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COMPLEX DIFFUSE RADIO EMISSION IN THE MERGING PLANCK ESZ CLUSTER A3411

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ABSTRACT

We present Very Large Array (VLA) radio and *Chandra* X-ray observations of the merging galaxy cluster A3411. For the cluster, we find an overall temperature of $6.4^{+0.6}_{-1.0}$ keV and an X-ray luminosity of $2.8 \pm 0.1 \times 10^{44}$ erg s⁻¹ between 0.5 and 2.0 keV. The *Chandra* observation reveals the cluster to be undergoing a merger event. The VLA observations show the presence of large-scale diffuse emission in the central region of the cluster, which we classify as a 0.9 Mpc size radio halo. In addition, a complex region of diffuse, polarized emission is found in the southeastern outskirts of the cluster along the projected merger axis of the system. We classify this region of diffuse emission as a radio relic. The total extent of this radio relic is 1.9 Mpc. For the combined emission in the cluster region, we find a radio spectral index of -1.0 ± 0.1 between 74 MHz and 1.4 GHz. The morphology of the radio relic is peculiar, as the relic is broken up into five fragments. This suggests that the shock responsible for the relic has been broken up due to interaction with a large-scale galaxy filament connected to the cluster or other substructures in the intracluster medium. Alternatively, the complex morphology reflects the presence of electrons in fossil radio bubbles that are re-accelerated by a shock.

Key words: galaxies: clusters: individual (A3411) – galaxies: clusters: intracluster medium – large-scale structure of universe – radiation mechanisms: non-thermal – X-rays: galaxies: clusters

Online-only material: color figures

1. INTRODUCTION

In a number of clusters, large-scale radio emission is found in the form of radio halos and relics (e.g., Feretti et al. 2012 and references therein). This large-scale diffuse emission indicates the presence of relativistic particles with Lorentz factors $\gamma \sim 10^4$, and μ Gauss magnetic fields in the intracluster medium (ICM). Giant radio halos and relics, i.e., those having Mpc-sizes, are exclusively found in merging galaxy clusters (e.g., Cassano et al. 2010; van Weeren et al. 2011b). This supports the picture that a fraction of the energy released during the merger event is channeled into the (re-)acceleration of particles to highly relativistic energies. In the presence of a magnetic field, these particles would then emit faint synchrotron radiation at radio wavelengths.

Giant radio halos have ~ 1 –2 Mpc sizes, are centrally located in clusters, display smooth morphologies, and are usually unpolarized (e.g., Feretti et al. 2001; Bacchi et al. 2003). Creating these large halos requires that local particle acceleration occurs over a large volume in the cluster (Jaffe 1977). Although the basic observational properties of halos have now been established (e.g., Feretti et al. 2012), the formation mechanism of halos is still being debated (e.g., Brunetti et al. 2008; Donnert et al. 2010a, 2010b; Macario et al. 2010; Brown & Rudnick 2011; Enßlin et al. 2011; Brunetti et al. 2012; Zandanel et al. 2012; Arlen et al. 2012). Scenarios for their origin include “primary” models, in which an existing electron population is re-accelerated by turbulence or shocks caused by recent cluster mergers (Brunetti et al. 2001; Petrosian 2001), and “secondary” models, in which relativistic electrons are continuously injected

into the ICM by inelastic collisions between cosmic rays and thermal ions (e.g., Dennison 1980; Blasi & Colafrancesco 1999; Dolag & Enßlin 2000; Miniati et al. 2001; Keshet & Loeb 2010). Combinations of both acceleration mechanisms have also been considered (Brunetti & Blasi 2005; Dolag et al. 2008; Brunetti & Lazarian 2011).

Giant radio relics are located in cluster outskirts, are usually polarized, and are elongated (e.g., Röttgering et al. 1997; Bagchi et al. 2011). They sometimes occur in pairs, located diametrically on opposite sides of the cluster center, so-called “double relics.” The currently favored scenario is that large relics are the signatures of electrons (re-)accelerated at shocks due to cluster mergers (e.g., Enßlin et al. 1998; Miniati et al. 2001; Giacintucci et al. 2008; Bonafede et al. 2009; van Weeren et al. 2010; Kale & Dwarakanath 2010). This should happen via the diffusive shock acceleration mechanism (e.g., Krymskii 1977; Axford et al. 1977; Bell 1978a, 1978b; Blandford & Ostriker 1978; Drury 1983; Blandford & Eichler 1987; Jones & Ellison 1991; Malkov & O’C Drury 2001; Kang & Ryu 2013). Cluster merger shocks typically have low Mach numbers ($\mathcal{M} \sim 1$ –3; e.g., Markevitch et al. 2002; Russell et al. 2010; Macario et al. 2011; Akamatsu & Kawahara 2013; Ogrea et al. 2013), and the efficiency with which such shocks can accelerate particles is unknown. Therefore, re-acceleration of pre-accelerated electrons in the ICM might be required to explain the observed brightness of relics, since re-acceleration is a more efficient mechanism for weak shocks (e.g., Markevitch et al. 2005; Giacintucci et al. 2008; Kang & Ryu 2011; Kang et al. 2012; Pinzke et al. 2013).

