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Scattering theory of plasmon-assisted entanglement transfer and distillation

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We analyze the quantum-mechanical limits to the plasmon-assisted entanglement transfer observed by Altewischer, van Exter, and Woerdman [Nature **418**, 304 (2002)]. The maximal violation S of Bell's inequality at the photodetectors behind two linear media (such as the perforated metal films in the experiment) can be described by two ratios τ_1 , τ_2 of polarization-dependent transmission probabilities. A fully entangled incident state is transferred without degradation for $\tau_1 = \tau_2$, but a relatively large mismatch of τ_1 and τ_2 can be tolerated with a small reduction of S . We predict that fully entangled Bell pairs can be distilled out of partially entangled radiation if τ_1 and τ_2 satisfy a pair of inequalities.

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The motivation for this work came from the recent remarkable demonstration by Altewischer, van Exter, and Woerdman of the transfer of quantum-mechanical entanglement from photons to surface plasmons and back to photons [1]. Since entanglement is a highly fragile property of a two-photon state, it came as a surprise that this property could survive, with little degradation, the conversion to and from the macroscopic degrees of freedom in a metal [2].

We present a quantitative description of the finding of Ref. [1] that the entanglement is lost if it is measured during transfer, that is to say, if the medium through which the pair of polarization-entangled photons is passed acts as a "which-way" detector for polarization. Our analysis explains why a few percent degradation of entanglement could be realized without requiring a highly symmetric medium. We predict that the experimental setup of Ref. [1] could be used to "distill" [3,4] fully entangled Bell pairs out of partially entangled incident radiation, and we identify the region in parameter space where this distillation is possible.

We assume that the medium is *linear*, so that its effect on the radiation can be described by a scattering matrix. The assumption of linearity of the interaction of radiation with surface plasmons is central to the literature on this topic [5–9]. We will not make any specific assumptions on the mode and frequency dependence of the scattering matrix, but extract the smallest number of independently measurable parameters needed to describe the experiment. By concentrating on model-independent results we can isolate the fundamental quantum-mechanical limitations on the entanglement transfer, from the limitations specific to any particular transfer mechanism.

The system considered is shown schematically in Fig. 1. Polarization-entangled radiation is scattered by two objects and detected by a pair of detectors behind the objects in the far field. The objects used in Ref. [1] are metal films perforated by a square array of subwavelength holes. The transmission amplitude $t_{\sigma\sigma' i}$ of object $i=1,2$ relates the transmitted radiation (with polarization $\sigma=H,V$) to the incident radiation (polarization $\sigma'=H,V$). We assume a single-mode incident beam and a single-mode detector (smaller than the coherence area) so that we require a set of eight transmission amplitudes $t_{\sigma\sigma' i}$ out of the entire scattering matrix (which also contains reflection amplitudes and transmission ampli-

tudes to other modes). The extension to a multimode theory (needed to describe some aspects of the experiment [1]) is left for a future investigation [10]. We do not require that the scattering matrix be unitary, so our results remain valid if the objects absorb part of the incident radiation.

The radiation incident on the two objects is in a known, partially entangled state and we wish to determine the degree of entanglement of the detected radiation. It is convenient to use a matrix notation. The incident two-photon state has the general form

$$|\Psi_{in}\rangle = a_{HH}^{in}|HH\rangle + a_{HV}^{in}|HV\rangle + a_{VH}^{in}|VH\rangle + a_{VV}^{in}|VV\rangle \quad (1)$$

The four complex numbers $a_{\sigma\sigma'}^{in}$ form a matrix

$$A_m = \begin{pmatrix} a_{HH}^{in} & a_{HV}^{in} \\ a_{VH}^{in} & a_{VV}^{in} \end{pmatrix} \quad (2)$$

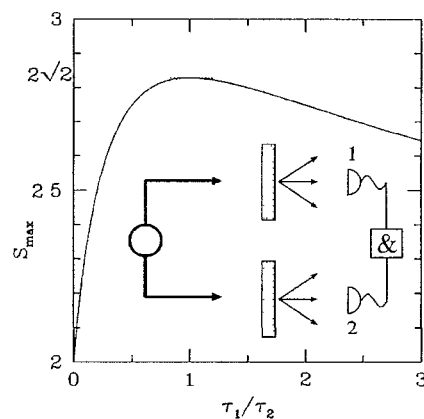


FIG. 1. Main plot: efficiency of the entanglement transfer for a fully entangled incident state as given by Eq. (14). The maximal violation S_{max} of Bell's inequality at the photodetectors is plotted as a function of the ratio $\tau_1/\tau_2 = T_{11} \Gamma_2 / T_{12} \Gamma_1$ of the polarization-dependent transmission probabilities. The inset shows schematically the geometry of the experiment of Ref. [1]. A pair of polarization-entangled photons is incident from the left on two perforated metal films. The photodetectors at the right, connected by a coincidence counter, measure the degree of entanglement of the transmitted radiation.

