf\beta$ EW's. LkH α 330 and HD 135344 B appear to have relatively low gas/dust ratios as compared to other transitional disks (Salyk et al. 2009). DoAr 44 might better be considered a

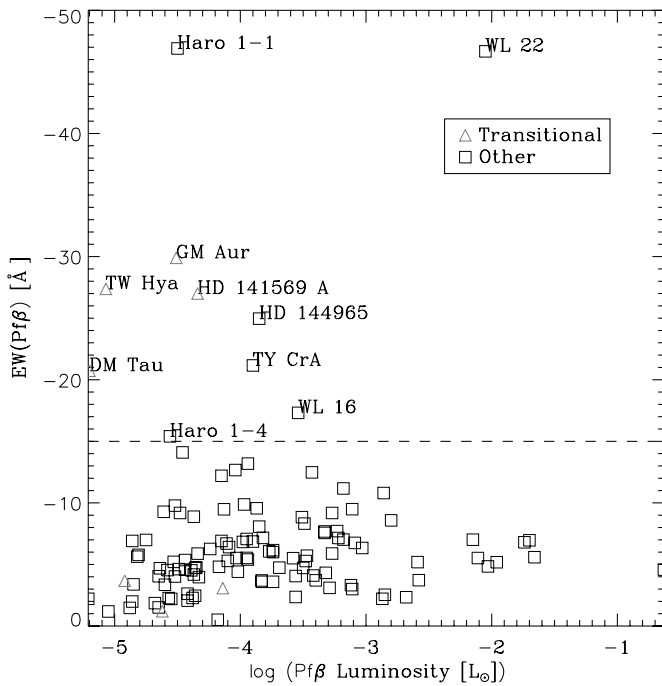


Figure 11. Pf β equivalent width (EW) plotted against Pf β line luminosity. Note that EW is not a good predictor for Pf β luminosity (or therefore accretion luminosity). The dashed line marks an EW of -15 \AA ; several transitional disks have extreme values of Pf β EW.

pre-transitional disk (Andrews et al. 2011)—a disk with an inner clearing but relatively high near-IR flux, that may be consistent with a small gap rather than a large clearing (Espaillat et al. 2007). And IRS 48 has a ~ 30 AU hole in its gas distribution (Brown et al. 2012). Therefore, high Pf β EW may be an indicator of only those transitional disks that have significantly cleared their inner regions of small dust grains but not gas. Thus, a measurement of Pf β EW might shed light on the process causing the inner disk depletion for any given disk, as different clearing scenarios predict different gas/dust ratios.

6.2.2. Transitional Disks in Comparison to Other Disks

Figure 12 shows the log of Pf β luminosity as a function of the log of stellar mass. Accretion luminosity and accretion rate are known to scale with stellar mass, and these results are no exception. Plotting Pf β luminosity as a function of stellar mass allows us to investigate potential outliers with high or low accretion rates.

Transitional disks have on average lower Pf β luminosities, at a marginally statistically significant level, with a mean log residual Pf β luminosity of -0.7 ± 0.3 . A low Pf β luminosity could represent a low accretion luminosity, or could result if the transitional disks have normal accretion luminosities but an anomalously low Pf β luminosity/accretion luminosity ratio. However, we have investigated the Pf β luminosity/accretion luminosity ratio for these targets and find that four of five transitional disks (DM Tau, GM Aur, HD 141569 A, and TW Hya) with literature accretion luminosities have Pf β luminosities even higher than would be predicted by Equation (1) (one, HD 135344 B, has a slightly but not significantly lower Pf β luminosity). Therefore, this study is in closer agreement with Najita et al. (2007) and Espaillat et al. (2012) who find slightly lower than average accretion rates for transitional disks, than with contrasting studies showing no difference between the two types of disks (Fang et al.