Only a few dozen clusters with diffuse radio emission are known today (Venturi et al. 2007, 2008; Giovannini et al. 2009; van Weeren et al. 2011b; Nuza et al. 2012; Bonafede et al. 2012).

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Table 1
VLA Observations

	VLA BnA	VLA DnC
Frequency (VLA band)	1365, 1435 MHz	1365, 1435 MHz
Bandwidth	2×50 MHz	2×50 MHz
Observation dates	2003 Oct 7	2004 Jun 14
Total on-source time (hr)	4.2	4.6
Beam size	$7''.0 \times 3''.9$	$48'' \times 33''$
rms noise ^a (σ_{rms})	$15 \mu\text{Jy beam}^{-1}$	$47 \mu\text{Jy beam}^{-1}$
Robust (briggs weighting)	-0.5	1.0

Note. ^a The rms noise is measured at the center of the images, between the radio halo and relic (see Figures 1 and 2).

This makes it difficult to study the halo and relic occurrence as a function of cluster properties, such as the dynamical state, in detail.

To identify both a larger number of clusters with radio halos and relics and a sample with more extreme clusters in terms of their dynamics and structural morphologies, we compared the radio emission as mapped by the NVSS (Condon et al. 1998) and Very Large Array (VLA) with *Chandra* X-ray images of the Planck ESZ clusters (Planck Collaboration et al. 2011a). Sunyaev–Zel’dovich surveys preferentially select massive clusters, including those in the early stages of formation (Planck Collaboration et al. 2011b, 2012).

Here we report on VLA, *Chandra* and Very Large Telescope (VLT) observations of the poorly studied cluster Abell 3411 (hereafter A3411) located at $z = 0.1687$ (Ebeling et al. 2002). It is listed in the Planck ESZ catalog as PLCKESZ G241.97+14.85 (Planck Collaboration et al. 2011a). Observations and data reduction are described in Section 2. The results are presented in Section 3, and this is followed by a discussion and conclusions in Sections 4 and 5.

Throughout this paper, we assume a Λ CDM cosmology with $H_0 = 71 \text{ km s}^{-1} \text{ Mpc}^{-1}$, $\Omega_m = 0.27$, and $\Omega_\Lambda = 0.73$. With the adopted cosmology, $1''$ corresponds to a physical scale of 2.85 kpc at $z = 0.1687$. All images are in the J2000 coordinate system.

2. OBSERVATIONS AND DATA REDUCTION

A3411 was observed with the VLA on 2003 October 7 in BnA array and on 2004 June 14 in DnC array (project AC0696). An overview of the observations and the resulting images is given in Table 1. The 1.4 GHz observations were carried out in single channel continuum mode with two intermediate frequencies with a bandwidth of 50 MHz. All polarization products (RR, RL, LR, and LL) were recorded.

The data were calibrated with the NRAO Astronomical Image Processing System package. The data were visually inspected for the presence of radio frequency interference, which was subsequently flagged. We then determined the gains for the calibrator sources. The fluxes for the primary calibrator sources (3C147 and 3C286) were set by the Perley & Taylor (1999) extension to the Baars et al. (1977) scale. The effective feed polarization parameters (the leakage terms or D-terms) were found by observing the phase calibrator (0902-142) over a range of parallactic angles and simultaneously solving for the unknown polarization properties of the source. The polarization angles were set using the polarized source 3C 286. For the R–L phase difference of 3C 286, we assumed a value of -66° . We transferred the calibrator solutions to the target and carried out several rounds of self-calibration to refine the calibration. Images were made with manually placed clean boxes and the

“briggs” weighting scheme (Briggs 1995; see Table 1), unless explicitly stated. All final images were corrected for the primary beam attenuation, unless noted otherwise.

The cluster was also observed using the *Chandra X-Ray Observatory* for 10 ks using the ACIS-I array. The data (Chandra ObsId 13378) were calibrated by applying the most recent calibration files to the level 1 events file. The events file was calibrated following the processing described in Vikhlinin et al. (2005). This included correction for position dependent charge transfer inefficiency, the application of gain and time-dependent gain maps to calibrate photon energies, and removal of counts from bad pixels and counts with a recomputed ASCA grade 1, 5, or 7. Periods of high local background were removed from the observation by examining the light curve and removing periods with a flux in the 5–10 keV band 1.2 times above the mean. This left a total exposure of 7.0 ks. For the image in Figure 1 (left panel), we extracted counts in the 0.5–3.0 keV band to maximize the ratio between source and background counts in the image. Pixels in the image were combined with a binning factor of 16.

VLT FORS1 images of A3411 were obtained from the ESO archive. These images were taken on 2003 December 31 with the V_{Bessel} , R_{Bessel} , and I_{Bessel} filters and only cover the central region of the cluster. In total, three images with an exposure time of 110 s each were available per filter. These images were bias-corrected, flat-fielded, and stacked using IRAF (Tody 1986, 1993).