Normalization of $|\Psi_{\text{in}}\rangle$ requires $\text{Tr} A_{\text{in}} A_{\text{in}}^\dagger = 1$, with ‘‘Tr’’ being the trace of a matrix

The four transmission amplitudes $t_{\sigma\sigma' i}$ of object $i=1,2$ form the matrix

$$T_i = \begin{pmatrix} t_{HH i} & t_{HV i} \\ t_{VH i} & t_{VV i} \end{pmatrix} \quad (3)$$

The transmitted two-photon state $|\Psi_{\text{out}}\rangle$ has matrix of coefficients

$$A_{\text{out}} = Z^{-1/2} T_1 A_{\text{in}} T_2^\dagger, \quad (4)$$

with normalization factor

$$Z = \text{Tr} (T_1 A_{\text{in}} T_2^\dagger) (T_1 A_{\text{in}} T_2^\dagger)^\dagger \quad (5)$$

(The superscript ‘‘t’’ denotes the transpose of a matrix)

We quantify the degree of entanglement in terms of the Clauser-Horne-Shimony-Holt parameter S [11], which measures the maximum violation of Bell’s inequality and was used in the experiment of Ref [1]. This parameter can be obtained from a decomposition of $|\Psi\rangle$ into a superposition of a fully entangled state (with weight \sqrt{P}) and a factorized state orthogonal to it [12,13]. The relation is

$$S = 2\sqrt{1+P^2}, \quad P^2 = 4\text{Det} A A^\dagger, \quad (6)$$

with ‘‘Det’’ being the determinant and $0 \leq P \leq 1$ (The concurrence [14] is identical to P). A fully entangled state has $P=1$, $S=2\sqrt{2}$, while a factorized state has $P=0$, $S=2$. The fully entangled state could be the Bell pair $(|HV\rangle - |VH\rangle)/\sqrt{2}$, or any state derived from it by a local unitary transformation ($A \rightarrow UAV$ with U, V arbitrary unitary matrices). The degree of entanglement $P_{\text{in}} = 2|\text{Det} A_{\text{in}}|$ of the incident state is given and we seek the degree of entanglement $P_{\text{out}} = 2|\text{Det} A_{\text{out}}|$ of the transmitted state. We are particularly interested in the largest P_{out} that can be reached by applying local unitary transformations to the incident state. This would correspond to the experimental situation that the polarizations of the two incoming photons are rotated independently, in order to maximize the violation of Bell’s inequality of the detected photon pair.

Before proceeding with the calculation we introduce some parametrizations. The Hermitian matrix product $T_i T_i^\dagger$ has the eigenvalue-eigenvector decomposition

$$T_1 T_1^\dagger = U \begin{pmatrix} T_{1+} & 0 \\ 0 & T_{1-} \end{pmatrix} U, \quad T_2 T_2^\dagger = V \begin{pmatrix} T_{2+} & 0 \\ 0 & T_{2-} \end{pmatrix} V \quad (7)$$

The matrices of eigenvectors U, V are unitary and the transmission eigenvalues $T_{i\pm}$ are real numbers between 0 and 1. We order them such that $0 \leq T_{i-} \leq T_{i+} \leq 1$ for each $i=1,2$. We will see that the maximal entanglement transfer depends only on the ratios $\tau_i = T_{i+}/T_{i-}$. This parametrization therefore extracts the two significant real numbers τ_1, τ_2 out of eight complex transmission amplitudes. The Hermitian matrix product $A_{\text{in}} A_{\text{in}}^\dagger$ has eigenvalues $\lambda_{\pm} = \frac{1}{2} \pm \frac{1}{2}(1 - P_{\text{in}}^2)^{1/2}$. These appear in the polar decomposition

$$U A_{\text{in}} V = e^{i\phi} \begin{pmatrix} u_+ & u_- \\ -u_- & u_+ \end{pmatrix} \begin{pmatrix} \sqrt{\lambda_+} & 0 \\ 0 & \sqrt{\lambda_-} \end{pmatrix} \begin{pmatrix} v_+ & v_- \\ -v_- & v_+ \end{pmatrix} \quad (8)$$

The phase ϕ is real and u_{\pm}, v_{\pm} are complex numbers constrained by $|u_{\pm}| = (\frac{1}{2} \pm u)^{1/2}$, $|v_{\pm}| = (\frac{1}{2} \pm v)^{1/2}$, with real $u, v \in (-\frac{1}{2}, \frac{1}{2})$. These numbers can be varied by local unitary transformations, so later on we will want to choose values which maximize the detected entanglement.

