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Citation

Lepp, S., Dalgarno, A., Dishoeck, E. F. van, & Black, J. H. (1988). Large molecules in diffuse interstellar clouds. Retrieved from <https://hdl.handle.net/1887/1984>

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LARGE MOLECULES IN DIFFUSE INTERSTELLAR CLOUDS

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Received 1987 August 31; accepted 1987 November 25

ABSTRACT

The effects of the presence of a substantial component of large molecules on the chemistry of diffuse molecular clouds are explored, and detailed models of the ζ Persei and ζ Ophiuchi clouds are constructed. The major consequence is a reduction in the abundances of singly charged atomic species. The long-standing discrepancy between cloud densities inferred from rotational and fine-structure level populations and from the ionization balance can be resolved by postulating a fractional abundance of large molecules of 1×10^{-7} for ζ Persei and 6×10^{-7} for ζ Ophiuchi. If the large molecules are polycyclic aromatic hydrocarbons (PAH) containing ~ 50 carbon atoms, they contain 1% of the carbon in ζ Persei and 7% in ζ Ophiuchi. Other consequences of the possible presence of PAH molecules are discussed.

Subject headings: interstellar: abundances — interstellar: molecules — molecular processes

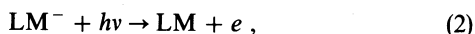
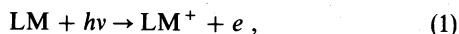
I. INTRODUCTION

Large molecules and in particular polycyclic aromatic hydrocarbons (PAH) have been suggested as the sources of the broad infrared emission features in nebulae (Duley and Williams 1981; Léger and Puget 1984; Allamandola, Tielens, and Barker 1985) and of the diffuse interstellar absorption bands (Léger and d'Hendecourt 1985; van der Zwet and Allamandola 1985; Crawford, Tielens, and Allamandola 1985). If present in dense clouds at the level suggested for nebulae, where the large molecules may contain 2% or 3% of the interstellar carbon, they modify the chemistry of dense clouds substantially (Lepp and Dalgarno 1988). If present in diffuse clouds at a comparable fractional abundance, large molecules will cause changes in the predicted compositions. We demonstrate here that the ratios of successive ionization stages of atomic systems that do not react with hydrogen are sensitive to the abundances of large molecules and that upper limits to and useful estimates of the fractional abundance of large molecules in diffuse clouds can be derived from measurements of the relative abundances of neutral and singly ionized constituents.

Models of the composition of the diffuse molecular clouds lying toward ζ Persei, ζ Ophiuchi, χ Ophiuchi, and σ Persei have been constructed by van Dishoeck and Black (1986). We supplement their chemistry by postulating the presence of a component of large molecules (LM) which we take to have a chemistry like that of the polycyclic aromatic hydrocarbons.

II. CHEMISTRY OF LARGE MOLECULES

We adopt a characteristic ionization potential for LM of 6.8 eV and a characteristic electron affinity of 3 eV (Omont 1986; Duley 1986; d'Hendecourt and Léger 1987). For the photoionization and photodetachment processes



we adopt cross sections of 10^{-16} cm^2 . The resulting unshielded interstellar photoionization and photodetachment rates are $1.5 \times 10^{-8} \text{ s}^{-1}$ and $6 \times 10^{-8} \text{ s}^{-1}$, respectively, for the standard interstellar radiation field (see Draine 1978; van Dishoeck and Black 1986). To calculate the dependence of the rates on depth into a cloud, we follow the procedures of Roberge, Dalgarno, and Flannery (1981) with their grain models 2 and 3. We fitted the depth dependence by a bi-exponential form (see van Dishoeck and Dalgarno 1984).

The distribution of large molecules among LM^+ , LM, and LM^- is determined by the balance between photoionization and photodetachment and recombination



and attachment



We adopted from Omont (1986) the recombination and attachment rate coefficients for molecules containing 50 atoms. The values are $5.7 \times 10^{-6} (100/T)^{1/2} \text{ cm}^3 \text{ s}^{-1}$ and $7.7 \times 10^{-7} \text{ cm}^3 \text{ s}^{-1}$, respectively, where T is the temperature.

Figure 1 shows the calculated abundances of LM^+ , LM, and LM^- with depth into a cloud of uniform density $n_{\text{H}} = 200 \text{ cm}^{-3}$, electron fraction $n_e/n_{\text{H}} = (1-3) \times 10^{-4}$ and temperature 50 K with visual extinction $A_V = 1$ for grain model 2 in the standard interstellar radiation field. In such a cloud, the concentrations of LM and LM^- are comparable and both are significantly larger than the concentration of LM^+ . The total electron density decreases with increasing large molecule abundance. This, in turn, causes the density ratios LM/LM^- and LM^+/LM to increase as illustrated in Figure 1.

