

A census of stellar mass in ten massive haloes at $z \sim 1$ from the GCLASS Survey

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ABSTRACT

Aims. We study the stellar mass content of massive haloes in the redshift range $0.86 < z < 1.34$, by measuring (1) the stellar mass in the central galaxy versus total dynamical halo mass; (2) the total stellar mass (including satellites) versus total halo mass; and (3) the radial stellar mass and number density profiles for the ensemble halo.

Methods. We use a K_s -band selected catalogue for the 10 clusters in the Gemini Cluster Astrophysics Spectroscopic Survey (GCLASS), with photometric redshifts and stellar masses measured from 11-band SED fitting. Combining the photometric catalogues with the deep spectroscopic component of GCLASS, we correct the cluster galaxy sample for interlopers. We also perform a dynamical analysis of the cluster galaxies to estimate the halo mass M_{200} for each cluster based on a measurement of its velocity dispersion.

Results. (1) We find that the central galaxy stellar mass fraction decreases with total halo mass and that this is in reasonable, quantitative agreement with measurements from abundance matching studies at $z \sim 1$. (2) The total stellar mass fractions of these systems decrease with halo mass, indicating that lower mass systems are more efficient at transforming baryons into stars. We find the total stellar mass to be a good proxy for total halo mass, with a small intrinsic scatter. When we compare these results from GCLASS with literature measurements, we find that the stellar mass fraction at fixed halo mass shows no significant evolution in the range $0 < z < 1$. (3) We measure a relatively high NFW concentration parameter $c_g \sim 7$ for the stellar mass distribution in these clusters, and debate a possible scenario for explaining the evolution of the stellar mass distribution from the GCLASS sample to their likely descendants at lower redshift.

Conclusions. The stellar mass measurements in the $z \sim 1$ haloes provided by GCLASS puts constraints on the stellar mass assembly history of clusters observed in the local Universe. A simple model shows that the stellar mass content of GCLASS can evolve in typical distributions observed at lower redshifts if the clusters primarily accrete stellar mass onto the outskirts.

Key words. galaxies: clusters: general – galaxies: evolution – galaxies: photometry

1. Introduction

One of the main objectives in the field of extragalactic astronomy is to understand the connection between galaxies and the distribution of the underlying dark matter. The growth of dark matter structures has been studied in large N -body simulations (e.g. [Springel et al. 2005](#); [Boylan-Kolchin et al. 2009](#); [Navarro et al. 2010](#)). From these simulations, the density profiles of collapsed structures have been found to be well represented by NFW-profiles ([Navarro et al. 1997](#)). These profiles are described by two parameters: the halo mass and the halo concentration parameter. The dependence of the concentration parameter c on the halo mass, formation time and redshift has been studied with N -body simulations (e.g. [Wechsler et al. 2002](#); [Neto et al. 2007](#); [Duffy et al. 2008](#); [Gao et al. 2008](#)). These have shown that c , for the dark matter, is higher for lower mass haloes, higher for haloes that collapse early, and higher for haloes at lower redshift.

How baryons affect the distribution of the dark matter is still under debate ([Dolag et al. 2009](#); [van Daalen et al. 2011](#); [Newman et al. 2013](#)). Baryons in the gas phase can cool and form stars at the bottom of the potential wells in the dark matter (sub-)haloes. The efficiency with which this happens depends on the properties

of the halo (see e.g. [Kravtsov & Borgani 2012](#); [Planelles et al. 2013](#)). To constrain the physics behind these processes, there are a number of key observables that can be exploited. In this paper we concentrate on three of these, which we introduce in turn below, and measure at $z \sim 1$ for a sample of ten cluster sized haloes.

First, to constrain the build up of stellar mass in central galaxies, we measure the stellar mass present in the central galaxies of GCLASS and compare it to direct measurements of their total halo masses. [Behroozi et al. \(2013\)](#) estimate the stellar mass in central galaxies versus total halo mass over a range of redshifts and halo masses in a *statistical* way using the abundance matching technique. In this technique observables, such as the stellar mass function and cosmic star formation history, are combined with merger trees from dark matter simulations to provide constraints on the processes that build up the stellar mass in central galaxies. The stellar content of central galaxies, or brightest cluster galaxies (BCGs) in the case of clusters, can grow by star formation in the galaxy itself or by merging with other galaxies. Given the significant growth of stellar mass in BCGs as a function of redshift ([Lin & Mohr 2004](#); [Lidman et al. 2012](#)), this build-up is likely to occur through mainly dry

mergers. However, the mass assembly has been shown (Lidman et al. 2012) to increase more slowly than is expected from semi-analytic models (De Lucia & Blaizot 2007), but in good agreement with more recent simulations (Laporte et al. 2013). Since the main halo also accretes matter while the central galaxy is building up its stellar content, studies have focussed on the relationship between those processes. The Behroozi et al. (2013) estimates at $z = 1.0$ cover a range of halo masses from $10^{11.3} < M_h/M_\odot < 10^{14.2}$, and are consistent with predictions from other abundance matching analyses (e.g. Moster et al. 2010, 2013). In general the highest central stellar mass fraction is found in haloes of around $10^{12} M_\odot$. By combining *direct* measurements of total mass and stellar mass in the same haloes, we test the results from abundance matching studies at $z \sim 1$.

Second, a key measurement for understanding the interplay between the growth of large scale structure and the formation and accretion of galaxies is to compare the total stellar mass as a function of halo mass. For a sample of groups selected at $0.1 < z < 1.0$ from COSMOS, Giodini et al. (2009) show that their stellar mass fraction is a decreasing function of halo mass. Similar results are found by Gonzalez et al. (2007, 2013), Andreon (2010) and Hilton et al. (2013) for samples of clusters around $z = 0.1$, $z < 0.1$ and $z = 0.5$, respectively. Given that the most massive haloes are expected to grow by accreting lower mass systems, which have a higher stellar mass fraction, one would naively expect the stellar mass fraction of massive haloes to grow with cosmic time, even in the absence of in situ star-formation processes. Consequently, measurements on the stellar mass fraction in these haloes are used to constrain the progenitor population that form the building blocks of these haloes (Balogh et al. 2008; McGee et al. 2009). Due to the major caveats in comparing measurements from different studies with inhomogeneous data and different analyses, the relation itself is hard to constrain observationally (Leauthaud et al. 2012a; Budzynski et al. 2013). So far little evolution with redshift has been found (Giodini et al. 2009; Lin et al. 2012).

Third, the spatial distribution of the stellar mass component of satellites within the main halo is intimately related to accretion processes, and eventually the growth of the central galaxy. While the sub-haloes in pure gravitational N -body simulations get destroyed by tidal disruptions, the galaxies that have formed inside of them are more resistive to those forces (e.g. Budzynski et al. 2012). The near-infrared luminosity and number density profiles are found to be described by NFW profiles for group-sized haloes (e.g. Giodini et al. 2009; Tal et al. 2013), and clusters (e.g. Carlberg et al. 1997b; Lin et al. 2004; Muzzin et al. 2007). Budzynski et al. (2012) measured the radial distribution of galaxies from the SDSS around Luminous Red Galaxies in a redshift range $0.15 < z < 0.4$, and found that this distribution is also well described by an NFW profile. However, they found that the concentration parameter c is lower for the galaxies than for the underlying dark matter. They found that the concentration is independent of mass, but that there is a hint of a mild dependence of the stellar mass concentration on redshift. A comparison of the radial stellar mass density distribution of clusters over a range of redshifts, linking high- z systems to their likely descendants, yields insights in the evolution of the galaxy distribution. In this study we extend the redshift baseline of these comparisons towards $z = 1$.

We perform the aforementioned key measurements in an unexplored combination of redshift and halo mass range using the GCLASS survey, which contains deep 11-band photometry and spectroscopy for 10 rich clusters at $0.86 < z < 1.34$. This paper builds further on the results presented in several papers on the

GCLASS sample. Muzzin et al. (2012, hereafter M12) present the spectroscopic sample, which is critical in this study to correct the photometrically selected galaxies by cluster membership. van der Burg et al. (2013, hereafter vdB13) measure the stellar masses of the galaxies in the sample and present their stellar mass function (SMF). We use the stellar masses estimated in this work for the current study. Lidman et al. (2012) identifies and studies the BCGs of GCLASS clusters as part of their analysis on the central galaxy stellar mass growth. The total GCLASS halo masses are estimated based on the velocity dispersions estimated in Wilson et al. (in prep.). To describe the masses of the clusters, we use R_{200} and M_{200} , which are defined as the radius at which the mean interior density is 200 times the critical density of the Universe, and the mass enclosed within this radius, respectively.