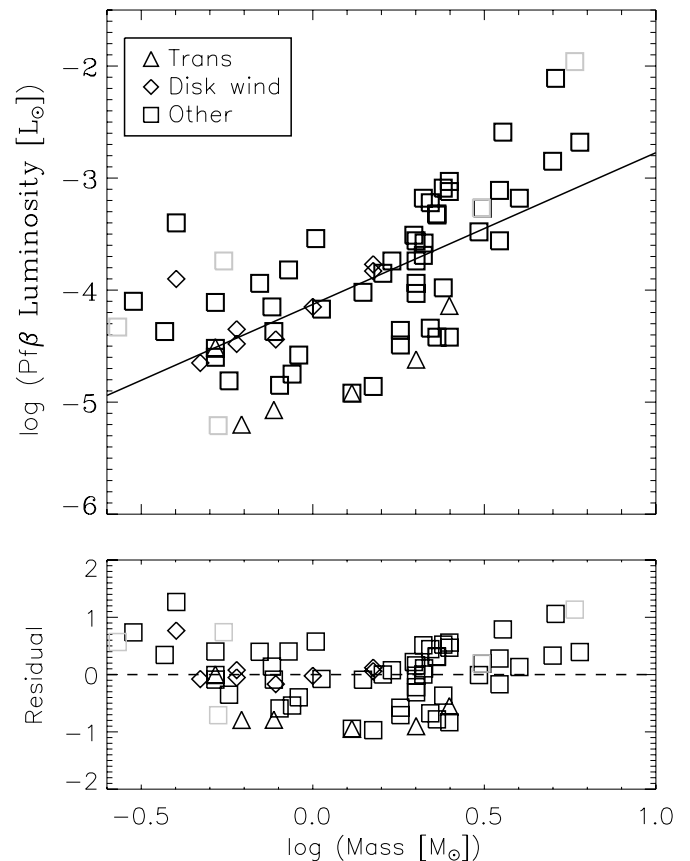


Figure 12. The top panel shows the log of Pf β line luminosity plotted against the log of stellar mass and best-fit correlation. The lower panel shows the fit residuals. Triangles and diamonds are transitional and disk wind targets, respectively; squares are other targets. Gray squares represent highly inclined disks.

2009; Sicilia-Aguilar et al. 2010). As suggested by Sicilia-Aguilar et al. (2010), the contrasting study results may be due to the different physical nature of transitional disks in different regions. However, it should be cautioned that the high Pf β EW's for transitional disks could produce an observational bias in our work, making it easier to detect a low accretion rate around a transitional, rather than classical, disk. Thus, our sample may be missing a number of classical disks with low accretion rates.

Since our sample derives from a large number of clusters of different ages, it is also worth asking whether we might measure spuriously low Pf β luminosities for transitional disks because they are derived from older clusters, and there is evidence that accretion rates decrease with stellar age (Hartmann et al. 1998). However, we find that the transitional disks are derived from a large number of the clusters sampled in this work, and that regardless of the cluster under consideration, transitional disks lie at the low end of the range of measured Pf β luminosities. In addition, the transitional disks are derived primarily from clusters with ages $\lesssim 3$ Myr (with the exception of TW Hya), and so are not biased toward older ages.

6.3. Accretion in Disk Wind Sources

Disk wind targets are a newly identified subset of disks defined by their single-peaked near-IR emission line flux and spectro-astrometric profiles, which may be evidence for the presence of a slow-velocity disk wind (Pontoppidan et al. 2011a; Bast et al. 2011). By comparing average accretion luminosities for disk wind and normal disk targets, Bast et al. (2011)

noted that disk wind targets may have higher than average accretion luminosities. However, Figure 12 shows that the $Pf\beta$ luminosities for these targets are actually typical for sources in our sample, implying no higher accretion luminosities. These targets also have typical ratios of $Pf\beta$ luminosity to derived accretion luminosity. And, the disk wind sources derive from a large number of clusters, implying no obvious age-related biases.

How can these two results be reconciled? We believe there are a few possible explanations. First, the disk wind sources in Bast et al. (2011) have somewhat higher stellar masses than their comparative sample of “normal” disks. Second, the disks in the comparative sample in Bast et al. (2011) have lower accretion luminosities at a given mass than disks in our sample. Since our full sample has not been “vetted” for disk wind targets, it is possible that our full sample includes unidentified disk wind targets that increase the average, or our sample may simply be biased toward higher accretion rates. Unfortunately, we cannot use the criteria in Bast et al. (2011) to classify all of our targets because of the relatively lower spectral resolution of NIRSPEC as compared to CRILES.