3. RESULTS

The *Chandra* ACIS observations of A3411 show a merging system, with the projected merger axis oriented SE-NW (see Figure 1). The cluster has a cometary shape with a well-defined subcluster core visible in the northwestern part of the system. Fainter X-ray emission is found surrounding the subcluster core and this emission seems to be part of a second, larger subcluster that likely extends beyond the northern boundary of the *Chandra* image. There is no clear surface brightness peak corresponding to the primary cluster core, which suggests the primary cluster has been disrupted by the collision with the subcluster, as has been the case for A2146 (Russell et al. 2012). The ICM temperature was fit using XSPEC with an absorbed single-temperature APEC model, which describes emission from collisionally-ionized diffuse plasma. The hydrogen column density was fixed using Dickey & Lockman (1990) at $5.39 \times 10^{20} \text{ cm}^{-2}$. The temperatures and metallicities were left free in the fit. From this we found a global temperature of $6.4^{+0.6}_{-1.0}$ keV. Since R_{500} partially extends beyond the *Chandra* field of view (FOV) we calculated the R_{500} -luminosity using a surface brightness profile. The (unabsorbed) X-ray luminosity is found to be $2.8 \pm 0.1 \times 10^{44} \text{ erg s}^{-1}$ between 0.5 and 2.0 keV within $R_{500} = 1.34 \text{ Mpc}$. This would correspond to a mass of about $M_{500} = 9 \times 10^{14} M_\odot$ (Vikhlinin et al. 2009). An overview of the cluster properties is given in Table 2.

The VLT FORS1 image is shown in Figure 1, right panel. A cD galaxy is visible just southeast of the X-ray peak and at about $10''$ south of a foreground star (the cD galaxy is marked in Figure 1, right panel). This galaxy, 2MASX J08415287-1728046, has a measured redshift of 0.16324 (Jones et al. 2009), consistent with the cluster redshift. The image displays a relatively large number of foreground stars as A3411 is located at a galactic latitude of $b \approx 15^\circ$.

The 1.4 GHz VLA continuum images are shown in Figures 2 and 3. A low-resolution image with the emission from compact