The structure of this paper is as follows. In Sect. 2 we present the GCLASS cluster sample, the available photometric and spectroscopic data, and give the results from a dynamical analysis to estimate the total halo masses. We also show how we obtain photometric redshifts and stellar mass estimates by summarizing the analysis from vdB13. We further show how the spectroscopic data are used to correct the full photometric catalogue for cluster membership. In Sect. 3 we compare the stellar mass in the central galaxies with their halo masses. In Sect. 4 we present results on the total stellar mass versus halo mass relation of the clusters. In Sect. 5 we show how the galaxies are distributed radially and compare this to the expected dark matter profiles for these systems. Further, we discuss a possible evolutionary model to connect the $z \sim 1$ measurements to their likely descendants at lower redshift. In each section we compare the results with the literature and discuss how they are affected by possible systematics. We summarise and conclude in Sect. 6.

All magnitudes we quote are in the AB magnitudes system and we adopt Λ CDM cosmology with $\Omega_m = 0.3$, $\Omega_\Lambda = 0.7$ and $H_0 = 70 \text{ km s}^{-1} \text{ Mpc}^{-1}$. For stellar mass estimates we assume the Initial Mass Function (IMF) from Chabrier (2003).

2. GCLASS data and analysis

The GCLASS cluster sample consists of 10 rich clusters in the redshift range $0.86 < z < 1.34$ selected with the red-sequence selection method (Gladders & Yee 2000) using the $z' - 3.6 \mu\text{m}$ colour from the 42 square degree SpARCS survey (Muzzin et al. 2009; Wilson et al. 2009). These 10 clusters, which are amongst the richest at $z \sim 1$ in this survey area, are described in M12, and can be considered as a fair representation of IR-selected rich clusters within this redshift range. It is always a question how representative a cluster sample is of the full distribution of massive haloes, as it is impossible to select a sample based on halo mass. Each selection method has potential biases, whether it is X-ray selected, SZ-selected or galaxy-selected. However, especially at the high-mass end of the distribution, these selection methods are unlikely to cause significant biases in favour of particular types of galaxy clusters. Specifically, as e.g. Blakeslee et al. (2003) and Mullis et al. (2005) show, X-ray and SZ-selected clusters also show significant over-densities of red-sequence galaxies. We discuss a possible selection bias further in Sect. 5.1. An overview of the GCLASS sample is given in Table 1.

The BCGs of these clusters have been identified and studied in Lidman et al. (2012). In general the identification of the BCGs is straightforward, being the brightest cluster member in the K_s -band, and we use the same identification as done in Lidman et al. (2012). In the cases of SpARCS-1051 and SpARCS-1634,

Table 1. The ten GCLASS clusters selected from SpARCS that form the basis of this study, with their dynamical properties.

Name ^a	z_{spec}	RA ^b J2000	Dec ^b J2000	σ_v ^c [km/s]	M_{200} ^d [$10^{14} M_{\odot}$]	R_{200} ^d [Mpc]	Spec- z Members
SpARCS-0034	0.867	00:34:42.06	-43:07:53.41	700^{+90}_{-150}	$2.4^{+1.0}_{-1.2}$	$0.9^{+0.1}_{-0.2}$	45
SpARCS-0035	1.335	00:35:49.70	-43:12:24.20	780^{+80}_{-120}	$2.5^{+0.9}_{-1.0}$	$0.8^{+0.1}_{-0.1}$	20
SpARCS-0036	0.869	00:36:45.03	-44:10:49.91	750^{+80}_{-90}	$2.9^{+1.0}_{-0.9}$	$1.0^{+0.1}_{-0.1}$	47
SpARCS-0215	1.004	02:15:24.00	-03:43:32.15	640^{+120}_{-130}	$1.7^{+1.1}_{-0.8}$	$0.8^{+0.2}_{-0.2}$	48
SpARCS-1047	0.956	10:47:33.43	57:41:13.30	660^{+70}_{-120}	$1.9^{+0.7}_{-0.9}$	$0.8^{+0.1}_{-0.2}$	31
SpARCS-1051	1.035	10:51:11.21	58:18:03.17	500^{+40}_{-100}	$0.8^{+0.2}_{-0.4}$	$0.6^{+0.1}_{-0.1}$	34
SpARCS-1613	0.871	16:13:14.63	56:49:29.95	1350^{+100}_{-100}	$16.9^{+4.0}_{-3.5}$	$1.8^{+0.1}_{-0.1}$	92
SpARCS-1616	1.156	16:16:41.32	55:45:12.44	680^{+80}_{-110}	$1.9^{+0.7}_{-0.8}$	$0.8^{+0.1}_{-0.1}$	46
SpARCS-1634	1.177	16:34:38.22	40:20:58.36	790^{+60}_{-110}	$2.9^{+0.7}_{-1.0}$	$0.9^{+0.1}_{-0.1}$	50
SpARCS-1638	1.196	16:38:51.64	40:38:42.91	480^{+50}_{-100}	$0.6^{+0.2}_{-0.3}$	$0.5^{+0.1}_{-0.1}$	44

Notes. ^(a) For full names we refer to Muzzin et al. (2012). ^(b) Coordinates of the BCGs, as identified by Lidman et al. (2012). ^(c) Velocity dispersions estimated by Wilson et al. (in prep.). ^(d) Dynamical properties estimated using the relation between σ_v and M_{200} from Evrard et al. (2008).

Lidman et al. (2012) found that the BCGs are off-set from the approximate cluster centre by about 250 kpc (projected).

The photometric data set consists of imaging in *ugrizJKs* and 4 IRAC channels for each cluster. For details on the data reduction, and a description of the catalogue, we refer to vdB13. In summary, the catalogues contain objects detected in the K_s -band, with Gaussian-weighted aperture fluxes in 11 filters to constrain the SEDs of the objects, and to separate stars from galaxies by combining their $u - J$ and $J - K$ colours. The depth of the images, and therefore the completeness of the catalogues, differs slightly from cluster to cluster. The median photometric completeness limit (80%), in terms of stellar mass, is $10^{10.16} M_{\odot}$ for the 10 clusters in the GCLASS sample.

Each cluster has substantial spectroscopic coverage provided by the GMOS instruments on Gemini North and Gemini South. The targets for spectroscopic follow-up were prioritized by their $3.6 \mu\text{m}$ flux and their projected cluster-centric distance, as explained in M12. The membership of the massive galaxies that constitute most of the stellar mass in the clusters are thus confirmed spectroscopically. Since the targeting completeness is well understood, we can use the sub-sample for which we have spectra to statistically correct the full catalogue for cluster membership. How this is done is outlined in vdB13 (Sect. 3.4), and expanded on in Sect. 2.3.

2.1. Total halo masses

Using the sample of spectroscopically identified cluster galaxies, totalling 457 members for 10 clusters, we perform a dynamical analysis to estimate masses for each cluster. From the line-of-sight velocity distributions, which show approximately Gaussian profiles, the velocity dispersions are measured (Wilson et al., in prep.) using standard methods such as the shifting gapper and the bi-weight estimator (Beers et al. 1990; Girardi et al. 1993; Fadda et al. 1996), see Table 1. Since we do not measure the velocity dispersion from dark matter particles but from subhaloes (or galaxies), several dynamical effects render this an imperfect tracer of the gravitational potential (e.g. Saro et al. 2013). In an attempt to take account of these biases (which also depend on the spectroscopic target selection), various scalings between the velocity dispersion and halo mass (M_{200}) have been proposed in the literature (e.g. Carlberg et al. 1997a; Evrard et al. 2008; Munari et al. 2013). These are of the form

$$\sigma_{\text{1D}} = A_{\text{1D}} \left[\frac{h(z) M_{200}}{10^{15} M_{\odot}} \right]^{\alpha} \text{ km s}^{-1}, \quad (1)$$

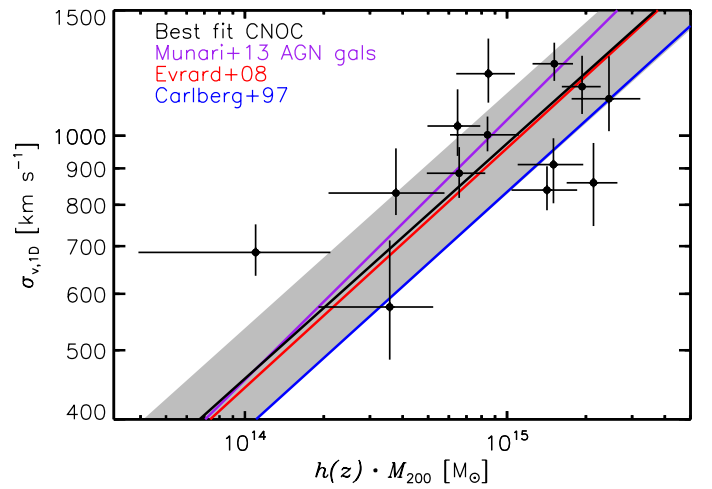


Fig. 1. Measured velocity dispersion versus halo mass ($h(z) \cdot M_{200}$). Data points are measurements on the CNOC sample. Lensing masses are from Hoekstra et al. (2012) (which are revised in Hoekstra et al., in prep.), whereas velocity dispersions are obtained from Borgani et al. (1999) and Girardi & Mezzetti (2001). Although there is a substantial amount of intrinsic scatter (grey region indicates $\pm 1\text{-}\sigma$ intrinsic scatter around the relation), the best fit to these data (black line) is very close to the Evrard et al. (2008) scaling relation (red).

where A_{1D} and α are parameters that are different for each study (Fig. 1).