The negative ions LM^- react rapidly with the positive ions X^+ in a process of mutual neutralization



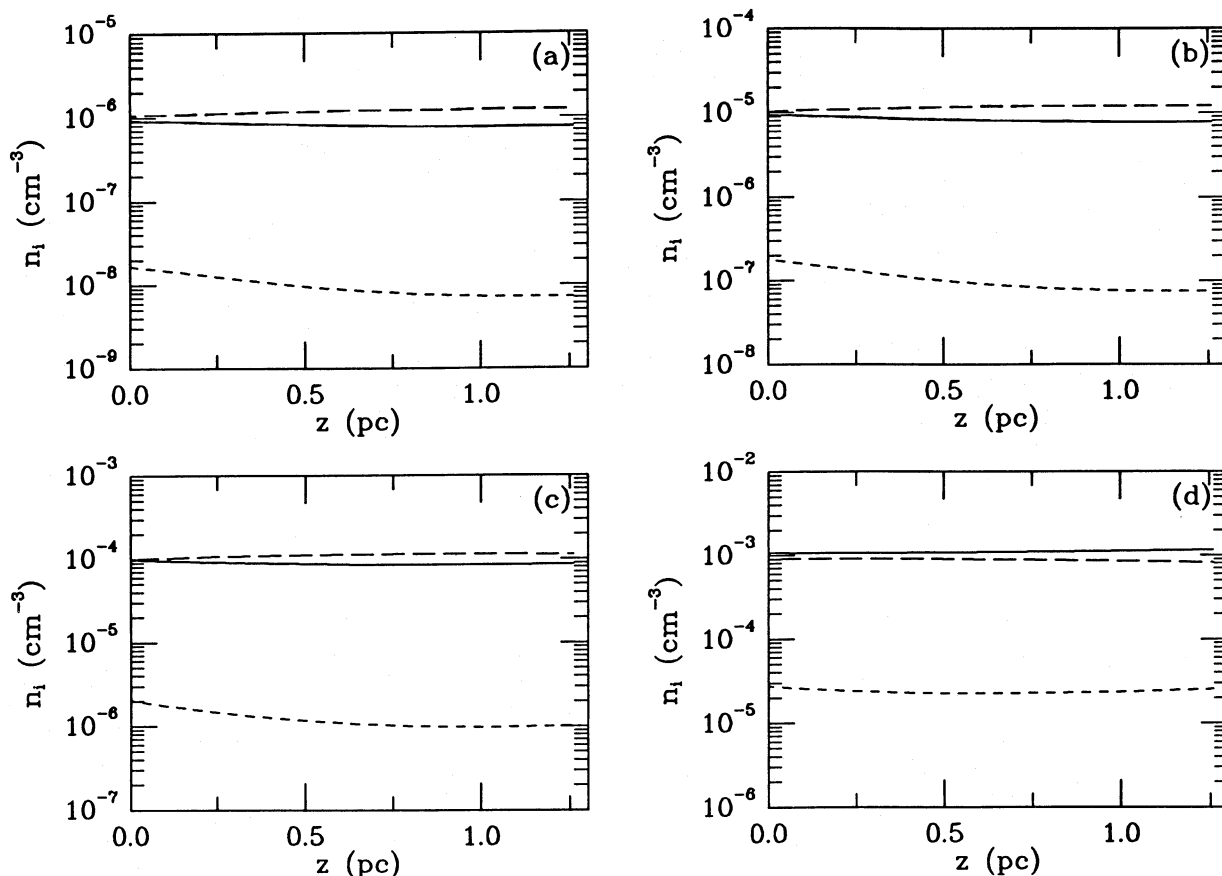


FIG. 1.—Depth variations of the concentrations of LM (solid curves), LM^- (long dashed curves), and LM^+ (short dashed curves) in a reference model of a diffuse cloud. This uniform model has $n_{\text{H}} = 200 \text{ cm}^{-3}$, $T = 50 \text{ K}$, and $A_{\text{V}} = 1.0 \text{ mag}$. Panels (a), (b), (c), and (d) show results for LM abundances of $x_{\text{LM}} = 10^{-8}$, 10^{-7} , 10^{-6} , and 10^{-5} , respectively. When $x_{\text{LM}} \geq 10^{-6}$, the effect of LMs on the overall ionization balance is large enough to produce a noticeable decrease in the electron density. This is accompanied by corresponding increases in the LM/LM^- and LM^+/LM^- density ratios.

(Omont 1986; Lepp and Dalgarno 1987). Positive ions also undergo charge transfer with neutral molecules



if the ionization potential of X is greater than that of LM. If it is less, the reaction



in which a hydrogen atom is extracted from LM is assumed to occur. The molecular ions then undergo dissociative recombination

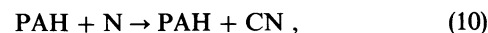
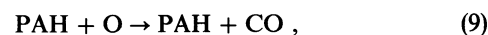


at a rate coefficient assumed to equal $10^{-6} T^{-1/2} \text{ cm}^3 \text{ s}^{-1}$.

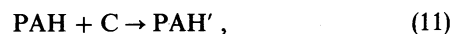
In our further discussions we assume that the large molecules are polycyclic aromatic hydrocarbons (PAH).