With these parametrizations a calculation of the determinant of A_{out} leads to the following relation between P_{in} and P_{out}

$$P_{\text{out}} = \frac{P_{\text{in}} \sqrt{\tau_1 \tau_2}}{(\tau_1 - 1)(\tau_2 - 1)} \left[\lambda_+ Q_+ + \lambda_- Q_- - 2\sqrt{\lambda_+ \lambda_-} \right. \\ \left. \times \left(\frac{1}{4} - u^2 \right)^{1/2} \left(\frac{1}{4} - v^2 \right)^{1/2} \cos \Phi \right]^{-1}, \quad (9)$$

$$Q_{\pm} = \left(u \pm \frac{1}{2} \frac{\tau_1 + 1}{\tau_1 - 1} \right) \left(v \pm \frac{1}{2} \frac{\tau_2 + 1}{\tau_2 - 1} \right) \quad (10)$$

The phase Φ equals the argument of $u_+ u_-^\dagger v_+ v_-$. To maximize P_{out} we should choose $\Phi = 0$.

We first analyze this expression for the case of a fully entangled incident state, as in the experiment of Ref [1]. For $P_{\text{in}} = 1$ one has $\lambda_+ = \lambda_- = 1/2$, and Eq (9) simplifies to

$$P_{\text{out}} = \frac{4\sqrt{\tau_1 \tau_2}}{(\tau_1 + 1)(\tau_2 + 1) + 4a(\tau_1 - 1)(\tau_2 - 1)}, \quad (11)$$

$$a = uv - \left(\frac{1}{4} - u^2 \right)^{1/2} \left(\frac{1}{4} - v^2 \right)^{1/2} \cos \Phi \quad (12)$$

Since $\tau_i \geq 1$ and $|a| \leq \frac{1}{4}$ we conclude that the degree of entanglement is bounded by $P_{\text{min}} \leq P_{\text{out}} \leq P_{\text{max}}$, with

$$P_{\text{min}} = \frac{2\sqrt{\tau_1 \tau_2}}{1 + \tau_1 \tau_2}, \quad P_{\text{max}} = \frac{2\sqrt{\tau_1 / \tau_2}}{1 + \tau_1 / \tau_2} \quad (13)$$

The maximum P_{max} can always be reached by a proper choice of the (fully entangled) incident state, so the maximal violation of Bell’s inequality is given by

$$S_{\text{max}} = 2\sqrt{1 + \frac{4\tau_1 / \tau_2}{(1 + \tau_1 / \tau_2)^2}} \quad (14)$$

The dependence of S_{max} on τ_1 / τ_2 is plotted in Fig. 1. Full entanglement is obtained for $\tau_1 = \tau_2$, hence for $T_{1+} T_{2-} = T_{1-} T_{2+}$. Generically, this requires either identical objects ($T_{1\pm} = T_{2\pm}$) or nonidentical objects with $T_{1+} = T_{1-}$. If $\tau_1 = \tau_2$ there are no which-way labels and entanglement fully survives with no degradation.

Small deviations of τ_1 / τ_2 from unity only reduce the entanglement to second order

$$S_{\text{max}} = 2\sqrt{2} \left[1 - \frac{1}{16} (\tau_1 / \tau_2 - 1)^2 + O(\tau_1 / \tau_2 - 1)^3 \right] \quad (15)$$

So for a small reduction of the entanglement one can tolerate a large mismatch of the transmission probabilities. In particular, the experimental result $S=2.71$ for plasmon-assisted entanglement transfer [1] can be reached with more than a factor two of mismatch ($S=2.71$ for $\tau_1/\tau_2=2.4$).

As a simple example we calculate the symmetry parameter τ_1/τ_2 for a Lorentzian transmission probability, appropriate for plasmon-assisted entanglement transfer [5–9]. We take

$$T_{i\pm} = \frac{\mathcal{T}\Gamma^2}{(\omega_0 - \omega_{i\pm})^2 + \Gamma^2}, \quad (16)$$

where ω_0 is the frequency of the incident radiation, Γ is the linewidth, and \mathcal{T} is the transmission probability at the resonance frequency $\omega_{i\pm}$. (For simplicity we take polarization-independent Γ and \mathcal{T} .) The transmission is through an optically thick metal film with a rectangular array of subwavelength holes (lattice constants $L_{i\pm}$). The dispersion relation of the surface plasmons is $\omega_{i\pm} = (1 + 1/\epsilon)^{1/2} 2\pi n c / L_{i\pm}$ [9], where ϵ is the real part of the dielectric constant and n is the order of the resonance, equal to the number of plasmon-field oscillations in a lattice constant. We break the symmetry by taking one square array of holes and one rectangular array (lattice constants $L_0 = L_{1+} = L_{2+} = L_{2-}$ and $L_1 = L_{1-}$). The lattice constant L_0 is chosen such that the incident radiation is at resonance. The symmetry parameter becomes

$$\frac{\tau_1}{\tau_2} = 1 + (2\pi)^2 \left(\frac{nl}{L_0} - \frac{nl}{L_1} \right)^2, \quad l = \frac{c}{\Gamma} \sqrt{\frac{\epsilon+1}{\epsilon}} \quad (17)$$