In order to determine which scaling relation gives the best halo mass estimate for the measured velocity dispersions in GCLASS, we consider a sample of clusters which were originally studied as part of the Canadian Network for Observational Cosmology (CNOC, Yee et al. 1996). A weak-lensing study has been performed for these systems, which provides for independent mass estimates (Hoekstra 2007; Hoekstra et al. 2012, revised in Hoekstra et al., in prep.). For 13 of the clusters in this sample, velocity dispersions have been measured from spectroscopic targets that were chosen in a similar way as the targets selected in the GCLASS sample (Borgani et al. 1999; Girardi & Mezzetti 2001). Figure 1 compares the weak-lensing masses (M_{200}) to the line-of-sight cluster velocity dispersions. We fit a linear relation in this logarithmic plane, while fixing the slope to $\alpha = \frac{1}{3}$, and allow for the presence of intrinsic scatter in the fit. The black line shows the best-fitting scaling relation to the data points ($A_{\text{1D}} = 972^{+60}_{-52} \text{ km s}^{-1}$), and we

find a significant amount of intrinsic scatter around this relation ($\log(\sigma_{\sigma_v/h(z) \cdot M_{200}}) = 0.07^{+0.03}_{-0.02}$ dex). The best-fitting scaling relation is very similar to the relation suggested by [Evrard et al. \(2008\)](#). To estimate halo masses of the GCLASS clusters, we therefore use the [Evrard et al. \(2008\)](#) scaling relation. This relation was also used by a recent dynamical study on the ACT cluster sample ([Sifón et al. 2013](#)), which simplifies a comparison with the results from this sample (e.g. [Hilton et al. 2013](#)) in the rest of this paper. Values of M_{200} and R_{200} are shown in Table 1. Statistical uncertainties are given (propagated from uncertainties on the velocity dispersion), but note that there is also a significant systematic uncertainty ($\sim 20\%$), corresponding to the choice of scaling relation, and indicated by the substantial amount of intrinsic scatter. Note that the R_{200} values have a smaller fractional uncertainty, since $R_{200} \propto M_{200}^{1/3}$.

2.2. Photometric redshifts and stellar masses

We estimate photometric redshifts for all galaxies in the K_s -band selected catalogue using the EAZY code ([Brammer et al. 2008](#)). In vdB13 we assessed the performance by comparing the photo- z estimates to spec- z measurements for the galaxies that have been observed spectroscopically. We found a scatter of $\sigma_z = 0.036$ in $\frac{\Delta z}{1+z}$, a negligible bias and fewer than 5% outliers.

After fixing the redshift for each object at its spec- z , or the photo- z when a spec- z is not available, we estimate stellar masses using FAST ([Kriek et al. 2009](#)). The stellar population libraries from [Bruzual & Charlot \(2003\)](#) are used to obtain the model SED that gives the best fit to the photometric data. We use a parameterization of the star formation history as $SFR \propto e^{-t/\tau}$, where the time-scale τ is allowed to range between 10 Myr and 10 Gyr. We also assume a [Chabrier \(2003\)](#) IMF, solar metallicity, and the [Calzetti et al. \(2000\)](#) dust law. For estimates on the stellar-mass completeness of the catalogues, we refer to vdB13. To approximate the statistical uncertainty on each stellar mass measurement, we perform 100 Monte-Carlo simulations in which we perturb the photometric aperture flux measurements within their estimated errors. Each realisation of the catalogue gives a slightly different SED fit, and therefore the mass-to-light ratio ($M_\star/L$) is different. We translate the obtained scatter in $M_\star/L$ into an approximate uncertainty on the stellar mass, after including uncertainties on the spectral templates themselves. We find typical uncertainties on $M_\star/L$ of 0.21 dex at $M_\star \sim 10^{10} M_\odot$, and 0.13 dex at $M_\star \sim 10^{11} M_\odot$.

2.3. Cluster membership correction

Figure 2 shows the distribution of total stellar mass contained in galaxies with a given stellar mass. The best-fitting Schechter function to the total galaxy population from vdB13 is also shown. The points with error bars are the measurements of the SMF presented there, and are integrated masses over the SMF in each bin, i.e. $\int_{M_{\min}}^{M_{\max}} \Phi(M) \cdot dM$, where $\Phi(M)$ represents the number density of galaxies as a function of stellar mass. With the characteristic mass of the Schechter function around $M_\star = 10^{11} M_\odot$, galaxies around this mass contribute most to the total stellar mass of the cluster. For the galaxies with stellar masses around $M_\star = 10^{11} M_\odot$, the completeness is high ($\geq 50\%$). For measurements within R_{500} the completeness is even higher. For that reason, the measurements of the total stellar mass of the clusters are based mostly on spectroscopic redshifts, and are robust with respect to how we correct the photometric sample for completeness.

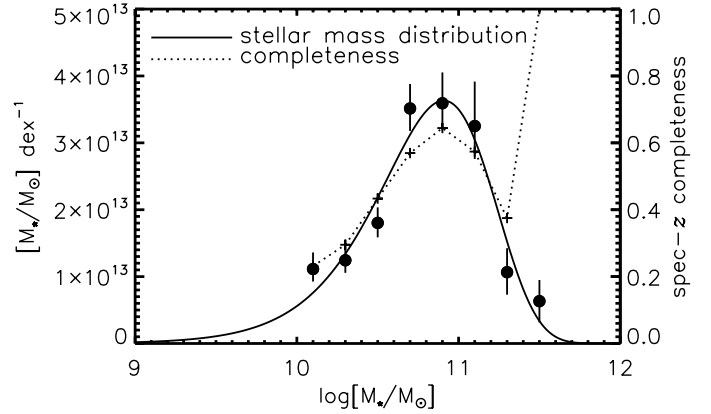


Fig. 2. Solid line: distribution of total stellar mass contained in galaxies with a given stellar mass. The points with error bars are the measurements of the SMF presented in [van der Burg et al. \(2013\)](#), but integrated over the mass bins. Dotted line: spectroscopic completeness for galaxies with projected distances from the BCG less than R_{200} . For the galaxies that constitute most of the stellar mass in the clusters, the spectroscopic completeness is high ($\geq 50\%$).

We use the limited number of galaxies in the fields that have been targeted spectroscopically to estimate the probability that a galaxy is part of the cluster for the objects that do not have a measured spectroscopic redshift. For objects with stellar masses exceeding $\sim 10^{10} M_\odot$ that were targeted, the success rate of obtaining a reliable spec- z is higher than 90% (M12). Given that the targeting prioritization is known (M12), we can correct the photometrically selected sample for cluster membership using the sub-sample of spectroscopic targets. To do this we take a similar approach as outlined in vdB13 (Sect. 3.4), but with a few adaptations.

The radial distance of each galaxy is rescaled to units of R_{200} , instead of physical distance. Then we measure for the cluster ensemble, in bins of radial distance and stellar mass, the fraction of correctly identified cluster galaxies based on their photo- z . Comparing this number to the total number of spec- z selected cluster members in this bin, we obtain membership correction factors that are used to correct the photometrically selected numbers for membership. The correction factors as a function of radial distance are shown in Fig. 3. The membership correction factors are a decreasing function of distance, since the clusters are less overdense further away from their cluster centre. The blue (red) points represent the population of star-forming (quiescent) galaxies. For the correction factors as a function of stellar mass we refer to vdB13 (Fig. 4).

To further improve the estimates on the total stellar mass associated with each cluster, we estimate the contamination by field galaxies for each individual cluster. This minor correction to the photo- z selected sample for each cluster is due to cosmic variance, slight differences in photometric redshift quality between the fields, and also the dependence of angular size associated with R_{200} on the cluster mass and redshift. To estimate this overdensity parameter we 1) apply the correction factors that we use on the photometric sample (e.g. Fig. 3 and vdB13 (Fig. 4)) on all spectroscopically targetted galaxies; then 2) use this to estimate the number of cluster members in this sample; and 3) divide the actual number of spectroscopic cluster members by the estimated number of cluster members to give the correction factor. This cluster overdensity parameter is by construction around 1.0 and ranges from 0.86 to 1.22 for the clusters in our sample.