The chemistry of PAH molecules is uncertain. We assume that the loss of atoms and molecules from the gas phase by attachment to PAH in any of its forms is small (Omont 1986). Duley and Williams (1986) have argued that reactions of PAH molecules with H, O, and N atoms are rapid and consequently destroy PAH on time scales less than 10^5 yr whereas Omont (1986) suggests that reactions occur more slowly, primarily at radical sites on PAHs from which H atoms have been ejected

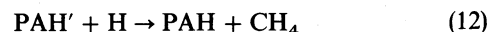
by photodissociation. We write the rate coefficients for



as $10^{-10} f_{\text{R}} \text{ cm}^3 \text{ s}^{-1}$ where f_{R} is an efficiency factor related to the fraction of radical sites on the periphery of the PAH (van Dishoeck and Black 1988). For diffuse clouds $f_{\text{R}} \approx 0.1$. We assume that reaction of atomic carbon with PAH forms a special C site that is not part of the lattice skeleton

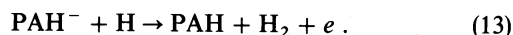


and that subsequent reactions with atomic hydrogen



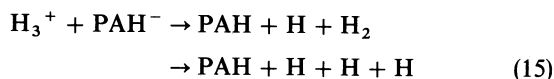
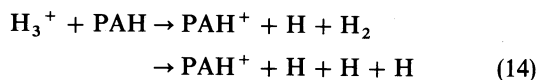
form CH_4 , which can be released by photolysis. The effective rate coefficient for the sequence (11) and (12) is $10^{-10} f_{\text{R}} (1 - f_{\text{R}})/4 \text{ cm}^3 \text{ s}^{-1}$ (van Dishoeck and Black 1988). We assume further that O, N, and C react with PAH^- to release the electron and liberate the saturated molecules H_2O , NH_3 , and CH_4 , respectively, with rate coefficients of $10^{-9} f_{\text{R}} (1 - f_{\text{R}})/m \text{ cm}^3 \text{ s}^{-1}$, where m is the number of hydrogen atoms in the liberated molecule. For reaction with H, we suppose that H_2 is

formed,



The effect of equation (13) may be subsumed into the efficiency factor y_f (van Dishoeck and Black 1986) for the formation of H_2 by recombination on grain surfaces. Note that if each of the reactions (9)–(12) were accompanied by the destruction of one hexagonal ring, the time scale for destruction of a typical 19 ring PAH molecule would be about 10^6 yr in the models. This lifetime is comparable to the longest equilibrium time scale for formation and destruction of H_2 , 5×10^6 yr in the centers of our cloud models; therefore the rates of these reactions represent the maximum rates that are in harmony with the assumption of steady state.

Because the dissociative recombination of H_3^+ and H_2D^+ is slow, we include the reactions



with total rate coefficients of respectively $4.4 \times 10^{-9} \text{ cm}^3 \text{ s}^{-1}$ and $1.3 \times 10^{-7} (10/T)^{1/2} \text{ cm}^3 \text{ s}^{-1}$ (Omout 1986) and equal branching ratios for the various channels. Similar reactions occur for H_2D^+ .

We have incorporated into the chemical models of van Dishoeck and Black (1986) the atoms Li, Na, Mg, Al, Si, P, S, K, Ca, Ti, Cr, Mn, Fe, Co, Ni, Zn, and Rb. They are all ionized by the interstellar radiation field and Ca is partly doubly ionized. The photoionization rates were calculated using the methods of Roberge, Dalgarno, and Flannery (1981). We adopted

recombination rate coefficients similar to those listed by Péquignot and Aldrovandi (1986). Charge transfer reactions of molecular ions with neutral metal atoms have been included (see Oppenheimer and Dalgarno 1974; Prasad and Huntress 1980). Detailed models have been worked out for the clouds toward ζ Persei and ζ Ophiuchi.

III. MODELS

For the ζ Persei cloud, we have explored the effects of large molecules on model F of van Dishoeck and Black (1986) which has a central density of 325 cm^{-3} , a central temperature of 30 K, and a radiation field enhanced by a factor of 3.5 for grain model 2. For the ζ Ophiuchi cloud, we have obtained results for models B and G. Model G has a central density of 350 cm^{-3} , a central temperature of 30 K, and a radiation field enhancement factor of 3.5 with grain model 2. Model B has a steeper density gradient with a central density of 800 cm^{-3} and a central temperature of 20 K.