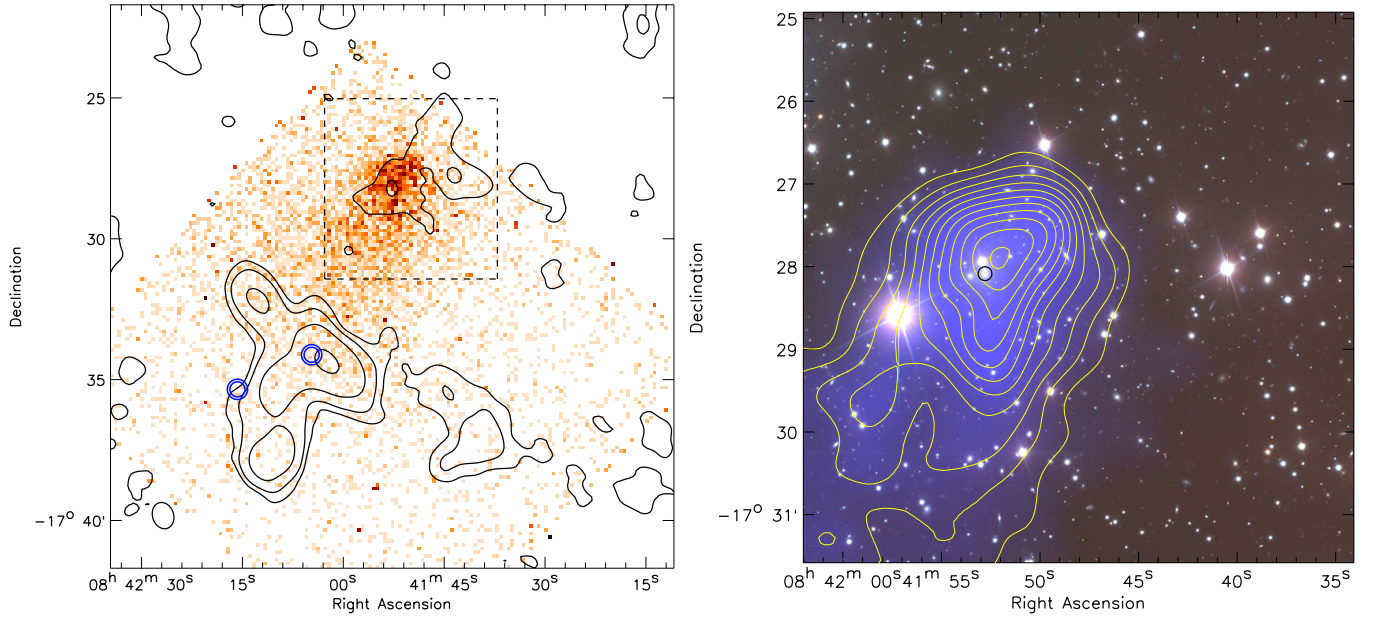


Figure 1. Left: *Chandra* ACIS 0.5–3.0 keV image of A3411. Radio contours are from the VLA 1.4 GHz DnC array image with natural weighting and compact sources subtracted. The resolution of the radio image is $57'' \times 45''$. Contour levels are drawn at $[1, 3, 9, 27, \dots] \times 3\sigma_{\text{rms}}$, with $\sigma_{\text{rms}} = 55 \mu\text{Jy beam}^{-1}$. The locations of two compact sources, labeled in Figures 2 and 3, are indicated with blue circles. The dashed box indicates the region shown in the right panel. Right: VLT FORS1 color image (using V_{Bessel} , R_{Bessel} , and I_{Bessel} filters) of the central part of the cluster. The *Chandra* ACIS 0.5–3.0 keV adaptively smoothed image is overlaid in blue, yellow contours are linearly spaced for the X-ray surface brightness. The cD galaxy 2MASX J08415287-1728046 is marked with a black circle.

(A color version of this figure is available in the online journal.)

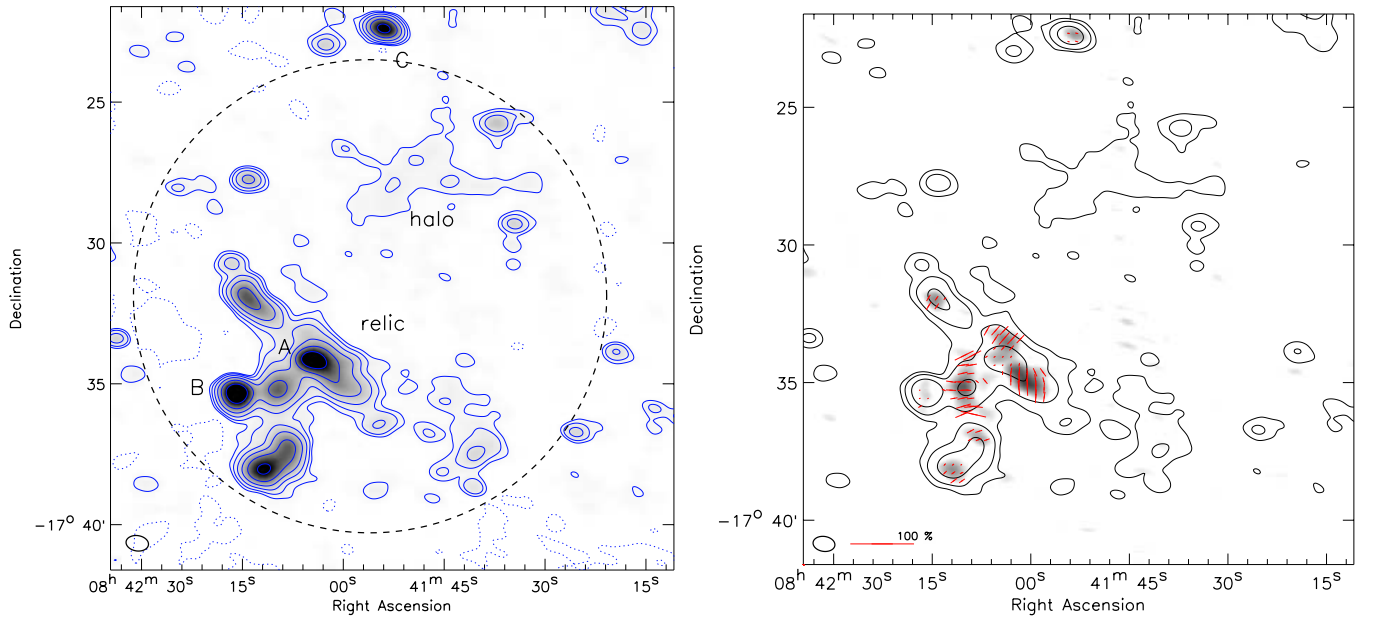


Figure 2. Left: VLA 1.4 GHz DnC array image. Contour levels are drawn at $[1, 2, 4, 8, \dots] \times 3\sigma_{\text{rms}}$ ($\sigma_{\text{rms}} = 47 \mu\text{Jy beam}^{-1}$). Negative $-3\sigma_{\text{rms}}$ contours are shown by the dotted lines. The dashed circle indicates the area where the flux densities were extracted for Figure 5. Right: VLA DnC array polarization map. The image is not corrected for the primary beam attenuation as the polarization vectors and polarization fraction are independent of the beam correction. Total polarized intensity is shown as a grayscale image. Red vectors depict the polarization E-vectors, their length represents the polarization fraction. A reference vector for a polarization fraction of 100% is shown in the bottom left corner. No vectors were drawn for pixels with a signal-to-noise ratio < 4 in the total polarized intensity image. Contour levels are drawn at $[1, 4, 16, 64, \dots] \times 3\sigma_{\text{rms}}$ ($\sigma_{\text{rms}} = 47 \mu\text{Jy beam}^{-1}$) and are from the Stokes I image.

(A color version of this figure is available in the online journal.)

sources subtracted is shown as an overlay on the *Chandra* image in Figure 1. For this, we first created an image of the compact sources using uniform weighing and an inner uv-range cut of $2 \text{ k}\lambda$. We then subtracted the clean components of this image from the uv-data and remade the radio image using natural weighting.