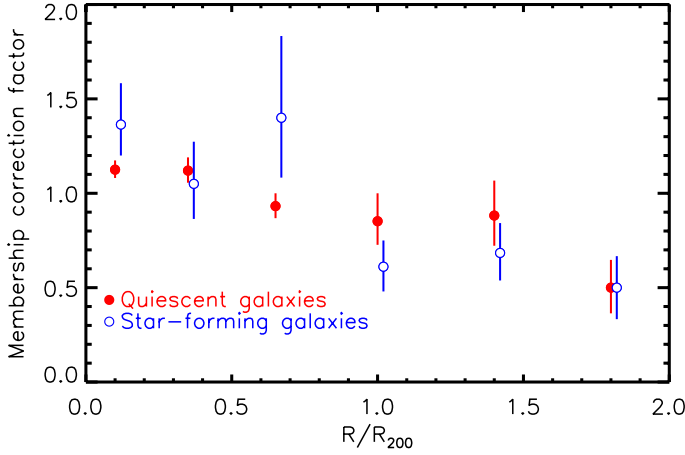


Fig. 3. Correction factors as a function of radius, scaled by R_{200} , for the cluster ensemble. Error bars represent uncertainties estimated from Monte-Carlo simulations. Further away from the projected centres, the correction factors go down because galaxies are increasingly more likely to be part of the field.

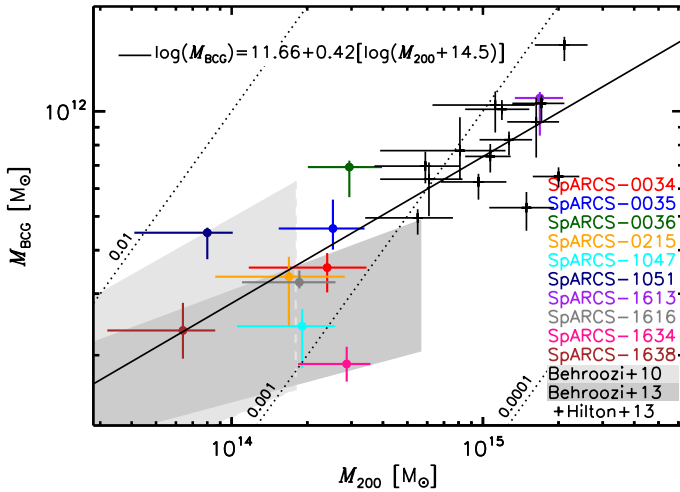


Fig. 4. BCG stellar mass versus total halo mass. Black dots show lines of constant stellar mass fractions of 0.0001, 0.001 and 0.01. '+'-signs show the results from Hilton et al. (2013). The relation is the best fit to the combined data set of Hilton et al. (2013) and the current study. Estimates from Behroozi et al. (2010) and Behroozi et al. (2013) are indicated by the shaded regions.

3. Central stellar mass versus halo mass

In Fig. 4 we show the stellar mass of the central galaxy as a function of dynamical halo mass. Stellar masses are measured based on M_*/L 's estimated with FAST, multiplied with the total flux in the K_s -band. Since BCGs generally have extended light profiles, their flux measured with SExtractor in Kron elliptical apertures is under-estimated. To account for the total flux of the BCGs in the K_s -band, we use GALFIT to fit Sérsic profiles to these galaxies. We make sure that we carefully mask any nearby satellite galaxies and perform 10 different fits where we convolve the profiles with different stars to approximate uncertainties due to the PSF. We compare the integrated flux in these Sérsic profiles with the SExtractor magnitudes in Table 2. The values show the median values and the maximum and minimum values for the 10 different GALFIT runs, after rejecting the highest and lowest value. The difference between the GALFIT and SExtractor measurements is typically about 0.2 mag, and

Table 2. K_s -band magnitudes for the BCGs identified in Lidman et al. (2012) for the GCLASS clusters.

Name	MAG_AUTO [mag _{AB}]	GALFIT [mag _{AB}]	GALFIT Sérsic - n	$M_{*,\text{BCG}}$ [$10^{11} M_\odot$]
SpARCS-0034	16.59 ± 0.01	$16.51^{+0.04}_{-0.03}$	$3.68^{+0.36}_{-0.37}$	$3.56^{+0.36}_{-0.54}$
SpARCS-0035	17.27 ± 0.01	$17.06^{+0.01}_{-0.01}$	$3.77^{+0.13}_{-0.13}$	$4.61^{+0.97}_{-0.60}$
SpARCS-0036	16.40 ± 0.01	$16.10^{+0.01}_{-0.01}$	$3.82^{+0.30}_{-0.37}$	$6.92^{+0.30}_{-1.25}$
SpARCS-0215	17.05 ± 0.01	$16.86^{+0.02}_{-0.02}$	$3.02^{+0.14}_{-0.13}$	$3.36^{+0.47}_{-0.94}$
SpARCS-1047	17.29 ± 0.01	$17.03^{+0.01}_{-0.03}$	$4.35^{+0.22}_{-0.41}$	$2.42^{+0.29}_{-0.58}$
SpARCS-1051	17.11 ± 0.02	$16.73^{+0.03}_{-0.04}$	$6.87^{+0.97}_{-0.93}$	$4.49^{+0.15}_{-0.72}$
SpARCS-1613	15.67 ± 0.01	$15.50^{+0.00}_{-0.01}$	$3.25^{+0.10}_{-0.15}$	$10.91^{+0.44}_{-2.40}$
SpARCS-1616	17.01 ± 0.01	$16.96^{+0.01}_{-0.02}$	$3.03^{+0.18}_{-0.15}$	$3.24^{+0.26}_{-0.13}$
SpARCS-1634	17.41 ± 0.01	$17.42^{+0.01}_{-0.01}$	$0.83^{+0.01}_{-0.01}$	$1.89^{+0.23}_{-0.21}$
SpARCS-1638	17.71 ± 0.02	$17.43^{+0.01}_{-0.01}$	$5.23^{+0.12}_{-0.12}$	$2.36^{+0.47}_{-0.40}$

Notes. The last column gives the stellar masses corresponding to the GALFIT total integrated magnitude, and the errors also include the statistical uncertainty on M_*/L .

depends mainly on the shape of the profile, which is described by the Sérsic parameter n . To obtain the total stellar masses of the central galaxies we multiply the total flux in the K_s -band with the M_*/L estimated using FAST, and include both the flux-error and the error on M_*/L (which is the dominant source of uncertainty).

Considering the GCLASS data in Fig. 4, we find mild evidence for a correlation between the BCG stellar mass and halo mass, with a Spearman rank coefficient $\rho = 0.49$. The fraction of mass contained in stellar form in the BCG is approximately 0.001 of the halo mass.

Behroozi et al. (2010) and Behroozi et al. (2013) estimated the stellar mass versus virial halo mass relation over a range of redshifts and halo masses using the abundance matching technique. At the high mass end we make a comparison between their estimates and our observations, which are based on direct measurements of the total halo masses and stellar masses of centrals in the same systems. We multiply the Behroozi halo masses by factor 1.11 to account for the difference between their virial halo masses and M_{200} (Bryan & Norman 1998). We show the Behroozi et al. (2010, 2013) prediction for $z = 1$ by the light (dark) shaded area in Fig. 4. Although the allowed areas are large because of including statistical and systematic uncertainties, the results from Behroozi et al. (2010) seem to be in better agreement with the GCLASS data than the results from Behroozi et al. (2013). What is different in both abundance matching studies is the specific treatment of intra-cluster light (ICL) in Behroozi et al. (2013). When a galaxy merger occurs in this new model, the stars associated with the satellite galaxy may either be deposited onto the central galaxy or be ejected into the ICL. Since Behroozi et al. (2013) estimate the ICL to be of a significant contribution to the total stellar mass at $z = 1$, this is potentially related to an under-prediction of the stellar mass in the central galaxies.

To increase the dynamic range in terms of cluster halo mass, in order to constrain the power-law slope of this relation, we compare our results to those from Hilton et al. (2013), which were obtained from a sample of ACT SZ-selected clusters. To be able to compare the results directly, we reduce the stellar masses estimated from Hilton et al. (2013) by 0.24 dex to account for differences in the adopted IMF. Note that Hilton et al. (2013) did not fit the SED of the BCG with a model to constrain M_*/L , but rather assumed a single burst stellar population that

has a formation redshift $z_f = 3$. For the purpose of estimating BCG stellar masses the difference between these approaches is small (<0.1 dex), because the BCGs contain relatively old stellar populations. The M_{200} measurements for this cluster sample are taken from Sifón et al. (2013). Figure 4 shows a clear relation between the BCG stellar mass and total halo mass from GCLASS and Hilton et al. (2013).

When we fit a slope to the combined set of data points, we have to account for intrinsic scatter in the relation to ensure that we do not give too much weight to precise measurements that are far off the mean relation. We follow the approach outlined in Hoekstra et al. (2011) to perform a three parameter fit to these data points. Besides the parameters describing the power-law relation, the intrinsic scatter is assumed to be described by a log-normal distribution, for which we fit the dispersion σ . The intrinsic scatter is best described by $\log(\sigma_{M_{\text{BCG}}|M_{200}}) = 0.12^{+0.03}_{-0.02}$ dex, and the best-fitting relation is $\log(M_{\text{BCG}}) = (11.66 \pm 0.03) + 0.42^{+0.06}_{-0.07} \cdot [\log(M_{200}) - 14.5]$. This relation is plotted in Fig. 4 and indicates that the BCG stellar mass fraction is lower for higher mass haloes. The fit shows that there is a significant amount of intrinsic scatter in the relation between central galaxy stellar mass and halo mass, which is consistent with the finding of e.g. Leauthaud et al. (2012b).