Figure 2 illustrates the variations of various species with depth into the ζ Persei model for a PAH abundance $x_{\text{PAH}} = n_{\text{PAH}}/n_{\text{H}} = 10^{-7}$, where $n_{\text{PAH}} = n(\text{PAH}) + n(\text{PAH}^-) + n(\text{PAH}^+)$, and $n_{\text{H}} = n(\text{H}) + 2n(\text{H}_2)$. PAH is mostly neutral at the edge of the cloud, whereas in the center of the cloud the PAH^- abundance is comparable to that of PAH. The PAH^+ abundance is an order of magnitude smaller. In comparing Figures 1 and 2, note that the relative amounts of PAH and PAH^- depend on the ratio of ultraviolet intensity to total density and that this ratio has a larger value in the ζ Persei model. Tables 1 and 2 list the resulting column densities for various species in models with PAH abundances ranging from 0 to 10^{-5} . In addition to the PAH reactions some reaction rate coefficients have been updated compared with the model of

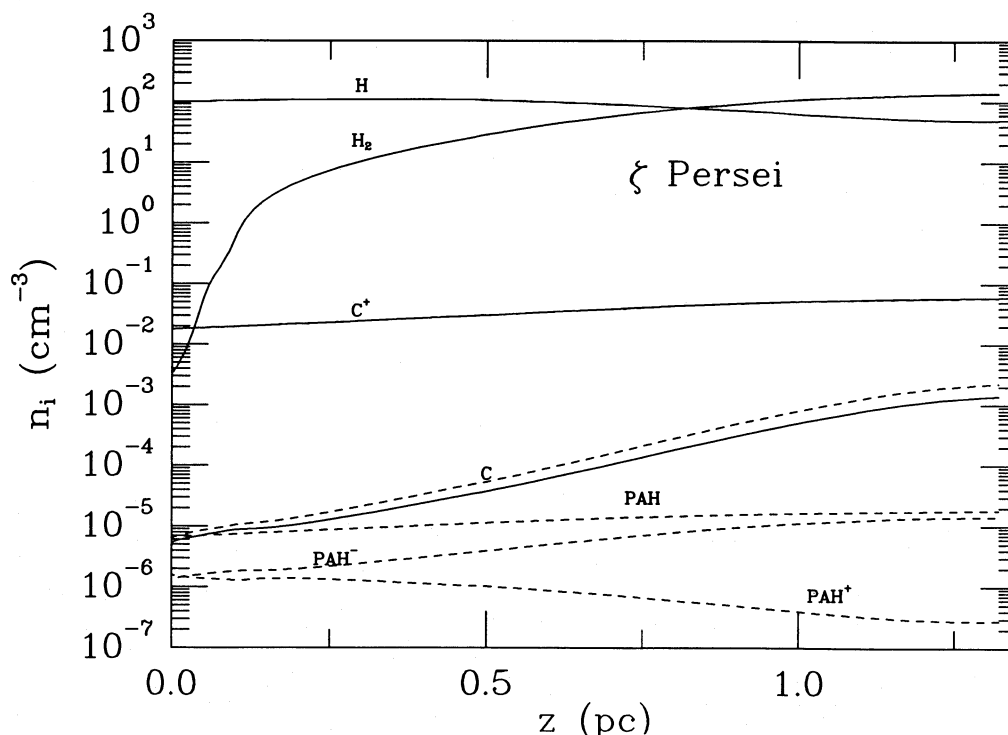


FIG. 2.—Depth variations of the concentrations of several species in model F of the ζ Persei cloud (see text). Solid curves show abundances for $x_{\text{PAH}} = 0$ while dashed curves indicate abundances for $x_{\text{PAH}} = 10^{-7}$. The presence of PAHs at the level of $x_{\text{PAH}} = 10^{-7}$ has no discernible effect on the abundance curves for H, H_2 , and C^+ .