The accretion rates for the disk wind targets, derived from this work, range from 6×10^{-9} to $2 \times 10^{-7} M_{\odot} \text{ yr}^{-1}$, with most accretion rates in the range of 10^{-8} to $10^{-7} M_{\odot} \text{ yr}^{-1}$. Thus, the disk wind sources are likely capable of supporting mass-loss rates of 10^{-10} to $10^{-9} M_{\odot} \text{ yr}^{-1}$, and in some cases $10^{-8} M_{\odot} \text{ yr}^{-1}$, as suggested by Pontoppidan et al. (2011a).

Although we find typical accretion rates for disk wind sources, it should be noted that one of the primary characteristics of disk wind targets is their high veiling at optical through IR wavelengths (e.g., Gahm et al. 2008). The seemingly contradictory observation of high veiling but normal accretion rates may be reconciled by the results of Gahm et al. (2008), who find that the optical veiling in such targets is not a good measure of disk accretion, as the derived optical veiling is affected by infilling of photospheric lines from line emission.

6.4. Accretion Rates for Embedded Disks

6.4.1. Classification of Evolutionary Stage

The process of star and planet formation is often deconstructed into four discrete evolutionary stages. The targets in our sample represent primarily, if not entirely, sources in stages I and II, where stage I sources are embedded protostars with still-collapsing envelopes, while stage II sources are revealed stars with circumstellar disks and little or no remaining envelope. For this work, the evolutionary class was determined primarily through examination of the M -band spectra themselves, with stage I disks defined as those showing CO ice absorption, and stage II disks those without. This offers the significant advantage that our entire sample can be easily classified using our own data set, with a single consistent criterion.

However, it is a known issue that edge-on stage II disks can masquerade as evolutionarily younger systems, especially with the use of spectral slope-based classification schemes (see, for example, the detailed discussion in Evans et al. 2009). Our classification scheme is not immune to this issue, as CO ices in edge-on disks can in principle produce absorption features. However, in practice, the production of strong ice absorption features requires a very specific geometry, with an inclination close to the opening angle of the flaring disk. Additionally, the upper layers of the disk that intercept the starlight are often too warm to contain significant quantities of CO ice

(Pontoppidan et al. 2005). Thus, it is perhaps not surprising that our classification scheme correctly identifies some high-inclination disks, including T Tau S and VV Ser, as stage II rather than stage I sources.

Nevertheless, we have made an effort to confirm stage I classifications whenever possible using additional diagnostics in the literature, including the use of spectral indices, bolometric temperatures, the presence of infall or outflow signatures in spectral lines, and the presence of extended envelopes as seen in millimeter-wave maps (e.g., Boogert et al. 2002; White & Hillenbrand 2004; Doppmann et al. 2005; Evans et al. 2009; van Kempen et al. 2009; Jørgensen et al. 2009; Herczeg et al. 2011). Our classifications are listed in Table 3, along with accompanying notes for sources with ambiguous properties.

6.4.2. Accretion Rates for Stage I Disks, and Comparison with Stage II Disks

Since stars are believed to accumulate most of their mass during their embedded stages, accretion rates for stage I disks were initially predicted to be significantly larger than for stage II disks. Smooth collapse models (Shu 1977) predicted stage I accretion rates similar to simple estimates—for example, accretion rates of order $\gtrsim 10^{-6} M_{\odot} \text{ yr}^{-1}$ are required to build a solar mass star over a $\sim 10^6$ year timescale. Accretion rates in stage II disks around solar-mass stars are typically near $10^{-8} M_{\odot} \text{ yr}^{-1}$ (see Table 5 and references therein), and thus stage I accretion rates would be predicted to be at least two orders of magnitude higher.