Note that our data do not allow for measurements of the intracluster light (ICL), and therefore the contribution of intracluster stars to the central stellar mass is neglected. Formally the measured values are therefore lower limits, but Burke et al. (2012) show that the contribution of intracluster stars to the total stellar mass at $z \sim 1$ is expected to be significantly smaller than at lower redshifts. In contrast, Behroozi et al. (2013) suggest a picture in which a significant fraction of the ICL has already been formed at $z \sim 1$. There is, however, slight tension between their statistical study and our observations of the stellar mass in the central galaxy (which is related to the build-up of the ICL component) in Fig. 4.

Lidman et al. (2012) measure the BCG stellar mass versus halo mass for a sample of 160 BCGs in the redshift interval $0.03 < z < 1.63$. Besides the different redshift range they study, their analysis is slightly different from ours. Lidman et al. (2012) constrain the M_*/L of the BCGs with J -, and K_s -band data and do not use GALFIT to probe the extended light profiles of the BCGs. The slope fitted by Lidman et al. (2012) is $M_{200} \propto M_{\text{BCG}}^{1.6 \pm 0.2}$. The reciprocal of this is consistent with our slope to within 2σ .

4. Total stellar mass versus halo mass

We make a comparison between the halo masses and the total stellar mass in the halo, including the satellites. We perform all measurements both within R_{200} and within R_{500} to provide a reference and facilitate the comparison with literature measurements. When necessary, we convert between R_{200} and R_{500} by applying the concentration parameter estimated from Duffy et al. (2008). For the mass and redshift range of the GCLASS clusters, Duffy et al. (2008) find a typical concentration of $c = 2.7$, which is consistent with a stacked weak-lensing measurement of clusters at $z \sim 1$ (Sereni & Covone 2013). Corresponding to this concentration parameter, we use the relationships $R_{500} = 0.632 \cdot R_{200}$ and $M_{500} = 0.631 \cdot M_{200}$.

For each cluster we sum the stellar mass contained in galaxies with a spectroscopic redshift consistent with the cluster that exceed the mass completeness limit of the cluster. The K_s -band flux limits were simulated for each cluster, and corresponding

Table 3. Total stellar masses projected onto spheres with radii R_{200} and R_{500} for the GCLASS clusters.

Name	$M_{200,*}$ [$10^{12} M_{\odot}$]	$M_{500,*}$ [$10^{12} M_{\odot}$]	$M_{500,*}(3.6 \mu\text{m})^a$ [$10^{12} M_{\odot}$]
SpARCS-0034	$2.40^{+0.16}_{-0.15}$	$2.10^{+0.14}_{-0.14}$	—
SpARCS-0035	$1.89^{+0.22}_{-0.20}$	$1.50^{+0.18}_{-0.16}$	$5.43^{+2.92}_{-1.90}$
SpARCS-0036	$3.30^{+0.16}_{-0.15}$	$2.74^{+0.14}_{-0.13}$	—
SpARCS-0215	$2.86^{+0.25}_{-0.23}$	$1.55^{+0.15}_{-0.14}$	—
SpARCS-1047	$1.45^{+0.15}_{-0.13}$	$0.94^{+0.09}_{-0.08}$	—
SpARCS-1051	$1.00^{+0.07}_{-0.07}$	$0.60^{+0.06}_{-0.06}$	—
SpARCS-1613	$7.35^{+0.60}_{-0.55}$	$5.68^{+0.42}_{-0.39}$	$18.72^{+9.65}_{-6.37}$
SpARCS-1616	$3.29^{+0.20}_{-0.19}$	$2.75^{+0.16}_{-0.15}$	$7.14^{+2.19}_{-1.68}$
SpARCS-1634	$1.88^{+0.13}_{-0.12}$	$1.38^{+0.11}_{-0.10}$	$3.37^{+2.58}_{-1.46}$
SpARCS-1638	$1.13^{+0.14}_{-0.13}$	$0.92^{+0.13}_{-0.11}$	$2.33^{+1.56}_{-0.93}$

Notes. ^(a) Taking the background subtracted flux in IRAC 3.6 μm and assuming the same M_*/L for every galaxy in each cluster, based on a single burst stellar population with $\tau = 0.1$ Gyr formed at $z_f = 3$.

stellar mass completeness limits were estimated and presented in vdB13 (Table 1). To this we add the photo- z selected sources that we correct for cluster membership using the method explained in Sect. 2.3, provided that their projected radii from the BCG are less than R_{200} (or R_{500}). Since the overdensity of the cluster with respect to the field is different for each cluster, as explained in Sect. 2.3, we correct the total stellar mass of the photometric sample with the cluster overdensity parameter for each cluster.

The stellar mass is now measured within a projected radius of R_{200} (or R_{500}), but to estimate the stellar mass fraction and be able to compare to results in the literature we have to deproject the stellar mass onto a sphere with radius r_{200} (or r_{500}), since the halo mass M_{200} (or M_{500}) is defined in that way. Assuming a concentration parameter $c = 2.7$ and integrating the NFW profile along the line of sight, we find that 74% of the mass in the cylinder also lies within the sphere with radius R_{200} (and 69% when we make this comparison for R_{500}). We therefore multiply the stellar mass estimates by a factor 0.74 (0.69 for R_{500}).

Since so far we only considered galaxies with stellar masses exceeding the mass completeness limits, we have to estimate the stellar mass contained in lower mass galaxies. We measured the Schechter parameters of the SMF in vdB13, and although these parameters were constrained by galaxies with stellar masses exceeding $10^{10} M_{\odot}$, we use the integral of this Schechter function for masses below the stellar mass completeness limits to correct for these lower mass galaxies. Figure 2 shows that the total stellar mass contained in low-mass galaxies is small. The percentage by which we correct the stellar mass depends on the stellar mass completeness and ranges from 4% for SpARCS-0035 to 25% for SpARCS-0036. Given the size of these correction factors, they do not have a significant effect on the results, especially because the depth in terms of stellar mass is independent of the redshift or halo mass of the clusters. Total stellar masses are listed in Table 3.

In Fig. 5 we show the total stellar mass versus total halo mass compared within R_{200} and R_{500} (left and right panels, respectively) for the GCLASS systems. Error bars in the vertical direction include statistical uncertainties on individual stellar mass measurements, and uncertainties on the estimated probabilities that a photometrically selected galaxy is part of the cluster. The latter uncertainty, which dominates, includes the error on the overdensity parameter for each cluster. The GCLASS data show

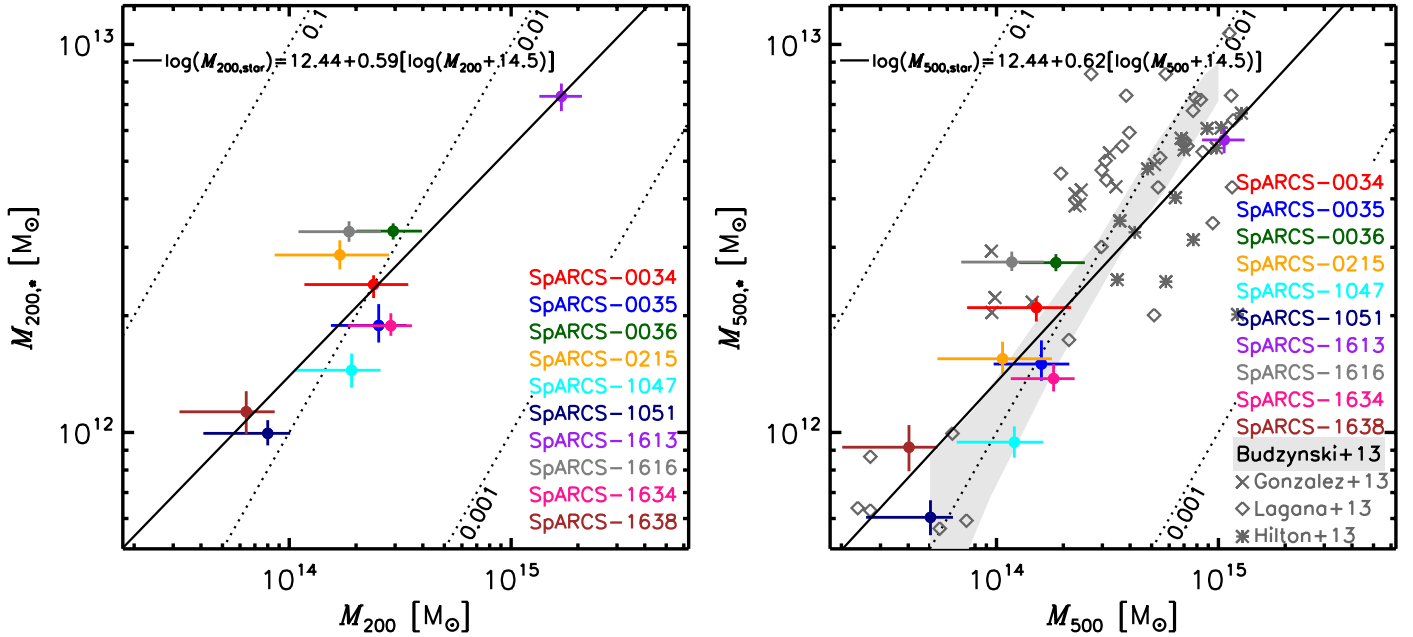


Fig. 5. Total stellar mass versus halo mass within a sphere with radius R_{200} and R_{500} (left and right panels, respectively). Error bars represent uncertainties on individual mass measurements and uncertainties on the membership correction for galaxies we do not have spectra for. Dotted lines show locations with constant stellar mass fractions. The literature measurements (right panel) are measured over a range of redshifts, and are based on different analyses. When possible, the data points are corrected for differences in IMF and M_*/L 's, as explained in the text.