TABLE 1
COMPUTED COLUMN DENSITIES^a IN THE ζ PERSEI AND ζ OPHIUCHI CLOUDS

Cloud	Model	x_{PAH}	δ_c	δ_b	C	C ⁺	CH	OH	HD	H ⁺	H ₃ ⁺	Notes
ζ Per	F	0	0.62	1.10	3.5(15)	4.7(17)	1.7(13)	4.8(13)	3.2(15)	...	2.1(14)	b
		0	0.62	1.10	3.8(15)	4.7(17)	1.5(13)	4.6(13)	5.2(15)	3.9(16)	2.1(14)	c
		1(-11)	0.40	1.10	2.3(15)	3.1(17)	9.9(12)	5.2(13)	5.7(15)	4.1(16)	2.1(14)	c
		1(-8)	0.40	1.10	2.5(15)	3.1(17)	9.9(12)	5.1(13)	5.4(15)	4.0(16)	2.1(14)	c
		3(-8)	0.40	1.10	2.8(15)	3.1(17)	1.1(13)	5.1(13)	5.0(15)	3.7(16)	2.1(14)	c
		5(-8)	0.40	1.10	3.0(15)	3.1(17)	1.1(13)	5.0(13)	4.6(15)	3.5(16)	2.1(14)	c
		1(-7)	0.40	1.10	3.7(15)	3.0(17)	1.2(13)	4.8(13)	3.7(15)	3.1(16)	2.1(14)	c
		1(-6)	0.40	1.10	1.5(16)	2.9(17)	3.2(13)	4.0(13)	3.7(14)	1.1(16)	1.9(14)	c
					(3.3 ± 0.4)(15)	(3 ± 1)(17)	(2.0 ± 0.1)(13)	(4.2 ± 0.5)(13)	(3.8 ± 1.4)(15)			d
ζ Oph	G	0	0.67	0.80	3.3(15)	4.2(17)	1.8(13)	4.5(13)	3.5(14)	...	1.6(14)	b
		0	0.67	0.80	4.3(15)	4.2(17)	1.5(13)	4.2(13)	6.1(14)	2.0(16)	1.5(14)	c
		1(-11)	0.23	0.30	7.3(14)	1.4(17)	5.4(12)	6.6(13)	2.6(14)	3.0(16)	1.6(14)	c
		1(-8)	0.23	0.28	7.8(14)	1.4(17)	5.6(12)	6.6(13)	2.3(14)	2.8(16)	1.6(14)	c
		1(-7)	0.23	0.39	1.2(15)	1.4(17)	6.9(12)	6.2(13)	2.2(14)	2.0(16)	1.6(14)	c
		5(-7)	0.23	0.75	3.2(15)	1.4(17)	1.3(13)	5.5(13)	2.1(14)	8.8(15)	1.5(14)	c
		6.2(-7)	0.21	0.87	3.2(15)	1.3(17)	1.3(13)	5.5(13)	2.2(14)	8.0(15)	1.5(14)	c
		1(-6)	0.23	1.00	5.5(15)	1.4(17)	2.0(13)	5.2(13)	1.8(14)	5.4(15)	1.5(14)	c
		1(-5)	0.23	3.21	3.0(16)	1.1(17)	1.1(14)	6.8(13)	1.6(14)	7.4(14)	1.0(14)	c
		0	0.70	0.80	3.1(15)	4.5(17)	2.1(13)	4.5(13)	2.9(14)	...	...	b
ζ Oph	B	1(-11)	0.22	0.26	6.7(14)	1.4(17)	5.7(12)	6.7(13)	2.1(14)	2.7(16)	1.3(14)	c
		1(-7)	0.22	0.39	1.3(15)	1.4(17)	6.8(12)	6.2(13)	1.9(14)	1.7(16)	1.3(14)	c
		1(-6)	0.22	1.00	6.6(15)	1.3(17)	1.7(13)	5.1(13)	1.6(14)	5.2(15)	1.2(14)	c
					(3.2 ± 0.6)(15)	(1 ± 0.5)(17)	(2.5 ± 0.1)(13)	(4.8 ± 0.5)(13)	(2.1 ± 0.2)(14)			d
Observed												

^a In cm⁻².

^b Cf. van Dishoeck and Black 1986. See also that paper for other parameters characterizing the models.

^c Present models including updated reaction rates.

^d See van Dishoeck and Black 1986 for references to the observations.

TABLE 2
COLUMN DENSITIES^a AND RATIOS OF COLUMN DENSITIES IN THE ζ PERSEI AND ζ OPHIUCHI CLOUDS

Cloud	Model	x_{PAH}	$N(\text{Si})$	$N(\text{Al})$	$N(\text{Zn})$	$N(\text{C}^+)/N(\text{C})$	$N(\text{Mg}^+)/N(\text{Mg})$	$N(\text{S}^+)/N(\text{S})$	$N(\text{Ca}^+)/N(\text{Ca})$	$N(\text{Fe}^+)/N(\text{Fe})$
ζ Per	F	1(-11)	1.1(13)	3.9(9)	3.6(10)	130	229	457	556	624
		1(-8)	1.1(13)	4.0(9)	4.0(10)	123	211	434	509	572
		3(-8)	1.3(13)	4.3(9)	4.7(10)	111	182	393	436	493
		5(-8)	1.4(13)	4.6(9)	5.4(10)	100	160	360	381	432
		1(-7)	1.8(13)	5.3(9)	7.2(10)	81.6	123	296	290	330
		1(-6)	8.0(13)	1.8(10)	3.9(11)	19.6	24.1	71.7	58.7	64.3
		Observed ^b				91	290	490	290	860
						63-135	66-1700	350-710	180-440	91- ∞
ζ Oph	G	1(-11)	1.9(11)	2.2(9)	2.0(10)	196	277	683	849	814
		1(-8)	2.0(11)	2.3(9)	2.3(10)	183	249	639	760	729
		1(-7)	3.5(11)	3.2(9)	4.7(10)	115	131	398	393	386
		5(-7)	1.0(12)	7.4(9)	1.5(11)	43.6	42.1	148	129	126
		6.2(-7)	1.1(12)	8.2(9)	1.8(11)	39.1	36.1	131	114	110
		1(-6)	1.8(12)	1.3(10)	2.8(11)	25.0	23.0	83.8	72.8	69.2
		1(-5)	1.4(13)	9.2(10)	2.0(12)	3.8	2.8	11.0	10.3	8.9
ζ Oph	B	1(-11)	2.1(11)	3.5(9)	2.0(10)	211	225	715	514	700
		1(-7)	4.6(11)	5.8(9)	5.4(10)	106	90.9	366	199	281
		1(-6)	2.7(12)	2.6(10)	3.4(11)	20.3	14.6	68.3	33.9	45.4
		Observed ^b	<3.6(12)	<1.6(10)	<1.2(12)	29	48	80	108	162
						13-260	17-67	52-124	98-118	40-280

NOTE.—In each case the heavy element abundances have been chosen to provide the best fit to observed column densities.