The length l is the propagation length of the surface plasmon [We have taken $c(1 + 1/\epsilon)^{1/2}$ for the plasmon group velocity, valid if ω_0 is not close to the plasma frequency [9].] Combining Eqs. (15) and (17) we see that the deviation of S_{max} from $2\sqrt{2}$ (the degradation of the entanglement) is proportional to the *fourth* power of the difference between the number of oscillations of the plasmon field along the two lattice vectors.

Turning now to the more general case of a partially entangled incident state, we ask the following question: Is it possible to achieve $P_{\text{out}}=1$ even if $P_{\text{in}}<1$? In other words, can one detect a $2\sqrt{2}$ violation of Bell's inequality after transmission even if the original state was only partially entangled? Examination of Eq. (9) shows that the answer to this question is *Yes*, provided τ_1 and τ_2 satisfy

$$\left| \ln \frac{\tau_1}{\tau_2} \right| \leq 2 \operatorname{arccosh}(P_{\text{in}}^{-1}) \quad \text{and} \quad \ln \tau_1 \tau_2 \geq 2 \operatorname{arccosh}(P_{\text{in}}^{-1}) \quad (18)$$

The allowed values of τ_1 and τ_2 lie in a strip that is open at one end, see Fig. 2. The boundaries are reached at $|u|=|v|=\frac{1}{2}$. The region inside the strip is reached by choosing both $|u|$ and $|v|<1/2$. For $P_{\text{in}}=1$ the strip collapses to the single line $\tau_1=\tau_2$, in agreement with Eq. (13).

The possibility to achieve $P_{\text{out}}=1$ for $P_{\text{in}}<1$ is an ex-

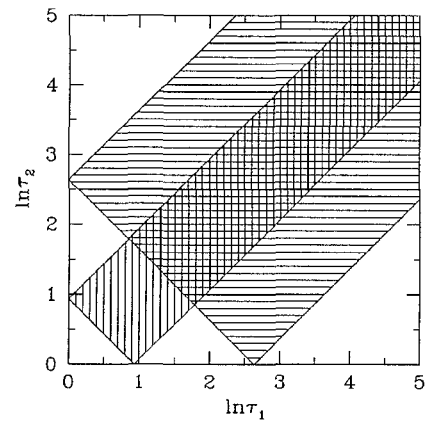


FIG. 2 The shaded strips indicate the values of $\ln \tau_1$ and $\ln \tau_2$ for which $P_{\text{out}}=1$ can be reached with $P_{\text{in}}=0.5$ (horizontally shaded) and $P_{\text{in}}=0.9$ (vertically shaded), in accordance with Eq. (18).

ample of distillation of entanglement [4]. The distillation method used here is the *Procrustean* method of Bennett *et al.* [3]. It requires only local linear filters (the metal films in our case) and classical communication (the coincidence counter). See Ref. [15] for an experimental realization and Refs. [16–20] for other distillation schemes. As it should, no entanglement is created in this operation. Out of N incoming photon-pairs with entanglement P_{in} , one detects NZ pairs with entanglement $P_{\text{out}} = P_{\text{in}} Z^{-1} \sqrt{T_1 + T_1 - T_2 + T_2}$, so that $NZ P_{\text{out}} \leq NP_{\text{in}}$.

In conclusion, we have shown that optical entanglement transfer and distillation through a pair of linear media can be described by two ratios τ_1 and τ_2 of polarization-dependent transmission probabilities. For fully entangled incident radiation, the maximal violation of Bell's inequality at the detectors is given by function (14) of τ_1/τ_2 which decays only slowly around the optimal value $\tau_1/\tau_2=1$. Distillation of a fully entangled Bell pair out of partially entangled incident radiation is possible no matter how low the initial entanglement, provided that τ_1 and τ_2 satisfy the two inequalities (18).

Our results provide a simple way to describe the experiment of Ref. [1] on plasmon-assisted entanglement transfer, in terms of two separately measurable parameters. By changing the square array of holes used in Ref. [1] into a rectangular array (or, equivalently, by tilting the square array relative to the incident beam), one can move away from the point $\tau_1=\tau_2=1$ and search for the entanglement distillation predicted here. The possibility of extracting Bell pairs by manipulating surface plasmons may have interesting applications in quantum information processing.

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