a clear correlation, with spearman coefficient $\rho = 0.65$ (within R_{200}), and $\rho = 0.62$ (within R_{500}).

We fit a power-law relation to the GCLASS data points, with the amount of intrinsic scatter as a free parameter, and described by a log-normal distribution with scatter σ . We find the following best-fitting parameters for the comparison within R_{200} ; $\log(\sigma_{M_{200,*}|M_{200}}) = 0.08^{+0.04}_{-0.05}$ dex, and the relation $\log(M_{200,*}) = (12.44 \pm 0.04) + (0.59 \pm 0.10) \cdot [\log(M_{200}) - 14.5]$. When we perform the fit to the data within R_{500} we find; $\log(\sigma_{M_{500,*}|M_{500}}) = 0.11^{+0.05}_{-0.04}$ dex, and the relation $\log(M_{500,*}) = (12.44^{+0.05}_{-0.06}) + (0.62 \pm 0.12) \cdot [\log(M_{500}) - 14.5]$. Both relations are shown in Fig. 5. The slope of the relation is consistent with the slope found by Lin et al. (2012), who measured it to be 0.71 ± 0.04 for a sample of redshift $z < 0.6$ clusters. The small amount of intrinsic scatter in the relation between total stellar mass and halo mass indicates that stellar mass is a good proxy for total halo mass (albeit with large measurement uncertainties on individual clusters), as was also suggested by Andreon (2012).

For 6 X-ray selected galaxy clusters at $z \sim 1$, Burke et al. (2012) show that the contribution of the ICL to the total J -band flux within R_{500} is about 1–4%. Since this contribution is much (factor ~ 2 –4) smaller than the contribution of the ICL at low- z , our measurements should be close to the actual mass in stars.

Given that this tight relation between total stellar mass and halo mass already exists at $z \sim 1$, and that the stellar mass fraction is decreasing with increasing halo mass, one would naively expect the stellar mass fraction of these massive haloes to increase towards lower redshifts. That is because the likely systems that will be consumed by these haloes are those with a high stellar mass fraction (McGee et al. 2009). In this simple picture the stellar mass fraction would increase, even in the absence of in-situ star formation. Given this naive expectation, it is therefore interesting to make a comparison of the stellar mass content of haloes at lower redshifts.

4.1. Comparison to other samples

We compare our measurements to others in the literature (mostly performed at low- z) for R_{500} , since this radius was used by most studies that estimate the halo masses with X-ray data. However, there are several important caveats to make before we can make a fair comparison. The adopted M_*/L is a major systematic uncertainty in any study and depends on the assumed IMF due to differences in the contribution of low mass stars to the total mass. We transform the results from other studies to the Chabrier IMF by subtracting 0.24 dex in mass for a Salpeter IMF, or adding 0.04 dex to the mass for a Kroupa IMF. The M_*/L depends on galaxy type, but due to the lack of multi-wavelength photometry, it is often assumed that all cluster galaxies are composed of the same stellar population. If one assumes an old stellar population (and therefore a high M_*/L), the mass of the late-type galaxies (and thus the cluster as a whole) is over-estimated. Such an effect will be more pronounced at higher redshift because of the higher number density of late-type galaxies in high- z clusters (M12, vdB13). We point out possible issues for each of the comparison samples below.

An obvious study to compare our results to is based on an SZ-selected cluster sample from the ACT, with a redshift range overlapping with GCLASS and a median redshift of $z = 0.50$ (Hilton et al. 2013). A complication is that Hilton et al. (2013) estimated cluster stellar masses based on the total IRAC 3.6 μm flux measured after a statistical background subtraction. Instead of fitting a M_*/L for each galaxy based on SED modelling, they assume a stellar population that is formed at $z_f = 3$, following a $\tau = 0.1$ Gyr single burst model and the Bruzual & Charlot (2003) stellar population synthesis model. To estimate the effects of these assumptions and see if this creates a bias, we follow the method described by Hilton et al. (2013) to obtain the background subtracted IRAC 3.6 μm flux within R_{500} for the 5 GCLASS clusters for which we have deep IRAC data (vdB13), and estimate the total stellar mass based on the described stellar

population. Table 3 compares these estimates with the total stellar mass in the clusters obtained by the full SED fitting analysis. The approach with a fixed M_*/L over-estimates the stellar mass in all clusters by at least a factor of 2, and this difference seems to be largest for the highest redshift cluster. This is consistent with the notion that the blue fraction, and therefore the fraction of galaxies with relatively low M_*/L , increases with redshift (cf. Butcher & Oemler 1978). It is also possible that the stellar population assumed by Hilton et al. (2013) has a formation redshift ($z_f = 3$) that is too high. After correcting the stellar masses from Hilton et al. (2013) to a Chabrier IMF, we divide them by an additional factor of 2 as an approximate correction for the M_*/L explained above. These data points are overplotted in Fig. 5 (right panel, $\times$ -symbols), and lie around the relation that is the best fit to the GCLASS data. Note that since we used the red-sequence selected GCLASS sample to measure this bias, the real bias might be even larger if the SZ-selected sample contains a lower fraction of quiescent galaxies.

To study a possible evolution in the stellar mass content of clusters we consider Laganá et al. (2013), who measure the stellar mass content in a sample of $z < 0.3$ clusters. Estimates for M_{500} are obtained from X-ray observations. To measure the total stellar mass from the available SDSS data, the galaxy population is separated between early-type and late-type galaxies using the $(u-i)$ colour. Exploiting the M_*/L from Kauffmann et al. (2003) in the i -band for these galaxy types, Laganá et al. (2013) estimate stellar masses. Since the Kauffmann et al. (2003) M_*/L 's are based on the Kroupa IMF, we subtract 0.04 dex to compare their results to ours, and overplot them in Fig. 5 (right panel, $\diamond$ -symbols).

Another nearby cluster sample is the one studied by Gonzalez et al. (2007), which is in the range $0.03 < z < 0.13$, and these measurements are revised in Gonzalez et al. (2013). In these studies, a single M_*/L was used for each galaxy, irrespective of their type. From a dynamical analysis of the SAURON project, they estimate the average M_*/L in the i -band, which they found to be lower than the M_*/L based on an assumed Salpeter IMF. We correct their M_*/L to a Chabrier IMF by subtracting 0.12 dex, and overplot the points from Gonzalez et al. (2013) in Fig. 5 (right panel, $\times$ -symbols). The stellar mass fractions they find are in approximate agreement with the stellar mass fractions of the GCLASS clusters, although they find a somewhat shallower slope of 0.52 ± 0.04 when they fit a relation to only their data set. Given that the fraction of red (with a large M_*/L) galaxies depends on halo mass, it is possible that this slope is biased due to the assumption of a single M_*/L for the sample.

To increase the dynamic range of the comparison samples, we make a comparison to the measurements from Budzynski et al. (2013), who measured the stellar mass fraction across a wide range of masses in the group and cluster regime from the SDSS. Their stacked measurement of over 20 000 optically selected systems at $0.15 < z < 0.4$ is shown by the shaded region in Fig. 5. Since their analysis is very similar to our, we do not have to correct their measurements for differences in e.g. M_*/L . Both the normalisation and their slope of 0.89 ± 0.14 are consistent with the relation we find for GCLASS. When they stack original SDSS images to measure the contribution from the ICL to the stellar mass in their sample, they find a slope that is even steeper.