While there is at least some evidence for greater accretion activity in stage I sources (e.g., Doppmann et al. 2005), measured accretion rates in stage I disks around low-mass stars are found to be near $10^{-8} M_{\odot} \text{ yr}^{-1}$ (e.g., White & Hillenbrand 2004; White et al. 2007)—two orders of magnitude lower than predicted. This discrepancy is related to the well-known “luminosity problem” in star formation (Kenyon et al. 1990), wherein measured bolometric luminosities (which derive primarily from the energy of gravitational contraction, and are therefore related to mass accretion) are lower than predicted from steady growth models. Instead, evidence is growing (e.g., Dunham et al. 2008; Enoch et al. 2009; Evans et al. 2009) that accretion may need to be episodic—primarily quiescent, with brief phases of intense accretion.

While direct measurements of accretion rates in stage I disks appear to support this hypothesis, such studies inevitably suffer from a bias toward more revealed disks, as the most embedded protostars are heavily extinguished, and cannot be easily studied with either optical or near-IR accretion tracers. Although it is difficult to quantify the many sample-selection biases here, extinction remains small in the M band even for high A_V , and our sample includes several heavily embedded protostars. As a point of comparison, the sample of young protostars in White & Hillenbrand (2004) all have $A_V < 20$, while 10 of 20 stage I disks in our study have $A_V \geq 20$. Thus, our study provides insight about accretion rates for the most embedded, and perhaps evolutionarily youngest, disks.

Figure 13 shows the empirical cumulative distributions of calculated $Pf\beta$ luminosities for stage I and stage II disks in our sample. There are small differences between the distributions of $Pf\beta$ line luminosities for the two samples, with slightly higher $Pf\beta$ luminosities in Stage II disks. Stage I disks have a mean $\log(L_{Pf\beta})$ of -3.7 and a standard deviation of 0.8, while stage II disks have a mean $\log(L_{Pf\beta}) - 3.8$ with a standard deviation of 0.9. However, a two-sample Kolmogorov–Smirnov

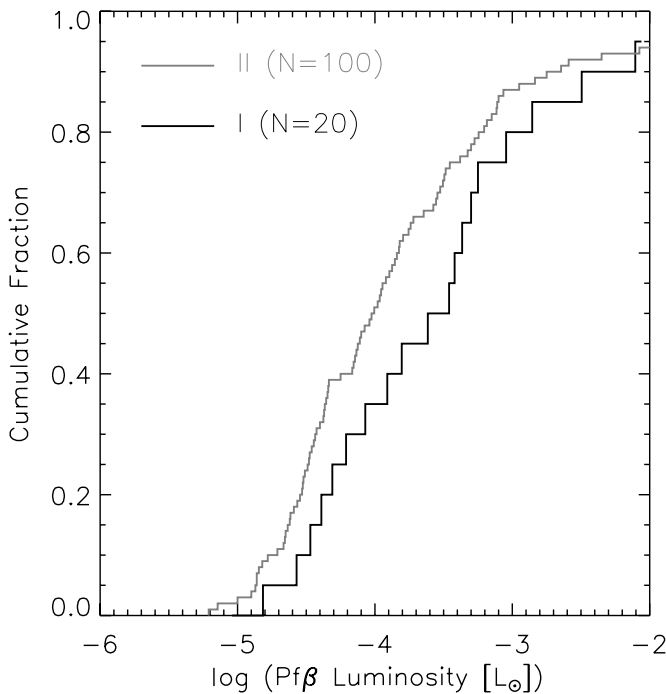


Figure 13. The cumulative distribution function of observed Pf β accretion luminosities for stage I and stage II disks.

test produces an associated probability of 30%. Therefore, the hypothesis that the stage I and stage II disks draw from the same distribution of accretion luminosities cannot be rejected, i.e., differences between the two distributions are not statistically significant.