^a In cm^{-2} .

^b See van Dishoeck and Black (1986) and Black (1987) for references to the observations.

van Dishoeck and Black (1986), including an enhanced CH photoionization rate and an improved treatment of the CO photodissociation rate (van Dishoeck and Black 1988). As the tables show for $x_{\text{PAH}} = 0$, these changes have little effect on the abundances of the species of interest here.

The atomic column densities and ratios presented in Tables 1 and 2 are the best available values from the literature. In the cases of Mg II, Fe I, Si I, and S I, published column densities have been adjusted for improved oscillator strengths. Table 2 lists for the ratios of column densities both the value and the extreme range of values permitted by the published uncertainties. The full effects of unresolved velocity structure on curve of growth analyses remain to be explored.

The most sensitive probes of PAH chemistry in diffuse clouds appear to be the abundance ratios of neutral atoms and their first ions. Atomic ions that recombine slowly with electrons and react inefficiently with H_2 can be removed more rapidly by PAHs when $x_{\text{PAH}} \geq 10^{-7}$. Although the C^+ column density is poorly determined from observations of highly saturated lines, it appears from Tables 1 and 2 that the neutral fraction of carbon is in better agreement with observation if PAH reactions (or some other agent of enhanced recombination) are at work in the ζ Persei and ζ Ophiuchi clouds. This suggestion is strengthened by the comparison of models with upper limits on neutral atomic species for which the corresponding ions are well-measured (Si, Al, Zn) and with measured ratios of column densities for neutral atom–first ion pairs (C, Mg, S, Ca, and Fe). If PAHs have the chemical behavior we have assumed, then PAH abundances larger than $\sim 5 \times 10^{-7}$ and $\sim 2\text{--}3 \times 10^{-6}$ are excluded by the data for the ζ Persei and ζ Ophiuchi clouds, respectively. The atomic data for the ζ Ophiuchi cloud are inconsistent with the models if $x_{\text{PAH}} < 10^{-7}$. For the ζ Persei cloud only the Ca^+/Ca ratio is sufficiently well measured to exclude the possibility that $x_{\text{PAH}} \approx 0$. Within the framework of the present models all of the data on atomic ionization in these clouds are reproduced well only if an additional recombining agent like the PAHs is

acting to enhance the neutral fractions. Specifically, the systematic trends in atomic abundances shown in Table 2 are in harmony with $x_{\text{PAH}} \approx 10^{-7}$ for the ζ Persei cloud and $x_{\text{PAH}} \approx 6 \times 10^{-7}$ for the ζ Ophiuchi cloud. Note that in all cases the heavy element abundances have been chosen to provide the best fit to the observed column densities. In particular, the gas-phase carbon abundance is reduced compared with the models of van Dishoeck and Black (1986), if PAHs are included. An uncertainty of the order of a factor of 2 applies to the inferred PAHs abundances. This analysis can be viewed alternatively not as a demonstration of the existence of LMs or PAHs in diffuse clouds but as a restatement of a long-standing issue concerning diffuse clouds: that diagnostic probes like rotational excitation in molecules and fine-structure excitation in atoms indicate lower densities than do the atomic ionization equilibria (cf. van Dishoeck and Black 1986).

For the generic 50-atom PAH assumed here, the inferred abundance in the ζ Persei cloud would account for 1% of the reference abundance of carbon in all forms ($[\text{C}/\text{H}] = 4.7 \times 10^{-4}$ in the Sun). The derived depletion factor $\delta_{\text{C}} = 0.40$ for the remainder of the gas-phase carbon still allows for almost 60% of the carbon to reside in grains. For the ζ Ophiuchi cloud 7% of the carbon is in PAHs, 21% in other gaseous forms, and 72% in grains.

Similar calculations can be performed for the σ Persei and χ Ophiuchi clouds. For the σ Persei cloud, a PAH abundance of a few times 10^{-7} would improve the comparison between the observational data on atomic species and the models. For the χ Ophiuchi cloud, the measurements are unfortunately too uncertain to make any inferences.

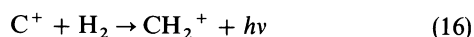
Our discussion of the ionization balance applies not only to PAHs, but to any other large molecule or small grain which is capable of neutralizing ions. The rate coefficients for charge exchange of ions with small grains of arbitrary shape are of the same order of magnitude as those adopted in this work with PAHs (Draine and Sutin 1987).