We note that there are caveats that arise when comparing different cluster samples, as was also pointed out by several other studies (e.g. Leauthaud et al. 2012a; Budzynski et al. 2013). Performing the analysis described by Hilton et al. (2013) on the

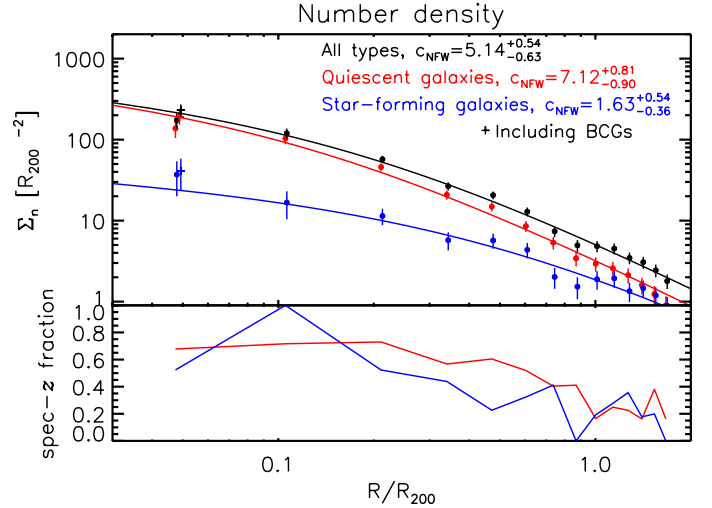


Fig. 6. Number density of galaxies with stellar masses $> 10^{10.2} M_\odot$ in the 10 GCLASS clusters as a function of radial distance. The total galaxy population (black) is separated between star-forming (blue) and quiescent (red) galaxies. Thick points show the membership-corrected number density, where the error bars represent the uncertainties that arise from membership correction. The points are fitted by projected NFW functions (lines), with different concentration parameters. The lower panel shows the fraction of galaxies in each bin with a spectroscopic redshift.

GCLASS data shows that there is a bias in the total stellar mass when a single M_*/L is assumed for all cluster galaxies, especially at high- z . This bias in the stellar mass can be larger than the evolution expected in the redshift range $0 < z < 1$. This shows that it is important to analyse the full SED of each galaxy to estimate its stellar mass. Thanks to the spectroscopic coverage of the GCLASS sample, which is more than 50% complete within R_{200} for the galaxies that dominate the total stellar mass content, membership assignment is relatively straight-forward. In other analyses, where a statistical background subtraction is performed, this can be a major uncertainty for individual systems. We attempted to correct for differences in the analyses between literature studies to be able to compare the total stellar mass fractions between different epochs. Within the uncertainties there seems to be a good agreement between the studies over this redshift range, showing that there is no significant evolution in the stellar mass fraction at fixed halo mass in the redshift range $0 < z < 1$. To tighten the constraints on a possible evolution of this relation, a large and more homogeneous dataset and analysis are required.

5. Radial stellar density distribution

Measurements of the evolution of the spatial galaxy number density and stellar mass density distributions are a key to understand how stellar mass accretes onto massive haloes. We perform these measurements in GCLASS by dividing the sample in radial bins. We do this by stacking the cluster ensemble at the location of the BCGs, and scaling the clusters by their respective R_{200} . We measure the area in each bin by masking the locations on the images that are contaminated by bright stars. Also, since we do not take the stellar mass of the BCGs into account in this study, we mask the location of the central galaxies since this location does not allow for the detection of typical cluster members.

The number density distribution is shown in Fig. 6, where in each radial bin the number of spec- z identified cluster members

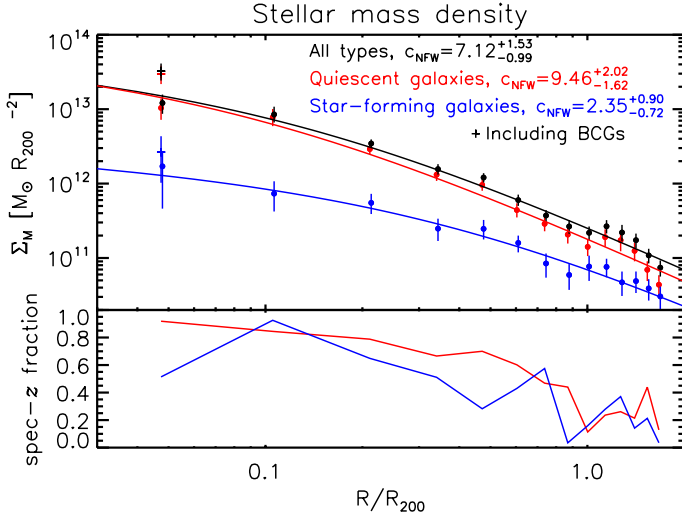


Fig. 7. Stellar mass density distribution in galaxies with stellar masses $> 10^{10.2} M_{\odot}$ of the composite cluster as a function of radial distance. Comparing these distributions to those shown in Fig. 6, we find that the stellar mass distributions are peaked more strongly than the number density distribution. That is an indication for mass segregation of quiescent galaxies in these systems. The *lower panel* shows the fraction of stellar mass in galaxies with a spectroscopic redshift.

Table 4. Best-fitting NFW parameters to the radial density distributions.

	c_{NFW}	$\chi^2/\text{d.o.f.}$
$\Sigma_{\text{M,all}}$	$7.12^{+1.53}_{-0.99}$	0.94
$\Sigma_{\text{n,all}}$	$5.14^{+0.54}_{-0.63}$	0.84
$\Sigma_{\text{M,quiescent}}$	$9.46^{+2.02}_{-1.62}$	1.07
$\Sigma_{\text{n,quiescent}}$	$7.12^{+0.81}_{-0.90}$	0.92
$\Sigma_{\text{M,star-forming}}$	$2.35^{+0.90}_{-0.72}$	0.36
$\Sigma_{\text{n,star-forming}}$	$1.63^{+0.54}_{-0.36}$	0.73

Notes. Reduced χ^2 values are given (14 degrees of freedom).

and the membership-corrected photo- z members with stellar masses exceeding $10^{10.2} M_{\odot}$ are combined. Errors on each point are a combination of Poisson sampling errors, and errors propagated from the membership correction which we estimated from a series of Monte-Carlo simulations. We used the area-weighted position to plot the data points in the horizontal direction. The “+”-signs show the innermost point including the BCGs. The bottom panel of Fig. 6 shows the spectroscopic targeting completeness as a function of radial distance, which shows that – as designed – the completeness is higher for objects near the projected cluster centres. Further away from the cluster centre, the errors that arise from membership estimates are dominant.

The radial distribution of stellar mass in the ensemble cluster is shown in Fig. 7. Besides Poisson counting errors and errors that arise from cluster membership corrections, the error bars include stellar mass measurement errors on individual galaxies. Compared to the number density profile, the spectroscopic targeting completeness is higher due to the selection of spectroscopic targets by their $3.6 \mu\text{m}$ flux.

We fit projected NFW (Navarro et al. 1997) profiles to the data points, excluding the BCGs, to be able to interpret the results in the context of the NFW concentration parameter. Using χ^2 minimization, taking account of the 2D annulus-shaped bins, we find the best fitting functions, which give good representations of the data (see the reduced χ^2 values in Table 4). We give the best-fitting concentration parameters and their marginalized

errors in the table for both the number density and the stellar mass density profiles. The best-fitting profiles are shown in the corresponding figures.

From both the number density and the stellar mass density profiles we find that the quiescent galaxy population is concentrated more strongly than the star-forming population, which is consistent with the view that the star-forming population is accreted more recently by the cluster (for a measurement at low- z , cf. Biviano et al. 2002).

We also find that the stellar-mass distribution of quiescent galaxies is concentrated more strongly than their number density profile, which is an indication that more massive galaxies are situated closer towards the cluster centres than lower mass galaxies. This is likely caused by dynamical friction of the cluster members, which is more efficient for massive galaxies. Note that this effect is observed without taking account of the BCGs.

5.1. Discussion

We measured the galaxy concentration parameters in the ensemble GCLASS cluster, and it may be that a subset of these systems is driving the concentration to this relatively high value. To investigate this we perform different stacks using subsets of the GCLASS sample. We separate the sample in 3 bins, and to make sure the statistics in each bin are sufficiently high, we rank order the clusters by total stellar mass and fill the bins by 6, 3, and 1 cluster(s), respectively. We find that the best-fitting stellar mass concentrations for these 3 ensembles are in the range $6.0 < c < 9.0$, and agree to within 2σ of their measurement errors. This suggests that the stellar mass in each of the GCLASS clusters is likely to be distributed with a concentration parameter around $c \sim 7$.