Unfortunately, accretion *rate* distributions cannot be compared directly, as the majority of the embedded disks in our sample do not have measured stellar masses or radii due to the relative difficulty of obtaining optical spectra for highly embedded stars. However, if average stellar masses and radii are similar for the two samples our results suggest that accretion rates are quite similar for stage I and stage II disks. Spectroscopic analyses of stage I protostars have suggested that stage I and stage II stars have similar average masses, while stellar radii are no more than a factor of ~ 2 larger for stage I stars (White & Hillenbrand 2004). For the small number (four) of stage I disks with measured stellar masses in our sample, the average stellar mass is $1.7 M_{\odot}$, as compared to an average of $1.8 M_{\odot}$ for our stage II sample. Assuming at most a factor of two difference in radius between stage I and II stars, and assuming no significantly different selection biases for our two subsamples, our results suggest that stage I accretion rates are at most factors of a few higher, on average, than they are for stage II disks.

We can directly measure the accretion rates of four stage I targets in our sample. These accretion rates range from 1×10^{-8} to $4 \times 10^{-7} M_{\odot} \text{ yr}^{-1}$. Thus, these targets have an accretion rate spread of more than an order of magnitude, spanning from a value typical for low-mass stars to something close to the $10^{-6} M_{\odot} \text{ yr}^{-1}$ required to build a solar mass star in 1 Myr. If we divide the currently measured stellar mass by \dot{M} , we can obtain a lower limit to the time required to form the star assuming a steady state accretion rate. We find timescales of 0.7, 6, 54, and 67 Myr for the four stage I sources—longer than estimated stage I lifetimes (Evans et al. 2009). Therefore, our data suggest that these sources must at times have had significantly higher accretion rates than currently observed.

Making some simple assumptions, we can also investigate accretion rates for the full stage I sample. For example, assuming $M_{\star} = 1 M_{\odot}$ and $R_{\star} = 2 R_{\odot}$ for the 16 stage I stars without measured stellar parameters, we find a mean and standard deviation of $\log(\dot{M}) = -7.1 \pm 0.7$. This value is at least an order of magnitude lower than the predictions from steady accretion models. Yet, with these assumptions, our sample does include two sources with accretion rates of $\sim 2 \times 10^{-6} M_{\odot} \text{ yr}^{-1}$. If our sample is representative of *all* stage I sources and homogeneous except for random changes in accretion rate due to episodic accretion, and taking into account a detection fraction of 60% (see Section 3), this would imply that the stage I sources spend $\sim 6\%$ of their lifetime with $\dot{M} > 10^{-6} M_{\odot} \text{ yr}^{-1}$. The large number (40%) of Pf β non-detections is further evidence for a wide range of accretion rates in stage I sources, again consistent with an episodic accretion scenario.

It is also interesting to note that no FU Orionis stars in our sample show detectable Pf β emission lines. The explanation for this is not known, but may have to do with the restructuring of the inner disk that occurs during high accretion-rate events (e.g., Zhu et al. 2009). In any case, while these sources may represent protostars with the highest rates of accretion (Hartmann & Kenyon 1996), they are necessarily excluded from this analysis due to their lack of Pf β emission, and thereby bias the remaining sample toward lower accretion rates.

7. CONCLUSIONS

In this work, we have introduced the use of H I Pf β as a mass accretion tracer for young stars with protoplanetary disks. H I Pf β offers several advantages over other accretion tracers, including being readily observed in heavily extinguished disks, requiring minimal correction for photospheric absorption and being commonly observed along with CO fundamental transitions. Using existing measurements of accretion from the literature, we derive a relationship between Pf β line luminosity and accretion luminosity, and show that this relationship can reproduce measured accretion rates with an accuracy similar to that of other commonly used tracers. Examining our large sample of accretion rates, we are further able to show that “disk wind” sources appear to have normal accretion rates, while our sample of transitional disks have slightly lower than average accretion rates.

We also examine the accretion rates of the stage I disks in our sample—a sample that includes significantly more embedded targets than in previous studies of this type. We find that stage I and stage II disks have statistically indistinguishable Pf β luminosities, implying similar accretion rates, and that the accretion rates of stage I disks are too low to build a stellar mass with quiescent accretion. Our results instead are consistent with both observational and theoretical evidence that stage I objects experience episodic, rather than quiescent, accretion.