Some caution is appropriate. The models described here

were constructed originally to reproduce as well as possible the available data on the abundance and excitation of H_2 , the abundances of other simple molecules, and the excitation of C_2 and of atomic fine structure levels. Apart from the PAH reactions, other factors which influence the ionization balance are the temperature and density structure in the cloud, and the strength of the incident radiation field. Although these parameters have been constrained to reproduce a wide variety of observations, they are uncertain by a factor of 2. The observed large molecular abundances suggest that it is unlikely that the density is much lower than the adopted values but it is possible that the density has been underestimated by a factor of 2 and the intensity of the radiation field somewhat overestimated. If so, the ionization balance would shift in favor of the neutral species, and the inferred upper limit on the PAH abundance would be lowered.

In order to examine the model-dependent effects of different density and temperature profiles, results are presented in Tables 1 and 2 for model B of the ζ Ophiuchi cloud, which has the highest central density of the models constructed by van Dishoeck and Black (1986). Within the factors of 2 uncertainty, the same LM/PAH abundance is inferred from the atomic ionization, as expected since all of the models were designed to have the same column-averaged ratios of ultraviolet absorption rates to gas density.

The PAHs have small effects on the abundances of observable molecules like CH, OH, and HD as long as $x_{\text{PAH}} \leq 10^{-7}$. For higher PAH abundances, there is a competition between the reduction in the concentrations of major ions that initiate molecule formation by reactions with PAH^- and the enhancement in the production of H_2O and CH_4 by the assumed reactions of C and O with PAH^- which lead to larger production rates of OH and CH by photodissociation. The predicted column densities of other molecules, such as H_2O and CN, approach or exceed the observed values only when $x_{\text{PAH}} > 10^{-6}$ for the ζ Ophiuchi cloud. When the depletion factor δ_C is lowered for $x_{\text{PAH}} \geq 10^{-7}$ in order to reproduce the column density of atomic carbon, the abundances of most carbon-bearing molecules are reduced as a result. A particularly interesting case is provided by the CH molecule. If the effects of PAHs on the ionization balance are taken into account but the reactions (11)–(12) of C and H with PAH and PAH^- to form CH_4 are not included in the models, then the CH column density falls below the observed values by factors of 2–4. The discrepancy could be remedied with a corresponding increase in the rate coefficient of radiative association



above the adopted value of $7 \times 10^{-16} \text{ cm}^3 \text{ s}^{-1}$; however, such a rate would begin to conflict with the upper limit of $1.5 \times 10^{-15} \text{ cm}^3 \text{ s}^{-1}$ (Luine and Dunn 1985) measured at 10 K. On the other hand, with the inclusion of the reactions with PAH and PAH^- , the production of CH_4 is large enough that its subsequent photodissociation provides a major source of CH and the CH abundance remains consistent with observations at $x_{\text{PAH}} \approx 10^{-7}$ – 10^{-6} without any other compensating changes in the carbon chemistry.

The abundances of oxygen-bearing species such as OH are not much affected by the presence of PAHs in the cloud, as Table 1 illustrates. Although the H^+ column density is reduced considerably by charge transfer for $x_{\text{PAH}} \geq 10^{-7}$, the H_3^+ column density remains nearly constant. The abundance of H_3^+ is governed primarily by cosmic-ray ionizations and reactions with oxygen-bearing neutral species. The principal initi-

ating reaction in the oxygen chain is the $O + H_3^+$ reaction to form OH^+ and the abundances of product molecules, including OH and CO, are not much changed.

In contrast, the HD column density decreases strongly with increasing PAH abundance mainly because of the large reduction of both the H^+ and D^+ concentrations through charge-transfer with PAHs, which reduces the rates of formation of HD by the $H^+ + D \rightarrow H + D^+$ reaction followed by $D^+ + H_2 \rightarrow HD + H^+$. The entries for ζ Persei in Table 1 illustrate the effect of PAH reactions on the HD column density for a fixed deuterium abundance ratio, $[D/H] = 1.5 \times 10^{-5} \delta_D$. The entries for ζ Ophiuchi show how the abundance must be increased with increasing x_{PAH} in order to reproduce the observed column density of HD. The analysis of the HD abundance involves two disposable parameters, δ_D and the cosmic-ray ionization rate. The latter is also constrained by the abundance of OH, for which the chemistry is rather well understood. For any deuterium abundance $\delta_D \leq 1.0$, the results of Table 1 and the observed HD column densities imply $x_{\text{PAH}} < 6 \times 10^{-7}$ for ζ Persei and $x_{\text{PAH}} < 10^{-6}$ for ζ Ophiuchi. These limits are virtually the same as those provided by the atomic data. Aside from the remaining uncertainties in the inferred cosmic-ray ionization rates, the implied deuterium abundances agree within 25% for $x_{\text{PAH}} = 10^{-7}$ in ζ Persei and $x_{\text{PAH}} = 6.2 \times 10^{-7}$ in ζ Ophiuchi. The HD abundance could also be increased in the models by lowering the line widths in the calculation of the HD photodissociation rates, but the HD column densities inferred from the measured equivalent widths would be increased and the discrepancy would persist.