Two relatively bright compact sources are visible to the SE of the cluster center in the higher resolution BnA array image. These sources are labeled A and B in Figures 2 and 3. We measure flux densities of 9.5 ± 0.5 and 16.1 ± 0.8 in the BnA array image for source A and B, respectively. Both sources have optical counterparts in the Digital Sky Survey

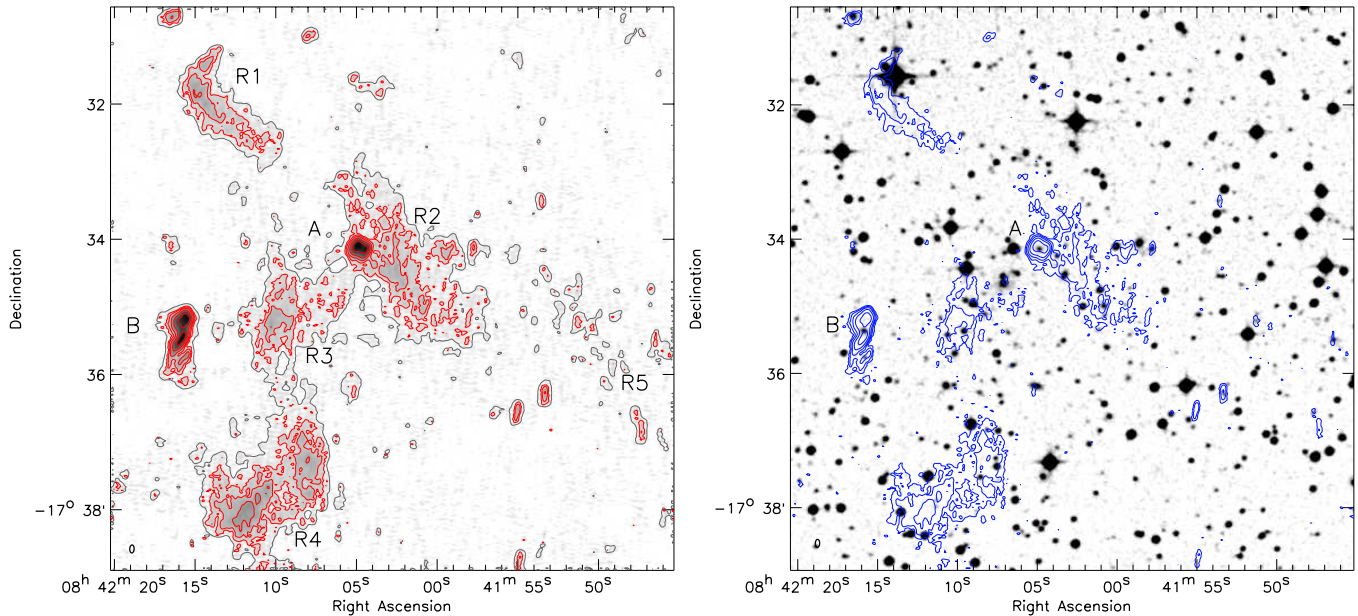


Figure 3. Left: VLA 1.4 GHz BnA array image of the relic. The red contour levels are drawn at $[1, 2, 4, 8, \dots] \times 4\sigma_{\text{rms}}$ ($\sigma_{\text{rms}} = 15 \mu\text{Jy beam}^{-1}$). The gray contours are drawn at $1.5\sigma_{\text{rms}}$ using a boxcar average of $7''$. The beam size is shown at the bottom left corner of the images. Right: optical DSS2 red image around sources *A* and *B* overlaid with radio contours from the VLA BnA array image. Contours levels are drawn at $[1, 2, 4, 8, \dots] \times 4\sigma_{\text{rms}}$. (A color version of this figure is available in the online journal.)

Table 2
Cluster Properties

Redshift	0.1687
$L_{X,R500}$ (erg s $^{-1}$, 0.5–2.0 keV a)	$2.8 \pm 0.1 \times 10^{44}$
Temperature (keV)	$6.4^{+0.6}_{-1.0}$
LLS relic, halo (Mpc)	1.9, 0.9
$P_{1.4\text{GHz}}$ relic, halo (W Hz $^{-1}$)	5.0×10^{24} , 4.6×10^{23}
$R_{\text{projected}}$ relic (Mpc)	1.3

Note. ^a Ebeling et al. (2002) reported $L_X = 3.8 \times 10^{44}$ from *ROSAT* in the 0.1–2.4 keV band (corrected for our adopted Λ CDM cosmology).