The calibration of Pf β presented here has allowed for a coherent comparison of accretion luminosity in 120 stars. While this work has focused exclusively on spectra from the NIRSPEC and CRIRES surveys with which we are associated, the correlation we derive can be extended to a much larger number of protoplanetary disk spectra, including some which may already reside in the NIRSPEC or CRIRES archive. We have also focused exclusively on disks with clearly detectable Pf β emission. However, with more extensive observations and focused analysis, our correlation can be extended to tenuous disks with low accretion rates—disks of utmost importance to our understanding of disk dissipation. Finally, as Pf β is

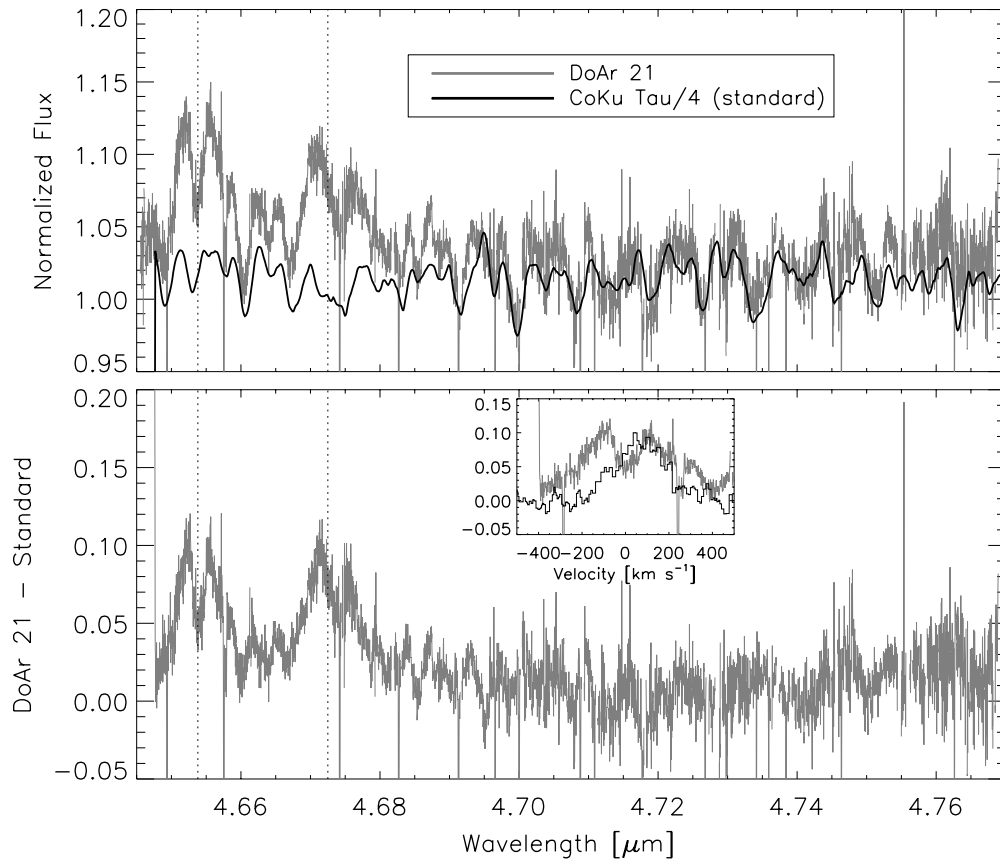


Figure 14. Top: spectrum of DoAr 21 (K1) observed with CRIRES (gray) and a rotationally broadened, veiled spectrum of CoKu Tau/4 (K3) observed with NIRSPEC (black). Dotted lines mark the locations of Pf β and H ϵ . Bottom: difference spectrum. Inset: difference spectrum centered at 4.6538 μm (gray) compared with observations of H α from Jensen et al. (2009).

observed contemporaneously with CO fundamental emission lines, the observation of Pf β can be used to account for the effects of accretion variability in searches for planet-induced CO line variability. We therefore expect the utility of Pf β as an accretion tracer to extend well beyond the work presented here.