This high concentration parameter for the stellar matter suggest that the stellar mass is concentrated more strongly than the dark matter is expected to be. For the GCLASS haloes Duffy et al. (2008) estimates a concentration parameter around $c = 2.7$ from simulations that only contain dark matter. Although this value is the median value for massive haloes at $z = 1$, the distribution of concentrations is found to be distributed by a log-normal distribution with a scatter $\sigma(\log(c)) = 0.15$. It is possible that the red-sequence selection method is biased towards systems with highly concentrated red-sequence galaxies. However, given the large difference in concentration between the stellar mass and dark matter, and the relatively small scatter in the distributions, it is unlikely that this difference is merely an effect of the selection method. Note that it is possible that the inclusion of baryonic physics in simulations will alter the dark matter distribution, as recent studies have suggested (e.g. van Daalen et al. 2011). This might bring the dark matter and stellar mass concentrations better in agreement. We checked that the results shown in Fig. 5 are only marginally affected if we change the concentration to $c = 7$.

The composite cluster sample is obtained after stacking the individual clusters on the locations of their BCGs. In some cases the identification of the BCG is ambiguous. For SpARCS-1051 and SpARCS-1634 the identified BCGs are separated by ~ 250 kpc from the approximate projected cluster centres. We test what the effect of possible mis-centring is on the concentration of the measured radial density profiles. We find that, if the intrinsic cluster profiles are described by a $c = 10$ NFW profile, and 10 clusters are stacked with a mis-centring sampled from a Gaussian distribution with $\sigma = 0.1 r_{200}$, the measured concentration would be $c = 7$. Any misalignment with the “true” cluster centre would result in a concentration that is biased

low. Given these tests, it is likely that the stellar mass is concentrated even more strongly than indicated by the NFW fits to the cluster ensemble.

5.2. Evolution towards lower redshift

From numerical simulations (Wechsler et al. 2002) we know that massive haloes are likely to grow by a factor of ~ 2.5 between $z = 1.0$ and $z = 0.3$. This suggests that the GCLASS cluster sample, with typical halo masses of $M_{200} \approx 2 \times 10^{14}$, is the likely progenitor population of the clusters observed in the CNOC survey (Yee et al. 1996; Carlberg et al. 1996), which have typical halo masses of $M_{200} \approx 7 \times 10^{14}$. The concentration of the underlying dark matter distribution is expected to increase by $\sim 10\%$ in this redshift interval (Duffy et al. 2008). Muzzin et al. (2007) measured the K -band luminosity and number density profiles for 15 of the CNOC1 clusters, and showed that the K -band luminosity distribution is well described by a projected NFW profile with concentration parameter $c = 4.28 \pm 0.57$. Although the luminosity in the K -band is a good proxy for the stellar mass, the mass-to-light ratio in this filter depends on galaxy type. Since we find a different distribution of stellar mass in quiescent and star-forming galaxies (Fig. 7), this suggests that the K -band luminosity profile differs from the stellar mass density profile. Indeed, if we scale the star-forming galaxies in GCLASS by a factor of 2 to account for the rough difference in M_*/L , we measure a luminosity profile with a concentration $c < 6$. Although the difference between GCLASS and CNOC1 is thus not as extreme, these results suggest that the dark matter and stellar mass density distributions evolve in distinct ways. This is also suggested by Budzynski et al. (2012), who based their study on a sample of groups and clusters in the redshift range $0.15 < z < 0.4$ from the SDSS. For this sample Budzynski et al. (2012) found that the concentration of the number density profile is lower than the dark matter prediction. There are several caveats, and possible explanations for the observed evolution of the stellar mass distribution.

First, since we do not take account of the stellar mass present in the central galaxies when fitting NFW profiles, accretion of galaxies onto the central galaxy might change the distribution of stellar mass in satellites, and therefore the concentration parameter, over time. Mergers play a dominant role in the build-up of stellar mass in BCGs (Lidman et al. 2013; Burke & Collins 2013). Massive galaxies that are close to the centre are expected to merge with the BCG on a relatively short timescale (Bildfell et al. 2012; Lidman et al. 2013), thereby rendering the BCG an increasingly statistically different population compared to cluster satellite galaxies. An indicator for this process is an increase in the luminosity gap between the BCG and the second BCG (e.g. Smith et al. 2010). However, given the shallow slope of the central stellar-halo mass relation (Sect. 3), BCGs are expected to grow only by a factor of 1.5 in the redshift range $1.0 > z > 0.3$ (see also Lidman et al. 2012). If the supply of this stellar mass growth is obtained from galaxies near the centre, the concentration parameter of the satellite galaxy population would go down.

We perform a simple simulation in which we reduce the stellar mass in satellite galaxies within $0.5 \cdot R_{200}$ in accordance with a BCG growth of a factor of 1.5, and this shows that this is not sufficient to explain the dramatic decrease in the concentration parameter (c decreases from 7.0 to 6.0). Nevertheless, it is possible that the build-up of the ICL component towards lower redshift plays a role in lowering the concentration parameter of stellar mass in satellites.

Second, as clusters get larger, the dynamical friction timescale of a galaxy with a given mass increases, so that it takes longer for galaxies to sink to the centre of the potential well. This is also hinted at when we compare the relation between central stellar and halo mass (Sect. 3), and between total stellar mass and halo mass (Sect. 4). Given that the latter slope is steeper, the fraction of stellar mass in satellite galaxies is higher in more massive haloes. It is possible that galaxies that are accreted onto the cluster at a later time are situated closer to the outskirts of the clusters due to the same process, and thus are less concentrated than the population that was accreted earlier.

We perform a simple test in which we increase the mass of the ensemble cluster by a factor 2.5 by adding stellar mass that is distributed following an NFW distribution with a given concentration. If we vary the concentration parameter of the population that we add, we find that, in order to end up with a concentration of $c = 4.0$ by $z = 0.3$ (i.e., similar to the concentration measured in CNOC), we have to add satellites with a concentration parameter of $c = 2.8$ to the stellar mass density distribution observed in GCLASS. This scenario could potentially explain the difference with the results from Budzynski et al. (2012), who find that at low- z the stellar mass is concentrated more strongly than the dark matter, and suggests that the stellar mass content mostly grows by accreting stellar mass onto the cluster outskirts.

6. Summary and conclusions

In this paper we provide three key measurements concerning the stellar content in 10 clusters at $z \sim 1$ from the GCLASS survey. GCLASS benefits from 11 band photometric coverage and deep spectroscopic coverage to provide a full census of stellar mass in cluster members down to about $M_* = 10^{10.2} M_\odot$. Combining these observations with measurements at lower redshifts we hope to provide constraints on the way baryons cool and form stars in galaxies in high density environments.

In Sect. 3 we presented a comparison of the central stellar mass with total halo mass, and found a correlation that suggests that the fraction of mass in the central galaxy is a decreasing function of halo mass, and about 0.001 for the mass range probed by GCLASS. We confirmed the trend predicted using abundance matching techniques, both in a qualitative as quantitative sense.

Section 4 showed a comparison of the total stellar masses (including satellites) with the dynamical halo masses, both within R_{200} and R_{500} . We found that the total stellar mass increases with halo mass, and that the fraction is around 0.01 for our sample and appears to decrease towards higher halo masses. A comparison of this relation with samples at other redshifts can yield insights on the way these systems accrete their stellar mass, but is difficult due to inhomogeneous sample selections and analyses. Especially inaccurate estimates on the stellar mass-to-light ratio are a source of confusion. After correcting the reference studies for differences in their analyses, we found no significant evolution with redshift in the stellar mass fraction at fixed halo mass.

In Sect. 5 we studied the radial number density and stellar mass density profiles of galaxies in the sample, and found that these are represented by projected NFW profiles. The stellar mass density distribution is concentrated more strongly than the galaxy number density distribution, which shows that more massive galaxies are situated closer to the cluster cores (i.e. mass segregation). The stellar mass density profile has an NFW concentration parameter ($c = 7$) that is significantly higher than the dark matter distribution is expected to be ($c = 2.7$) from numerical simulations. Comparison of the concentration parameter

with the CNOC1 survey at $z = 0.3$ suggests that the stellar mass concentration should decrease towards lower redshift. A simple simulation showed that stellar mass growth of the BCG alone is not enough to explain the evolution between GCLASS and CNOC1, and that the clusters are likely to accrete more stellar mass on the cluster outskirts as they grow by a factor of 2.5 in total mass from $z = 1$ to $z = 0.3$. Also the build-up of the ICL can play a role in the observed evolution.

We note that comparisons of our results with other studies are complicated due to inhomogeneous samples and different analyses. In order to draw firm conclusions regarding the evolution of the baryonic content, and in particular stellar mass, observational data need to be homogenized. We have also seen that the assumption of a single stellar mass-to-light ratio is inadequate to measure the total stellar mass content of galaxy clusters. Rather, one should fit the full Spectral Energy Distributions (SEDs) to estimate the stellar masses for individual galaxies. Moreover, since the galaxies with a high mass-to-light ratio are generally more concentrated in the cluster cores, measurements of the K -band luminosity profile and stellar mass density profile should not be taken as equivalent measurements. Thanks to the advance of large optical and near-infrared imaging facilities over the past decades, these multi-wavelength data are relatively easy to obtain, so that we will soon expect to be able to compare consistent stellar mass measurements with full SED fitting over a large redshift baseline.

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