(DSS) image (see Figure 3 right panel for an overlay of radio contours on a DSS image). No redshifts are available for these counterparts. With the Two Micron All Sky Survey *Ks*-band magnitudes ($\text{mag}_{Ks} = 15.1$ for both *A* and *B*; Skrutskie et al. 2006), we estimate redshifts of ~ 0.3 for both galaxies using the *K*-*z* relation for massive elliptical galaxies from Willott et al. (2003). The uncertainty in this estimate is about 0.1, so either these galaxies are cluster members or they are located behind the cluster. Source *A* is resolved into two components that are separated by about $8''$. The morphology of source *B* is more complex and it has a largest angular extent of $53''$. At the distance of A3411, $53''$ corresponds to a physical extent of about 150 kpc. Merging galaxy clusters often host disturbed and/or tailed radio sources (e.g., Bridle et al. 1979), so it seems likely that at least source *B* is associated with A3411.

The only other galaxy with a published redshift in the A3411 field, besides the central cD galaxy, is 2MASX J08415393-1722206. This galaxy is associated with the bright compact radio source *C* at the far north in Figure 2 (left panel) and is likely a cluster member with $z = 0.162569$ (Jones et al. 2009). We measure a flux density of 9.1 ± 0.5 mJy for the source in the BnA array image.

The radio images display a region of complex radio emission to the SE of the cluster, along the major axis of the extended

X-ray emission. The diffuse emission to the SE consists of five elongated components (labeled R1–R5, Figure 3 left), with a total extent of $11'$ in the NE-SW direction, corresponding to 1.9 Mpc at the distance of A3411. Four of those components can be seen in the higher-resolution BnA array image (Figure 3 left). The fifth fainter component (R5) extends further to the SW (see Figure 1 left). The source (the part around *A*) is located at a projected distance ($R_{\text{projected}}$) of 1.3 Mpc from the X-ray peak of the cluster. We measure an integrated flux of 53 ± 3 mJy for the combined emission from R1 to R5 in the image with the compact sources subtracted (Figure 1 left panel). This corresponds to a radio power of $P_{1.4\text{GHz}} = 5.0 \times 10^{24}$ W Hz $^{-1}$. In addition, we detect polarized emission from this diffuse source at the 10 to 25% level (see Figure 2 right). The emission cannot be classified as the lobes of an FR-I type radio source (Fanaroff & Riley 1974) related to either source *A* or *B*, since the counterparts of *A* and *B* are located in or behind the cluster, making the source too large for an FR-I type radio galaxy. Instead, we classify this complex radio source as a radio relic due to its size, location with respect to cluster center, lack of optical counterparts, and polarization properties.

In the central part of the cluster, diffuse and faint extended radio emission is found which we classify as a radio halo. The halo appears to be elongated along the major axis of the ICM and has a physical extent of ~ 0.9 Mpc. We measure a total flux in the halo region of 4.8 ± 0.5 mJy which corresponds to $P_{1.4\text{GHz}} = 4.6 \times 10^{23}$ W Hz $^{-1}$. Most of the halo is only detected at $\sim 3\sigma_{\text{rms}}$. The actual halo power might be somewhat higher as the total extent of the halo could be larger. The radio power we find is consistent with the X-ray luminosity–radio power (L_X – $P_{1.4}$) correlation for radio–halo clusters (e.g., Liang et al. 2000; Enßlin & Röttgering 2002; Cassano et al. 2006).

3.1. Radio Spectral Index

The radio relic is also visible in an image from the 352 MHz Westerbork In the Southern Hemisphere (WISH) survey