The postulated reactions of atomic species with PAHs to form molecules such as CH_4 and NH_3 open a pathway for the synthesis of complex molecules in diffuse clouds. In particular, the predicted abundances of related species such as H_2CO , CH_3 , and C_2H_2 are enhanced by several orders of magnitude if PAHs are included. In somewhat thicker and denser clouds than those considered here, the predicted amounts of H_2CO may be observable.

Regardless of the interpretation of the ad hoc chemistry introduced here in the guise of large molecules, the model analysis provides an internally consistent set of semiempirical corrections for unseen stages of ionization in diffuse clouds. This is crucial for the determination of gas-phase abundances from observations of a rare form of an element such as Li, Na, K, or Rb that is expected to exist predominantly in an unobservable stage of ionization. In Table 3 the element abundances deduced from the best available observations and the best-fitting ionization models are collected. The results in Table 3 for depletions are in qualitative agreement with values elsewhere in the literature. The inferred depletions are more severe for ζ Ophiuchi than for ζ Persei, with the exceptions of Ti and Mn. We stress that detailed ionization models are required to correct element abundances for unseen stages of ionization. Empirical corrections are often used in the literature; for example, the observed ratio of calcium column densities, $N(\text{Ca}^+)/N(\text{Ca})$, is taken to characterize the ratio of the rates of ionization Γ and recombination α

$$N(\text{Ca}^+)/N(\text{Ca}) \approx n(\text{Ca}^+)n(\text{Ca}) = \Gamma(\text{Ca})/[\alpha(\text{Ca})n(e)] \quad (17)$$

With adjustments for different photoionization cross sections and recombination rate coefficients, this ratio is applied to the ionization balance in lithium or sodium, for example, for which only the neutral species are observed. This procedure is strictly applicable only if the ratios of local densities do not vary significantly with depth. Our detailed models for ζ Persei show

TABLE 3
INFERRED GAS-PHASE ABUNDANCES

ELEMENT X	ζ OPH		ζ PER	
	n_X/n_H	δ_X	n_X/n_H	δ_X
PAH	6 (-7)		1(-7)	
Li	3.5(-10)	1.6(-1)	6.2(-10)	2.8(-1)
C	9.7(-5)	2.1(-1)	1.9(-4)	4.0(-1)
Na	7.6(-7)	3.6(-1)	2.1(-6)	1.0
Mg	5.0(-6)	1.2(-1)	5.5(-6)	1.2(-1)
Al	7.6(-9)	2.5(-3)	...	...
Si	8.2(-7)	1.9(-2)	2.6(-5)	6.1(-1)
P	2.2(-8)	7.8(-2)	3.8(-4)	1.4(-1)
S	8.3(-6)	4.9(-1)	1.2(-5)	6.9(-1)
K	1.0(-8)	7.8(-2)	3.7(-8)	2.9(-1)
Ca	6.0(-10)	2.7(-4)	8.5(-10)	3.9(-4)
Ti	2.3(-10)	1.9(-3)	5.6(-11)	4.7(-4)
Mn	1.5(-8)	5.9(-2)	8.6(-9)	3.3(-2)
Fe	3.2(-8)	7.5(-4) ^a	3.6(-7)	8.3(-3)
Ni	9.9(-9)	5.8(-3)	...	...
Zn	1.4(-8)	3.6(-1)	...	...
Rb	<7.2(-11)	<1.8(-1)	<7.5(-11)	<1.9(-1)

^a Fe abundance adjusted to Fe I column density.

that the local ratio of concentrations, $n(\text{Li})/n(\text{Na})$, increases by a factor of 2.6 from the edge to the center of the cloud, while the density ratio, $n(\text{Ca}^+)/n(\text{Ca})$, decreases from 4800 at the edge to 160 at the center. This example of depth dependent ionization suggests that empirical ionization corrections based on ratios of column densities can be inaccurate. Based on the present ionization models of ζ Persei and ζ Ophiuchi, the abundance of interstellar lithium in the gas phase is found to be 3–5 times larger than the primordial abundance $[\text{Li}/\text{H}] = (1.1 \pm 0.4) \times 10^{-10}$ derived by Spite and Spite (1982). The depletion factors for lithium in Table 3 refer to the abundance, $[\text{Li}/\text{H}] = 2.2 \times 10^{-9}$, in carbonaceous chondrites.

This work has been supported by grant AST-86-17675 to Harvard University by the National Science Foundation, Division of Astronomical Sciences, by the NASA Astrophysical Theory Program through grant NAGW-763 to the University of Arizona, and by the National Science Foundation through grant RII-86-20342 to Princeton University. E. F. v D. and J. H. B. acknowledge the hospitality of the Institute for Advanced Study, Princeton, where part of this work was carried out, and its support through NSF grant PHY 86-20266. They also enjoyed stimulating discussions with B. T. Draine.

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