C.S. thanks Casey Deen and Chris Sneden for providing an introduction to MOOG. C.S. also acknowledges the financial support of the NOAO Leo Goldberg Fellowship program and the University of Texas at Austin McDonald Observatory Harlan J. Smith Postdoctoral Fellowship. E.v.D. is supported by EU A-ERC grant 291141. This publication makes use of data products from the Wide-field Infrared Survey Explorer, which is a joint project of the University of California, Los Angeles, and the Jet Propulsion Laboratory/California Institute of Technology, funded by the National Aeronautics and Space Administration. Some of the data presented herein were obtained at the W. M. Keck Observatory, which is operated as a scientific partnership among the California Institute of Technology, the University of California and the National Aeronautics and Space Administration. The Observatory was made possible by the generous financial support of the W. M. Keck Foundation.

APPENDIX

THE SPECIAL CASE OF DoAr 21

DoAr 21 is known to be an unusual disk, with little or no reported excess emission below 7 μm and weak or non-existent accretion, but asymmetric extended emission beyond ~ 100 AU, and line emission from polycyclic aromatic hydrocar-

bons (PAHs) and H $_2$ (Jensen et al. 2009 and references therein). Observations of DoAr 21 with the Very Long Baseline Array have revealed that DoAr 21 is a 5 mas (0.6 AU) binary (Loinard et al. 2008). Detailed studies of evolved sources like DoAr 21 can provide insight into the process of disk dissipation, as it may represent a case of a disk at the latest stages of dissipation.

DoAr 21 is the only disk in our sample that shows both low but detectable Pf β equivalent width, and low veiling, meaning that the underlying photosphere must be accounted for in the calculation of Pf β line flux. In Figure 14, we show the CRIRES spectrum of DoAr 21 (K1) along with a NIRSPEC spectrum of a photospheric template star—CoKu Tau/4 (K3)—rotationally broadened and veiled to match the observed spectrum of DoAr 21 longward of 4.68 μm . Although DoAr 21 and CoKu Tau/4 do not have the same spectral sub-type, we find only small differences in spectra across a range of observed spectral types (G2–M0), and both are young stars with similarly lowered surface gravities compared to main-sequence stars.

The bottom of Figure 14 shows the difference between the observed and template spectra, which reveals strong and double-peaked Pf β and H ϵ emission lines. As an inset in Figure 14, we compare Pf β to an observed spectrum of H α from Jensen et al. (2009). The H α line shown here is in fact the difference between H α lines observed on two different nights, and therefore represents the variable, perhaps flare-related, component to H α . We find that the Pf β emission line is broader and more double-peaked than the H α line, suggesting that the Pf β may not have a flare origin, but instead have an accretion origin. (It should be noted, however, that the H ϵ line flux appears unusually high compared to other disks in our sample.)

We also note that the best fit to the observed DoAr 21 spectrum requires a veiling of ~ 2.6 , and that the spectrum is inconsistent with zero veiling, consistent with the results of Salyk et al. (2009). Assuming the continuum flux is emitted entirely by grains at 1500 K with an opacity of $500 \text{ cm}^2 \text{ g}^{-1}$ (appropriate for sub-micron sized grains), the mass in small grains is $\sim 10^{19} \text{ kg}$, or $\sim 10^{-4} M_{\odot}$. This is almost certainly a lower limit, as there is likely a range of grain sizes and a range of grain temperatures, with 1500 K being the approximate maximum temperature for solid silicates. An alternative way to explain finite veiling would be for the stellar disk to contain hot spots with featureless emission spectra. However, even in an extreme case of 10^4 K hot spots, 54% of the stellar disk would need to be covered by hot spots to explain the observed veiling. In contrast, results from photometric variability studies (Bouvier et al. 1995) find hot spot temperatures of no more than a few hundred to a few thousand K hotter than the stellar temperatures, and 0.5%–40% disk covering fractions. Thus, DoAr 21 appears to retain both a gas and dust disk at small radii.

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