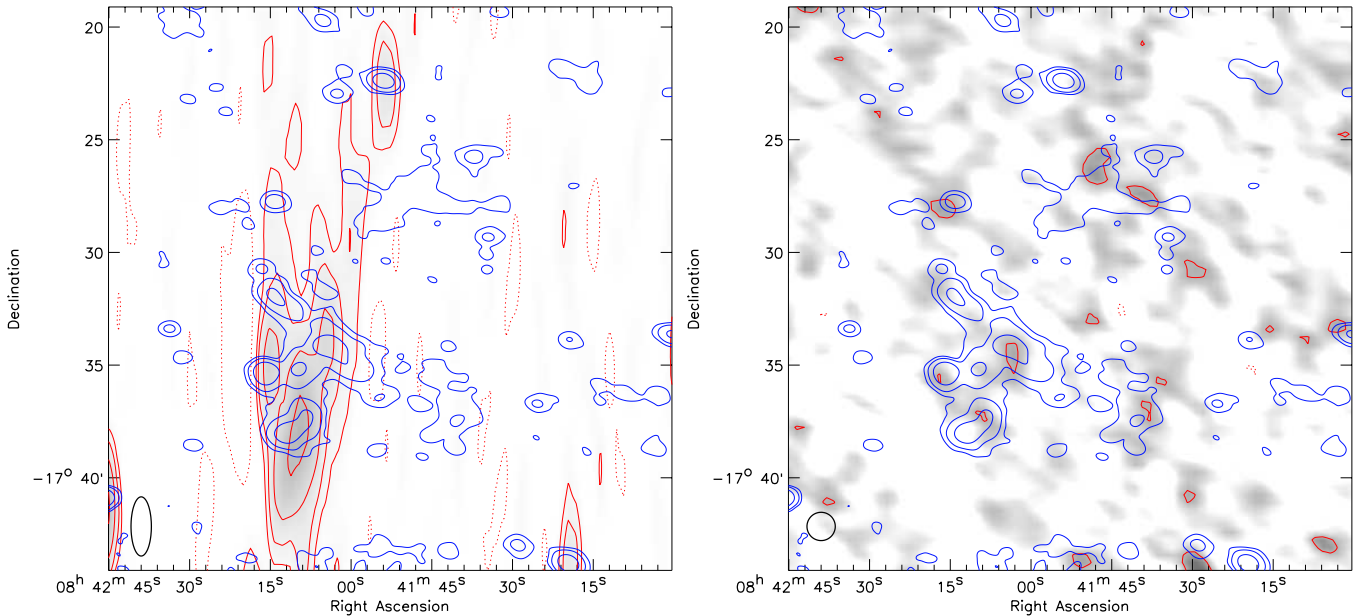


Figure 4. Left: 352 MHz WISH image. The red contour levels are drawn at $[1, 2, 4, 8, \dots] \times 4.5 \text{ mJy beam}^{-1}$. Negative $-3\sigma_{\text{rms}}$ contours are shown by the dotted lines. The image has a very elongated beam shape with a size of $157'' \times 54''$. The blue contours are from the VLA DnC array image and drawn at levels of $[1, 2, 4, 8, \dots] \times 0.141 \text{ mJy beam}^{-1}$. Right: VLSS 74 MHz image of the cluster covering the same region as the 352 MHz WISH image. The red contours are drawn at 0.2 (solid) and -0.3 (dotted) Jy beam^{-1} . The grayscales display the range between 0 and 0.7 Jy beam^{-1} (white to black). The image has a circular beam size of $75''$.

(A color version of this figure is available in the online journal.)

(De Breuck et al. 2002; see Figure 4, left panel). Some positive residuals are also seen in the reprocessed 74 MHz VLSS image (Lane et al. 2012; Figure 4, right panel), but none of these are above the $3\sigma_{\text{rms}}$ level. We convolved the VLA DnC, VLSS and WISH images to a common resolution of $157'' \times 80''$. Since it is not possible to properly take into account the flux of the compact sources in the WISH and VLSS images we summed the flux over the entire cluster area (Figure 2, right panel). The radio relic flux is expected to dominate the flux contribution, since it is the brightest source at 1.4 GHz and should have a spectral index steeper than the compact sources (mainly sources A and B). At 1.4 GHz, the relic flux contributes 66% to the total flux measured in the cluster area.

From a power-law fit through the flux measurements at 74, 352,⁸ and 1.4 GHz we find a spectral index $\alpha = -1.0 \pm 0.1$ (see Figure 5), typical for bright radio relics (e.g., Clarke & Enßlin 2006; van Weeren et al. 2012).

4. DISCUSSION

The sharp northwestern outer edge of the subcluster core, visible in the *Chandra* image, suggests that this subcluster fell in from the SE. The edge itself could be a cold front, although deeper *Chandra* observations are needed to prove that. The presence of a large radio relic, likely tracing a shock, and the X-ray morphology imply that we currently view the merger after core passage of the subclusters. However, during a head-on binary cluster merger event, two outward propagating shock waves should form along the merger axis (e.g., Roettiger et al. 1999; Ricker & Sarazin 2001). Based on the X-ray morphology, a counter shock wave is expected to be traveling in front of the visible core in Figure 1. This region is mostly outside the FOV

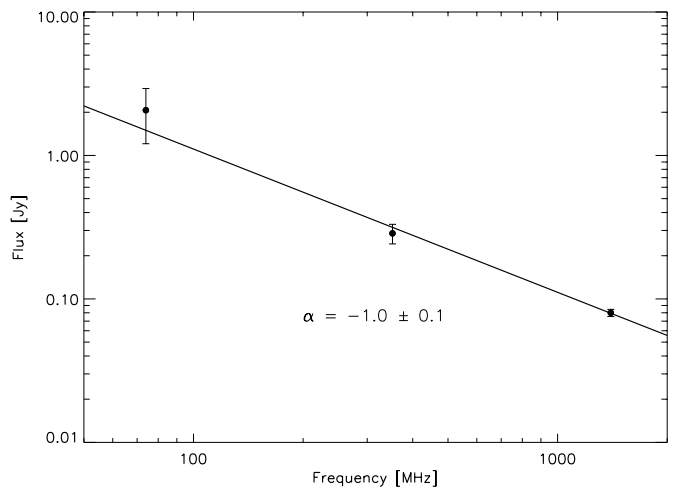


Figure 5. Radio flux measurements of the combined flux (diffuse + compact sources) in the entire A3411 cluster region. The flux measurements were taken from the 74 MHz reprocessed VLSS survey, the 352 MHz WISH survey, and the 1.4 GHz VLA DnC array image. The errors are based on the map noise, scaling with the square root of the number of beam areas, and absolute flux calibration uncertainties of 10% for the VLSS, 15% for WISH, and 5% for the VLA measurement. We note that the reprocessed VLSS survey is on the Scaife & Heald (2012) scale so the 74 MHz flux could be a little overestimated compared to the WISH and NVSS flux measurements (see Scaife & Heald 2012).

of the *Chandra* observation. The VLA observations cover the entire region up to 2 Mpc NW of the subcluster core but no relic is found there. This could indicate that (1) the Mach number of this proposed counter shock is not high enough for efficient particle acceleration (e.g., Hoefl & Brüggén 2007; Kang et al. 2012). Simulations indicate that for cluster mergers with mass ratios $\gtrsim 3$ very little radio emission is generated by the shock wave in front of the less massive subcluster (van Weeren et al. 2011a). From simulations it is also found that the radio emission

⁸ We took the empirical flux correction factor of 0.83, listed in De Breuck et al. (2002), into account.

in clusters is highly dependent on the merger state, varying on time scales of a few hundred million years (Skillman et al. 2011). (2) The lack of a radio relic could also imply there is no supply of pre-accelerated fossil electrons available to be re-accelerated by the shock (Russell et al. 2011) because of the lack of nearby radio galaxies. Some of the halo emission is located NW of the well-defined subcluster core. A third possibility therefore is that we actually observe relic emission from the counter shock but it is very faint and partly mixed with halo emission. Because the NW cluster region is at the edge of the *Chandra* FOV and the radio emission is only visible at the 3σ level, it is currently difficult to determine the precise nature of this radio emission. Deep polarization observations could distinguish between radio halo and relic emission since relics are usually polarized at the 10%–20% level or more.

The A3411 radio relic is located at $R_{\text{projected}} = 1.3$ Mpc, a typical distance for radio relics. Few relics are found at distances beyond 2 Mpc from the cluster center (van Weeren et al. 2011b; Vazza et al. 2012). If we assume that the A3411 relic is located at a distance of 2 Mpc or less from the cluster center, then we can constrain the line of sight of the merger geometry. The shock waves are then seen under an angle of $<50^\circ$ from edge-on and the merger axis is located under an angle of $>40^\circ$ along the line of sight.

The relic in A3411 has a very complex morphology compared to some of the other clusters with single relics like A521 (Giacintucci et al. 2008), A746 (van Weeren et al. 2011b), and A2744 (Govoni et al. 2001). The morphology of A3411 could have been the result of the shock front interacting with substructures or a large-scale galaxy filament from the cosmic web connected to the cluster (e.g., Paul et al. 2011; van Weeren et al. 2011a). Numerical simulations also show that the morphology of relics in clusters can be quite complex (Hoeft et al. 2011; Vazza et al. 2012; Skillman et al. 2011) and detailed observations of the single relics in, for example, A2256, the Coma Cluster, and 1RXS J0603.3+4214 also reveal more intricate morphologies (e.g., Giovannini et al. 1991; Clarke & Enßlin 2006; van Weeren et al. 2012). The broken morphology of the relic could also be caused by fluctuations in the magnetic field strength across the relic, which again could have been caused by varying conditions along the shock front.

Another possibility is that the shock re-accelerates relativistic electrons trapped in several fossil radio bubbles from previous episodes of active galactic nucleus activity or the relic traces the lobes of radio galaxies whose plasma has been shredded and re-accelerated. This would imply that electrons are not directly accelerated from the thermal pool (Markevitch et al. 2005) but were already pre-accelerated. The fossil radio bubbles could for example be related to radio galaxy *B* and/or *A*. What is problematic is that it might be very difficult to distinguish between shock acceleration and shock re-acceleration on the basis of the spectral properties of relics.

The dynamical state of the cluster is not well known due to the lack of a sufficient number of redshift measurements. For example, it is unclear if the cD galaxy near the peak of the X-ray emission (2MASX J08415287-1728046; see Figure 1 right panel) is associated with the well-defined subcluster core. Cluster merger events can decouple the baryonic matter component from the dark matter and galaxies (e.g., Clowe et al. 2006; Shan et al. 2010; Dawson et al. 2012) due to ram pressure stripping. If this subcluster undergoes ram pressure stripping and the subcluster moves toward the NW, then we would expect the cD galaxy to be located ahead of the subcluster's core, which

is not the case. Therefore it is possible that the cD galaxy is associated with the primary subcluster surrounding the smaller subcluster core.

5. CONCLUSIONS

We have presented *Chandra* and VLA observations of the Planck ESZ cluster A3411. The *Chandra* image reveals the cluster to be undergoing a merger event. The VLA observations display a central 0.9 Mpc radio halo and a one-sided polarized 1.9 Mpc relic with a complex unusual morphology to the SE of the cluster. For the combined emission in the cluster area (diffuse emission + compact sources), we find a radio spectral index of -1.0 ± 0.1 between 74 MHz and 1.4 GHz, consistent with a bright radio relic dominating the flux contribution. The morphology of the relic could be related to a complex-shaped shock surface due to the presence of substructures in the ICM, or reflect the presence of electrons trapped in several distinct fossil radio bubbles. This would then suggest that the synchrotron emitting electrons are re-accelerated by the shock and not accelerated directly from the thermal pool. Another possibility is that the complex morphology of the relic is caused by fluctuations in the magnetic field strength across the relic, for example caused by varying shock properties related to substructures.

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